



**Taking String
Theory to the
Standard
Model and
Beyond**

Esben
Mølgaard

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String theory
D-Branes

My Work

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Gauge groups in
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The model

Beyond

Conclusions

Taking String Theory to the Standard Model and Beyond

Esben Mølgaard

Niels Bohr Institute
University of Copenhagen

May 2, 2011



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Massless open string states



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NS-sector



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Massless open string states

NS-sector Vector boson $b_{-1/2}^{\mu}|p\rangle_{NS}$
R-sector



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NS-sector	Vector boson	$b_{-1/2}^{\mu} p\rangle_{NS}$
R-sector	Chiral spinor	$ p\rangle_R$



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NS-sector	Vector boson	$b_{-1/2}^{\mu} p\rangle_{NS}$
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The theory is spacetime supersymmetric.



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The theory is spacetime supersymmetric.
It is consistent only in 10 dimensions.



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The theory is spacetime supersymmetric.

It is consistent only in 10 dimensions.

The effective dimension can be reduced through compactification.



D-branes

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- Introduced through T-duality



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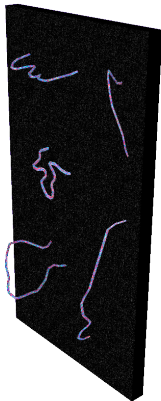
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- Introduced through T-duality
- Objects on which Dirichlet strings are stuck





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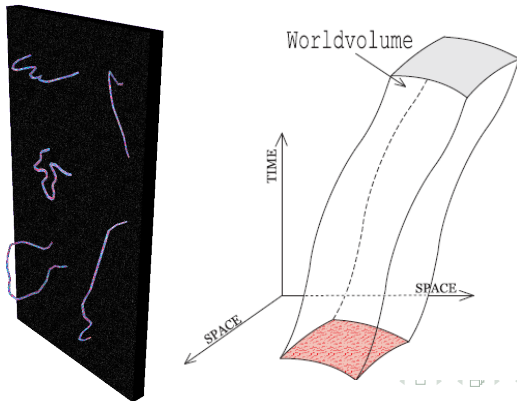
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- Introduced through T-duality
- Objects on which Dirichlet strings are stuck
- Dynamical objects in their own right





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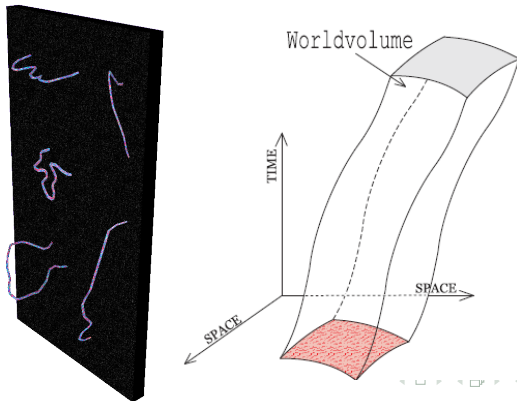
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- Introduced through T-duality
- Objects on which Dirichlet strings are stuck
- Dynamical objects in their own right
- Many possible configurations





Changes from magnetization

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- Single string stretched between two D-branes that are magnetized in the 6 compact directions.



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NS-Sector

R-sector



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Ground state mass

$$M^2 = \frac{1}{2\alpha'} \sum_{a=1}^3 \nu_a$$

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10d chirality broken



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4d chirality remains



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10d chirality broken

4d chirality remains

Supersymmetry is broken



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Chan-Paton indices give $U(N)$ symmetry



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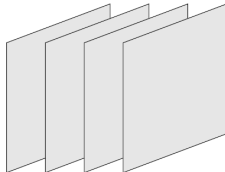
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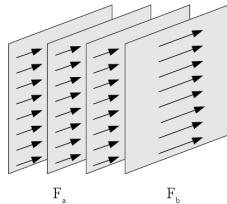
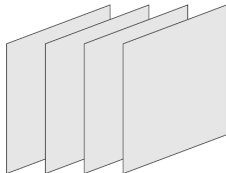
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Chan-Paton indices give $U(N)$ symmetry



Different magnetizations breaks this to $U(K) \times U(L)$.



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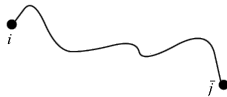
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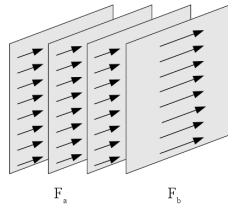
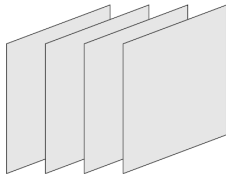
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Chan-Paton indices give $U(N)$ symmetry



Different magnetizations breaks this to $U(K) \times U(L)$.
R-strings transform in the bifundamental representation $(\mathbf{K}, \bar{\mathbf{L}})$.



The Standard Model reproduced

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To get all the particles of the standard model, we need 4 stacks



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To get all the particles of the standard model, we need 4 stacks

$$U(3)_{\text{baryonic}} \times U(2)_{\text{left}} \times U(1)_{\text{right}} \times U(1)_{\text{leptonic}}$$



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Magnetization gives rise to Landau levels and thus generations

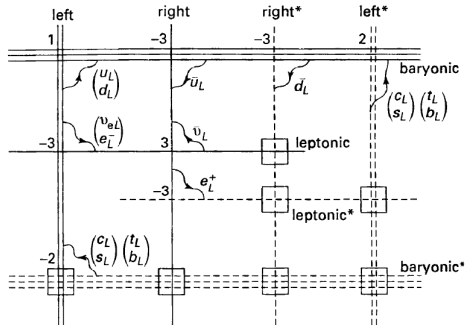


The Standard Model reproduced

To get all the particles of the standard model, we need 4 stacks

$$U(3)_{\text{baryonic}} \times U(2)_{\text{left}} \times U(1)_{\text{right}} \times U(1)_{\text{leptonic}}$$

Magnetization gives rise to Landau levels and thus generations
In the T-dual picture, these are intersection numbers,





Beyond the Standard Model

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- Dark matter candidate in sterile neutrinos.



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- Dark matter candidate in sterile neutrinos.
- $\mathcal{N} = 1$ supersymmetry between scalars and spinors



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- Dark matter candidate in sterile neutrinos.
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- $B_{1/2-\nu_a}^{a\dagger} |p\rangle_{NS}, \quad M^2 = \frac{1}{\alpha'} \left[\frac{1}{2} \sum_{b=1}^3 \nu_b - \nu_a \right]$



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- Technicolor through a new stack of branes



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- Technicolor through a new stack of branes
- $U(N)_{TC} \times SU(3) \times SU(2) \times U(1)$



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- $SU(3) \times SU(2)_{left} \times SU(2)_{right} \times U(1)$



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Using the proper brane configuration and dimensional compactification, it is possible to construct something that is very close to the Standard Model.



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Using the proper brane configuration and dimensional compactification, it is possible to construct something that is very close to the Standard Model.

The model in question makes testable predictions that go beyond what it was designed to do.



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The model in question makes testable predictions that go beyond what it was designed to do.

It is possible to extend the model to go even further beyond the Standard Model.