Flux compactifications and geometry in string theory

Peter Patalong

String theory group, Arnold-Sommerfeld-Center

July 26th, 2011

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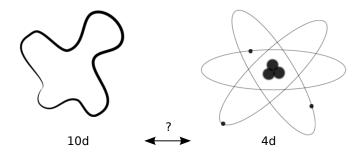
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Geometry Flux Compactification

Geometry?

A problem...



... and its meaning

Structure of space-time?

- "Ordinary" space-time for the 4-dimensional part
- "Small" space for the 6-dimensional remain

... and its meaning

Structure of space-time?

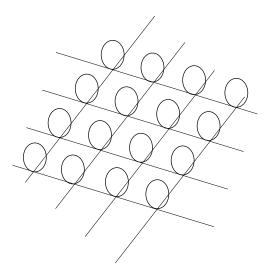
- "Ordinary" space-time for the 4-dimensional part
- "Small" space for the 6-dimensional remain

Mathematically: a manifold. What kind of manifold?

First guess: a product

 $M_{10} = M_4 \times M_6$

Example: $\mathbb{R}^2 \times S^1$



The geometry of M_4

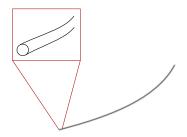
- General relativity: a four-dimensional, smooth, connected Lorentzian manifold
- Cosmology: homogeneous and isotropic universe
 - \Rightarrow Maximally symmetric space
 - $\Rightarrow {\sf Constant\ curvature}$
 - \Rightarrow M_4 is either Minkowski, AdS or dS

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Ok, let's do that!

- Find a solution to the theory:
 - Solve the equations of motion
 - Here: type II string theory in its low-energy approximation (supergravity)
 - Take the easiest setup: a solution that gives 4d vacuum
 - ▶ Expectation value of the fermion fields: < (fermion) >= 0
- Preserve supersymmetry:

$$\delta_{\varepsilon}(\text{fermion}) = (\text{boson}) \stackrel{!}{=} 0$$

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► Supersymmetry ⇒ equations of motion (luckily!)

Equations, equations!

Variation of the gravitino field:

$$\delta_{\varepsilon}\Psi_{\mu} = \nabla_{\mu}\varepsilon \stackrel{!}{=} 0$$

Can be decomposed:

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Equations, equations!

Variation of the gravitino field:

$$\delta_{\varepsilon}\Psi_{\mu}=\nabla_{\mu}\varepsilon\stackrel{!}{=}0$$

- Can be decomposed:
 - On M_4 : $\mathcal{R}_4 = 0$
 - On M_6 : $\nabla_m \eta = 0$
 - \Rightarrow Internal manifold has covariantly constant spinor
 - \Rightarrow Strong geometrical constraint: reduced holonomy

Holonomy?

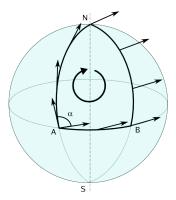


Image: Wikimedia Commons

The result: a Calabi-Yau manifold

- Holonomy reduced to SU(3), Ricci-flat ($Ric_{\mu\nu} = 0$)
- Examples: Torus T^2 (in 2d), T^4 and K3 (in 4d)

Geometry Flux Compactification

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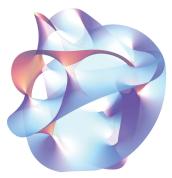


Image: Wikimedia Commons

Not good

A Calabi-Yau geometry is not satisfactory:

- ► *M*₄ is bound to be Minkowski.
- There are moduli.
- There is too much supersymmetry.

Anything else?

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- Recall the first guess: a product manifold $M_{10} = M_4 \times M_6$
- Can we allow for more? Yes, a warped product!

$$\mathrm{d}s^2 = e^{2A(y)}g_{\mu\nu}\mathrm{d}x^{\mu}\mathrm{d}x^{\nu} + g_{mn}\mathrm{d}y^m\mathrm{d}y^n$$

Size of M_4 depends on the position in M_6 .

Susy does not like warping

• The internal components of $abla_{\mu} \varepsilon \stackrel{!}{=} 0$ become

$$k\mathcal{R}_4+(\nabla A)^2=0$$

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• M_6 compact $\Rightarrow \nabla A = 0$ \Rightarrow no warping, M_4 again Minkowski. Geometry Flux Compactification

Flux?

Type II closed string spectrum:

$$\begin{array}{ll} \text{NS-NS} & \mathbf{8}_{v}\otimes\mathbf{8}_{v}=\mathbf{35}\oplus\mathbf{28}\oplus\mathbf{1} & g_{\mu\nu}, \ \mathbf{B}_{\mu\nu}, \ \Phi \\ \text{R-R} & \mathbf{8}_{s}\otimes\mathbf{8}_{s/c}=\text{p-forms} & \text{RR-fields} \\ \text{NS-R} & \mathbf{8}_{v}\otimes\mathbf{8}_{s/c}=\mathbf{8}_{s/c}\oplus\mathbf{56}_{s/c} & \Psi_{\mu}, \ \lambda \\ \text{R-NS} & \mathbf{8}_{s}\otimes\mathbf{8}_{v}=\mathbf{8}_{s}\oplus\mathbf{56}_{s} & \Psi_{\mu}', \ \lambda' \end{array}$$

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RR-fluxes:

Tensor decomposition:

$$\begin{aligned} \mathbf{8}_s \otimes \mathbf{8}_c &= C^{(1)} \oplus C^{(3)} & \text{IIA} \\ \mathbf{8}_s \otimes \mathbf{8}_s &= C^{(0)} \oplus C^{(2)} \oplus C^{(4)} & \text{IIB} \end{aligned}$$

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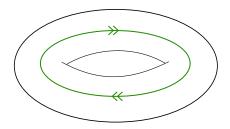
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• *H*-flux: H = dB

Flux in M_6

- Take electrodynamics as inspiration: field strength $F^{(2)}$
- Fluxes can wrap cycles in the internal manifold.



Fluxes and geometry I

- Flux lines \rightarrow energy density
- Decreasing volume \rightarrow increasing density \rightarrow "pressure"
- Mathematically: Derive equations of motion from action and find for example

$$\mathcal{R}_{6} + rac{1}{2}g_{s}^{2}|F_{2}|^{2} + rac{3}{2}\left(g_{s}^{2}|F_{0}|^{2} - |H|^{2}
ight) = 0$$

Fluxes and geometry II

Supersymmetry variation reloaded:

$$\delta_{\varepsilon}\Psi_{\mu} = \left(\nabla_{\mu} + \frac{1}{8}\mathcal{P}H_{\mu}\right)\varepsilon + \frac{1}{16}e^{\Phi}\sum_{n}\frac{1}{(2n)!}F_{\mu}\mathcal{P}_{n}\varepsilon$$

Fluxes and geometry II

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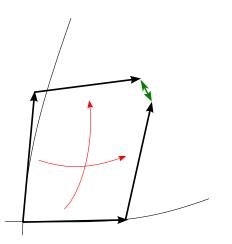
Decompose derivative of the spinor,

$$abla_{\mu} arepsilon = (ext{torsion classes})
eq 0$$

 \Rightarrow Fluxes induce torsion!

Geometry Flux Compactification

Torsion?



Torsion!

Classification of M_6 with respect to its torsion:

Torsion!

Classification of M_6 with respect to its torsion:

- Complex
- Symplectic
- Half-flat
- Special Hermitean
- Nearly Kähler
- Almost Kähler
- Kähler
- Conformal Calabi-Yau
- Calabi-Yau

Fluxes and geometry III

Geometry also affects the choice of the fluxes:

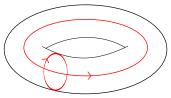
Dirac quantisation condition

$$\int_{\Sigma_p} F^{(p)} = N_p \in \mathbb{Z}$$

Values are determined by integrals over cycles

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 Number of different cycles is determined by topological features of M₆



Compactification?

$$\begin{aligned} \int d^{10}x \ (10d) & \to \int d^4x \int d^6x \ (4d) \cdot (6d) \\ &= (number) \cdot \int d^4x \ (4d) \end{aligned}$$

A type II example:

$$\mathcal{S}=rac{1}{2\kappa^2}\int \mathrm{d}^{10}x \; e^{-2\phi}\sqrt{|g|}\left(\mathcal{R}+4|\mathrm{d}\phi|^2-rac{1}{2}|\mathcal{H}|^2
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▶ Dilaton: $e^{-\phi} \rightarrow e^{-\phi^{(0)}} e^{-\varphi}$ and $\sigma = \rho^{3/2} e^{-\varphi}$

Integrating out the internal dimensions:

$$S = M_4^2 \int d^4 x \sqrt{|g_{\mu\nu}|} \left(\mathcal{R}_4 + kin + \sigma^{-2} \rho^{-1} \mathcal{R}_6 - \frac{1}{2} \sigma^{-2} \rho^{-3} |\mathcal{H}|^2 \right)$$

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 - make even more deformations possible and
 - give rise to masses for moduli.

Take-away message

Fluxes \rightarrow Internal geometry \rightarrow 4-dimensional theory