# Non-geometric spacetimes in string theory 

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- coarse grained space
- spin foams (loop quantum gravity)
- string theory?


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## overview

(1) motivation
(2) non-commutative and non-associative spacetimes
(3) strings in (bosonic) backgrounds
(4) from geometric to non-geometric fluxes

## uncertainty relations of spacetime

- non-commutative spacetime $\rightarrow$ physical minimal areas

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\left[\hat{x}^{\mu}, \hat{x}^{v}\right]=i \theta^{\mu v} \quad \Rightarrow \quad \Delta x^{\mu} \cdot \Delta x^{v} \geq \theta^{\mu v}
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(similar) example: charged point particle in a magnetic field $\vec{B}$

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- non-associative spacetime: $\rightarrow$ physical minimal volumes
- Jacobi identity (infinitesimal version of associativity - blackboard):

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\operatorname{Jac}(\hat{A}, \hat{B}, \hat{C})=[\hat{A},[\hat{B}, \hat{C}]]+[\hat{B},[\hat{C}, \hat{A}]]+[\hat{C},[\hat{A}, \hat{B}]]=0
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- violation of Jacobi identity $\rightarrow$ higher uncertainty relation, e.g.

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\operatorname{Jac}\left(\hat{x}^{1}, \hat{x}^{2}, \hat{x}^{3}\right)=i R^{123} \neq 0 \quad \rightarrow \Delta x^{1} \cdot \Delta x^{2} \cdot \Delta x^{3} \geq R^{123}
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How do these emerge in string theory?

## strings in (bosonic) backgrounds

- single string picture:
- string moving in/coupling to a background spacetime generalisation of point particle moving in curved space
- different background fields, among others: metric $G_{\mu \nu}$, Kalb-Ramond field $B_{\mu \nu}$ : (higher) gauge field


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- different background fields, among others: metric $G_{\mu v}$, Kalb-Ramond field $B_{\mu v}$ : (higher) gauge field
- modifications in relation string $\leftrightarrow$ (target) spacetime
- compact extra dimensions: winding modes of string
- finite string length
- not necessarily ordinary manifolds:
interplay between metric $G$ and other fields
- dualities, e.g. $T$-duality $-R \leftrightarrow \frac{1}{R}$


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- open strings with $B$-field [Seiberg, Witten 99]

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- $\rightarrow$ 'non-commutative Yang-Mills theory':
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# example 2: from geometric to non-geometric fluxes I 

[Shelton, Taylor, Wecht 05]
chain of $T$-dualities:

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\mathbf{H}_{a b c} \quad \xrightarrow{T_{1}} \quad \mathbf{f}_{a b}^{c} \quad \xrightarrow{T_{2}} \quad \mathbf{Q}_{c}^{a b} \xrightarrow{T_{3}} \quad \mathbf{R}^{a b c}
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- start with $\mathbf{H}$-flux on flat 3-torus $T^{3}: \mathbf{H}_{123}=\partial_{[1} B_{23]}=h \mathbb{R}^{3}$ with $x^{i} \sim x^{i}+1+$ gauge field $\rightarrow$ geometric


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- $T$-duality along $x^{1}$ :
- 'twisted torus': $\left(x^{1}, x^{2}, x^{3}\right) \sim\left(x^{1}-h x^{2}, x^{2}, x^{3}+1\right)$
- similar to a Lie group: vector fields $e_{i}^{a}$

$$
\partial_{[i} e_{j]}^{c}+\mathbf{f}^{c}{ }_{a b} e_{i}^{a} e_{j}^{b}=0 \rightarrow \text { geometric } \mathbf{f} \text {-flux, here } \mathbf{f}^{3}{ }_{12}=h
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$\rightarrow$ 'globally non-geometric'
- for closed strings in $\mathbf{Q}$-flux: $\left[\hat{x}^{a}, \hat{x}^{b}\right] \sim \mathbf{Q}_{c}{ }^{a b} w^{c}$
- (formal) T-duality along $x^{3}: \theta^{12}=h \tilde{x}^{3}$
- $\tilde{x}^{3}$ 'winding' coordinate
- no description in terms of standard coordinates
$\rightarrow \quad$ 'locally non-geometric'

$$
\mathbf{R}^{a b c}=\tilde{\partial}\left[c \theta^{a b]} \longrightarrow \text { non-geometric } \mathbf{R} \text {-flux, here } \mathbf{R}^{123}=h\right.
$$

- for closed strings in R-flux: $\operatorname{Jac}\left(\hat{x}^{a}, \hat{x}^{b}, \hat{x}^{c}\right) \sim \mathbf{R}^{a b c}$


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- closed strings: non-geometric Q- and R-flux describe non-commutative resp. non-associative spacetimes.
- string dualities: different backgrounds, same physics new view on spacetime and its mathematical description

