#### Teilchenphysik mit höchstenergetischen Beschleunigern (Higgs & Co)



#### 8. Precision Tests of the Standard Model

12.12.2016

Prof. Dr. Siegfried Bethke Dr. Frank Simon



Max-Planck-Institut für Physik (Werner-Heisenberg-Institut)

### **Overview**

- The Standard Model Structure, Motivation
- Vector boson properties
  - Z decay & width
  - W, Z production
  - W mass
  - W width
  - Triple Gauge couplings
- Topics of future lectures in the framework of the Standard Model:
  - QCD (Lecture 8)
  - Higgs (Lectures 9 & 10)
  - Top quark (Lecture 12)



• The SM describes our visible Universe by a (reasonably small) set of particles:



- The SM describes our visible Universe by a (reasonably small) set of particles:
  - The particles that make up matter: Spin 1/2 Fermions

Elementary Particles					
	Generation				
	1	2	3		
Quarks	u d	C S	t b		
Leptons	v <sub>e</sub> <b>e</b>	ν <sub>μ</sub> μ	ν <sub>τ</sub> τ		



- The SM describes our visible Universe by a (reasonably small) set of particles:
  - The particles that make up matter: Spin 1/2 Fermions
  - ... and the force carriers: Spin 1 Vector bosons

Elementary Particles				Elen	nentary Forces	
	1	Generatio	n   3		exchange boson	relative strength
Quarks	u	С	t	Strong	g	1
	d	S	b	elmagn.	γ	1/137
Leptons	v <sub>e</sub>	V <sub>µ</sub>	ν <sub>τ</sub>	Weak	W±, Z <sup>0</sup>	<b>10</b> -14
	е	μ	τ	Gravitation	G	<b>10</b> <sup>-40</sup>



- The SM describes our visible Universe by a (reasonably small) set of particles:
  - The particles that make up matter: Spin 1/2 Fermions
  - ... and the force carriers: Spin 1 Vector bosons

Elementary Particles				Elen	nentary Forces	
	1	Generatio 2	<b>n</b> 3		exchange boson	relative strength
Quarks	u d	C	t b	Strong	g	1
Lontone	V <sub>e</sub>	V	۲ ۷ т	elmagn. Weak	γ W±, Z <sup>0</sup>	1/137 <b>10</b> -14
Leptons	е	μ	τ	Gravitation	G	<b>10</b> <sup>-40</sup>

... plus the Higgs particle as a consequence of the mechanism to generate mass



- The SM describes our visible Universe by a (reasonably small) set of particles:
  - The particles that make up matter: Spin 1/2 Fermions
  - ... and the force carriers: Spin 1 Vector bosons

Elementary Particles				Elen	nentary Forces	
	1	Generatio	n 3		exchange boson	relative strength
Quarks	u	С	t	Strong	g	1
	a	S	D	elmagn.	γ	1/137
Leptons	v <sub>e</sub>	ν <sub>μ</sub>	ν <sub>τ</sub>	Weak	W±, Z <sup>0</sup>	<b>10</b> <sup>-14</sup>
	е	μ	τ	Gravitation	G	<b>10</b> <sup>-40</sup>

... plus the Higgs particle as a consequence of the mechanism to generate mass

Underlying theories:

QCD

QED / weak interaction

→ electroweak unification (GSW)



### The Success of the Standard Model

 The Standard Model was developed in the 1970s following experimental observations (at that point only three quarks were known, the charm discovery followed shortly thereafter)



### The Success of the Standard Model

- The Standard Model was developed in the 1970s following experimental observations (at that point only three quarks were known, the charm discovery followed shortly thereafter)
- It:
  - describes the unified electroweak interactions and the strong force with gauge invariant quantum field theories
  - is extremely successful in consistently and precisely describing all particle reaction observed to date
  - provides a consistent (yet incomplete) picture of the evolution of the early universe
    -> particle cosmology



### **The Standard Model - Combining Theories**





*Teilchenphysik mit höchstenergetischen Beschleunigern:* WS 16/17, 08: Standard Model

Frank Simon (fsimon@mpp.mpg.de)

### **The Standard Model - Combining Theories**



The Standard Model of Particle Physics



### "Rediscovering" The Standard Model at LHC





#### "Rediscovering" The Standard Model at LHC





 The electroweak part of the SM is based on the gauge group SU(2) x U(1)



- The electroweak part of the SM is based on the gauge group SU(2) x U(1)
- This gives rise to the gauge bosons W<sup>+</sup>, W<sup>-</sup>, Z for SU(2) and  $\gamma$  for U(1)



- The electroweak part of the SM is based on the gauge group SU(2) x U(1)
- This gives rise to the gauge bosons W<sup>+</sup>, W<sup>-</sup>, Z for SU(2) and  $\gamma$  for U(1)
- Left-handed fermion fields transform as doublets under SU(2) right handed fermions as singlets (no coupling of right-handed fermions to W; V-A structure of the weak interaction)



- The electroweak part of the SM is based on the gauge group SU(2) x U(1)
- This gives rise to the gauge bosons W<sup>+</sup>, W<sup>-</sup>, Z for SU(2) and  $\gamma$  for U(1)
- Left-handed fermion fields transform as doublets under SU(2) right handed fermions as singlets (no coupling of right-handed fermions to W; V-A structure of the weak interaction)
- There are three fermion families



- The electroweak part of the SM is based on the gauge group SU(2) x U(1)
- This gives rise to the gauge bosons W<sup>+</sup>, W<sup>-</sup>, Z for SU(2) and  $\gamma$  for U(1)
- Left-handed fermion fields transform as doublets under SU(2) right handed fermions as singlets (no coupling of right-handed fermions to W; V-A structure of the weak interaction)
- There are three fermion families
- A complex scalar Higgs field is added for mass generation through spontaneous symmetry breaking to give mass to the gauge bosons and fermions -> Gives rise to one physical neutral scalar particle, the Higgs boson



- The electroweak part of the SM is based on the gauge group SU(2) x U(1)
- This gives rise to the gauge bosons W<sup>+</sup>, W<sup>-</sup>, Z for SU(2) and  $\gamma$  for U(1)
- Left-handed fermion fields transform as doublets under SU(2) right handed fermions as singlets (no coupling of right-handed fermions to W; V-A structure of the weak interaction)
- There are three fermion families
- A complex scalar Higgs field is added for mass generation through spontaneous symmetry breaking to give mass to the gauge bosons and fermions -> Gives rise to one physical neutral scalar particle, the Higgs boson
- The electroweak SM describes in lowest order ("Born approximation) processes such as  $f_1f_2 \rightarrow f_3f_4$  with only 3 free parameters:  $\alpha$ ,  $G_f$ ,  $sin^2\theta_W$



### **Testing the Standard Model**

- mainly physics with
  - electroweak gauge bosons (W, Z, γ)
  - top quarks (-> lecture 9)
  - with hadron jets (QCD) (-> lecture 7)



### **Testing the Standard Model**

- mainly physics with
  - electroweak gauge bosons (W, Z, γ)
  - top quarks (-> lecture 9)
  - with hadron jets (QCD) (-> lecture 7)
- measurements of
  - production cross sections
  - masses
  - decay rates / widths
  - decay asymmetries
  - gauge bosons couplings (WW, Wγ, WZ, ZZ, Zγ)



### **Motivations for these Tests**

 Since the establishment of the Standard Model, one main goal of particle physics has been (and still is) to test its predictions as a consistency check, and to look for cracks



### **Motivations for these Tests**

- Since the establishment of the Standard Model, one main goal of particle physics has been (and still is) to test its predictions as a consistency check, and to look for cracks
- Search for deviations from the SM:
  - properties, production and decay of gauge bosons are sensitive to the particle content and to various particle properties, and are modified by new physics



 $\Rightarrow$  used to place indirect limits on the Higgs mass based on M<sub>top</sub> and M<sub>W</sub>



### **Motivations for these Tests**

- Since the establishment of the Standard Model, one main goal of particle physics has been (and still is) to test its predictions as a consistency check, and to look for cracks
- Search for deviations from the SM:
  - properties, production and decay of gauge bosons are sensitive to the particle content and to various particle properties, and are modified by new physics



 $\Rightarrow$  used to place indirect limits on the Higgs mass based on M<sub>top</sub> and M<sub>W</sub>

- Use well-understood SM processes to measure luminosity at LHC
- Precisely define SM backgrounds in the search for new physics



### The Z Boson in e<sup>+</sup>e<sup>-</sup> Annihilation

 A short excursion to e<sup>+</sup>e<sup>-</sup> Annihilation (covered in somewhat greater detail in the Summer)







• A key measurement at the Z resonance: The total decay width





• A key measurement at the Z resonance: The total decay width

Given by:  $\Gamma_{Z} = \Gamma_{ee} + \Gamma_{\mu\mu} + \Gamma_{\tau\tau} + \Gamma_{had}$   $+ \Gamma_{veve} + \Gamma_{v\mu\nu\mu} + \Gamma_{v\tau\nu\tau}$   $= 3 \Gamma_{II} + \Gamma_{had} + N_{v} \Gamma_{vv}$ 





• A key measurement at the Z resonance: The total decay width

Given by:  $\Gamma_{Z} = \Gamma_{ee} + \Gamma_{\mu\mu} + \Gamma_{\tau\tau} + \Gamma_{had}$   $+ \Gamma_{veve} + \Gamma_{v\mu\nu\mu} + \Gamma_{v\tau\nu\tau}$   $= 3 \Gamma_{II} + \Gamma_{had} + N_{v} \Gamma_{vv}$ 

The partial width into visible final states can be directly measured





This precision can not be reached at hadron colliders - LEP input used for calibration at LHC



**Teilchenphysik mit höchstenergetischen Beschleunigern:** WS 16/17, 08: Standard Model  A key measurement at the Z resonance: The total decay width

Given by:  $\Gamma_{Z} = \Gamma_{ee} + \Gamma_{\mu\mu} + \Gamma_{\tau\tau} + \Gamma_{had}$   $+ \Gamma_{veve} + \Gamma_{v\mu\nu\mu} + \Gamma_{v\tau\nu\tau}$   $= 3 \Gamma_{II} + \Gamma_{had} + N_{v} \Gamma_{vv}$ 

The partial width into visible final states can be directly measured



This precision can not be reached at hadron colliders - LEP input used for calibration at LHC

• A key measurement at the Z resonance: The total decay width

Given by:  $\Gamma_{Z} = \Gamma_{ee} + \Gamma_{\mu\mu} + \Gamma_{\tau\tau} + \Gamma_{had}$   $+ \Gamma_{veve} + \Gamma_{v\mu\nu\mu} + \Gamma_{v\tau\nu\tau}$   $= 3 \Gamma_{II} + \Gamma_{had} + N_{v} \Gamma_{vv}$ 

The partial width into visible final states can be directly measured

The SM makes a clean prediction for  $\Gamma_{vv}$ - from the measured cross section and total width the number of (light) neutrinos can be determined

$$N_v = 2.984 \, \pm \, 0.008$$



### Production (and Decay) of Gauge Bosons at LHC

• For precision measurements: hadronic final states can not be used due to dominating QCD background





### Production (and Decay) of Gauge Bosons at LHC

• For precision measurements: hadronic final states can not be used due to dominating QCD background



... but also t/u channel processes such as





# Production (and Decay) of Gauge Bosons at LHC

• For precision measurements: hadronic final states can not be used due to dominating QCD background



 theoretical uncertainties mainly due to quark structure of the proton: PDF uncertainties



### Z Production at LHC

• Candidate Z->e+e-





### **Z** Production at LHC

Candidate Z->µ<sup>+</sup>µ<sup>-</sup>





## W Production at LHC

•  $W^- \rightarrow \mu^- \nu$  candidate





# W Production at LHC

• W<sup>+</sup> -> e<sup>+</sup>v candidate





### Z Production at LHC with high Pileup



Z-> µµ
 ... with 20 additional vertices





Arbizit



- Measurement of Cross Sections:
  - Z selection:
    - one lepton with "tight" selection (high energy, isolation, unambiguous ID)
    - second lepton with more relaxed criteria





- Measurement of Cross Sections:
  - Z selection:
    - one lepton with "tight" selection (high energy, isolation, unambiguous ID)
    - second lepton with more relaxed criteria
    - W selection:
      - one lepton with "tight" selection
      - missing transverse energy / momentum





- Measurement of Cross Sections:
  - Z selection:
    - one lepton with "tight" selection (high energy, isolation, unambiguous ID)
    - second lepton with more relaxed criteria
  - W selection:
    - one lepton with "tight" selection
    - missing transverse energy / momentum
  - Determination of cross section corrections to event numbers:
    - trigger efficiency (data)
    - reconstruction efficiency (MC, data)
    - luminosity

$$\sigma_{Z} = \frac{N}{\int Ldt \cdot Br(Z^{0} \rightarrow e^{+}e^{-}) \cdot \varepsilon_{ee}}$$











*Teilchenphysik mit höchstenergetischen Beschleunigern:* WS 16/17, 08: Standard Model

Frank Simon (fsimon@mpp.mpg.de)

20





*Teilchenphysik mit höchstenergetischen Beschleunigern:* WS 16/17, 08: Standard Model

Frank Simon (fsimon@mpp.mpg.de)

20





*Teilchenphysik mit höchstenergetischen Beschleunigern:* WS 16/17, 08: Standard Model

Frank Simon (fsimon@mpp.mpg.de)

20



• "Best results" typically in the Muon channel



#### W and Z Production at the Tevatron





### W and Z Measurements at the LHC



- Measured cross sections corrected for efficiency and acceptance
- Higher cross section for W<sup>+</sup> than for W<sup>-</sup>: Due to valence quark content of protons: uud - higher probability to make a W<sup>+</sup>



### W and Z Production at the LHC



• Combined with Tevatron results to illustrate evolution with energy



#### Measuring the Mass of the W Boson

• Measurement of the mass from the transverse momentum distribution of the lepton and of the neutrino (inferred from lepton and hadronic system)



$$\vec{P}_T^\nu = -(\vec{P}_T^l + \vec{U})$$

• Reconstruct transverse mass:

$$M_T = \sqrt{(E_T^l + E_T^{\nu})^2 - (\vec{P}_T^l + \vec{P}_T^{\nu})^2}$$



### Measuring the Mass of the W Boson

• Measurement of the mass from the transverse momentum distribution of the lepton and of the neutrino (inferred from lepton and hadronic system)



$$\vec{P}_T^{\nu} = -(\vec{P}_T^l + \vec{U})$$

• Reconstruct transverse mass:

$$M_T = \sqrt{(E_T^l + E_T^{\nu})^2 - (\vec{P}_T^l + \vec{P}_T^{\nu})^2}$$

- Compare measured M<sub>T</sub> distribution to simulated distributions with different W mass assumptions ("template fit")
- Requires excellent understanding of momentum and energy scale and resolution



### Measuring the Mass of the W Boson



- Best measurement from Tevatron
- Combination of CDF and D0:  $M_W = 80.387 \pm 0.016 \; \text{GeV}$
- World average with LEP:  $M_W = 80.385 \pm 0.015$  GeV



### The Impact of the W Mass Measurement



 W mass measurement (together with top mass) provides indirect constraints on Higgs mass in the Standard Model









First LHC W mass measurement from ATLAS to be announced tomorrow!





*Teilchenphysik mit höchstenergetischen Beschleunigern:* WS 16/17, 08: Standard Model

Frank Simon (fsimon@mpp.mpg.de)

### Measuring the Width of the W Boson

- The tail of the  $M_T$  distribution is sensitive to the total width of the W boson:
  - Events with M<sub>T</sub> > M<sub>W</sub> are due to detector resolution effects and due to the finite width - the resolution contribution to this falls faster than the width contribution, allowing an accurate measurement of the width





### The Width of the W - Summary of Results



 Excellent agreement with the Standard Model



#### **Triple Gauge Couplings**



• In the SM: Space-like diagrams are = 0 if two of the three bosons are identical



#### **Triple Gauge Couplings**



- In the SM: Space-like diagrams are = 0 if two of the three bosons are identical
- BSM: May contribute to triple Gauge couplings in non-standard ways

$$\sum_{i=1}^{n} \sum_{i=1}^{n} \sum_{i$$



### **Triple Gauge Couplings**



- In the SM: Space-like diagrams are = 0 if two of the three bosons are identical
- BSM: May contribute to triple Gauge couplings in non-standard ways

$$\sum_{i=1}^{n} \sum_{i=1}^{n} \sum_{i$$

- SM: Time-like diagrams with two identical bosons in the final state are allowed
  - NB No triple gauge coupling! SM background to TGC measurements





### **Measurement of WW Production**



- Looking for W->lv: Best separation from background
- Cleanest signal: events w/o jets - one additional jet also considered

#### **Measurement of ZZ Production**





### **Measurement of ZZ Production**



So far all observations consistent with SM expectations



### **Double Vector Boson Production - CMS Summary**



• Overall excellent agreement with SM expectations - Consistent for 7, 8 and 13TeV



### Summary

- The (electroweak) Standard Model combines QED and the weak interaction theory to describe electromagnetic and weak interactions - based on the Gauge Group SU(2) x U(1)
- It has been extremely successful in describing all observations to date
- Its predictions are tested by measurements of
  - masses
  - cross-sections (and production asymmetries not covered)
  - decay widths
  - triple gauge couplings particularly sensitive to New Physics
- The Tevatron provides the most precise W mass measurement to date global uncertainty 15 MeV – LHC might ultimately go to 5 MeV
  - requires very precise understanding of detectors and excellent control of all systematics - first ATLAS results expected tomorrow!



## Summary

- The (electroweak) Standard Model combines QED and the weak interaction theory to describe electromagnetic and weak interactions - based on the Gauge Group SU(2) x U(1)
- It has been extremely successful in describing all observations to date
- Its predictions are tested by measurements of
  - masses
  - cross-sections (and production asymmetries not covered)
  - decay widths
  - triple gauge couplings particularly sensitive to New Physics
- The Tevatron provides the most precise W mass measurement to date global uncertainty 15 MeV - LHC might ultimately go to 5 MeV
  - requires very precise understanding of detectors and excellent control of all systematics - first ATLAS results expected tomorrow!

#### Next Lecture: Top Physics, F. Simon 19.12.2016



### Schedule

1.	Introduction	17.10.
2.	Accelerators	24.10.
	no lecture	31.10.
3.	Particle Detectors I	07.11.
4.	Particle Detectors II	14.11.
5.	Trigger, Data Acquisition, Computing	21.11.
6.	Monte Carlo Generators and Detector Simulation	28.11.
7.	QCD, Jets, Proton Structure	05.12.
8.	Tests of the Standard Model	12.12
9.	Top Physics	19.12.
	Christmas	
10.	Higgs Physics I	09.01.
11.	Higgs Physics II	16.01.
12.	Physics beyond the SM	23.01.
13.	LHC Outlook & Future Collider Projects	30.01
	no lecture	06.02

