On couplings to matter in bimetric theory

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by

Marvin Lüben

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Why massive graviton?

We have theories for massless and massive particles of different spin s

	massless	massive
spin-0	Klein-Gordon	Klein-Gordon
spin-1	Maxwell	Proca
spin-2	GR	?

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GR is the unique theory (unitary and Lorentz-invariant) for a massless spin-2 field (in 4 dim) [Lovelock '71]

GR is a fully non-linear theory of gravity

	massless	massive
linear	lin. gravity	Fierz-Pauli 1939
non-linear	Einstein 1915	?

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 Coupling FP to matter problematic already for solar system tests (vDVZ discontinuity) and requires non-linear completion (Vainshtein mechanism) [van Dam&Veltman '70; Zakharov '70; Vainshtein '72]

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non-linear	Einstein 1915	Massive (bi)gravity

 Coupling FP to matter problematic already for solar system tests (vDVZ discontinuity) and requires non-linear completion (Vainshtein mechanism) [van Dam&Veltman '70; Zakharov '70; Vainshtein '72]

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 - d \neq 4 dimensions
 - higher-order derivatives of the metric tensor
 - non-locality
 - Add new DoF's: scalars, vectors, tensors

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Introduction to bigravity

Linear theory & Boulwere-Deser ghost

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Linear theory & Boulwere-Deser ghost

Consider small perturbations about Minkowksi: g_{μν} = η_{μν} + h_{μν}
 Lagrangian of lin. gravity + Fierz-Pauli mass term

$$\mathcal{L} = -rac{1}{4} h_{lphaeta} \mathcal{E}_{\mu
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Propagates a ghost with mass $m^2 \sim a^{-1}$. Fierz-Pauli tuning: a = 0. [Fierz&Pauli '39] Consider small perturbations about Minkowksi: $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$ Lagrangian of lin. gravity + Fierz-Pauli mass term

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- Propagates a ghost with mass $m^2 \sim a^{-1}$. Fierz-Pauli tuning: a = 0. [Fierz&Pauli '39]
- BUT: tuning does not persist non-linearly and the ghost will reappear at higher order! [Boulwere&Deser '72]

Bigravity

Non-linear action for massless & massive graviton [deRham, Gabadadze, Tolley

$$\mathsf{S} = -\frac{m_g^2}{2}\int \mathsf{d}^4 x \Big(\sqrt{g}\mathsf{R}(g)$$

^{&#}x27;10; Hassan&Rosen '11]

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$$S = -\frac{m_g^2}{2} \int d^4x \Big(\sqrt{g}R(g) - 2m^2\sqrt{g}V(g,f) + \alpha^2\sqrt{f}R(f) \Big)$$

- Symmetric under $g_{\mu\nu} \leftrightarrow f_{\mu\nu}$
- Potential is symmetric and involves square-root matrix

$$V(g, f) \supset \sqrt{g^{-1}f}, \ \sqrt{g}V(g, f) = \sqrt{f}V(f, g)$$

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Symmetric under
$$g_{\mu\nu} \leftrightarrow f_{\mu\nu}$$

Potential is symmetric and involves square-root matrix
 $V(a, f) \supset \sqrt{a^{-1}f}, \sqrt{a}V(a, f) = \sqrt{f}V(f, a)$

To which metric should matter couple? Physical metric?

Matter couplings

Matter couplings

Singly and doubly coupled

$$S_{\mathsf{m}} = \int \mathsf{d}^4 x \sqrt{g} \mathcal{L}_{\mathsf{m}}(g;\partial\phi,\phi)$$

Simplest choice: matter couples to one metric (singly-coupled BG)

- (classically) ghost-free [Hassan&Rosen '11]
- matter loops do not detune the potential, but generate coupling between *f* and φ. BD ghost reappears above strong coupling scale

[deRham et.al. '14, Heisenberg '15]

Nice phenomenology, but gradient instabilities

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- Nice phenomenology, but gradient instabilities
- P Trivial extension: two independent matter sectors ϕ and χ
- Same matter couples to both metrics: $\phi = \chi$
 - BD ghost at low energies [Yamashita et.al. '14; deRham et.al. '14]
 - A matter field must not have two kinetic terms, but even then matter loops bring back the BD ghost at unacceptable low scales

Effective composite metric

$$S_{\rm m} = \int {\rm d}^4 x \sqrt{h} \mathcal{L}_{\rm m}(h;\partial\phi,\phi)$$

Matter couples to eff. metric, composed out of g and f [deRham et.al. '14; Heisenberg '14 & '15]

$$h_{\mu\nu} = a^2 g_{\mu\nu} + 2ab \, g_{\mu\lambda} (\sqrt{g^{-1}f})^{\lambda}_{\ \nu} + b^2 f_{\mu\nu}$$

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- For computations, switch to trimetric description



Summary



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- More than 1 dynamical metric: Physical metric?
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- What's next?
 - Other ghost-free symmetric matter couplings?
 - Study phenomenology, in particular in high energy environments (early universe, BH)