

# High precision predictions for a new suite of Lund-based observables

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# The Lund-plane: central tool for pQCD









## **Definition of Lund-based observables**

#### INPUT

#### Anti-k<sub>t</sub> jet

#### Hemisphere in e<sup>+</sup>e<sup>-</sup>





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#### INPUT

Anti-k<sub>t</sub> jet

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#### Lund-based observables: resuming tion vs=data 140 fb<sup>-1</sup>

#### **Primary Lund-plane density CMS** [CMS Collab arXiv:2312.16343] 138 fb<sup>-1</sup> (13 TeV) 0.5 AK8 jets p<sub>\_</sub><sup>jet</sup> > 700 GeV, |y<sub>jet</sub>| < 1.7 0.45 $1.584 < \ln(k_/GeV) < 2.084$ .... 0.4 4.87 < k<sub>τ</sub> < 8.03 GeV Data 0.35 0.3 (Lifson, Salam, Soyez) 0.25 0.2 0.15 0.1 Pred./Data 1.4 1.2 0.8<sup>E</sup> 0.5 1.5

State-of-the-art calculations describe data within 10-20% precision































### Lund-based observables: resummations vs parton showers



New application of resummation calculations: test of parton showers

[Ferrario Ravasio et al. PRL 131 (2023) 16, 161906]





### Lund-based observables: resummations vs parton showers

#### See Alexander's talk next



[PanScales Collaboration, in preparation]







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Note: definition for colour singlet events in pp given in [van Beekveld, ASO et al. JHEP 11 (2022) 020]











C/A reclustered jet

Note: definition for colour singlet events in pp given in [van Beekveld, ASO et al. JHEP 11 (2022) 020]

$$S \beta$$

$$V(\{\tilde{p}\}, k) = d_{\ell} \left(\frac{k_t^{(\ell)}}{Q}\right)^{a_{\ell}} e^{-b_{\ell}\eta^{(\ell)}} g_{\ell}(\phi)$$

OUTPUT

#### **Reminiscent of soft-and**collinear behaviour of event shapes

[Banfi, Salam and Zanderighi JHEP 03 (2005) 073] [See also Berger et al. PRD 68 (2003) 014012 ]









Note: definition for colour singlet events in pp given in [van Beekveld, ASO et al. JHEP 11 (2022) 020]

This talk: two new Lund-based observables [Dasgupta et al. PRL 125 (2020) 5, 052002]







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#### OUTPUT



### Plan for this talk



Work in progress with: Melissa van Beekveld (NIKHEF), Luca Buonocore (CERN), Basem El-Menoufi (Monash U.), Silvia Ferrario Ravasio, Pier Francesco Monni (CERN) and Gregory Soyez (IPhT)

#### Sneak peek at pheno with PanScales-Pythia interface

[PanScales Collab. arXiv:2312.13275]



[Banfi, Salam and Zanderighi JHEP 03 (2005) 073]

The cumulative cross section for an (rIRC safe) observable v reads

$$\Sigma(v) = \exp\left(-\int [dk] |M(k)|^2\right) \times \sum_{m=1}^{\infty} \frac{1}{2}$$
virtual corrections

# $\sum_{m=1}^{\infty} \frac{1}{m!} \prod_{i=1}^{m} \int [dk_i] |M(k_i)|^2 \Theta(v - v(k_1, \dots, k_m))$

real emissions





[Banfi, Salam and Zanderighi JHEP 03 (2005) 073]

$$\Sigma(v) = \exp\left(-\int [dk] |M(k)|^2\right) \times \sum_{m=1}^{\infty} \frac{1}{2}$$
virtual corrections

$$\sum_{m=1}^{\infty} \frac{1}{m!} \prod_{i=1}^{m} \int [dk_i] |M(k_i)|^2 = \sum_{n=1}^{\infty} \frac{1}{n!} \int [dk_i] |M(k_i)|^2 = \sum_{n=1}^{\infty}$$

- The cumulative cross section for an (rIRC safe) observable v reads
  - $\sum_{n=1}^{\infty} \frac{1}{m!} \prod_{i=1}^{m} \int [dk_i] |M(k_i)|^2 \Theta(v v(k_1, ..., k_m))$ real emissions
- Introduce a slicing parameter satisfying  $\epsilon \ll 1, \ln 1/\epsilon \ll \ln 1/\nu$ 
  - $\frac{1}{n!} \prod_{i=1}^{n} \int_{\varepsilon v} [dk_i] |M(k_i)|^2 \text{ resolved}$  $\infty$  1  $\frac{n+k}{k}$   $\epsilon v$  $[dk_i]|M(k_i)|^2$  unresolved i=n+1 \*





[Banfi, Salam and Zanderighi JHEP 03 (2005) 073]

rIRC safety implies that unresolved emissions don't contribute to v

$$\Sigma(v) = \exp\left(-\int [dk] |M(k)|^2 (1-e^{-1}) \left[\frac{dk}{m!}\right] + \sum_{m=1}^{\infty} \frac{1}{m!} \prod_{i=1}^{m} \int_{e^v} [dk_i] |M(k_i)|$$

At NLL, we can further expand the no-emission probability

 $\exp\Big(-\int [dk]|M(k)|^2(1-\Theta(\epsilon v -$ 

 $\Theta(\epsilon v - v(k))$  virtual & unresolved

 $|^2 \Theta(v - v(k_1, \dots, k_m))$  real emissions

$$-v(k)$$
) = exp $\left(-R(v) - R' \ln 1/\epsilon\right)$ 



[Banfi, Salam and Zanderighi JHEP 03 (2005) 073]

So that the cumulative cross-section can be written as

$$\Sigma(v) = \exp\left(-\int [dk] |M(k)|^2 [(1 - e^{-R \ln 1/\epsilon})] \sum_{m=1}^{\infty} \frac{1}{m!} \prod_{i=1}^{m} \int_{\epsilon v} [dk_i] \right)$$

$$\int [dk] |M(k)|^2 \sim \int \frac{dk_t}{k_t} \alpha_s^{\text{CMW}}(k_t) |M(k_t)|^2 \sim \alpha_s(k_t) \frac{dk_t}{k_t} d\eta$$

- $\Theta(v v(k))] \end{pmatrix} \times$  Sudakov radiator
- $||M(k_i)|^2 \Theta(v v(k_i))$  Transfer function
- Different log counting for Sudakov and Transfer function ( $\epsilon v < v(k_i) < v$ )
  - $(k_t)dzP(z)$  CMW + hard-coll. splitting

Ensemble of independent softcollinear gluons



## **NLL results for Lund observables**

[Banfi, Salam and Zanderighi JHEP 03 (2005) 073]

So that the cumulative cross-section can be written as

$$\Sigma(v) = e^{-R_{sc}(v) - \frac{1}{2}}$$
Analytic

The NLL results for the Lund observables are: [Dasgupta et al. PRL 125 (2020) 5, 052002]

$$\Sigma(M_{\beta}) = e^{-1}$$
$$\Sigma(S_{\beta}) = e^{-1}$$

# $-R_{hc}(v) \times \mathcal{F}_{NII}$ Numerical (in general) VTIC

 $-R^{\mathrm{NLL}}(M_{\beta})$   $-R^{\mathrm{NLL}} \qquad e^{-\gamma_{E}R'}$ 



# NNLL resummation (ARES approach) in a nutshell

[Banfi, McAslan, Monni, Zanderighi JHEP 05 (2015) 102] [Banfi, El-Menoufi and Monni JHEP 01 (2019) 083]

The interplay between real and virtual emissions notably more delicate

$$\Sigma_{\text{NNLL}}(v) = e^{-R_{sc}(v) - R_{hc}(v)} \\ \times \left[ \mathscr{F}_{\text{NLL}} \left( 1 + \frac{\alpha_s(Q)}{2\pi} H^1 + \frac{\alpha_s(Q_{hc})}{\pi} C^1 \right)_+ \frac{\alpha_s(Q)}{\pi} \delta \mathscr{F}_{\text{NNLL}} \right]$$

 $n \equiv \hat{}$ 

•  $e^{-R_{hc}(v)}, C_1$ : end-point of DGLAP splitting function

•  $H_1$ : finite part of one-loop virtual corrections

where, schematically, the physical origin of each NNLL correction is •  $e^{-R_{sc}(v)}$ :  $\alpha_s^{\text{phys}} = \alpha_s(1 + \sum_{n=1}^{2} \left(\frac{\alpha_s}{2\pi}\right)^n K^n)$  with  $K^1 = K^{\text{CMW}}$ 



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where, schematically, the physical origin of each NNLL correction is



•  $\delta \mathcal{F}_{NNLL}$  : one single emission with NLL-like kinematics

wide-

soft, commensurate angle to other emission  $\delta \mathcal{F}_{correl}, \delta \mathcal{F}_{clust}$ 



### Some remarks for Lund observables at NNLL

We recycled a few ingredients from previous works

• F<sub>NLL</sub> analytic and computed in [Dasgupta et al. PRL 125 (2020) 5, 052002]

and computed the remaining NNLL corrections





# • $e^{-R_{sc}(v)-R_{hc}(v)}$ $\left(1+\frac{\alpha_s(Q)}{2\pi}H^1+\frac{\alpha_s(Q_{hc})}{\pi}C^1\right)$ analytic and computed in [Banfi, El-Menoufi and Monni JHEP 01 (2019) 083]



#### **Cross-checks against Event2 and PanScales showers**

#### **Fixed-order check**



#### **All-orders check**

#### Full agreement between analytic calculation and numerics



## Plan for this talk







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### Sneak peek at pheno with PanScales-Pythia interface

[PanScales Collab. arXiv:2312.13275]



### Sensitivity to non-perturbative corrections @LEP energies



Significantly smaller hadronisation corrections for  $M_{\beta}$ 



### Sensitivity to non-perturbative corrections @FCs energies



Perturbative regime clearly extended when going to higher energies



## Comparison between showers @LEP energies



Physical value of  $\alpha_s(M_Z)$  for PG<sub>0</sub>. Missing shower uncertainties



### Comparison between showers @FCs energies



Similar results for other  $\beta$  values. Plots with uncertainties coming soon





### Summary and outlook

Presented NNLL resummation for two new Lund observables

$$S_{\beta} = \sum_{i \in \text{declust}} \frac{k_{t,i}}{Q} e^{-\beta|\eta_i|} \qquad M_{\beta} = \max_{i \in \text{declust}} \frac{k_{t,i}}{Q} e^{-\beta|\eta_i|}$$

- non-perturbative corrections

•  $M_{\beta}$  has a particularly simple resummation structure and small sensitivity to

 Future directions: extension to hadron collisions (both globally and inside a jet), matching of the resummed predictions, systematic pheno study @LHC

