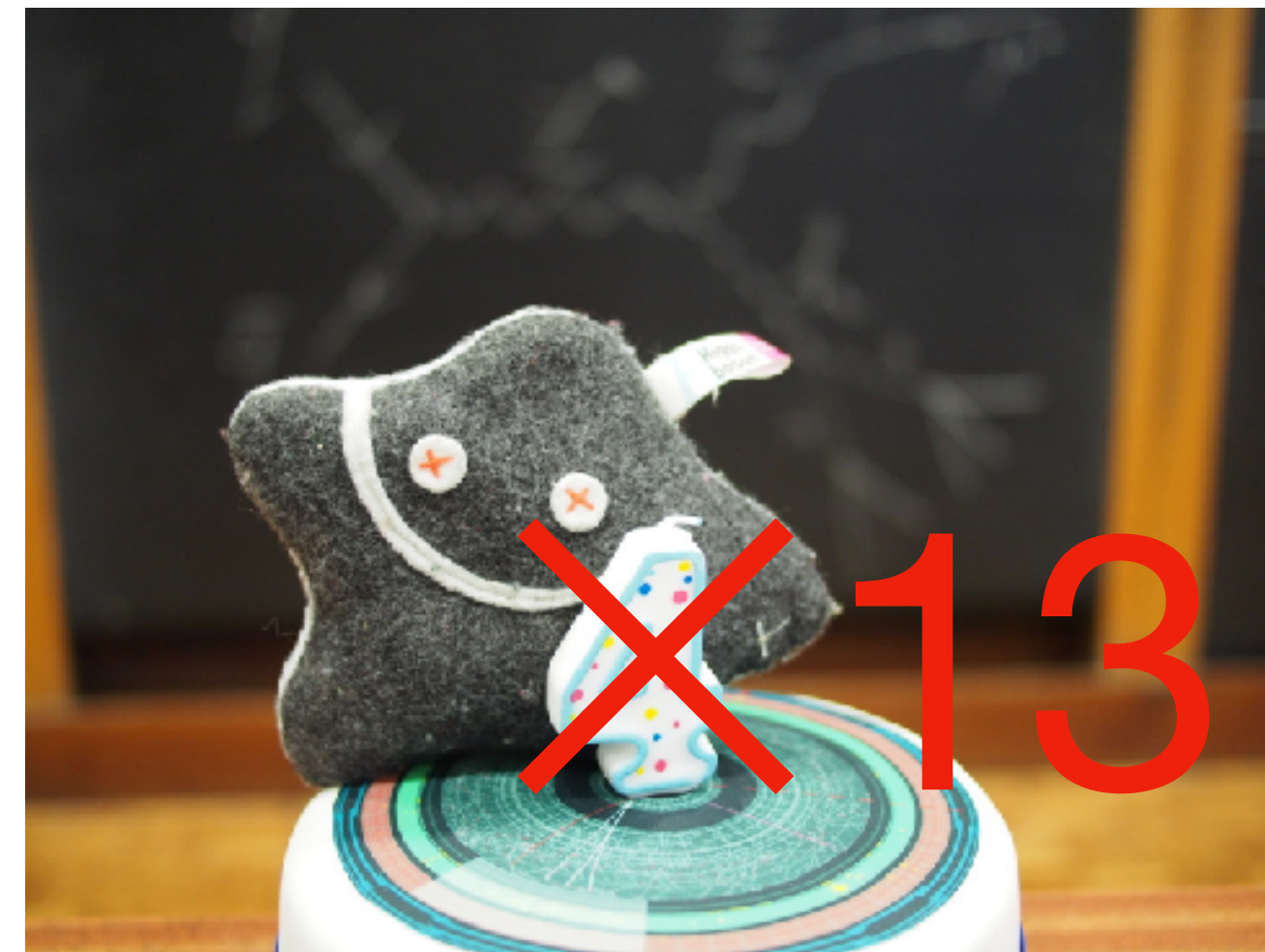


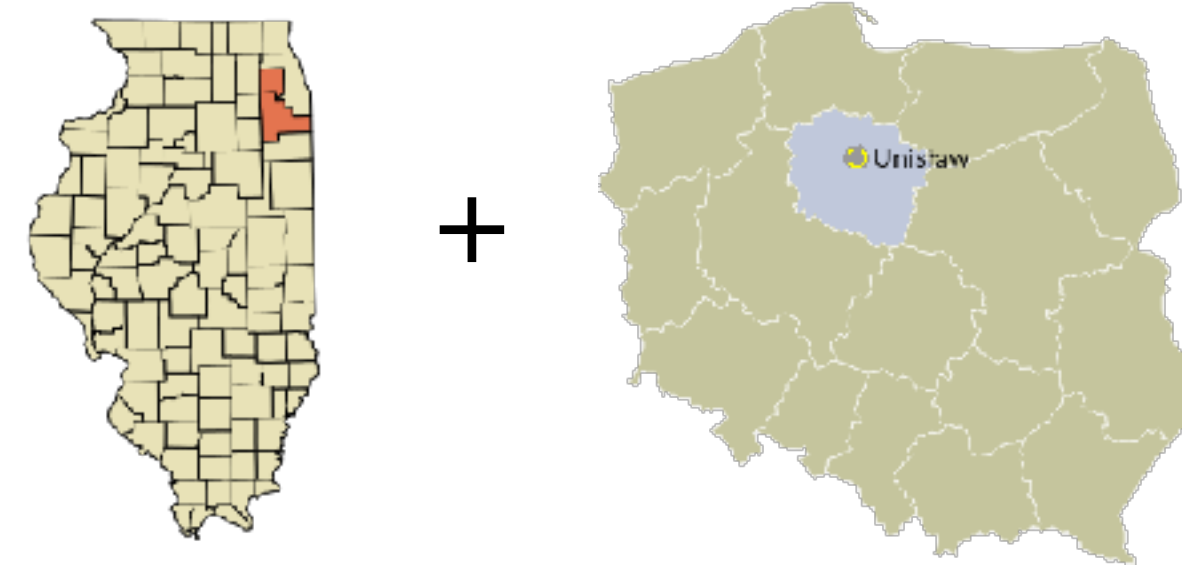
Phenomenology of boosted b-jets in di-Higgs processes at the LHC

Anna Przybyl



A little bit about me...

Anna Przybyl



- From Naperville, IL
- Studied at UIUC
 - Researched and also worked
- Currently finishing my master's in Uppsala, Sweden
- I enjoy ceramics and bouldering



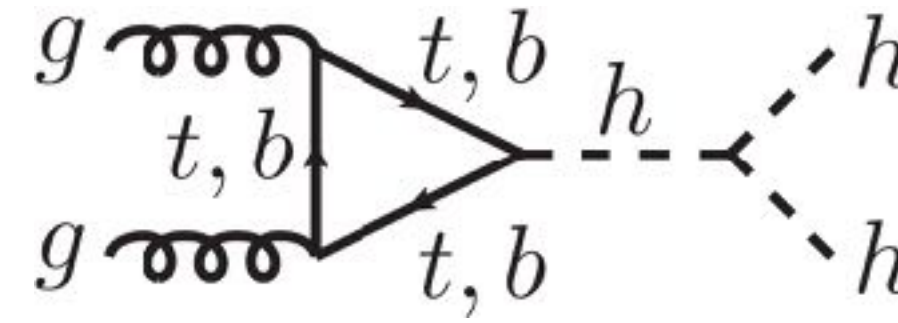
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URBANA-CHAMPAIGN



UPPSALA
UNIVERSITET

Di-Higgs in SUSY

MSSM + NMSSM



Physical Higgses:

h, H, A, H^\pm

$h, H_1, H_2, A_1, A_2, H^\pm$

- SUSY is symmetry between fermions and bosons
 - Higgs sector must be extended, various scenarios with $m_h = 125$ GeV
- In di-Higgs:
 - Colored squarks enter loops
 - Extra Higgses (i.e. possible resonances)
- MSSM extends SM via additional Higgs doublet, NMSSM extends MSSM via additional Higgs singlet

The goal of the thesis

Di-Higgs in SUSY models

$$gg \rightarrow hh \rightarrow b\bar{b}b\bar{b}$$

most probable!

Using benchmark points within simplified models from earlier papers [1,2]:

1. Optimize jet reconstruction

- Slimjet vs. fatjet reconstruction

- anti- k_T (R) vs. variable R (minR, maxR, ρ)

$$R_{\text{eff}} = \frac{\rho}{p_T}$$

2. Come up with kinematic cut to distinguish signal (NP diagrams + NP diagram interferences) from background (SM di-Higgs and QCD background)

[1] Stefano Moretti et al. “Deconstructing squark contributions to di-Higgs production at the LHC”. In: Physical Review D 107.11 (June 2023). issn: 2470-0029. doi: 10.1103/physrevd.107.115010. url: <http://dx.doi.org/10.1103/PhysRevD.107.115010>.

[2] Stefano Moretti et al. Deconstructing resonant Higgs pair production at the LHC: effects of coloured and neutral scalars in the NMSSM test case. 2025. arXiv: 2506.09006 [hep-ph]. url: <https://arxiv.org/abs/2506.09006>.

The benchmark points

Di-Higgs in SUSY models

$$gg \rightarrow hh \rightarrow b\bar{b}b\bar{b}$$

most probable!

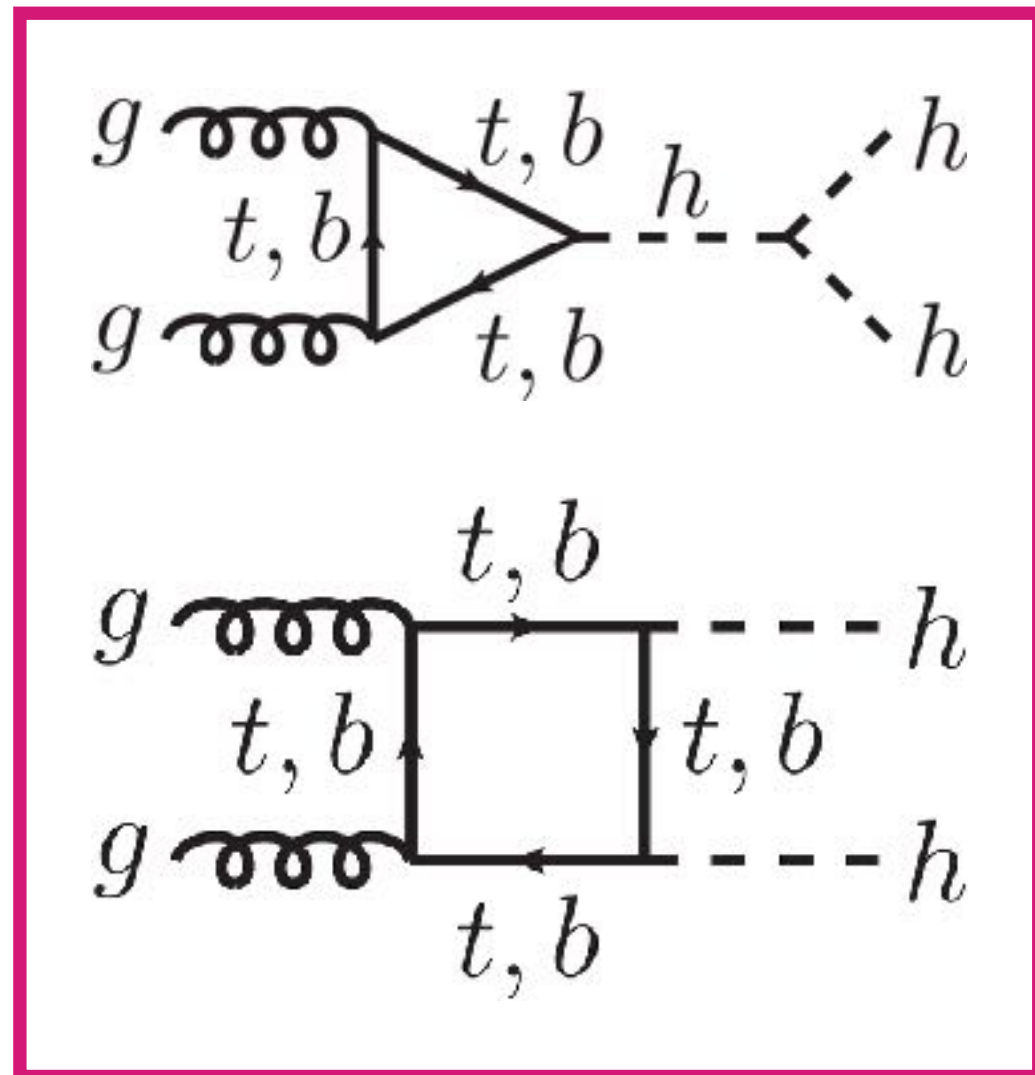
Three BPs:

- MSSM non-resonant case
- NMSSM resonant cases (350 GeV resonance and 800 GeV resonance)

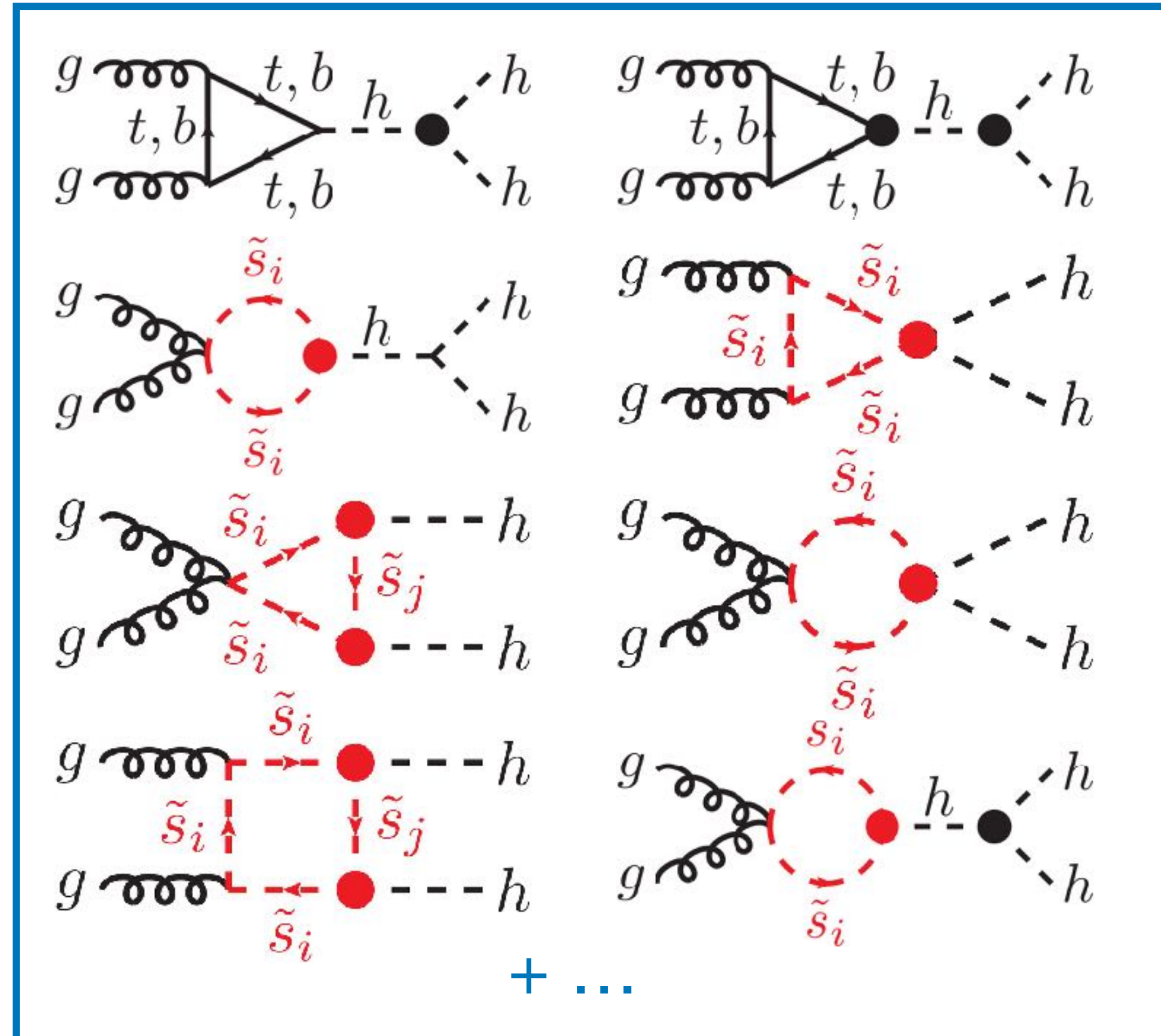
Simplified SUSY models written (limited number of squarks) → appropriate parameter spaces scanned → BPs chosen

Feynman diagrams

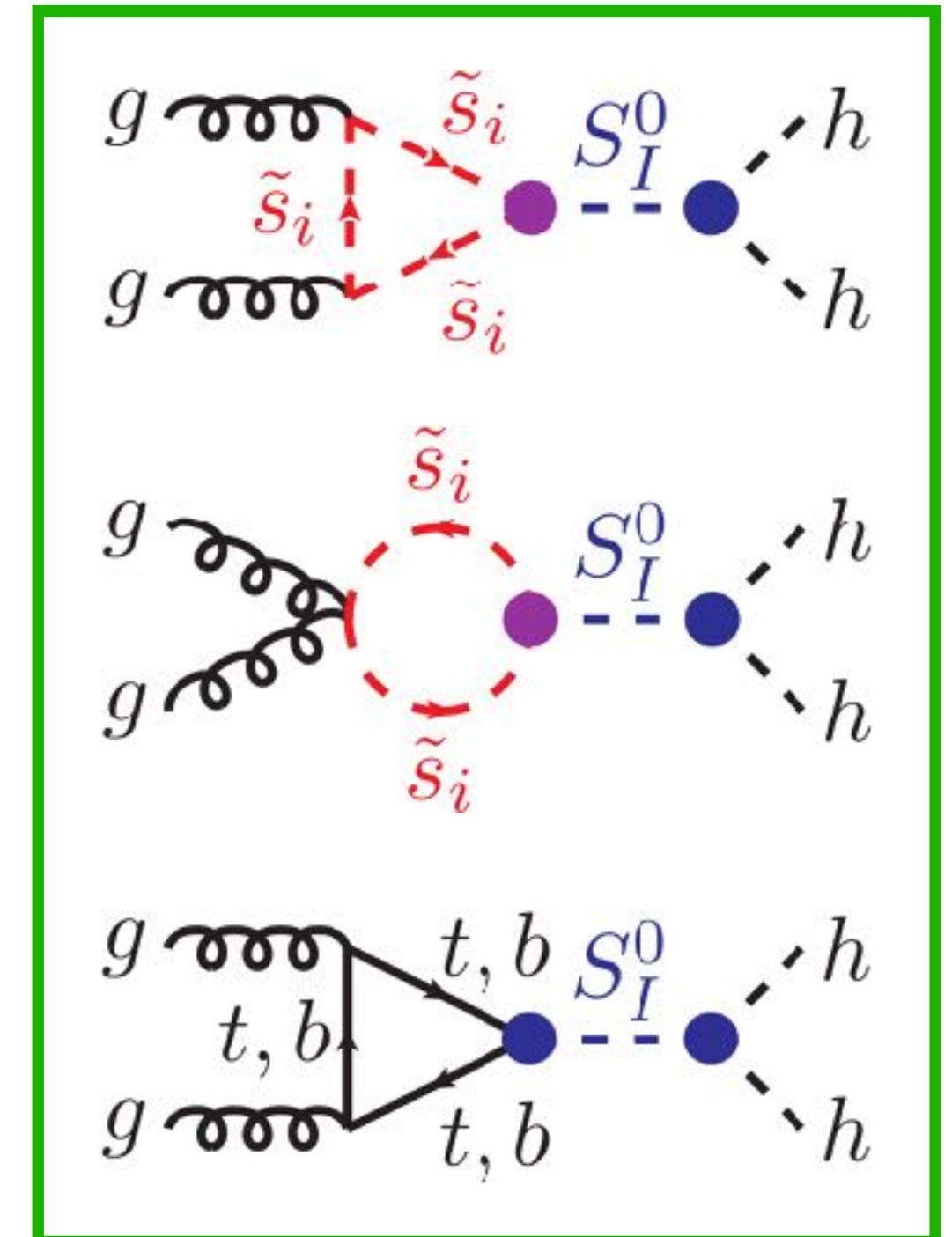
Di-Higgs processes in **SM** + **MSSM** + **NMSSM**



+

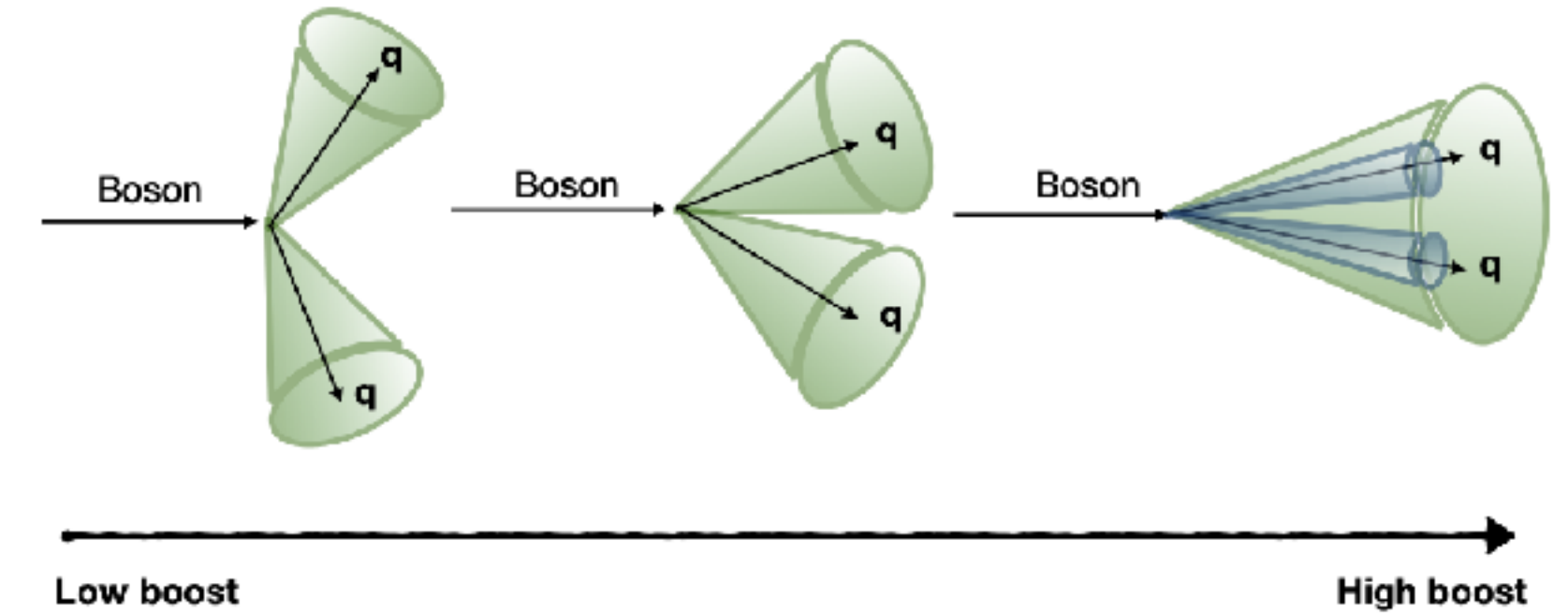


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Jet reconstruction

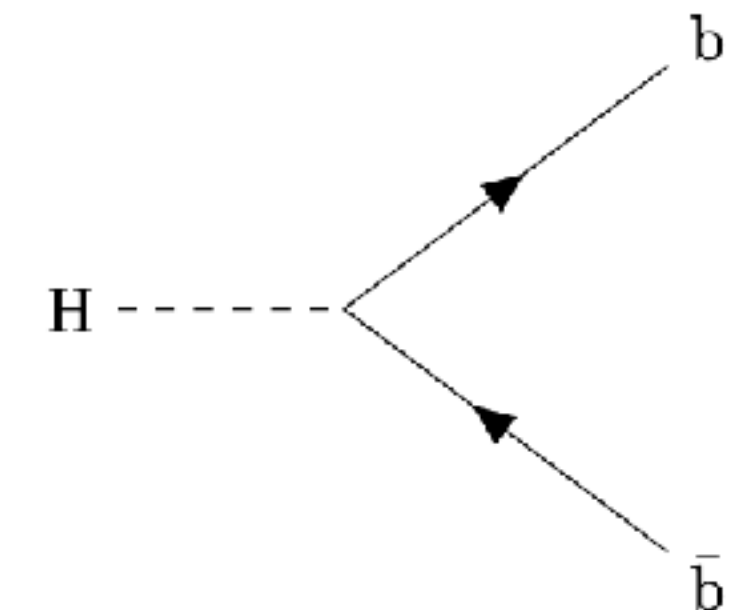
Slimjet/fatjet analysis, jet algorithms...



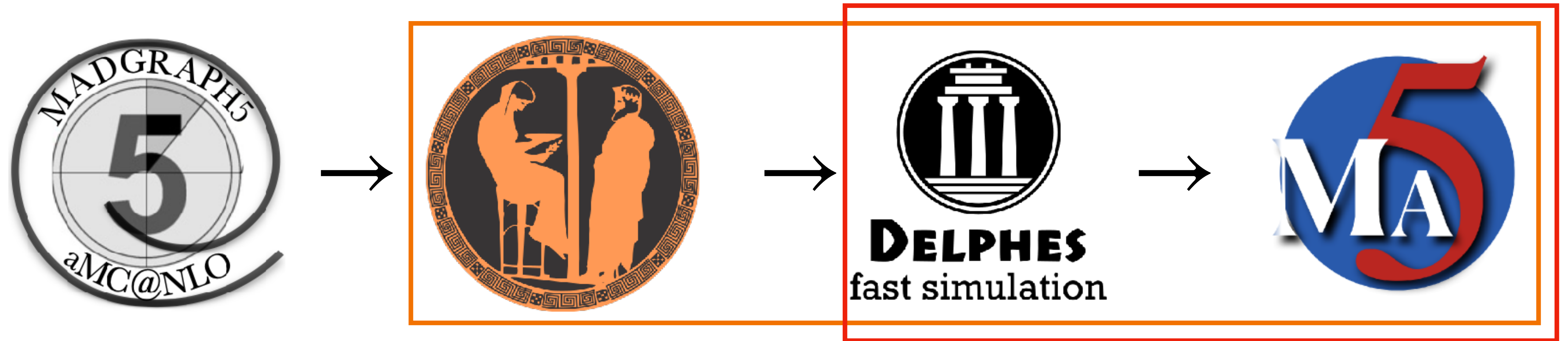
- Slimjet reconstruction:
 - Typical “resolved” case, reconstructs 4 separate b-jets with small R
- Fatjet reconstruction:
 - Boosted b-jets are collimated together, makes sense to reconstruct them together as one “fatjet” with higher R

- Jet algorithms:

- anti- k_T and variable R work in the same way, but variable R uses an effective radius



Workflow and analysis process



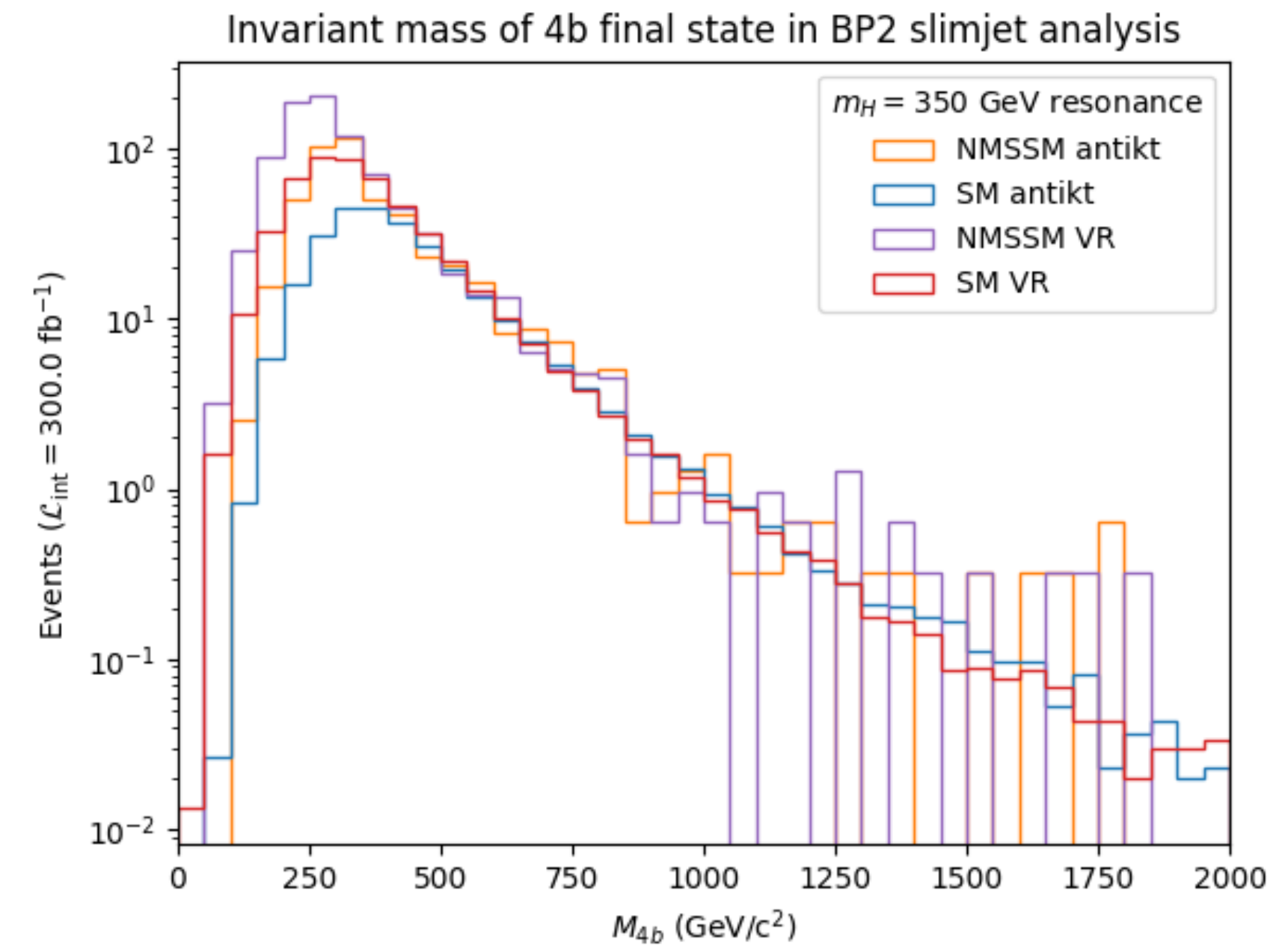
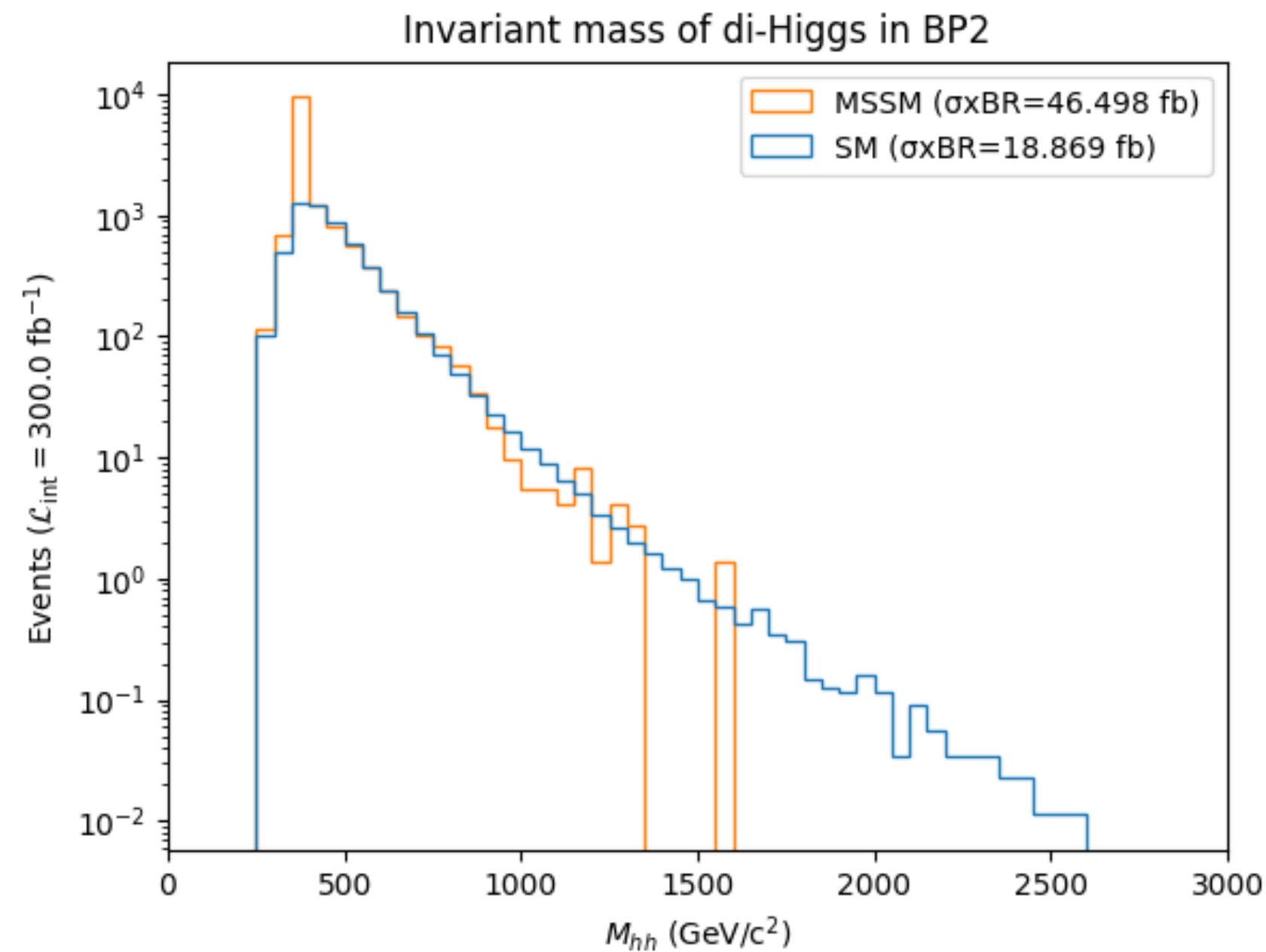
1. Jet optimization (R, ρ) done using Delphes
2. Higgs tagging done (dijets/fatjets closest to Higgs mass are labeled as “1”, the next ones are labeled as “2”)
3. Kinematic cuts placed on THT, dijet/fatjet mass and p_T using MadAnalysis5

BP2 — 350 GeV resonance

Jet optimization

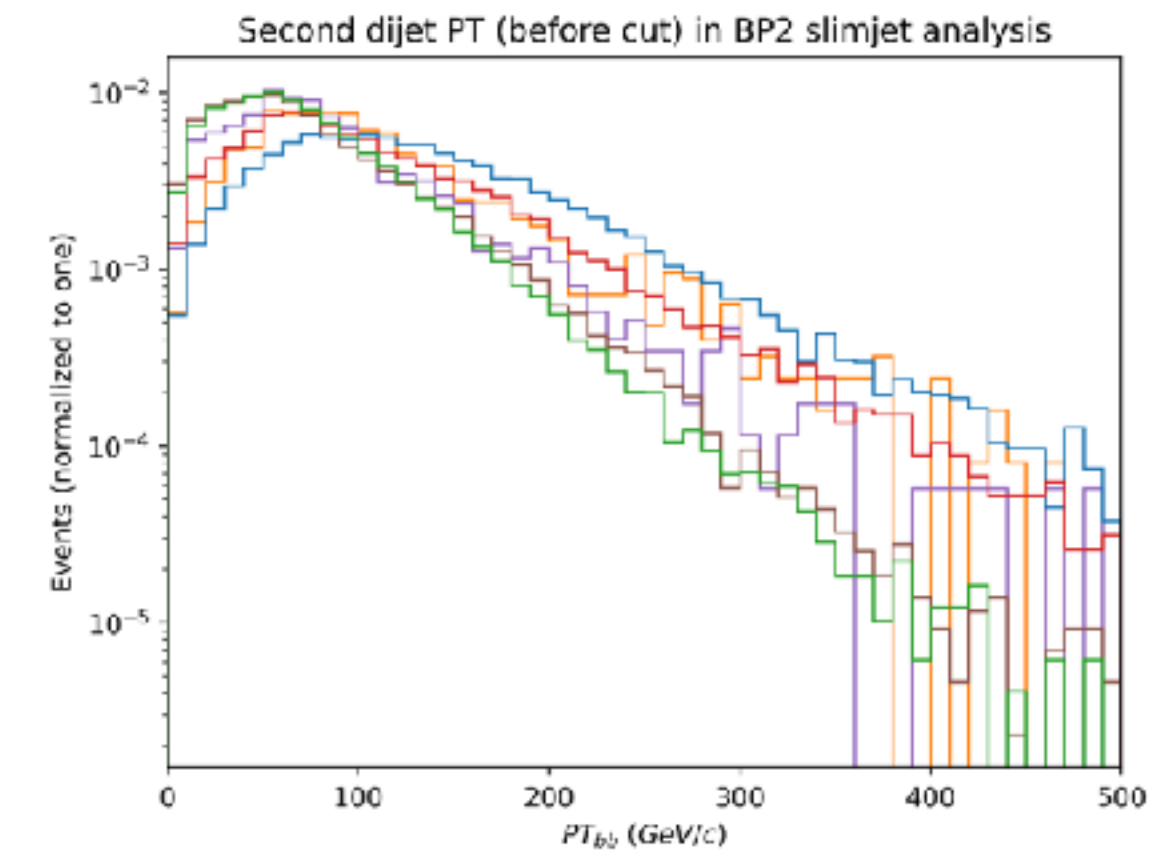
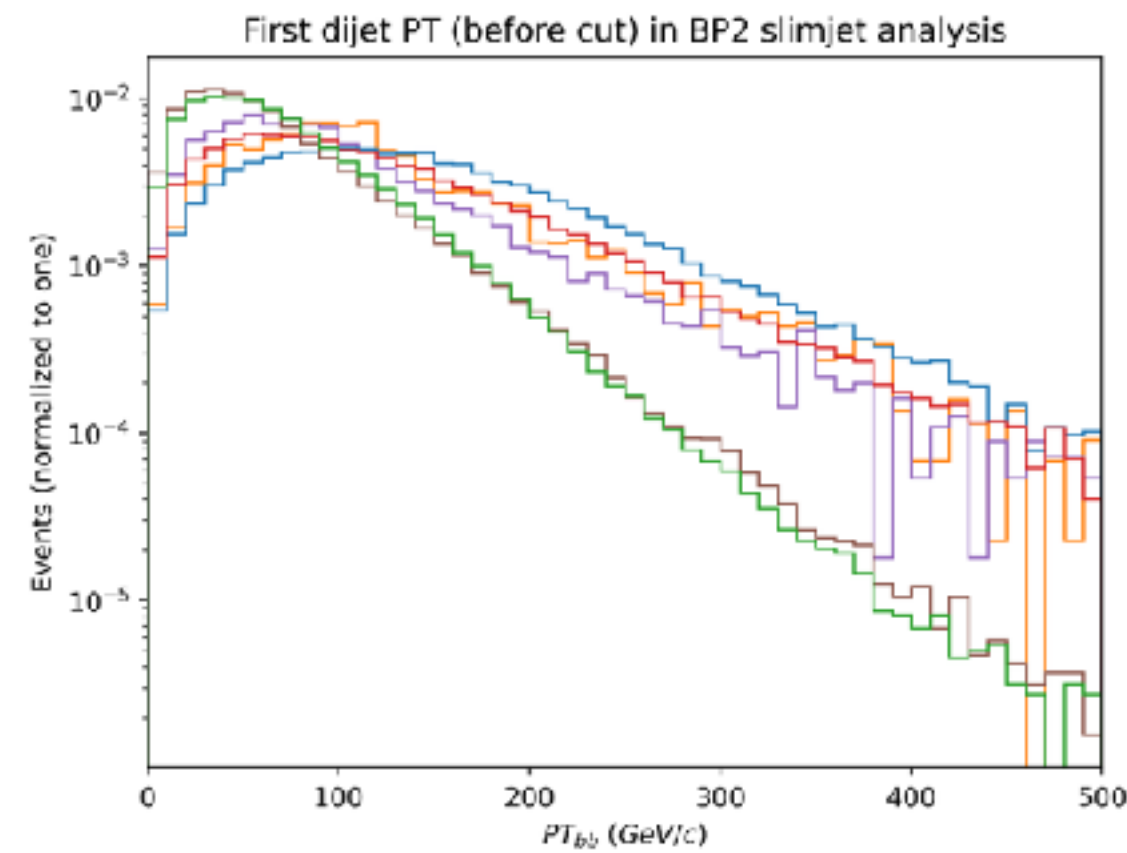
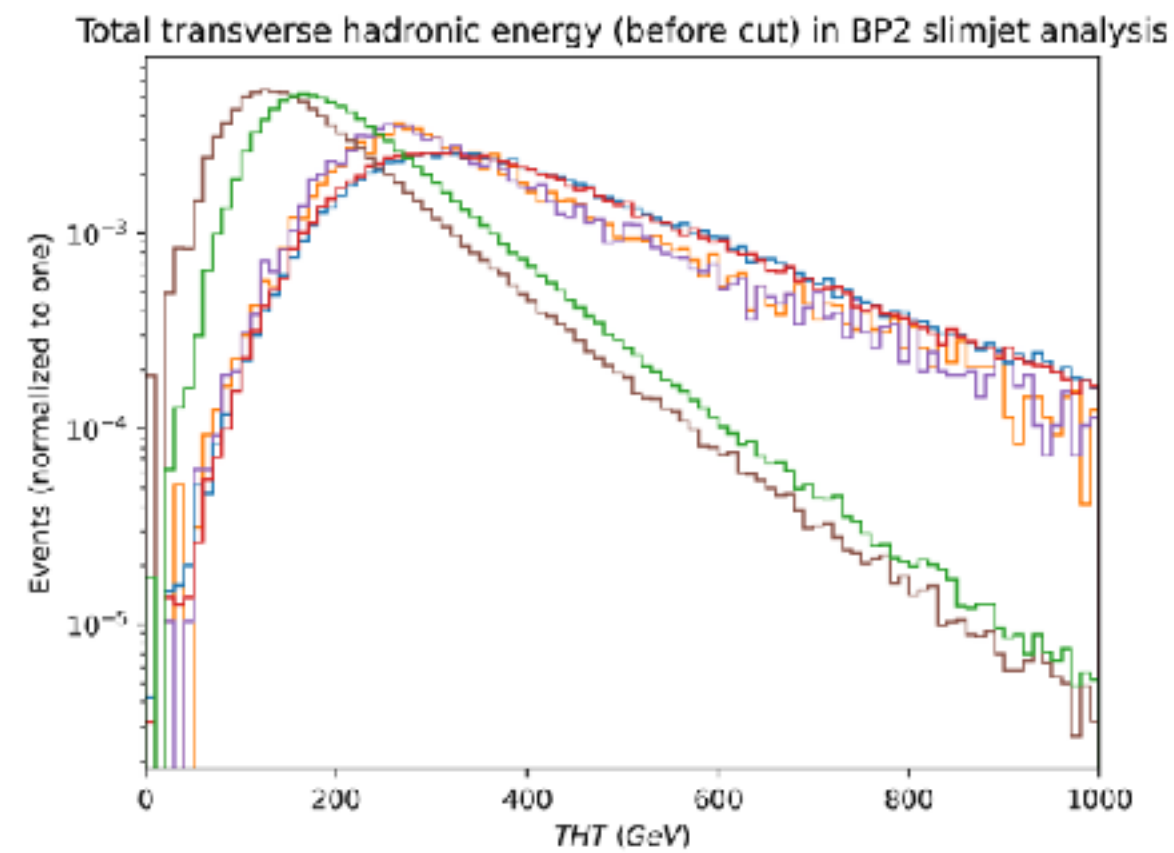
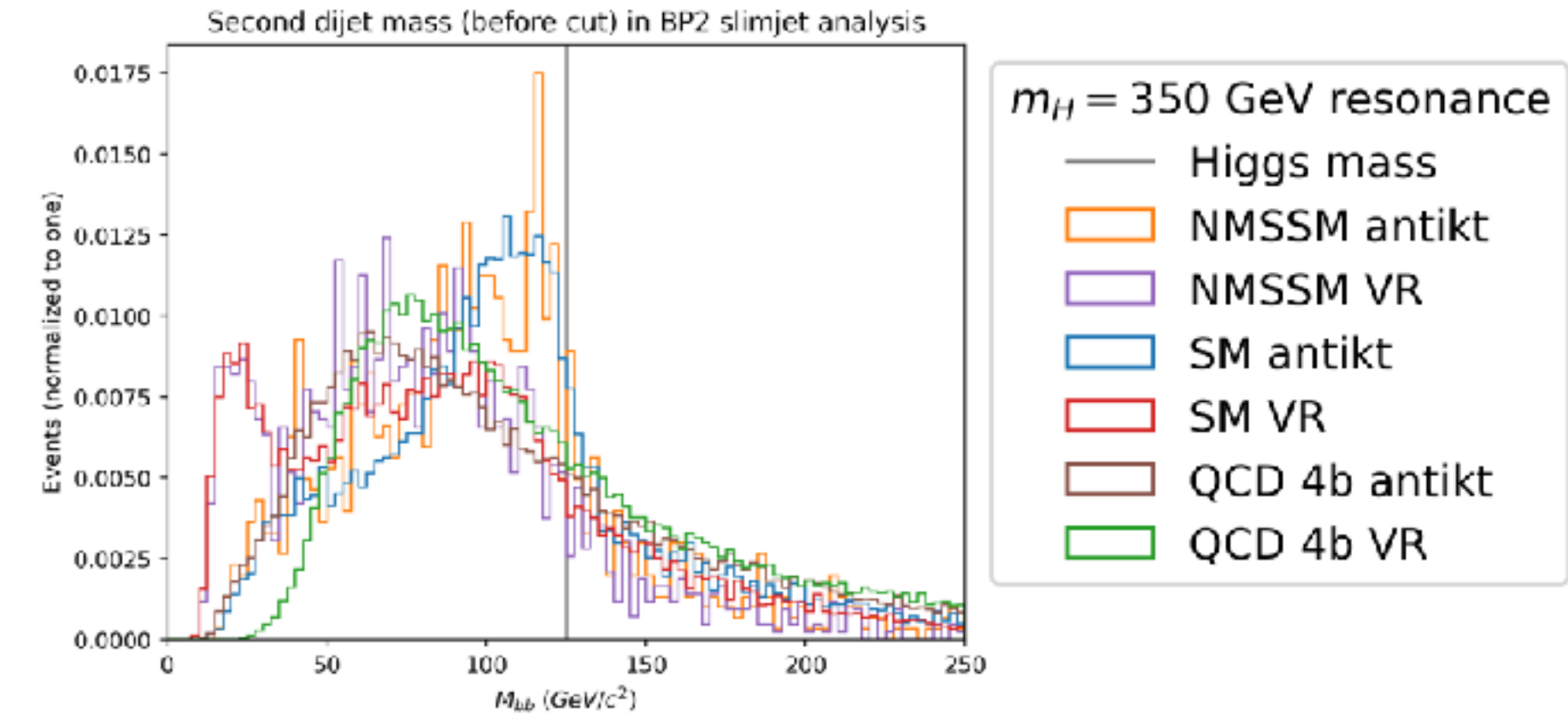
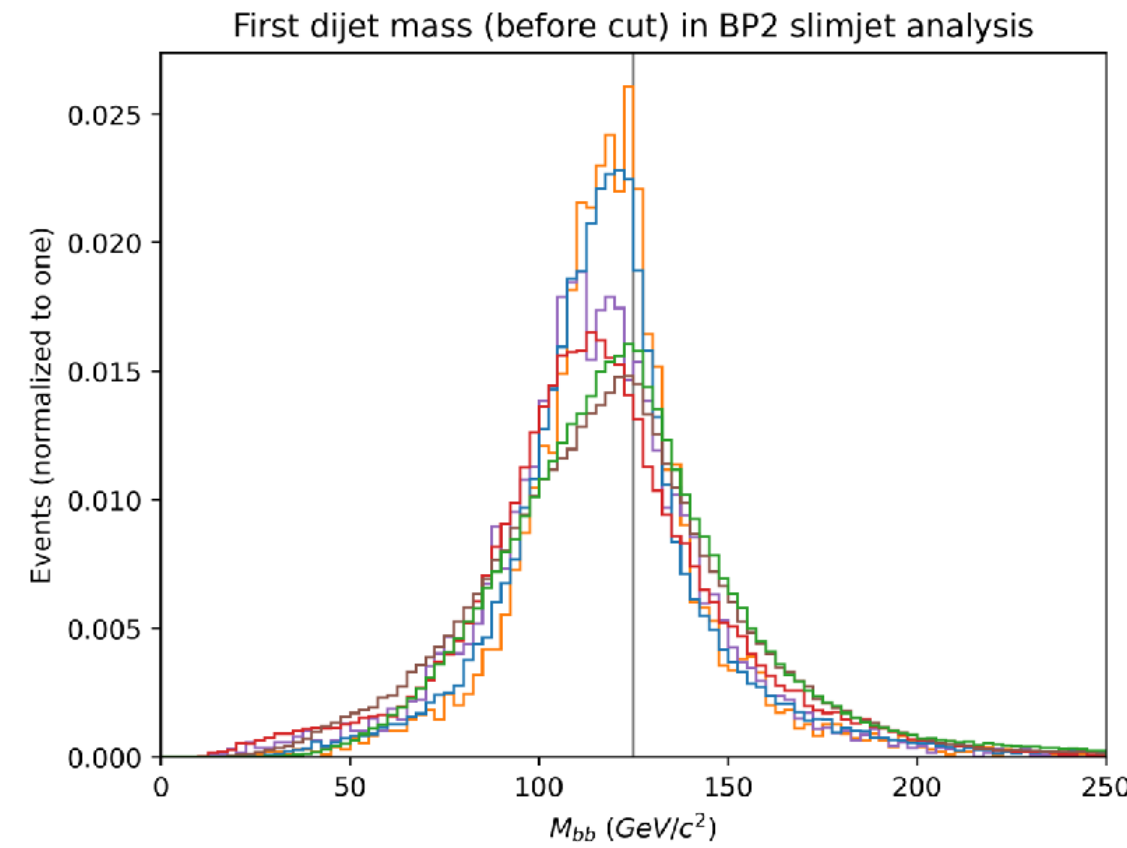
	$\sigma_{\text{before decay}}$ (fb)	$\sigma_{\text{after decay}}$ (fb)
SM	18.869	5.492
NMSSM	46.498	10.6

	slimjet
R	0.4
minR	0.2
maxR	0.5
ρ	10



BP2 — 350 GeV resonance

Kinematic cut
(including QCD background)



	THT	Dijet 1 mass	Dijet 1 pT	Dijet 2 mass	Dijet 2 pT
If $N(b) > 2$	$THT > 300$	$110 < m < 140$	$pT < 175$	—	—
If $N(b) > 3$	$THT > 300$	$110 < m < 140$	$pT < 175$	$m < 150$	$pT < 175$

BP2 — 350 GeV resonance (4b)

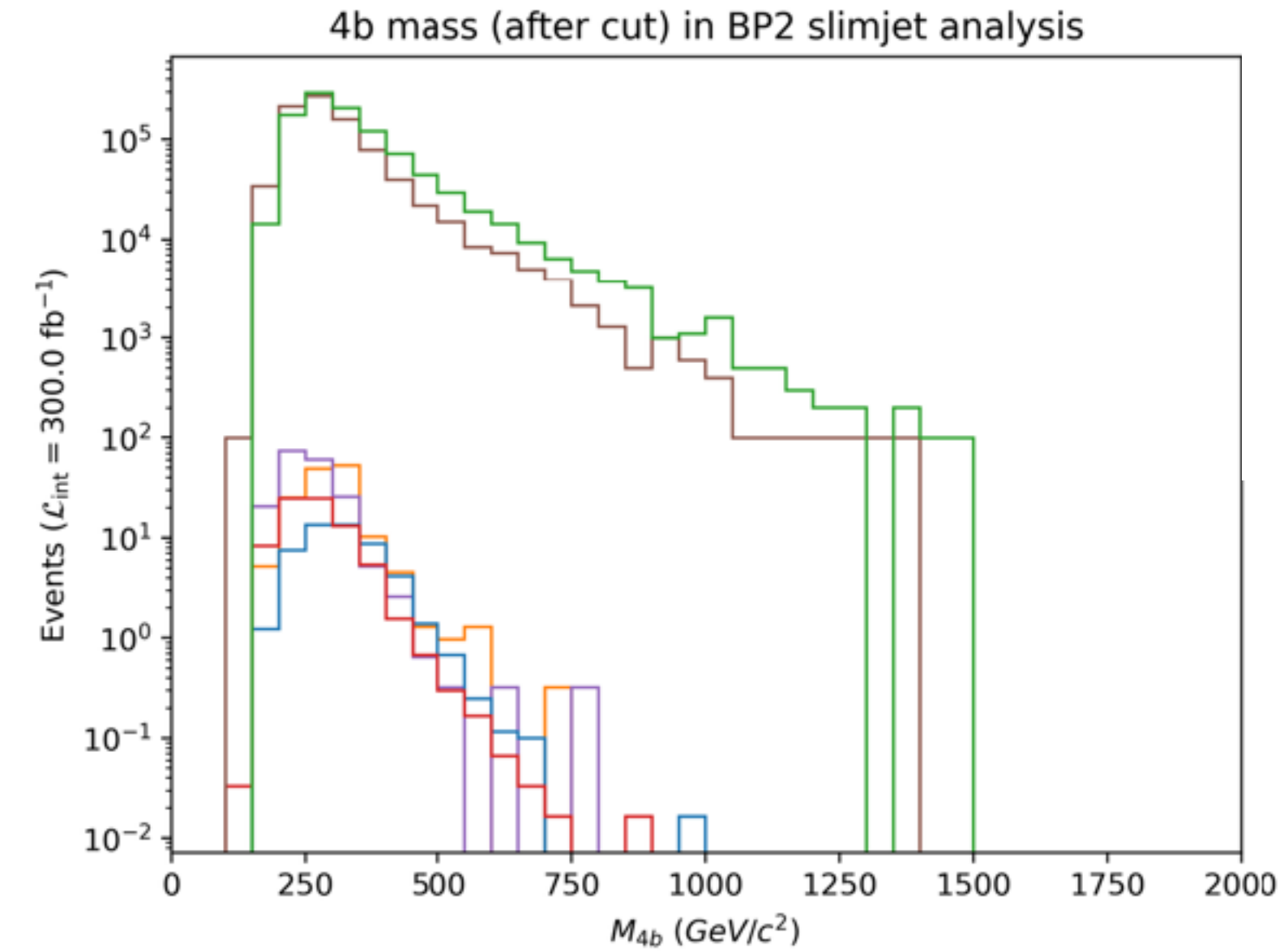
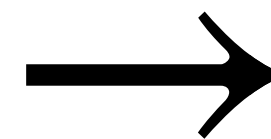
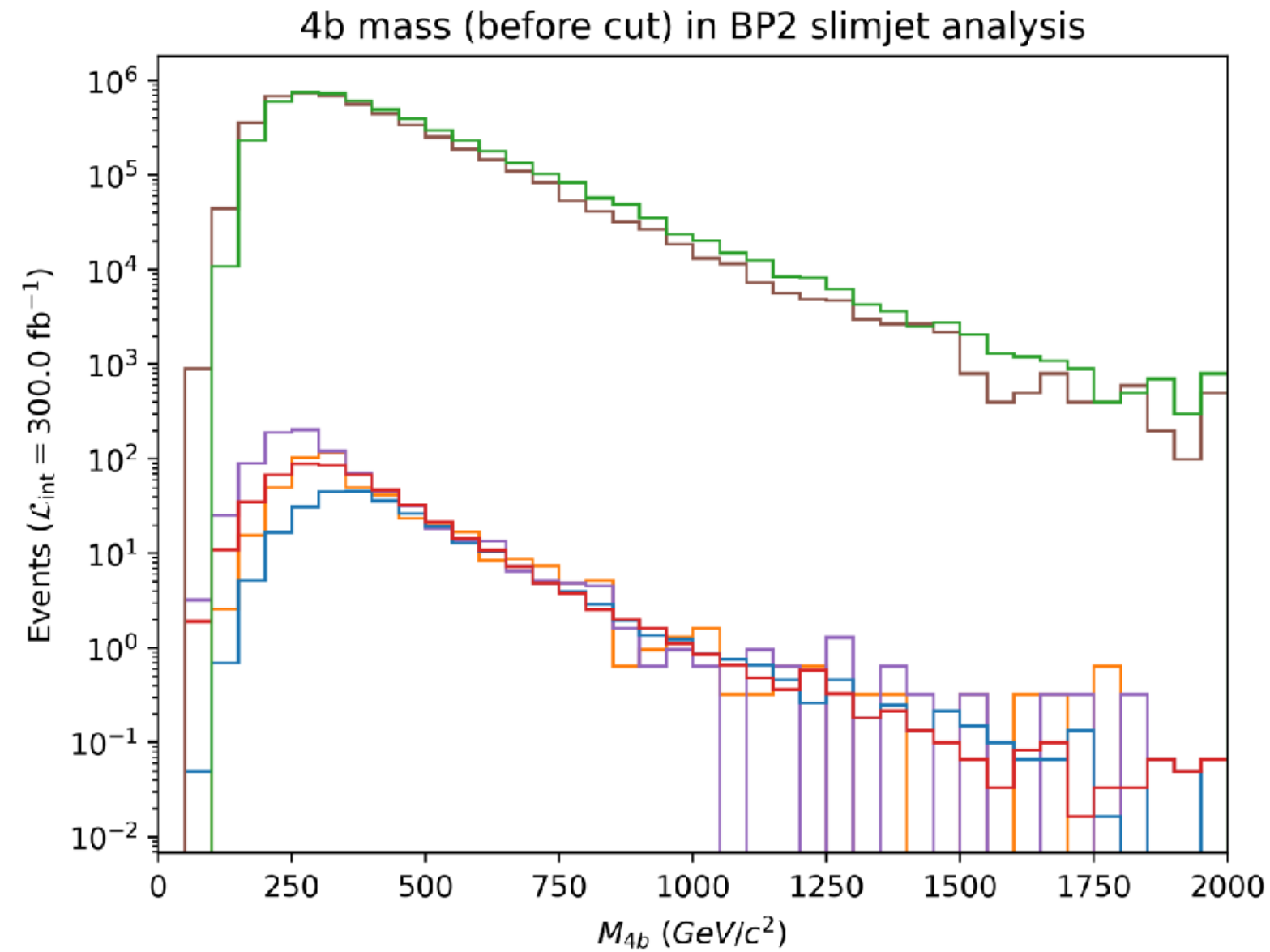
Applying the cut

SR: $M_{4b} < 500 \text{ GeV}$

$\sigma_{\text{QCD}} = 192 \text{ pb}$

$m_H = 350 \text{ GeV}$ resonance

- ▭ NMSSM antikt
- ▭ NMSSM VR
- ▭ SM antikt
- ▭ SM VR
- ▭ QCD 4b antikt
- ▭ QCD 4b VR



SR events before cut (4b)	anti-kT	Variable R
BSM di-Higgs	402	778
SM di-Higgs	206	435
Signal (BSM-SM)	196	343
QCD 4b	3867732	3821103
S/(\sqrt{B})	0.0999	0.175

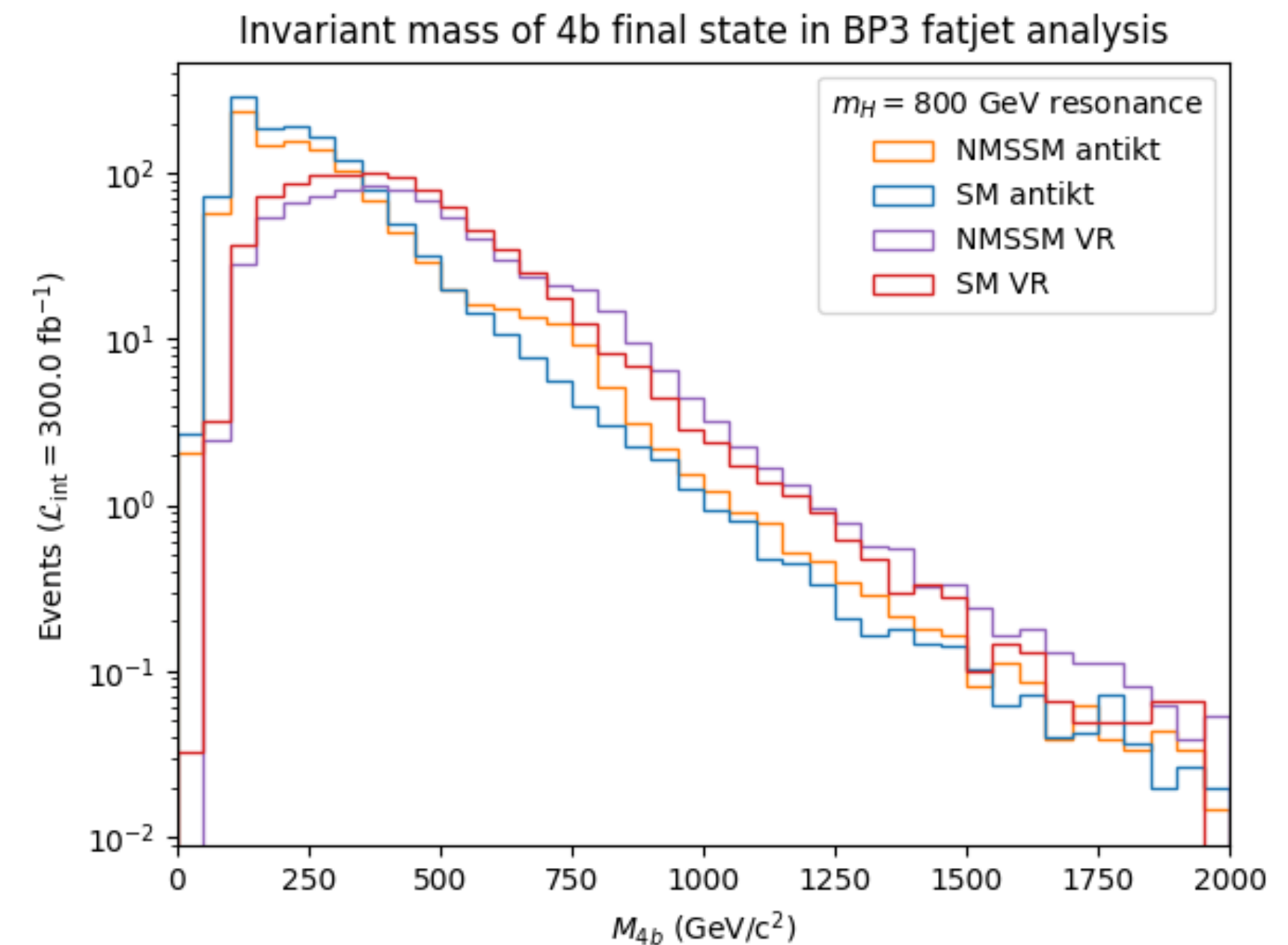
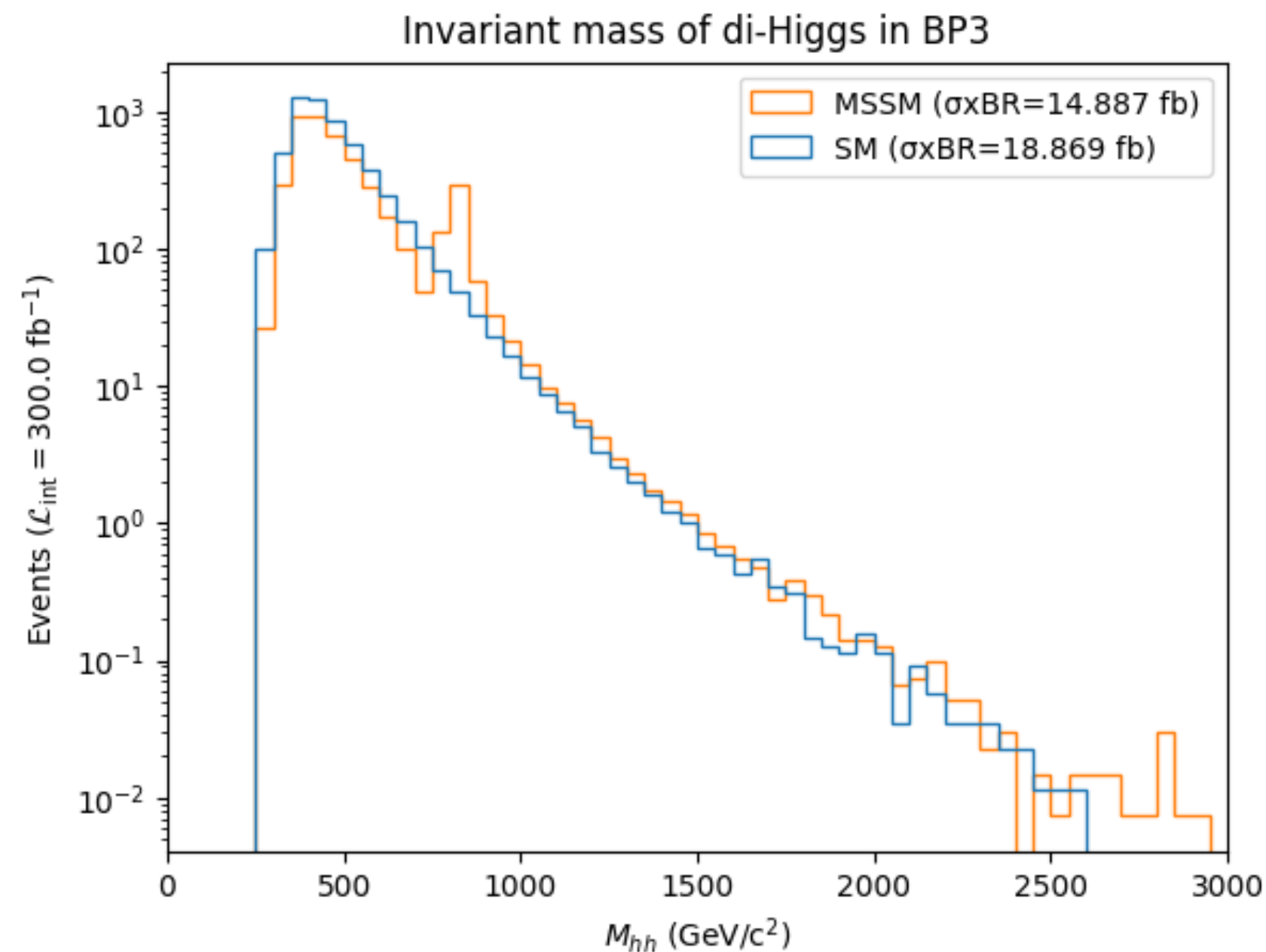
SR events after cut (4b)	anti-kT	Variable R
BSM di-Higgs	148	189
SM di-Higgs	50	78
Signal (BSM-SM)	98	111
QCD 4b	816552	924165
S/(\sqrt{B})	0.109	0.115

BP3 — 800 GeV resonance

Jet optimization

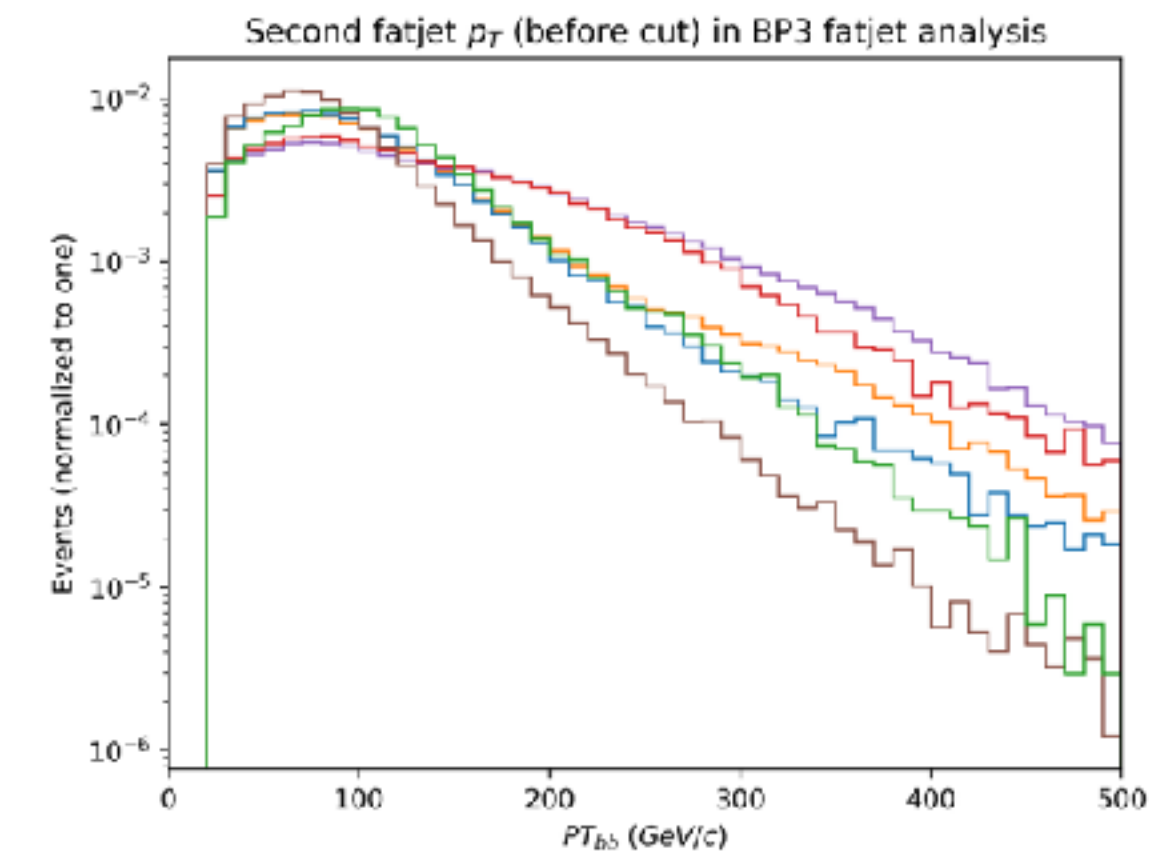
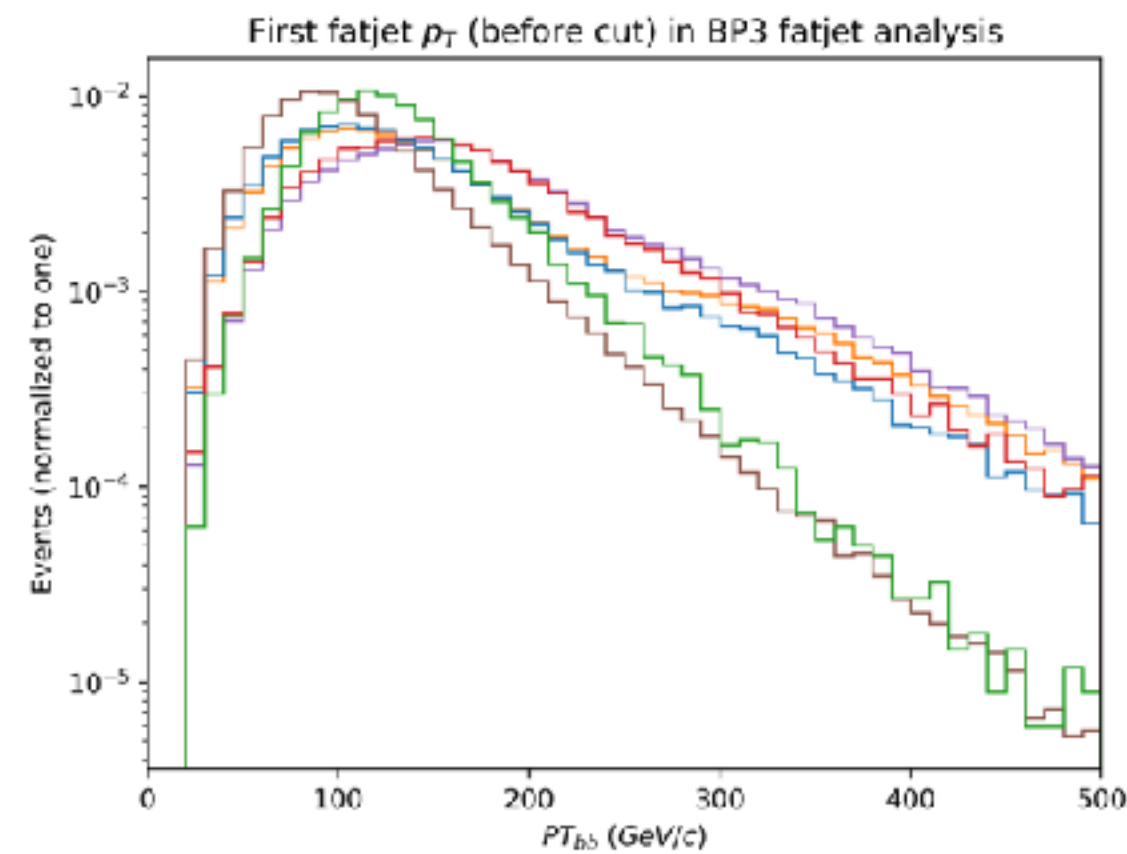
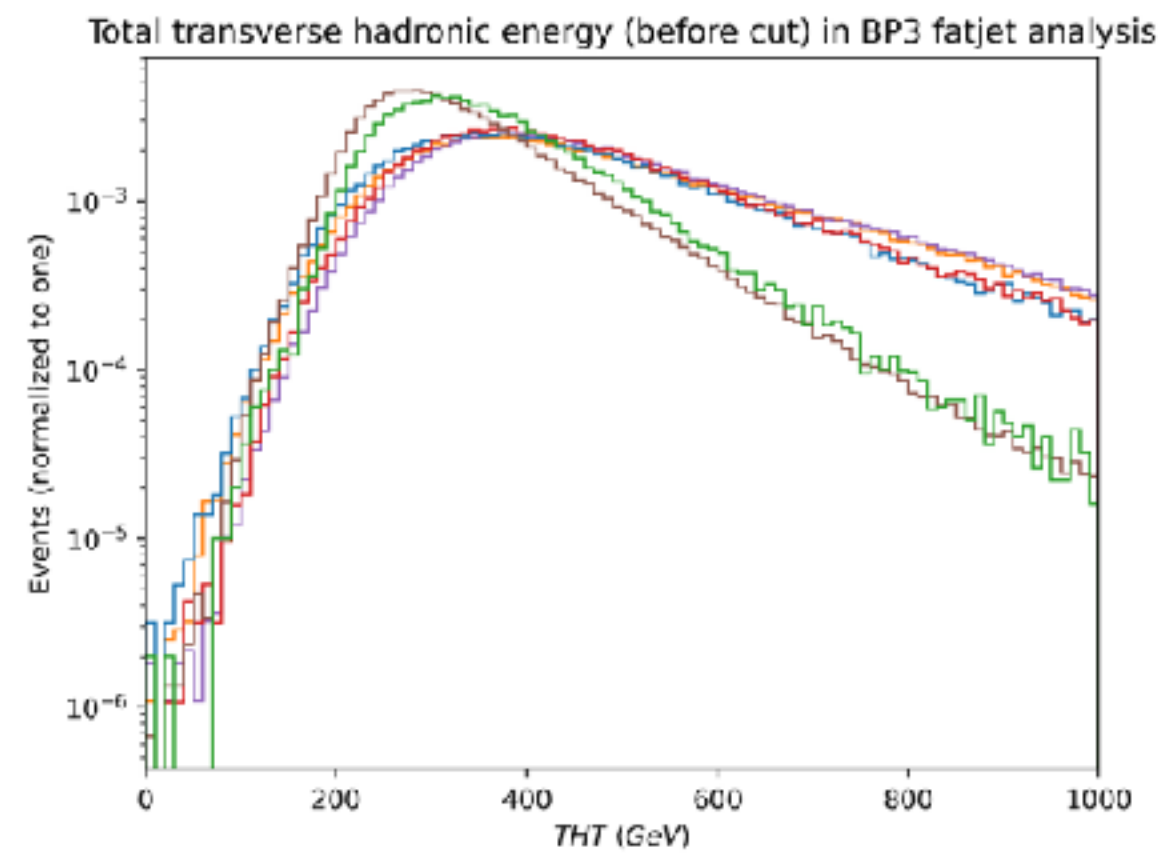
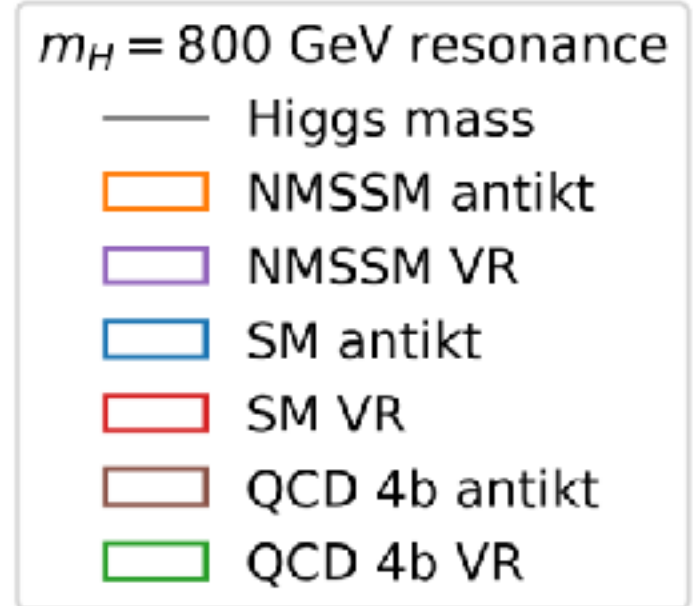
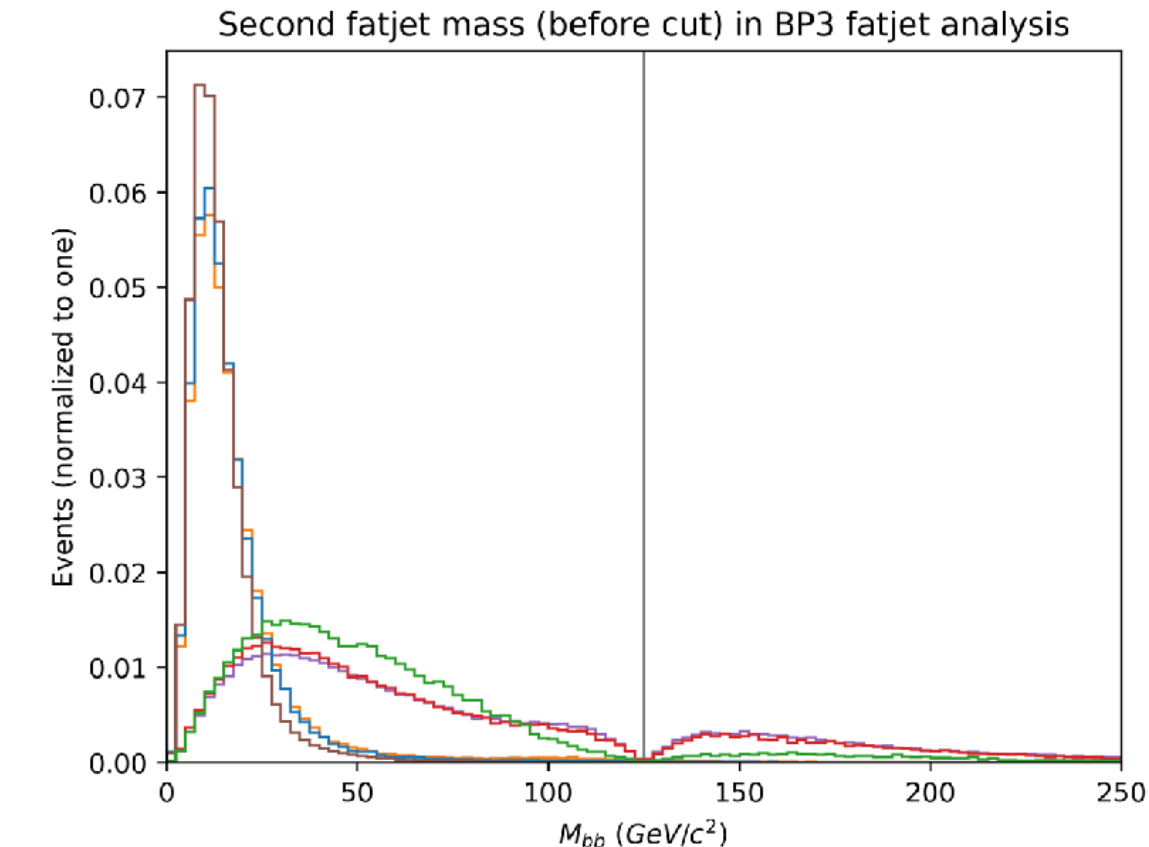
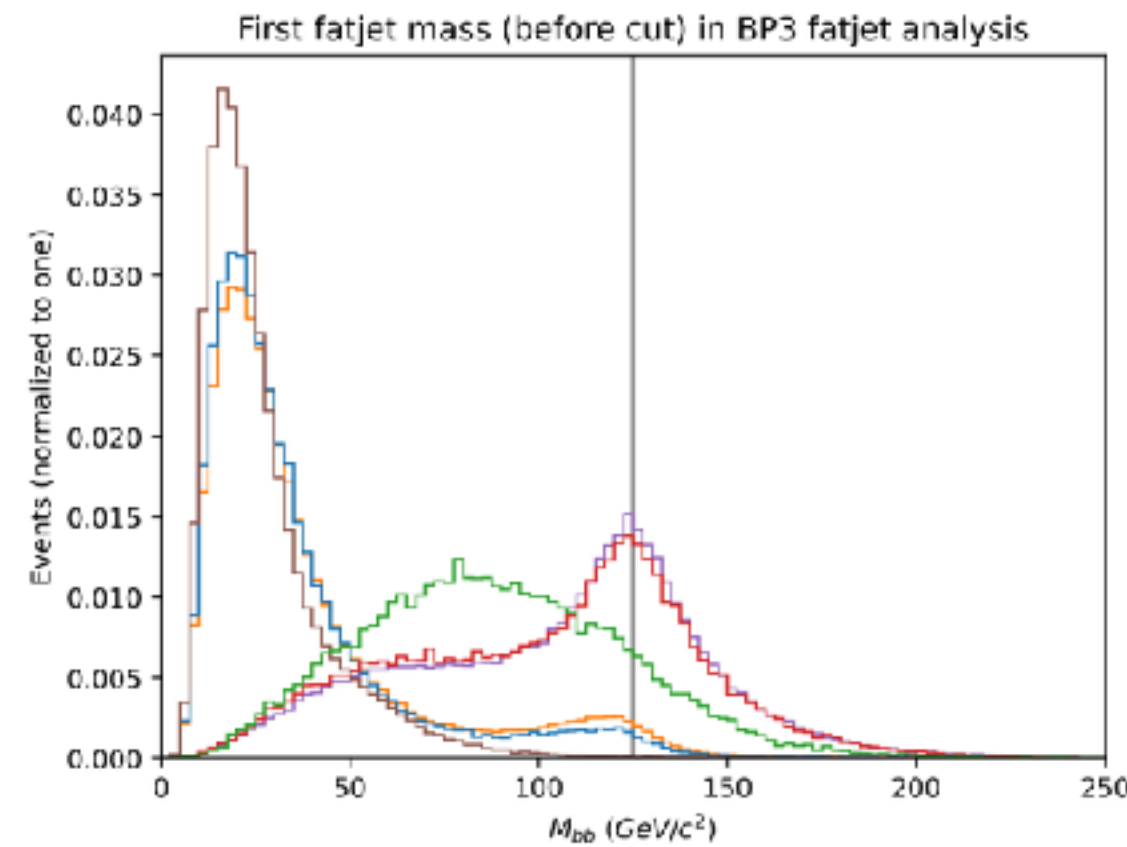
	$\sigma_{\text{before decay}}$ (fb)	$\sigma_{\text{after decay}}$ (fb)
SM	18.869	5.492
NMSSM	14.887	4.84

	fatjet
R	0.8
minR	0.6
maxR	1.6
ρ	400



BP3 — 800 GeV resonance

Kinematic cut
(including QCD background)



	THT	Fatjet 1 mass	Fatjet 1 p_T	Fatjet 2 mass	Fatjet 2 p_T
If $N(b) > 1$	$THT > 300$	$100 < m < 150$	$250 < p_T$	$M < 150$	$250 < p_T$

All units in GeV.

SR (in GeV): $600 < M_{4b} < 1200$

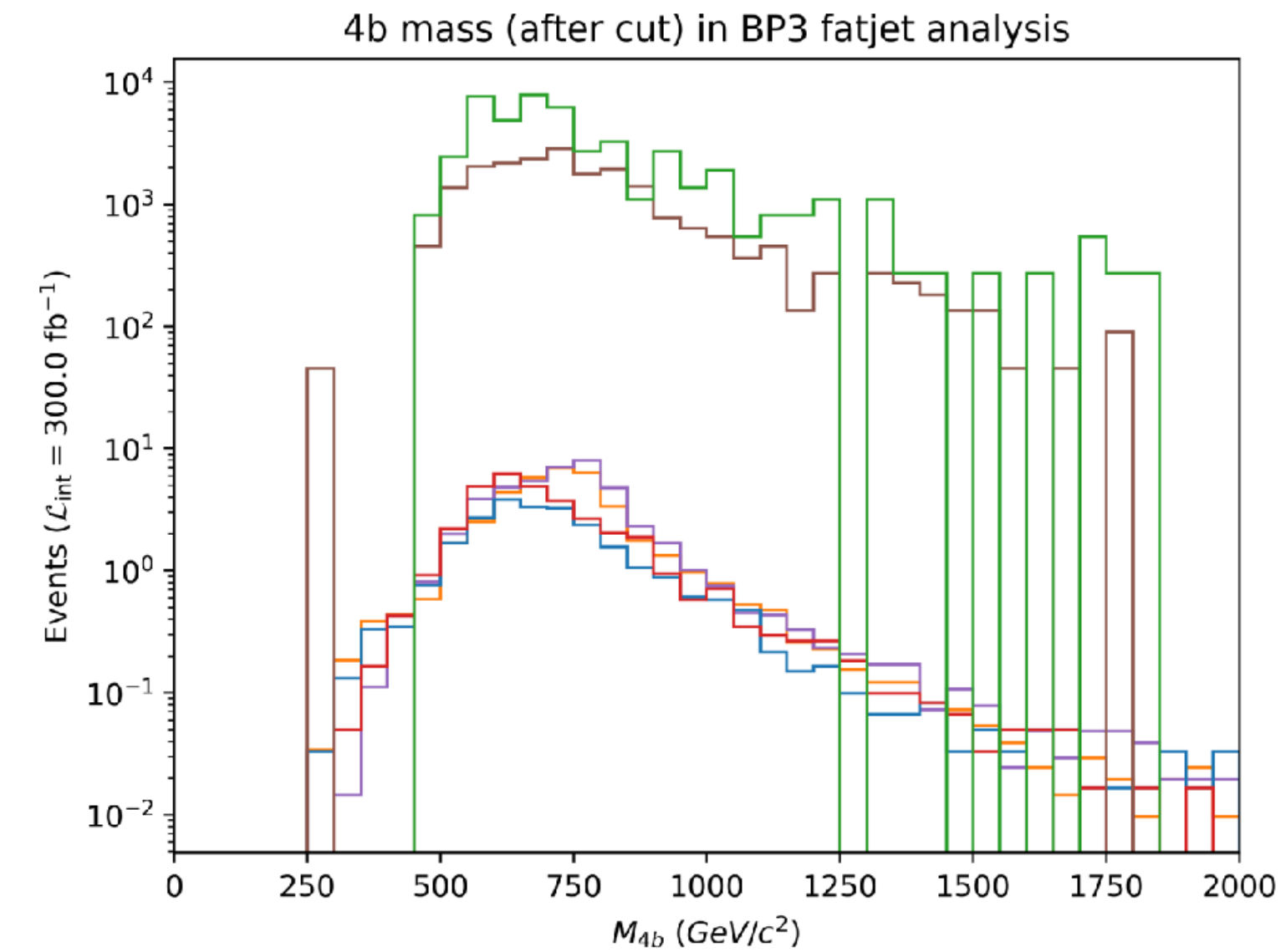
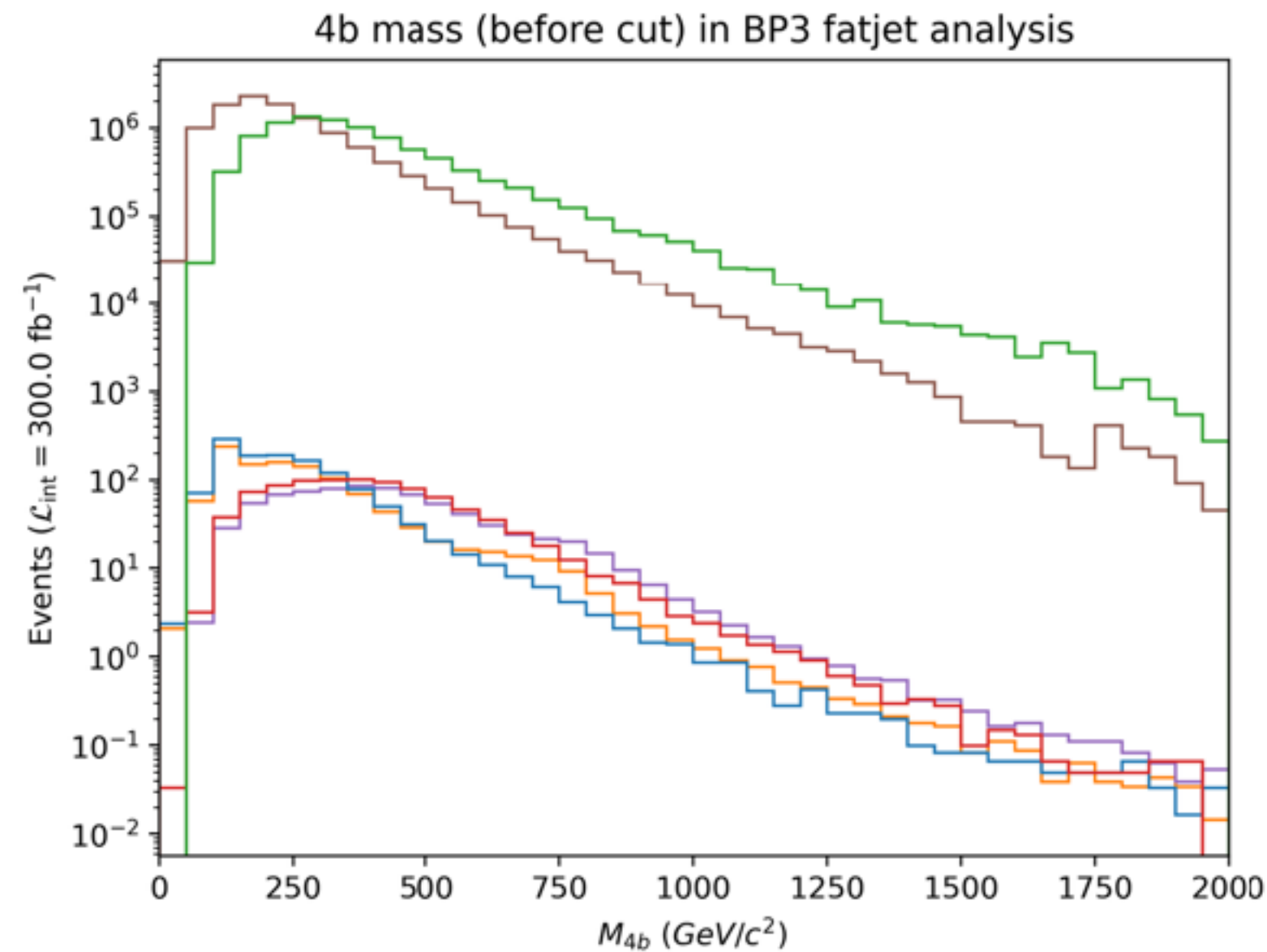
$\sigma_{\text{QCD}} = 45 \text{ pb}$

BP3 — 800 GeV resonance

Applying the cut

$m_H = 800 \text{ GeV}$ resonance

- ▭ NMSSM anti-kt
- ▭ NMSSM VR
- ▭ SM anti-kt
- ▭ SM VR
- ▭ QCD 4b anti-kt
- ▭ QCD 4b VR

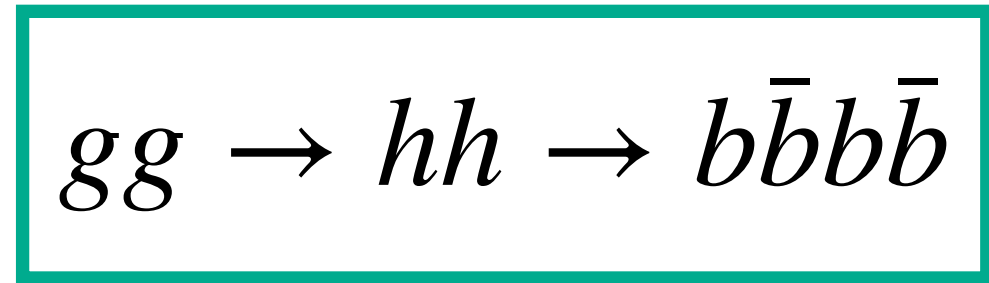


SR events before cut (4b)	anti-kT	Variable R
BSM di-Higgs	66	140
SM di-Higgs	39	120
Signal (BSM-SM)	27	20
QCD 4b	383012	1122262
$S/(\sqrt{B})$	0.0426	0.0188

SR events after cut (4b)	anti-kT	Variable R
BSM di-Higgs	33	37
SM di-Higgs	18	24
Signal (BSM-SM)	15	13
QCD 4b	15401	34247
$S/(\sqrt{B})$	0.118	0.0668

Conclusion

Di-Higgs in SUSY models



most probable!

- Variable R (generally) reconstructs more signal events
 - But it also reconstructs more QCD events
 - Smaller R works much better in excluding random background
- Kinematic cuts on the BPs decrease QCD background by ~ an order of magnitude
 - SM di-Higgs and BSM di-Higgs have similar kinematics
- Future work: investigating/placing cuts on jet substructure would also be a good idea, cuts as selection cuts for ML study

Thank you!

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Matteo Cacciari, Gavin P. Salam, and Gregory Soyez. “FastJet user manual: (for version 3.0.2)”. In: *The European Physical Journal C* 72.3 (Mar. 2012). issn: 1434-6052. doi: 10.1140/epjc/s10052-012-1896-2. url: <http://dx.doi.org/10.1140/epjc/s10052-012-1896-2>.

Eric Conte, Benjamin Fuks, and Guillaume Serret. “MadAnalysis 5, a user-friendly framework for collider phenomenology”. In: *Computer Physics Communications* 184.1 (Jan. 2013), 222–256. issn: 0010-4655. doi: 10.1016/j.cpc.2012.09.009. url: <http://dx.doi.org/10.1016/j.cpc.2012.09.009>.

The simplified MSSM model

- Based on superpotential

$$W_{\text{MSSM}} = y_u \bar{Q} H_u U^c + y_d \bar{Q} H_d D^c + y_\ell \bar{L} H_d E^c + \mu H_u H_d$$

- Introduces two stop squarks
- Lagrangian induced by NP:

$$\mathcal{L}_{\text{NP}}^{\tilde{q}} = -(\lambda^{\text{SM}} + \kappa_{hhh})vh^3 - \frac{1}{\sqrt{2}}(y_t^{\text{SM}} + \kappa_{htt})h\bar{t}t + vh(\tilde{s}_1^* \tilde{s}_2^*) \begin{pmatrix} \kappa_{h\tilde{s}\tilde{s}}^{11} & \kappa_{h\tilde{s}\tilde{s}}^{12} \\ \cdot & \kappa_{h\tilde{s}\tilde{q}}^{22} \end{pmatrix} \begin{pmatrix} \tilde{s}_1 \\ \tilde{s}_2 \end{pmatrix} + hh(\tilde{s}_1^* \tilde{s}_2^*) \begin{pmatrix} \kappa_{hh\tilde{s}\tilde{s}}^{11} & \kappa_{hh\tilde{s}\tilde{s}}^{12} \\ \cdot & \kappa_{hh\tilde{s}\tilde{s}}^{22} \end{pmatrix} \begin{pmatrix} \tilde{s}_1 \\ \tilde{s}_2 \end{pmatrix}$$

Squark terms

Modified couplings

Higgs masses in the MSSM

Physical
CP-even
states:
h and *H*

- Upper bound on lighter Higgs:

$$m_{h,0}^2 \leq m_Z^2 \cos^2 2\beta, \text{ where } \tan \beta = \frac{\langle H_u^0 \rangle}{\langle H_d^0 \rangle}$$

- Light Higgs must be lower than Z boson mass (91.2 GeV)
- Large loop corrections needed for $m_h = 125$ GeV
 - Can come from top-stop sector \rightarrow large $\tan \beta$
- Resonant di-Higgs production is suppressed because at large $\tan \beta$, heavy Higgs mass must be higher than a TeV

The simplified NMSSM model

- Based on superpotential

$$W_{\text{NMSSM}} = y_u Q H_u U^c + y_d Q H_d D^c + y_\ell L H_d E^c + \lambda S H_u H_d + \frac{\kappa}{3} S^3$$

- Introduces two stop squarks, two sbottom squarks, one neutral scalar
- Lagrangian terms induced by NP:

Modified couplings

$$\mathcal{L}_M = -(\lambda_{\text{SM}} + \kappa_{hhh})vh^3 - \frac{1}{\sqrt{2}}(y_f^{\text{SM}} + \kappa_{hff})h\bar{f}f$$

Squark terms

$$\mathcal{L}_{\tilde{s}} = \sum_i (\kappa_{h\tilde{s}\tilde{s}}^{ii}vh\tilde{s}_i^*\tilde{s}_i + \kappa_{hh\tilde{s}\tilde{s}}^{ii}hh\tilde{s}_i^*\tilde{s}_i) + \left(\sum_{i>j} \kappa_{h\tilde{s}\tilde{s}}^{ij}vh\tilde{s}_i^*\tilde{s}_j + \text{h.c.} \right)$$

Scalar terms

$$\mathcal{L}_S = \sum_I \kappa_{Shh}^I v S_I^0 hh + \kappa_{Sff}^I S_I^0 \bar{f}f$$

Scalar-squark terms

$$\mathcal{L}_{\tilde{s}S} = \sum_{I,i} \kappa_{S\tilde{s}\tilde{s}}^{Ii} v S_I^0 \tilde{s}_i^* \tilde{s}_i$$

Higgs masses in the NMSSM

- Upper bound on lighter Higgs:

$$m_{h,0}^2 \leq m_Z^2 \left(\cos^2 2\beta + \frac{2\lambda^2}{g^2 + g'^2} \sin^2 2\beta \right)$$

- Light Higgs can be heavier than Z
- Less constraints than MSSM
- At $1 \lesssim \tan \beta \lesssim 3$ and large λ , $m_h = 125$ GeV without loop corrections

Benchmark points

Di-Higgs in SUSY models

In MSSM

- Large squark loops at large $\tan \beta$

In NMSSM

- Large resonance effects:

Large λ and small $\tan \beta$

- Large squark effects:

Small λ and high $\tan \beta$

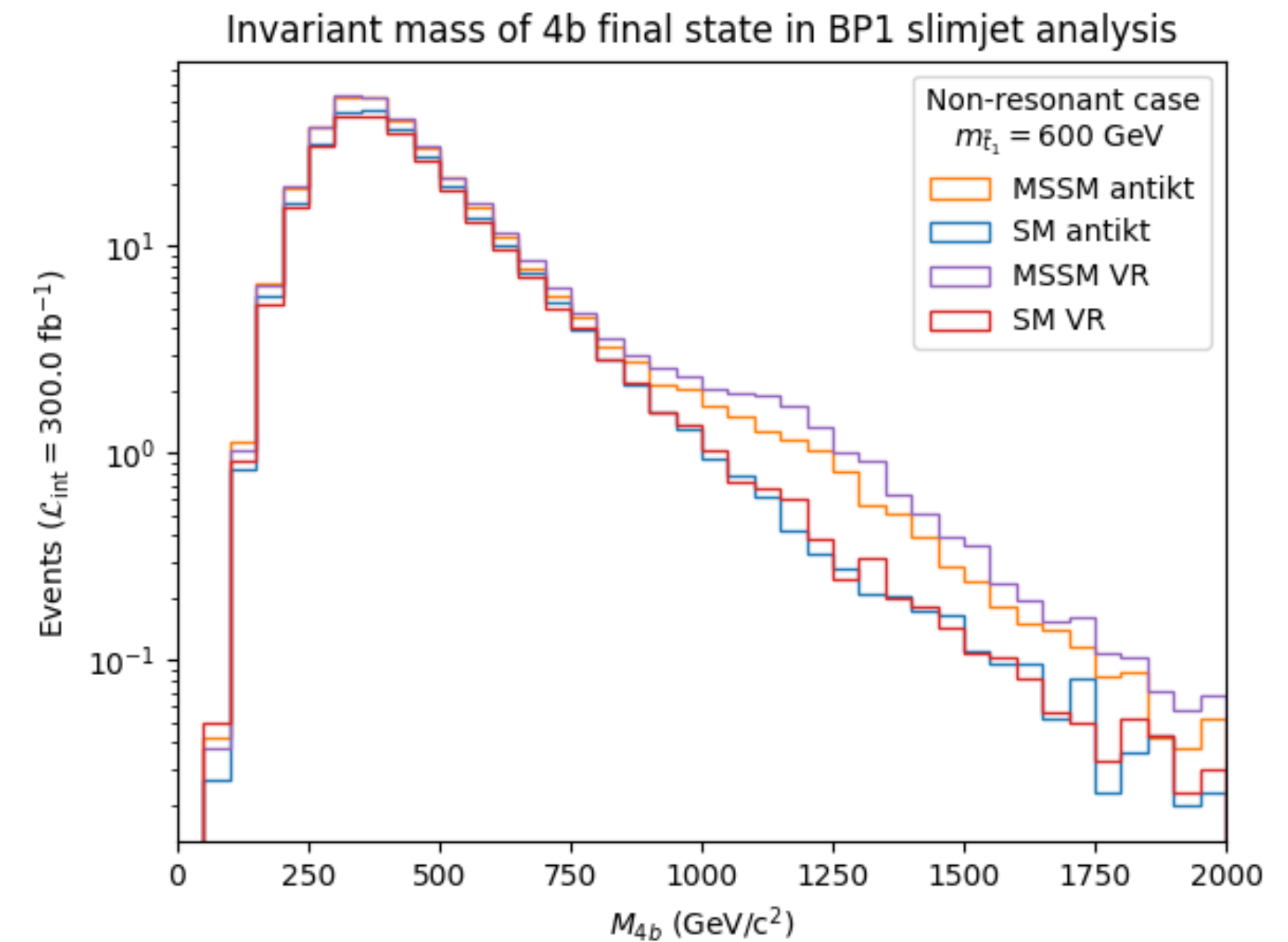
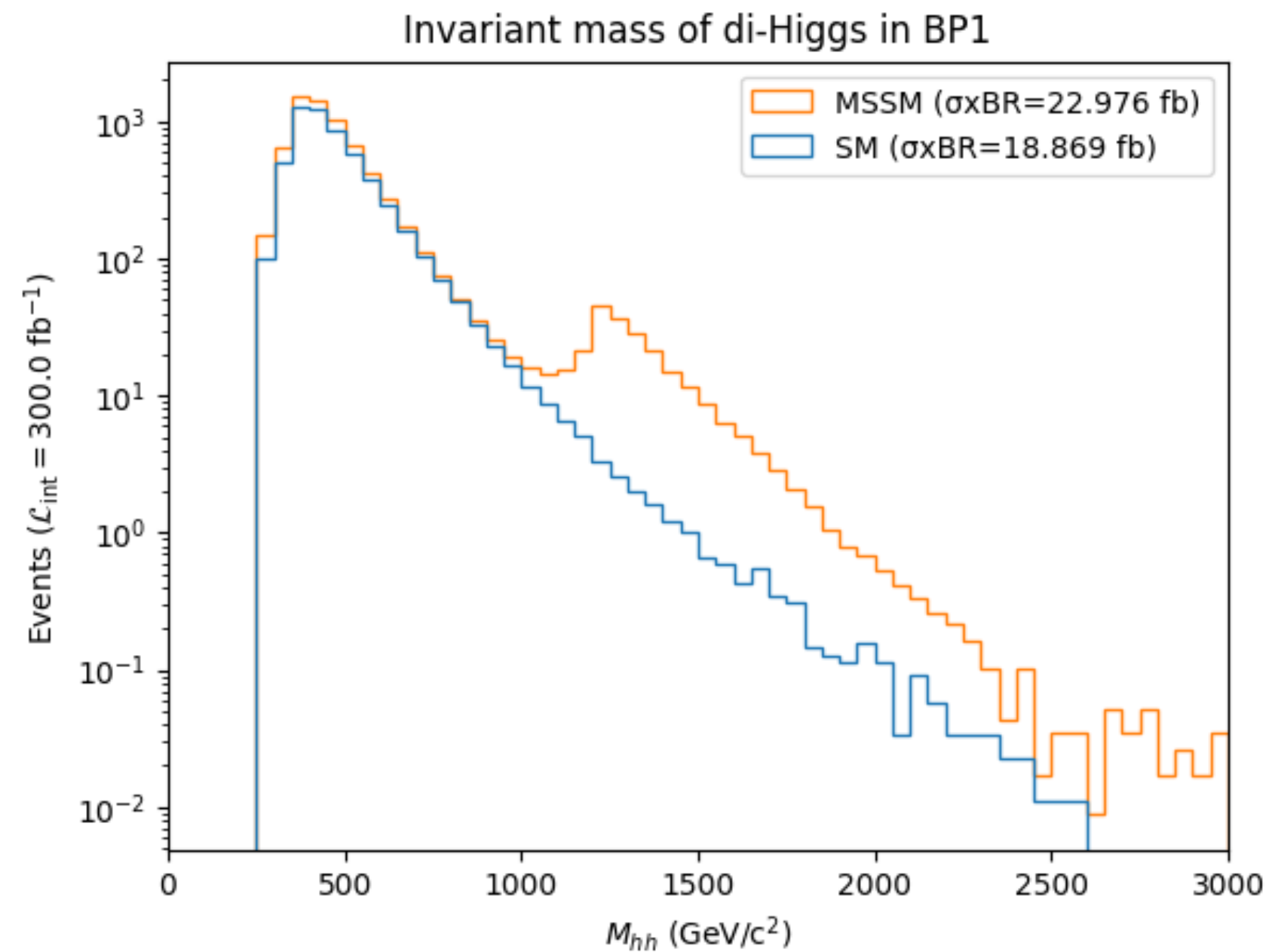
	BP1	BP2	BP3
production mode	non-resonant	resonant	resonant
$m_{\tilde{t}_1}$ (in GeV)	600.6	1601.5	602.4
$m_{\tilde{t}_2}$ (in GeV)	1301.0	1951.8	1398.9
$m_{\tilde{b}_1}$ (in GeV)	—	1852.7	1365.0
$m_{\tilde{b}_2}$ (in GeV)	—	1951.8	1623.0
$m_{S_1^0}$ (in GeV)	—	354.6	800.0
$\tan \beta$	45.4	30	2.5
λ	—	0.043	0.7
reconstruction	slimjet and fatjet	slimjet	fatjet

BP1 — non-resonant slimjet

Jet optimization

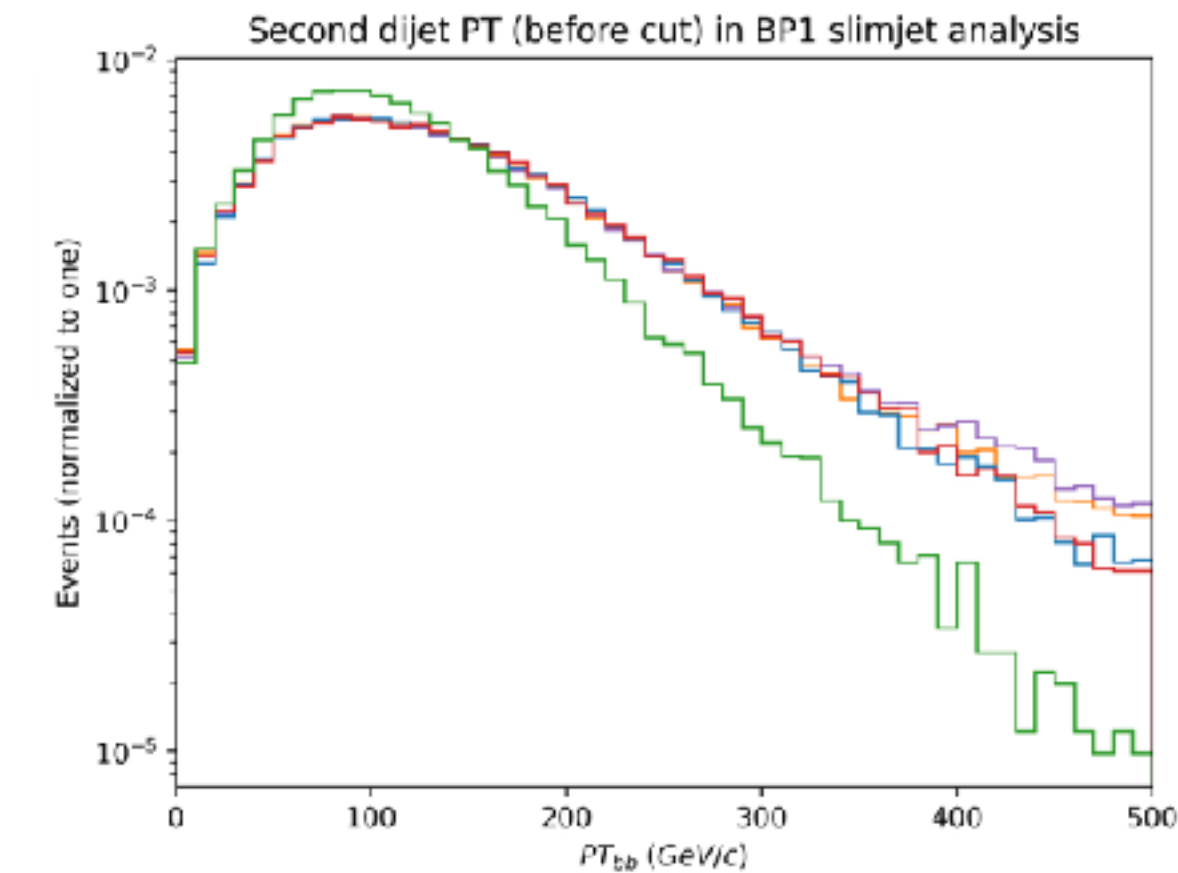
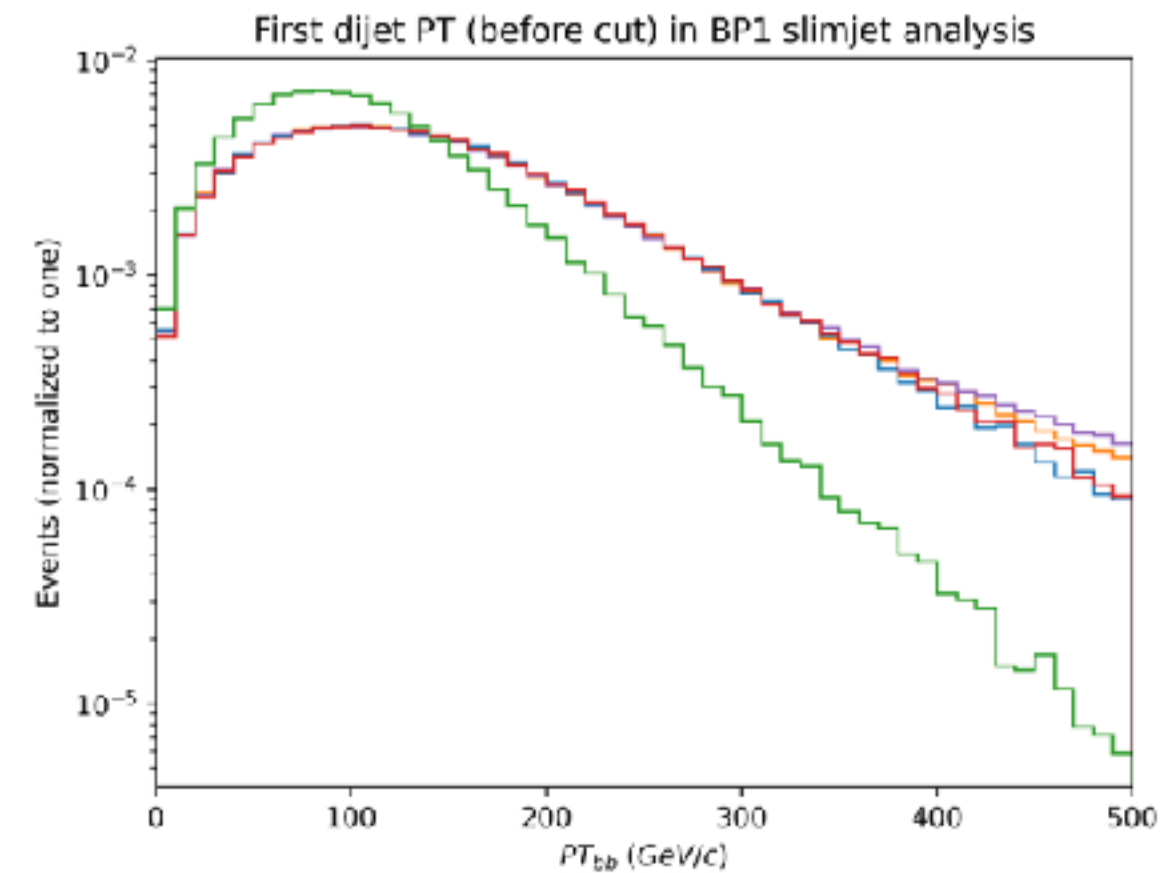
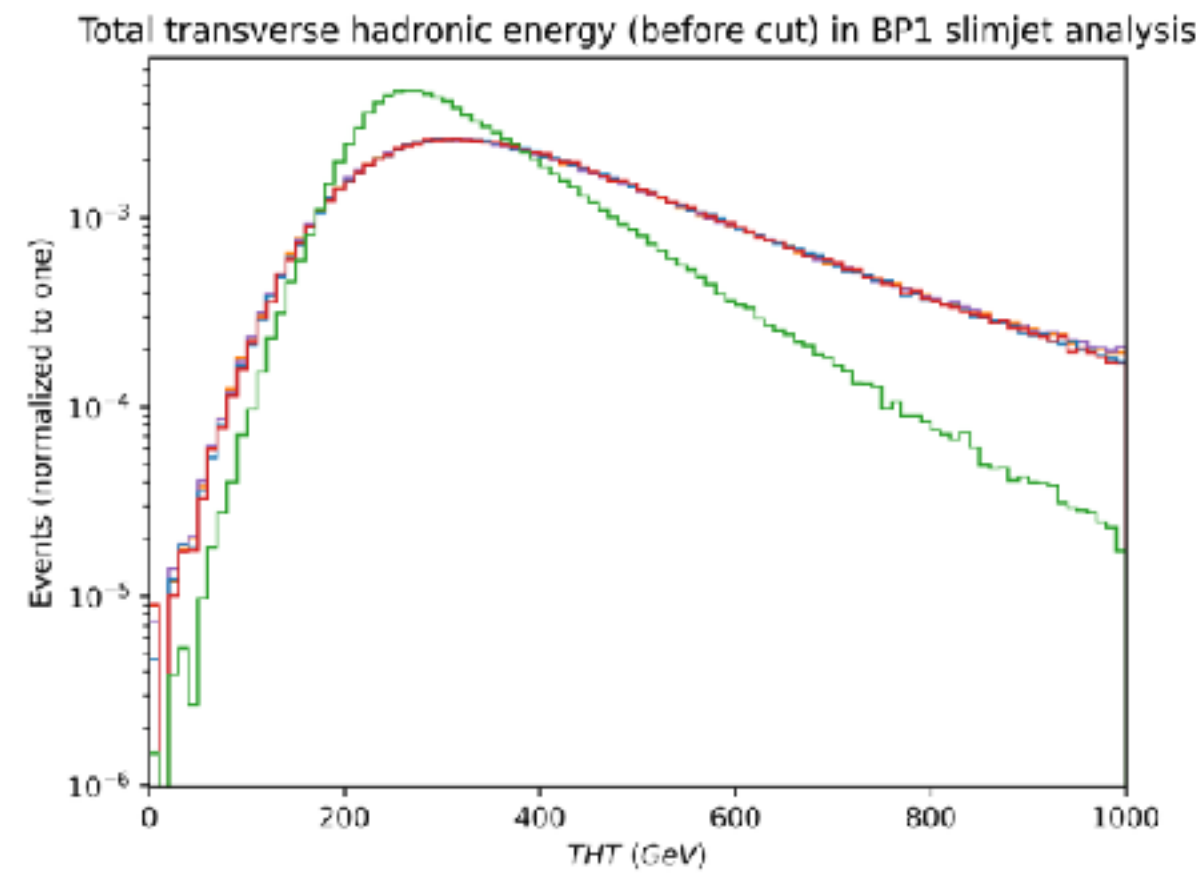
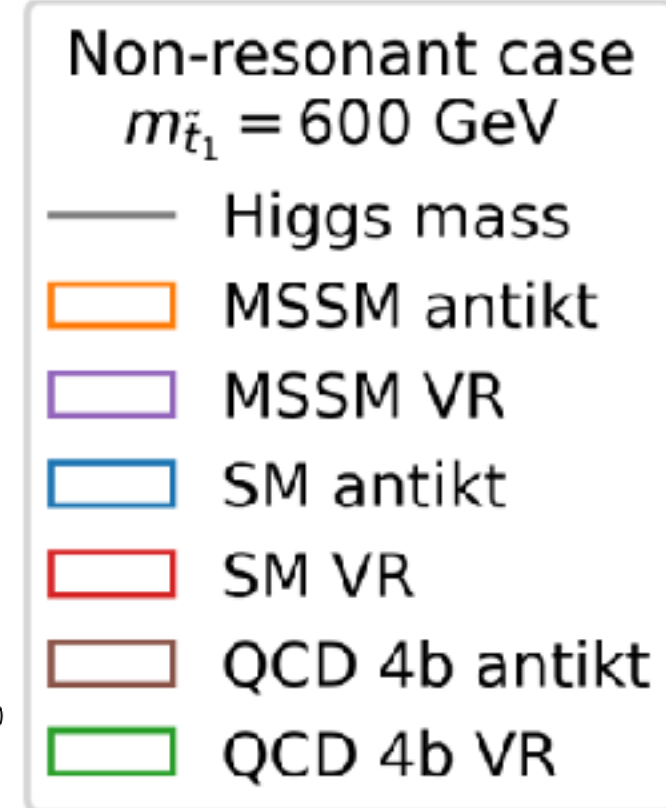
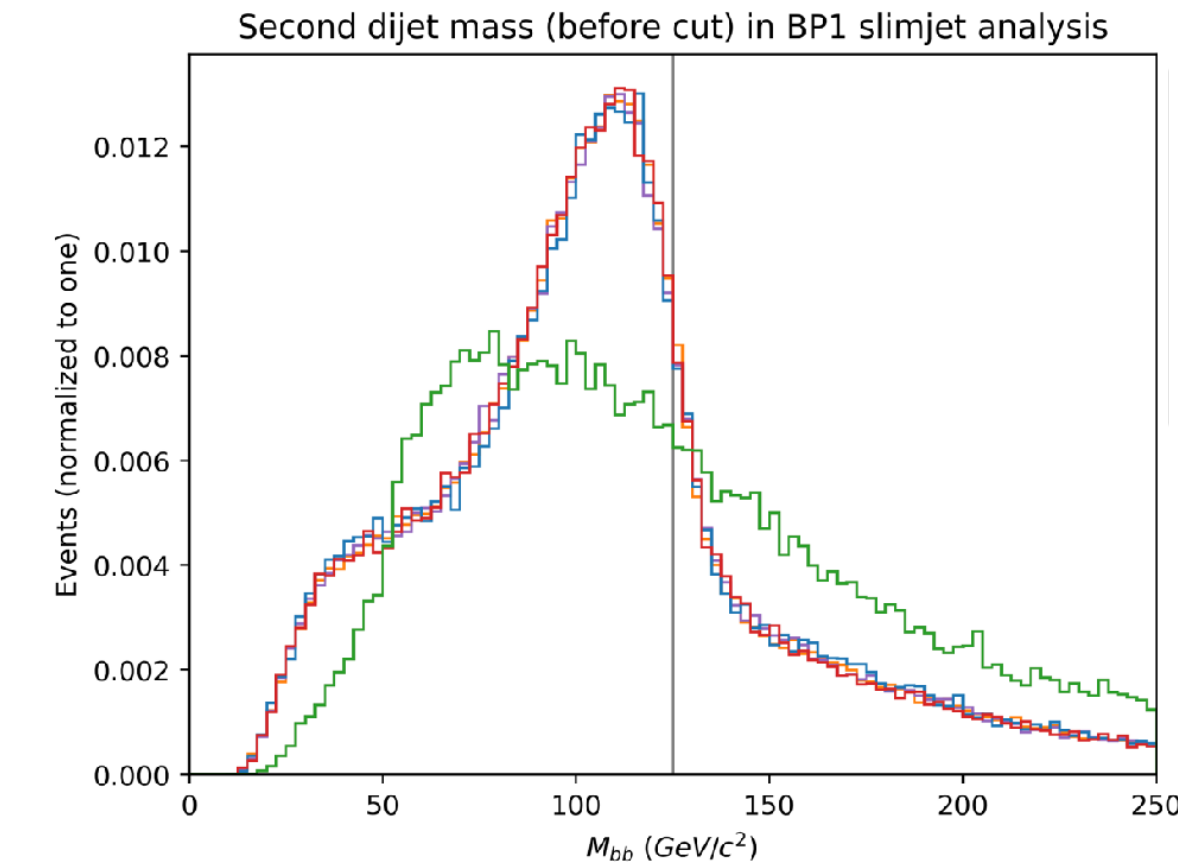
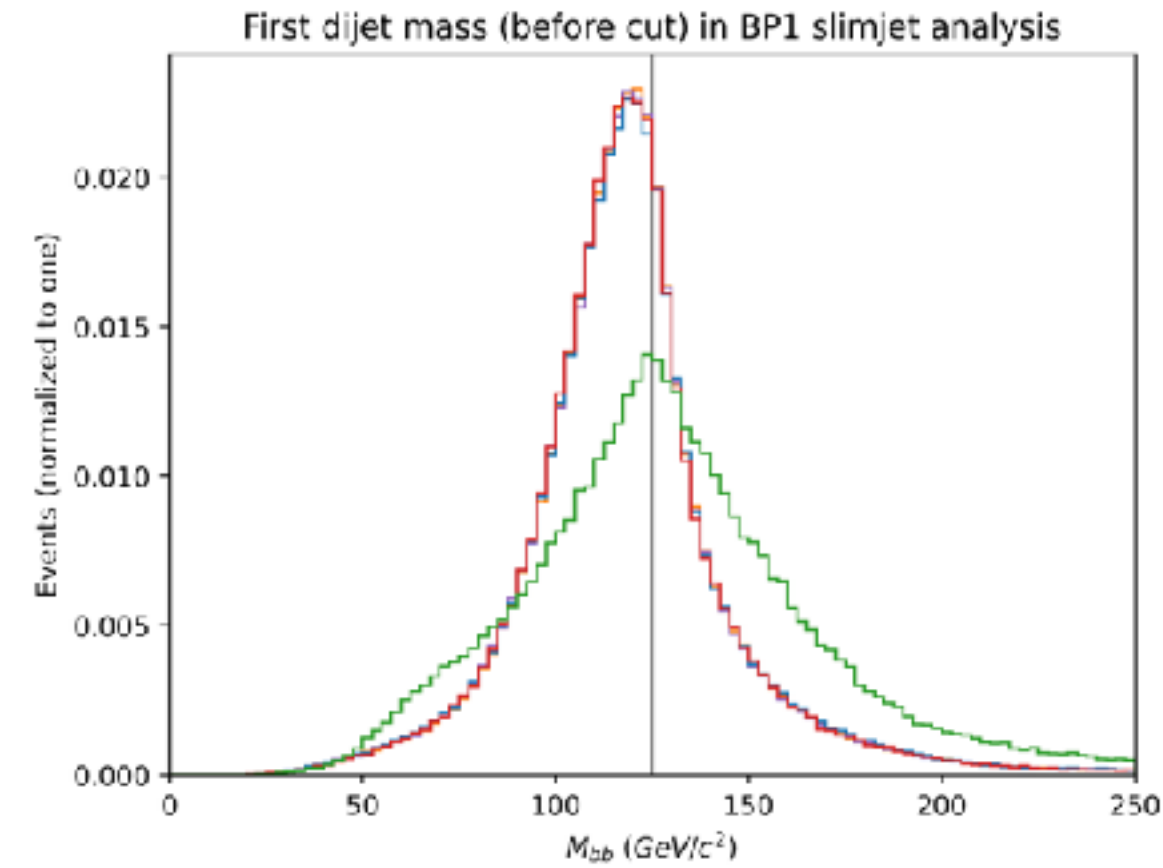
	$\sigma_{\text{before decay}}$ (fb)	$\sigma_{\text{after decay}}$ (fb)
SM	18.869	5.492
MSSM	22.976	6.697

	slimjet
R	0.4
minR	0.2
maxR	0.4
ρ	100



BP1 — non-resonant slimjet

Kinematic cut
(including QCD background)



	THT	Dijet 1 mass	Dijet 1 pT	Dijet 2 mass	Dijet 2 pT
If $N(b) > 2$	$THT > 300$	$110 < m < 140$	$200 < pT$	—	—
If $N(b) > 3$	$THT > 300$	$110 < m < 140$	$200 < pT$	$m < 150$	$200 < pT$

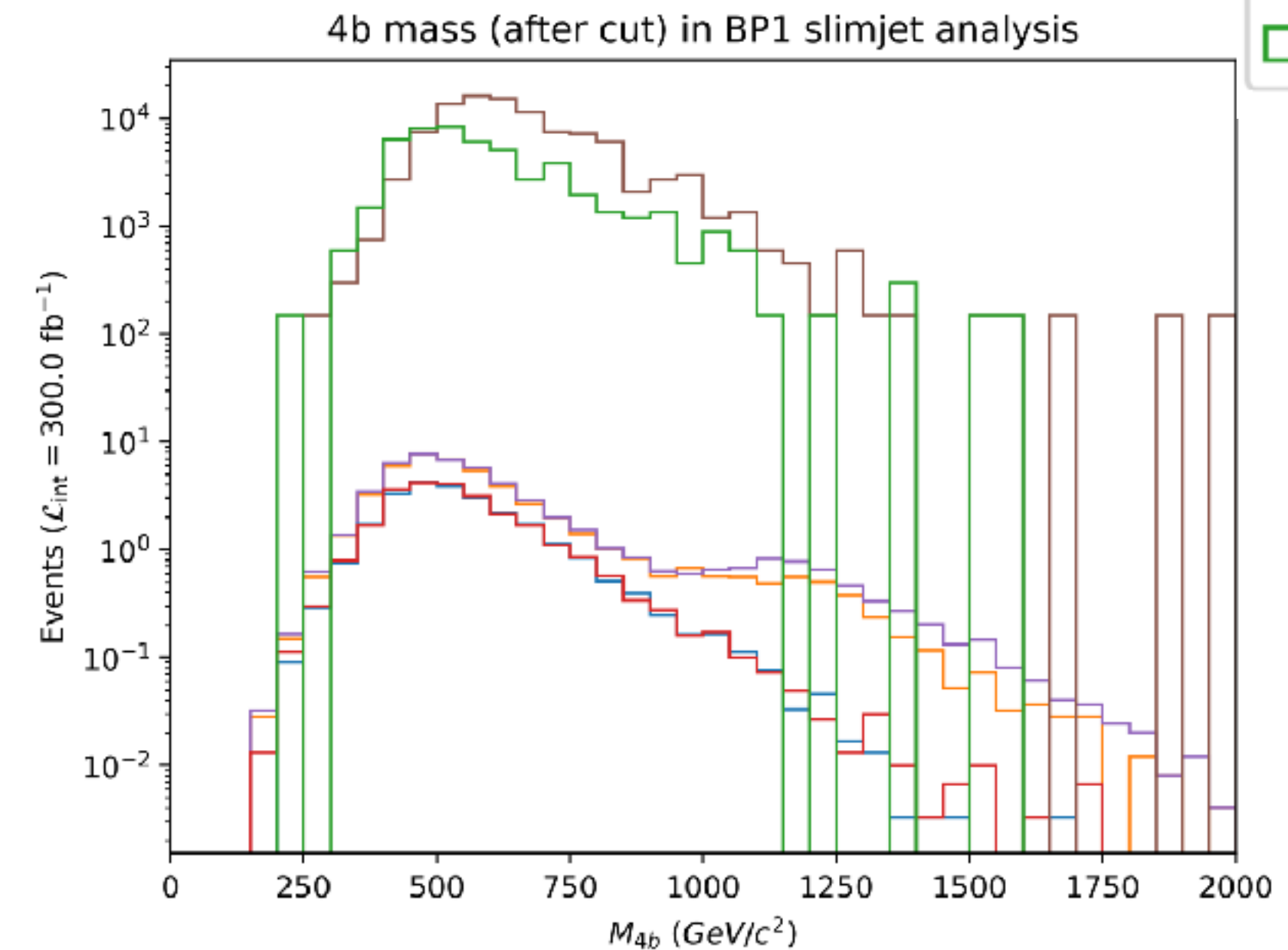
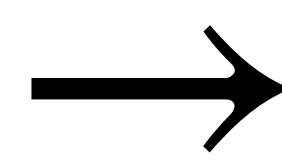
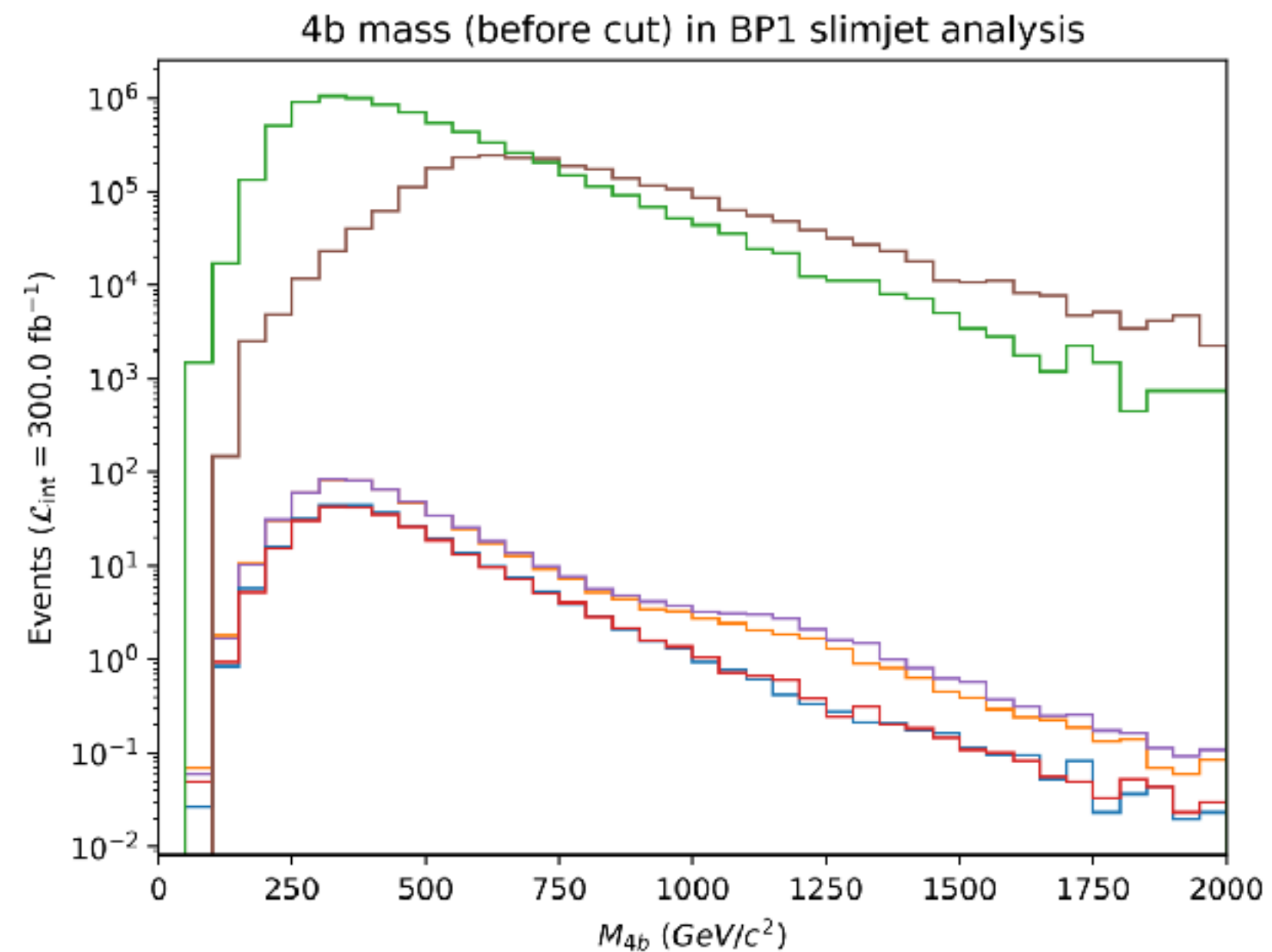
BP1 — non-resonant slimjet (4b)

Applying the cut

SR: $M_{4b} > 800 \text{ GeV}$

$\sigma_{\text{QCD}} = 45 \text{ pb}$

- Non-resonant case
 $m_{\tilde{t}_1} = 600 \text{ GeV}$
- ▭ MSSM antikt
 - ▭ MSSM VR
 - ▭ SM antikt
 - ▭ SM VR
 - ▭ QCD 4b antikt
 - ▭ QCD 4b VR



SR events before cut (4b)	anti-kT	Variable R
BSM di-Higgs	33	40
SM di-Higgs	13	13
Signal (BSM-SM)	20	27
QCD 4b	994459	519268
$S/(\sqrt{B})$	0.0204	0.0380

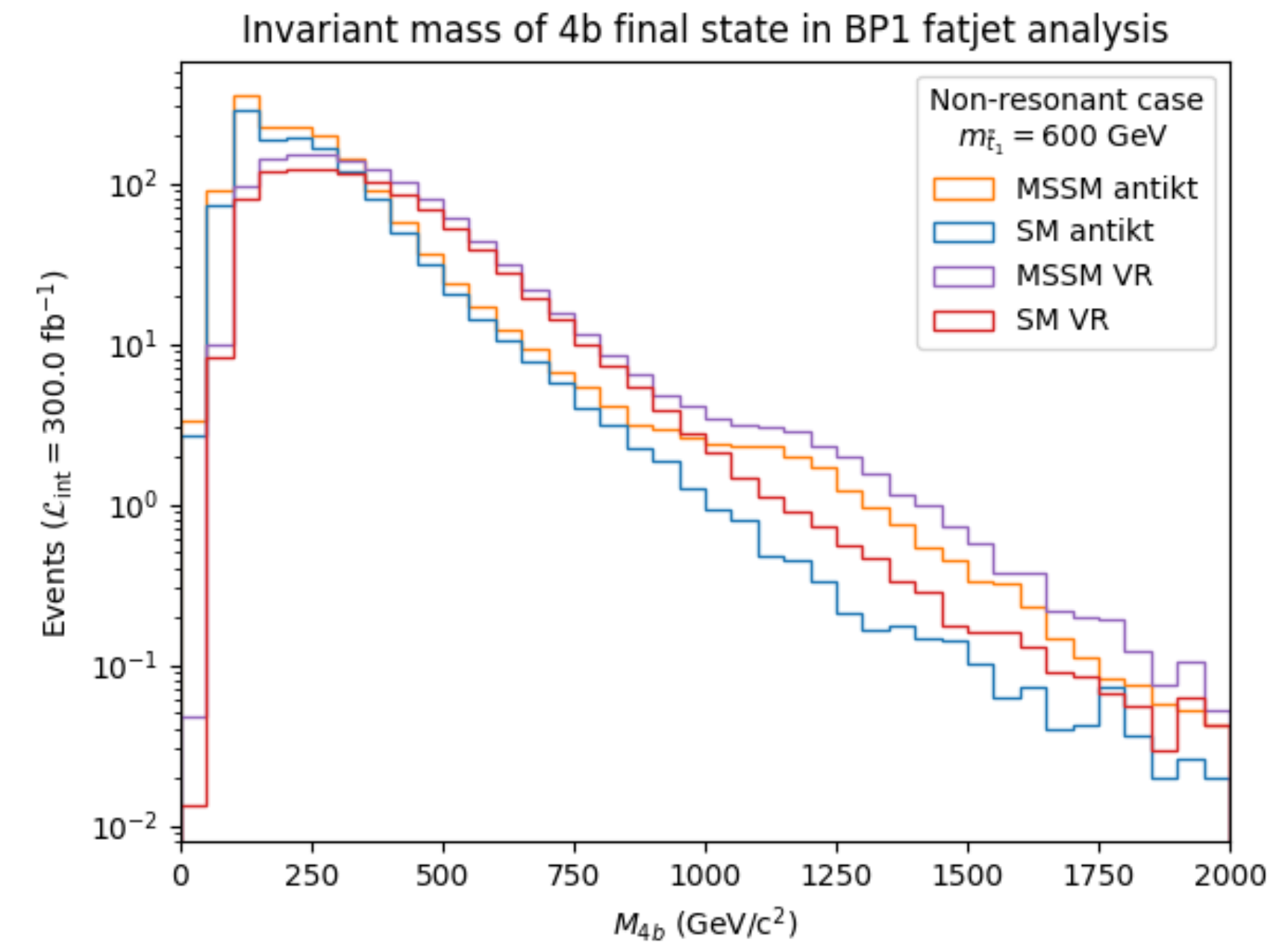
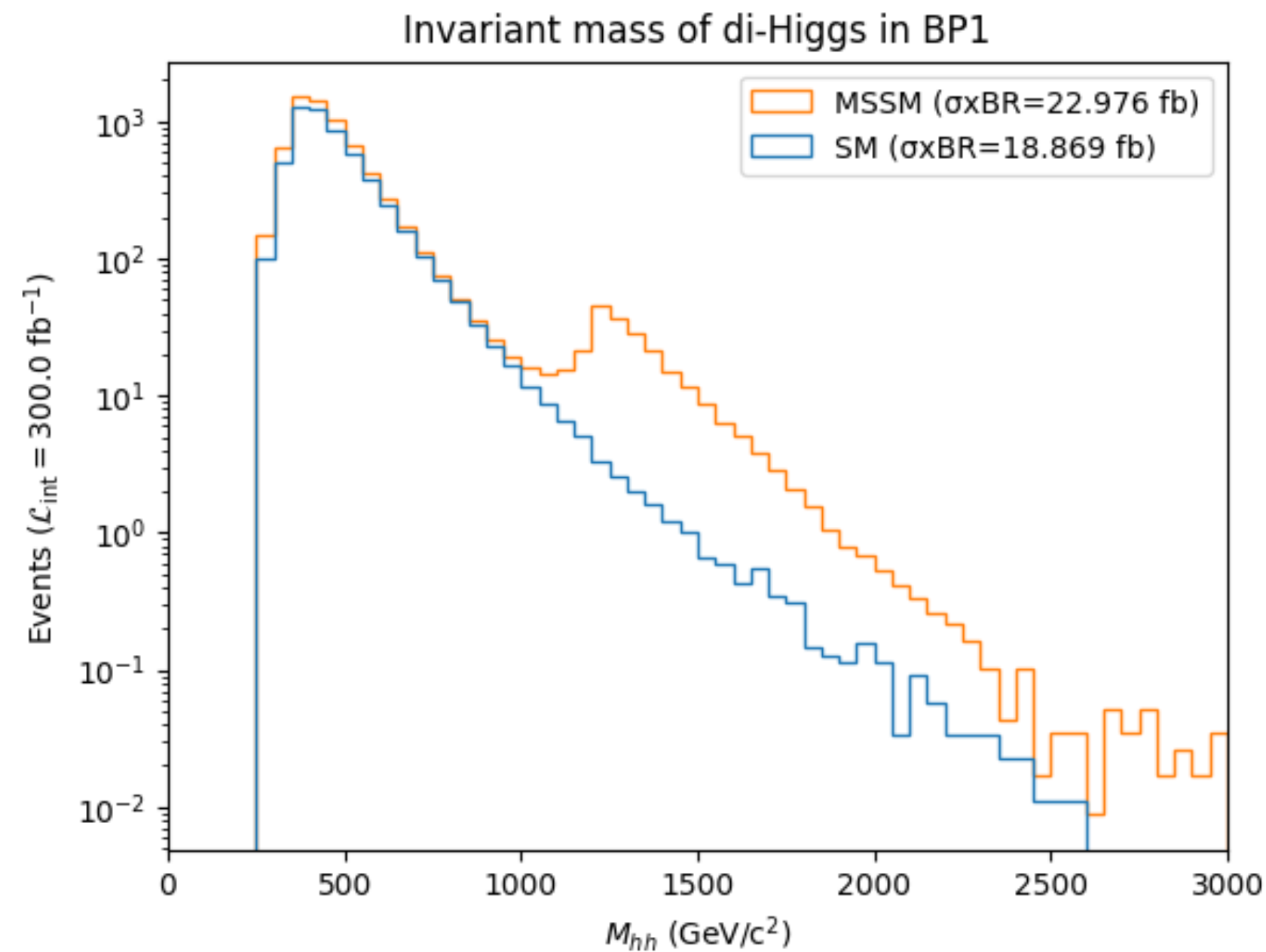
SR events after cut (4b)	anti-kT	Variable R
BSM di-Higgs	7	8
SM di-Higgs	2	2
Signal (BSM-SM)	5	6
QCD 4b	19144	51120
$S/(\sqrt{B})$	0.0367	0.0813

BP1 — non-resonant fatjet

Jet optimization

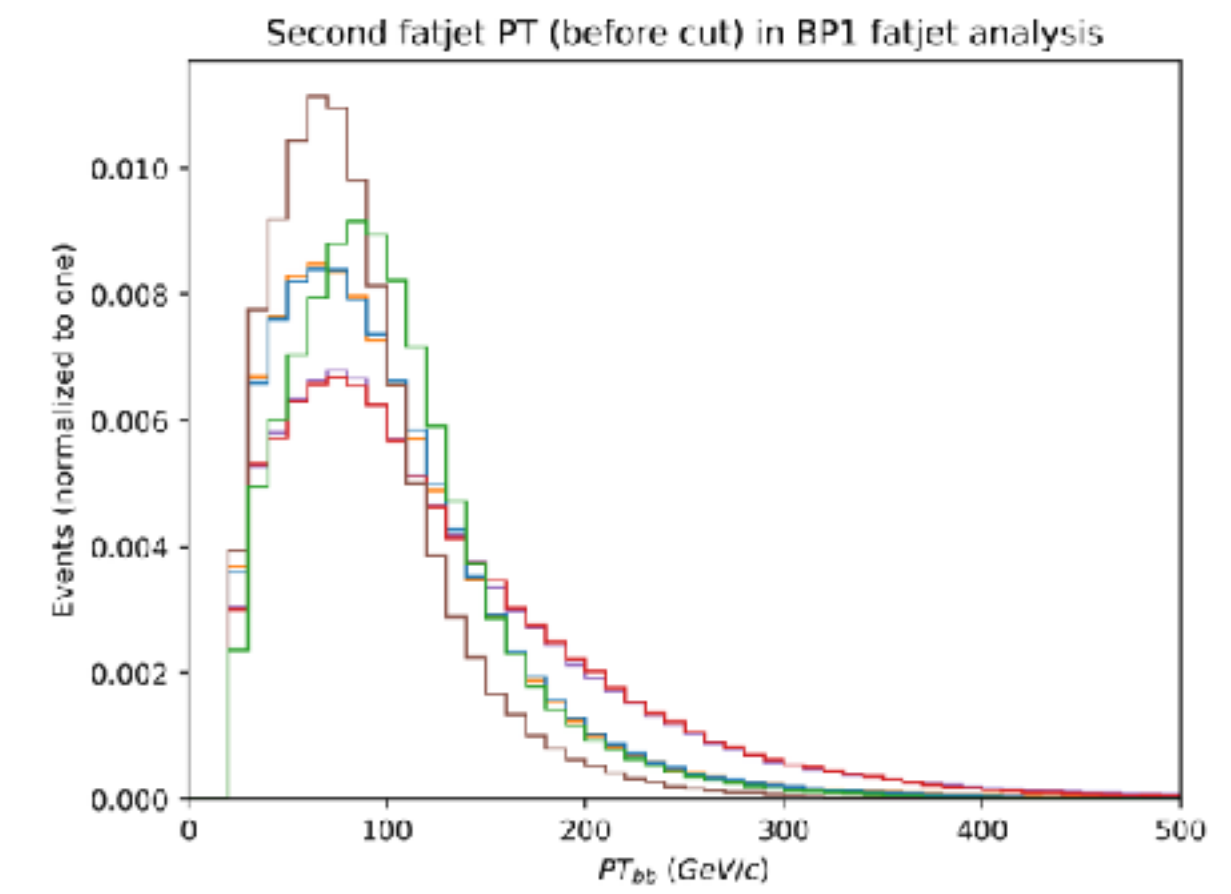
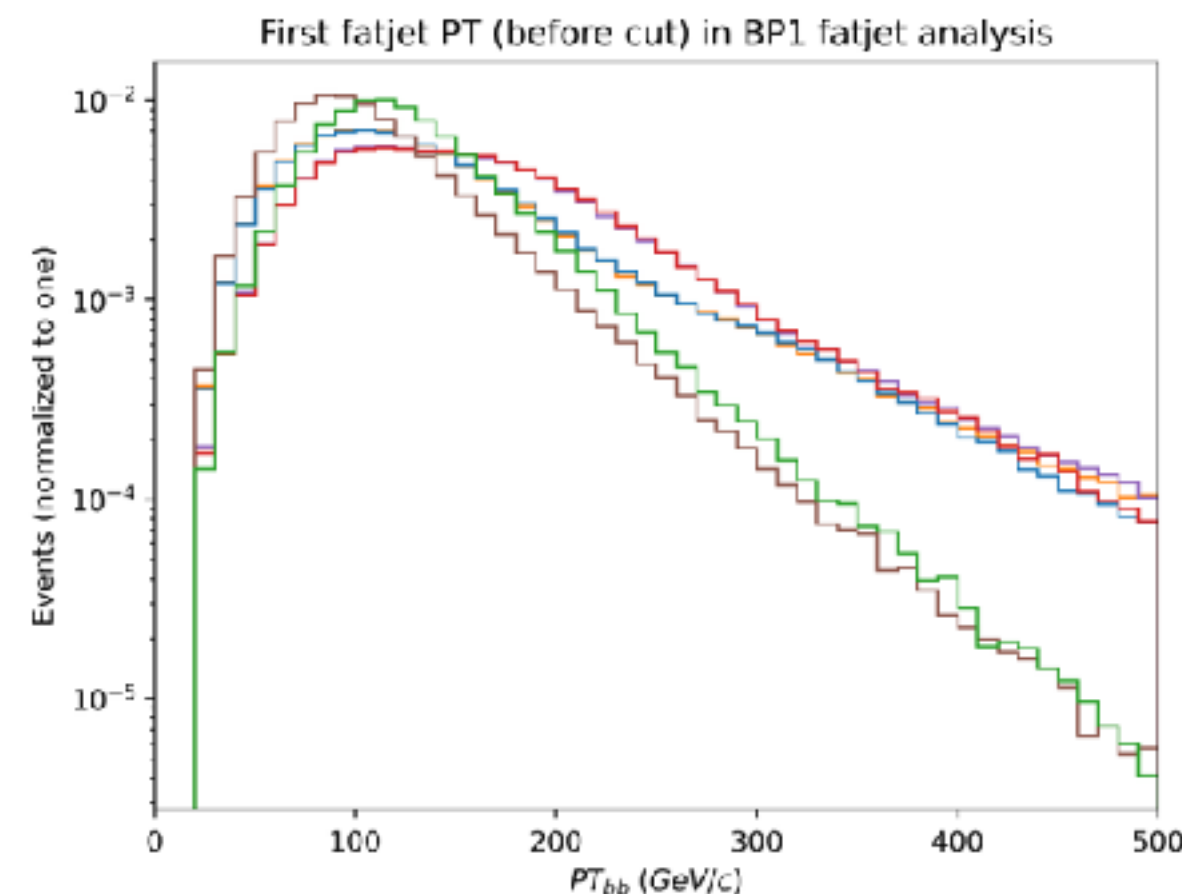
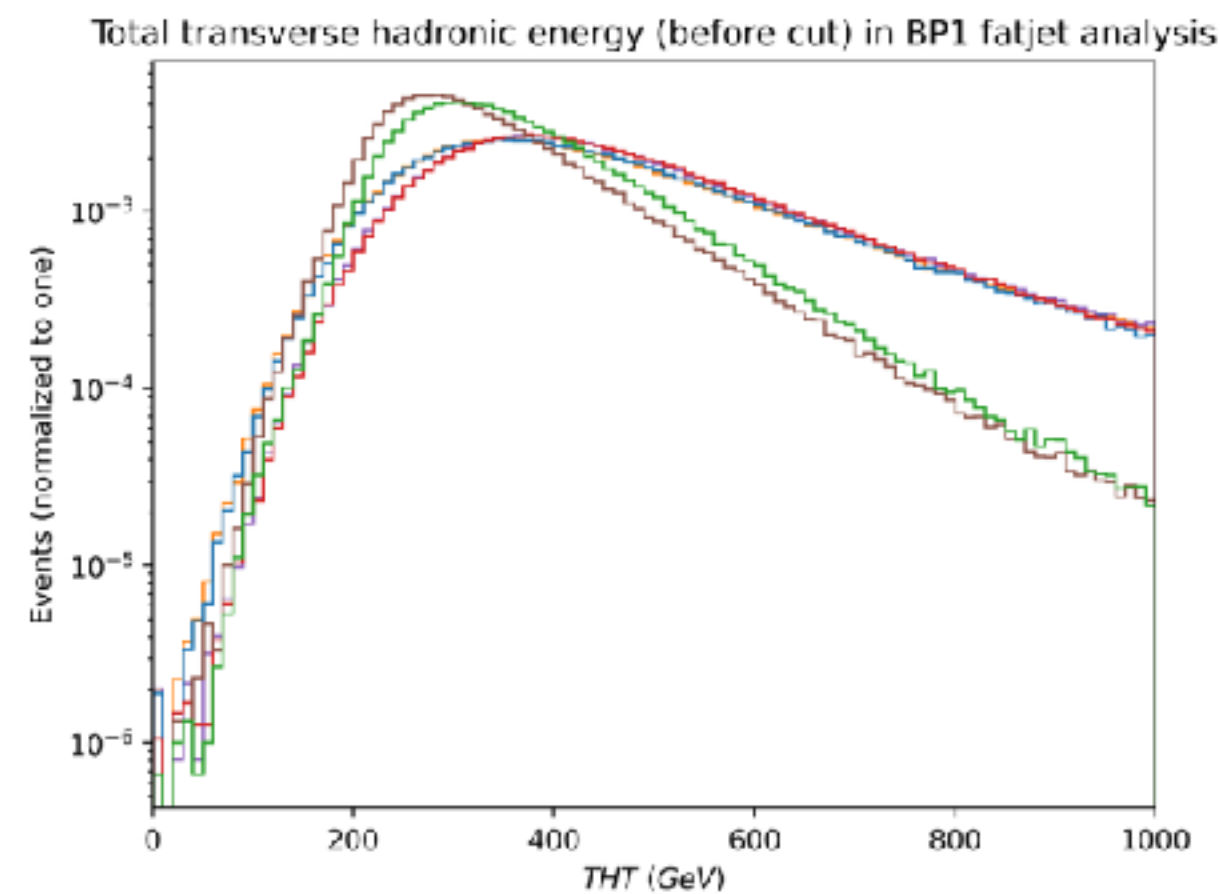
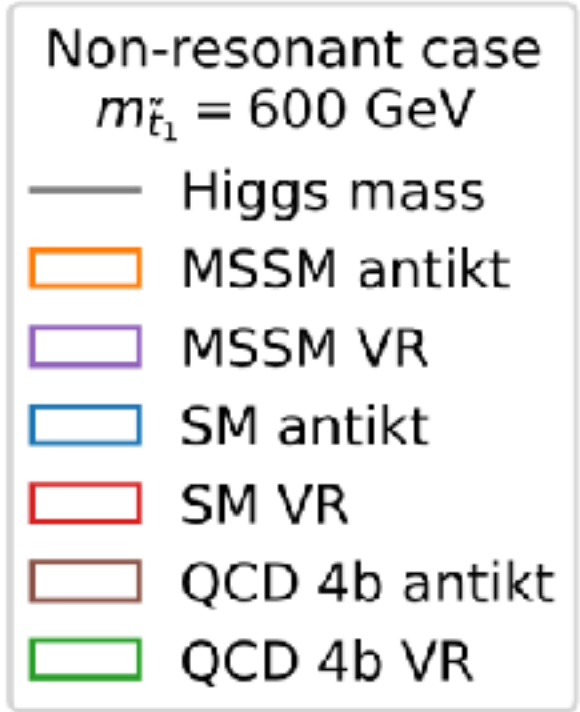
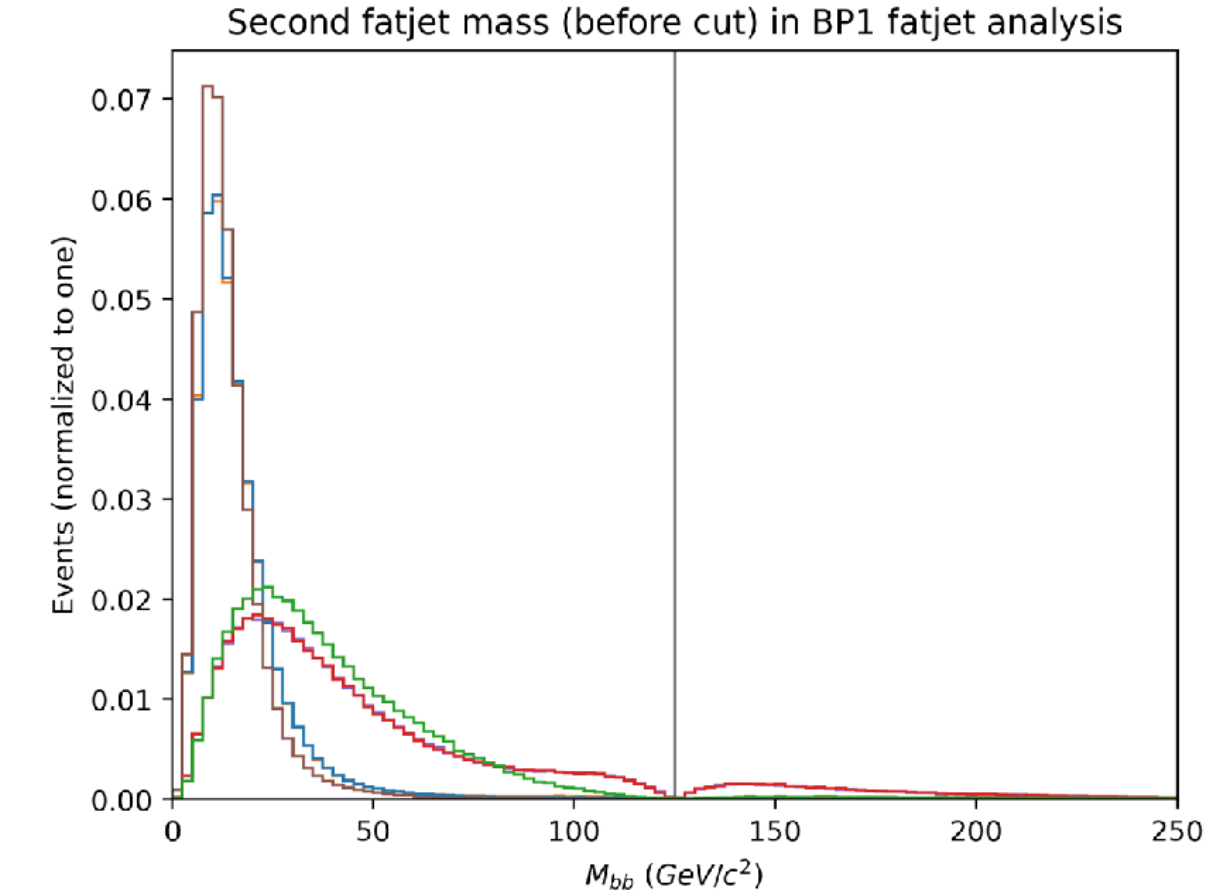
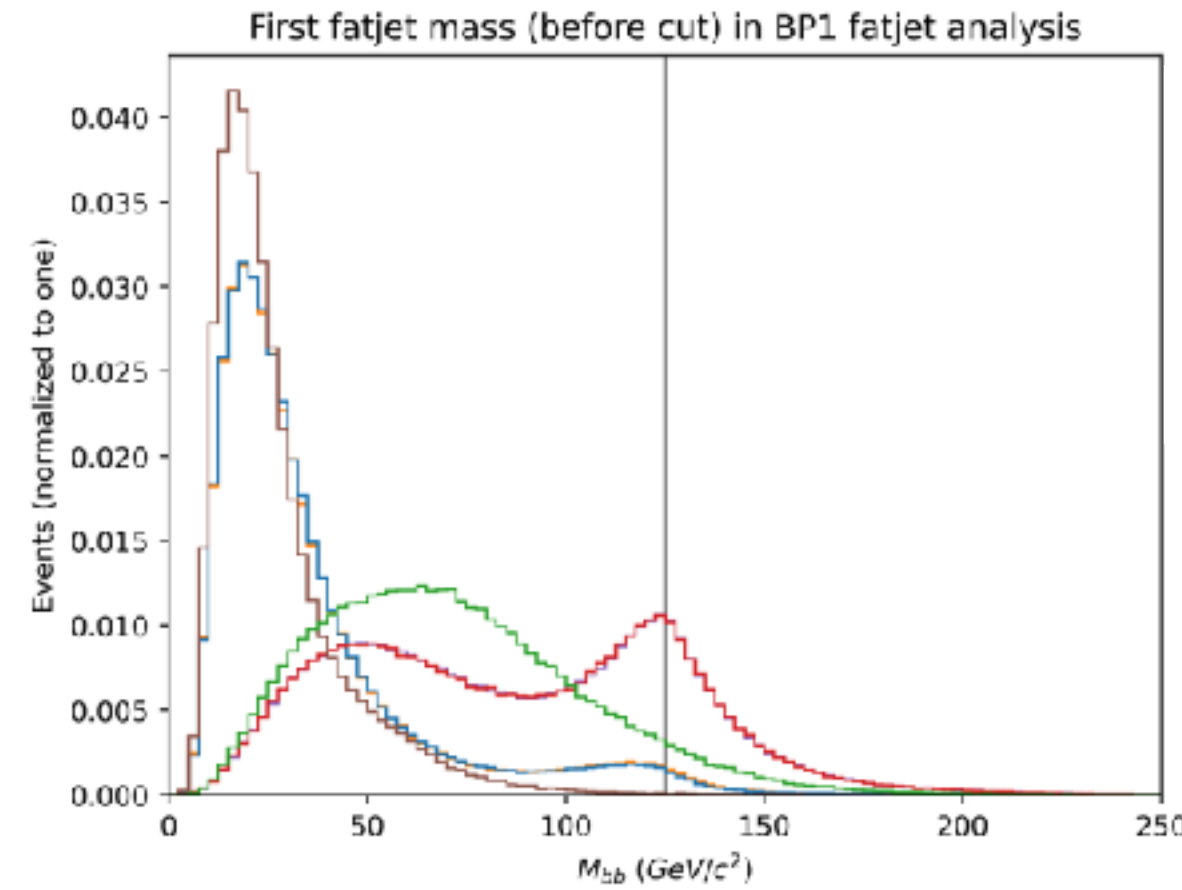
	$\sigma_{\text{before decay}}$ (fb)	$\sigma_{\text{after decay}}$ (fb)
SM	18.869	5.492
MSSM	22.976	6.697

	fatjet
R	0.8
minR	0.6
maxR	1.4
ρ	450



BP1 — non-resonant fatjet

Kinematic cut
(including QCD background)



	THT	Fatjet 1 mass	Fatjet 1 pT	Fatjet 2 mass	Fatjet 2 pT
If $N(b) > 1$	$THT > 300$	$90 < m < 200$	—	$90 < m < 200$	—

BP1 — non-resonant fatjet

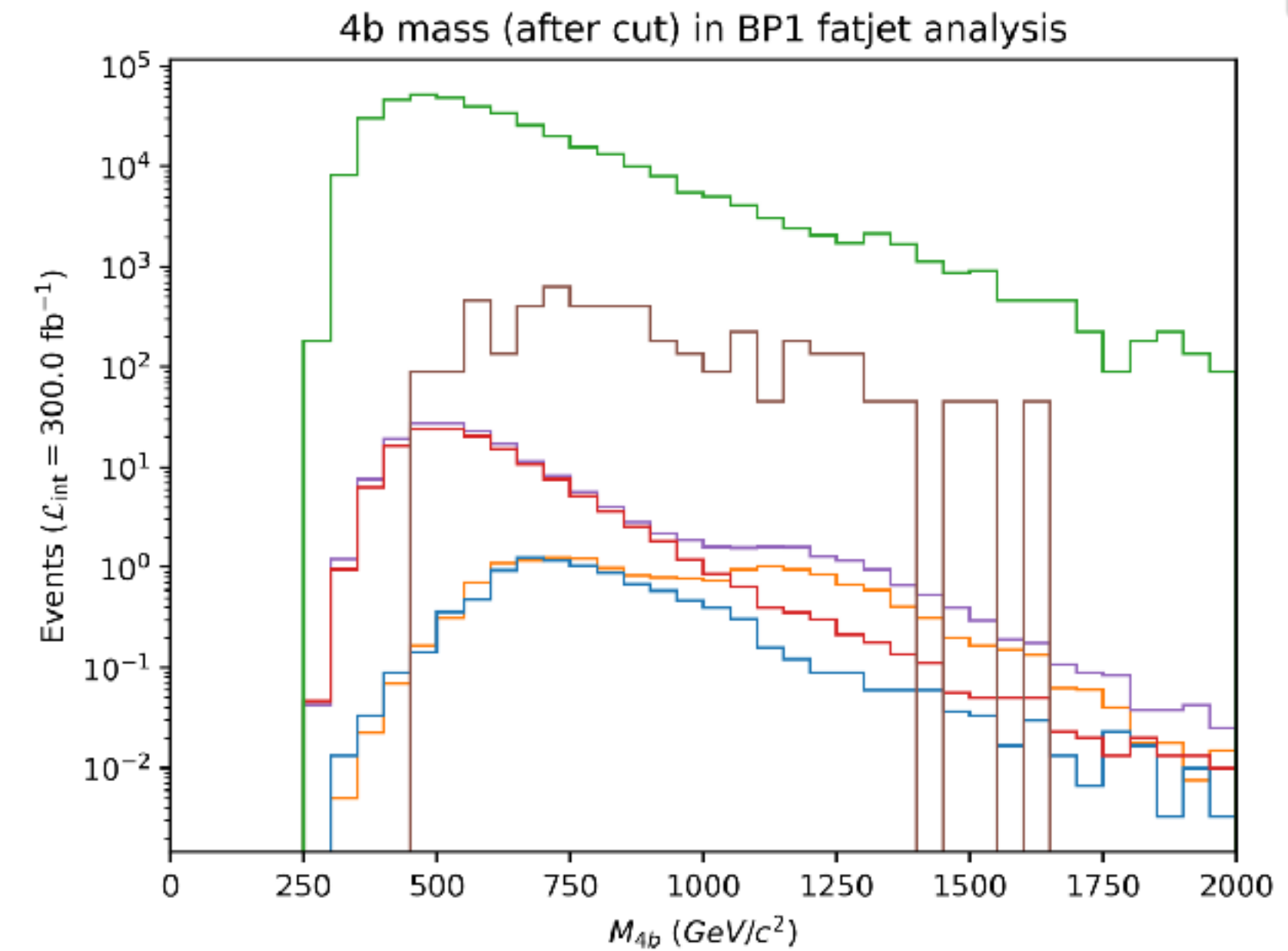
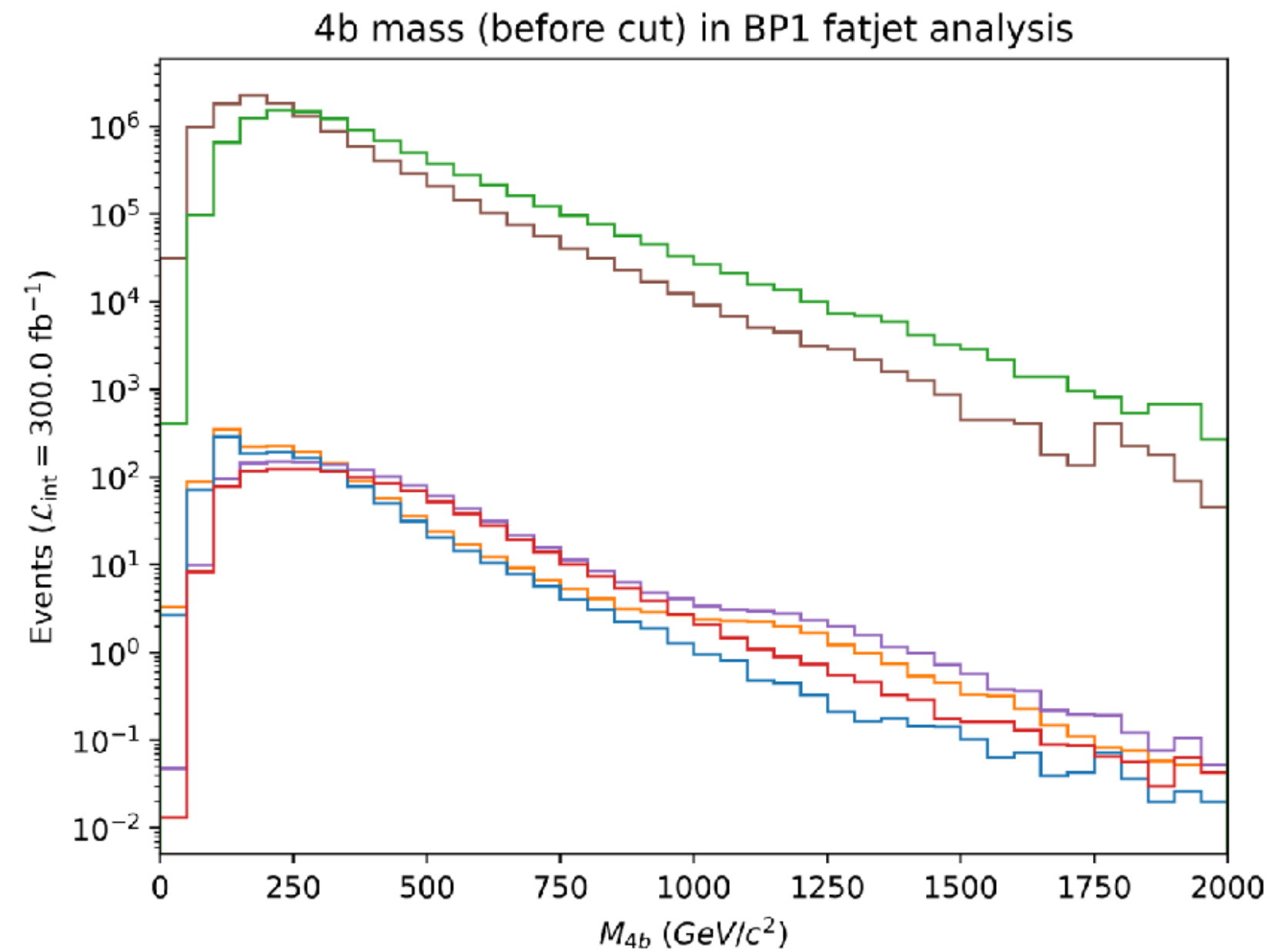
Applying the cut

SR: $M_{4b} > 800 \text{ GeV}$

$\sigma_{\text{QCD}} = 45 \text{ pb}$

Non-resonant case
 $m_{\tilde{t}_1} = 600 \text{ GeV}$

- ▭ MSSM anti-kt
- ▭ MSSM VR
- ▭ SM anti-kt
- ▭ SM VR
- ▭ QCD 4b anti-kt
- ▭ QCD 4b VR



SR events before cut (4b)	anti-kT	Variable R
BSM di-Higgs	29	47
SM di-Higgs	13	28
Signal (BSM-SM)	16	19
QCD 4b	123941	337937
$S/(\sqrt{B})$	0.0455	0.0325

SR events after cut (4b)	anti-kT	Variable R
BSM di-Higgs	11	23
SM di-Higgs	4	13
Signal (BSM-SM)	7	10
QCD 4b	2175	64371
$S/(\sqrt{B})$	0.141	0.0421