Anisotropic Hydrodynamics from Gauge/Gravity Duality

Hansjörg Zeller

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based on arXiv:1011.5912 in collaboration with J. Erdmenger and P. Kerner

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Gauge/Gravity Duality:

- novel method to study strongly coupled theories
- simple calculations to determine the dynamics of these systems
- prediction $\frac{\eta}{s} = \frac{1}{4\pi}$, good agreement with measurements of Quark-Gluon-Plasma (RHIC)

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Anisotropic Systems:

- more than one distinct shear mode
- deviation of $\frac{\eta}{s}$ from $\frac{1}{4\pi}$ in holographic systems \Rightarrow step towards applying gauge/gravity duality to real-world systems

- Gauge/Gravity Duality
- 2 Hydrodynamics an Overview
- 3 Hydrodynamics in a Holographic Setup with Anisotropic Spacetime
- 4 Conclusion

Gauge/Gravity Duality

- classical gravity in asymptotically AdS space
- dual field theory on the conformally flat boundary of the AdS space



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strongly coupled \Leftrightarrow weakly coupled	
finite temperature \Leftrightarrow black holefinite chemical potential \Leftrightarrow $\langle A_t \rangle \neq 0$ spontaneous breaking of global symmetries \Leftrightarrow hairy black holes	

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- constitutive equations

$$T^{\mu\nu} = T^{\mu\nu}_{eq.} + \Pi^{\mu\nu}$$
 and $J^{\mu} = J^{\mu}_{eq.} + \Upsilon^{\mu}$
with $\Pi_{ij} \sim \eta \left(\partial_i u_j + \partial_j u_i \right)$

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- isotropic systems:
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- transversely isotropic systems:
 2 independent shear viscosities



What is the Shear Viscosity?



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• using $s \sim T^3$

$$\Rightarrow \frac{\eta}{s} \sim \frac{1}{\lambda^2}$$

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$$rac{\eta}{s} = rac{1}{4\pi} \simeq 0.079$$



 λ

Mateos

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 - conjecture by Kovtun, Son and Starinets:

$$\frac{\eta}{s} \geq \frac{1}{4\pi}$$

for all substances found in Nature

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Son et al.

- setup: *SU*(2) Einstein-Yang-Mills theory in (4 + 1)-dimensional asymptotically *AdS* space
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Erdmenger et al.

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- \Rightarrow rotational symmetry spontaneously broken

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- *SO*(2) rotational symmetry in *yz*-plane (i.e. transversely isotropic background)
- \Rightarrow 2 distinct shear viscosities: η_{yz} and η_{xy}



colour coding: $\frac{\eta_{yz}}{s} = \frac{1}{4\pi}$ and $\frac{\eta_{xy}}{s}$ depending on the amount of symmetry breaking (blue > red > green)

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Conductivity

 $J^{y} = \sigma^{yy} E_{y}$



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Outlook:

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- \Rightarrow enables us to write down the field theory
 - zero temperature behaviour of transport coefficients

Thank you!

• SU(2) Einstein-Yang-Mills theory:

$$S = \frac{1}{2\kappa_5^2} \int d^5 x \left[R - \Lambda - \frac{\alpha^2}{2} F^a_{MN} F^{aMN} \right]$$

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• metric ansatz:

$$ds^{2} = -N(r)\sigma(r)^{2}dt^{2} + \frac{r^{2}}{f(r)^{4}}dx^{2} + r^{2}f(r)^{2}(dy^{2} + dz^{2}) + \frac{1}{N(r)}dr^{2}$$

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• equations of motion:

$$\begin{aligned} R_{MN} - 4g_{MN} &= \alpha^2 \left(T_{MN} - \frac{1}{3} T_P^P g_{MN} \right) \,, \\ \text{with:} \quad T_{MN} &= F_{PM}^a F^a {}^P_N - \frac{1}{4} F_{PQ}^a F^a {}^{PQ} g_{MN} \\ \nabla_M F^{a \, MN} &= -\varepsilon^{abc} A_M^b F^{c \, MN} \end{aligned}$$

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• real-time retarded Green's functions

$$\begin{split} G^{\mu\nu,\rho\sigma}(\omega,\vec{k}) &= \int \mathrm{d}^{4}x \, \mathrm{e}^{-ik_{\mu}x^{\mu}}\theta(t) \, \langle [T^{\mu\nu}(t,\vec{x}), T^{\rho\sigma}(0)] \rangle \\ G^{\mu\nu,i\rho}(\omega,\vec{k}) &= \int \mathrm{d}^{4}x \, \mathrm{e}^{-ik_{\mu}x^{\mu}}\theta(t) \, \langle [T^{\mu\nu}(t,\vec{x}), J^{i\rho}(0)] \rangle \\ G^{i\mu,j\nu}(\omega,\vec{k}) &= \int \mathrm{d}^{4}x \, \mathrm{e}^{-ik_{\mu}x^{\mu}}\theta(t) \, \langle [J^{i\mu}(t,\vec{x}), J^{j\nu}(0)] \rangle \end{split}$$

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• transport coefficients using Kubo formulas:

$$\eta_{xy} = -\lim_{\omega \to 0} \frac{1}{\omega} \operatorname{Im} \left(G^{xy, xy}(\omega, 0) \right)$$
$$\sigma_{yy} = -i \frac{G^{3y, 3y}}{\omega}$$