

Hot Field Theory

Outline

- 1) Motivation
- 2) Thermal Field Theory
 - 2.1) Vacuum Field Theory
 - 2.2) The Density Operator
 - 2.3) Green's Functions at finite T
 - 2.4) Imaginary Time Formalism
- 3) Hard Thermal Loops and Resummation
 - 3.1) The HTL Limit
 - 3.2) HTL Resummation and Quasiparticles
- 4) Applications: Physical Interpretation of Discontinuities

Previously

- TFT : average over interactions w/ bath
- vacuum : $\mathcal{G}, \Gamma \rightarrow$ via $\mathcal{M} \rightarrow$ Feynman Rules
 \rightarrow Propagator $\langle 0 | \Gamma \phi_x \phi_y | 0 \rangle$

$$\text{tr} \langle 0 | A | 0 \rangle \xrightarrow{T \gg 0} \text{tr} (\rho A)$$

- Thermal propagator (scalar)

$$i \Delta_F^{T \gg 0}(x-y) = \int \frac{d^3 k}{(2\pi)^3} \frac{1}{2\omega_k} \left\{ \underbrace{[\underbrace{1}_{T=0} + \underbrace{n_B(\omega_k)}_{\text{induced creation}}]}_{\text{induced creation}} e^{-ik(x-y)} + \underbrace{n_B(\omega_k)}_{\text{absorption @ } x} e^{ik(x-y)} \right\}$$

$\beta = \frac{1}{T}$

- Imaginary Time Formalism :

$$t \rightarrow t = -i\tau \quad (\tau \in \mathbb{R}) \quad (\text{for } \tau = \beta \quad \Delta(\tau) = \Delta(\tau + n\beta))$$

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$$i \Delta_F^{T \gg 0} = i T \sum_{k_0 = 2\pi i n T} \int \frac{d^3 k}{(2\pi)^3} \frac{i}{k^2 - m^2} e^{-ik \cdot x}$$

$$e^{-\beta H} = e^{-itH}$$

$$\beta = \bar{c} = iT$$

Feynman Rules: • As in vacuum

$$\bullet \int \frac{dk_0}{2\pi} \rightarrow i T \sum_{k_0 = 2\pi i n T}$$

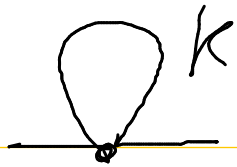
Fermions: $i S_F^{T \gg 0}(k) = i \frac{k \not{+} m}{k^2 - m^2}$ $k_0 = (2n+1)\pi iT$

Mixed Rep.: $\Delta(\bar{c}, \omega_k) = -T \sum_{k_0} e^{-k_0 \tau} \Delta(k)$

$$\Delta(k) = - \int_0^\beta d\tau e^{k_0 \tau} \Delta(\bar{c}, \omega_k) = \frac{1}{k^2 - m^2}$$

$$\Delta(\bar{c}, \omega_k) = \frac{1}{2\omega_k} \left\{ [1 + n_B(\omega_k)] e^{-i\omega_k \tau} + n_B(\omega_k) e^{i\omega_k \tau} \right\}$$

Our first diagram :



$$\mathcal{L}_{\text{int}} = -g^2 \phi^4$$

$$-i\Pi = \frac{1}{2} (-i 4! g^2) i \mathcal{T} \sum_{k_0} \int \frac{d^3k}{(2\pi)^3} i \Delta(k)$$

$$\text{Use } \mathcal{T} \sum_{k_0} e^{k_0 \tau} = \mathcal{T} \sum_{n=-\infty}^{\infty} e^{2\pi i n \tau} = \mathcal{T} \delta(\tau) = \delta(\tau)$$

$$\Rightarrow \mathcal{T} \sum_{k_0} \Delta(k) = - \int_0^{\beta} d\tau \underbrace{(\mathcal{T} \sum_{k_0} e^{k_0 \tau})}_{\delta(\tau)} \Delta(\tau, \omega_k) = \Delta(\tau=0, \omega_k)$$

$$\Rightarrow \Pi = 12 g^2 \int \frac{d^3k}{(2\pi)^3} \Delta(0, \omega_k)$$

$$= 6 g^2 \int \frac{d^3k}{(2\pi)^3} \frac{1}{\omega_k} [1 + 2n_B(\omega_k)]$$

\downarrow $\tau=0$ \downarrow $\tau > 0$

$$\Pi = \underbrace{\Pi_{\tau=0}}_{\downarrow \text{UV-divergent} \Rightarrow \text{Renormalisation}} + \underbrace{\Pi_{\tau > 0}}_{\rightarrow \text{finite because } n_B \text{ damps the integral}}$$

\Rightarrow Thermal does not introduce new UV-divergences

Massless case: $m=0$

$$\boxed{\Pi} = G g^2 \int_0^{\infty} \frac{dk \, 4\pi k^2}{(2\pi)^3} \frac{1}{k} \frac{1}{e^{k/T} - 1} \stackrel{x = \frac{k}{T}}{=} \frac{G}{\pi^2} g^2 T^2 \int_0^{\infty} \frac{x}{e^x - 1} = \boxed{g^2 T^2}$$

$$\zeta(2) = \frac{\pi^2}{6}$$

3) Hard Thermal Loop Resummation

Bare thermal propagators :

- IR divergences
- gauge dependent result
- miss contributions to lower order from higher order diagrams

CURE: HTL Resummation (Braaten, Pisatski 90s)

Separate :

- $k \sim T$ hard bare propagators
- $k \sim gT$ soft HTL resummed propagators

$g \ll 1$

Scalar Field :

$$\begin{array}{c} k \sim gT \\ \text{---} \bullet \end{array} = \text{---} + \overset{\sim T}{\text{---} \bigcirc \text{---}} + \text{---} \bigcirc \bigcirc \text{---} + \dots = \text{---} \leftarrow \bigcirc \text{---} \bullet$$

$$i\Delta^* = i\Delta + i\Delta(-i\pi)i\Delta + \dots$$

$$\begin{aligned} \Delta^* &= \Delta + \Delta\pi\Delta + \Delta\pi\Delta\pi\Delta + \dots = \Delta \sum_{n=0}^{\infty} (\pi\Delta)^n \\ &= \Delta \frac{1}{1-\pi\Delta} = \frac{1}{\Delta^{-1} - \pi} \\ &= \frac{1}{k^2 - m^2 - \pi} \end{aligned} \quad \text{convergent } \pi\Delta < 1$$

Recall from vacuum QFT: Setting $\Delta^{-1} = 0$ gives energy-momentum relation of a "free" particle
(dispersion relation, on-shell condition)

$$\Delta^{-1} = k^2 - m^2 \stackrel{!}{=} 0 \Rightarrow k_0 = \left(\frac{+}{-}\right) \omega_k \quad \omega_k = \sqrt{k^2 - m^2}$$

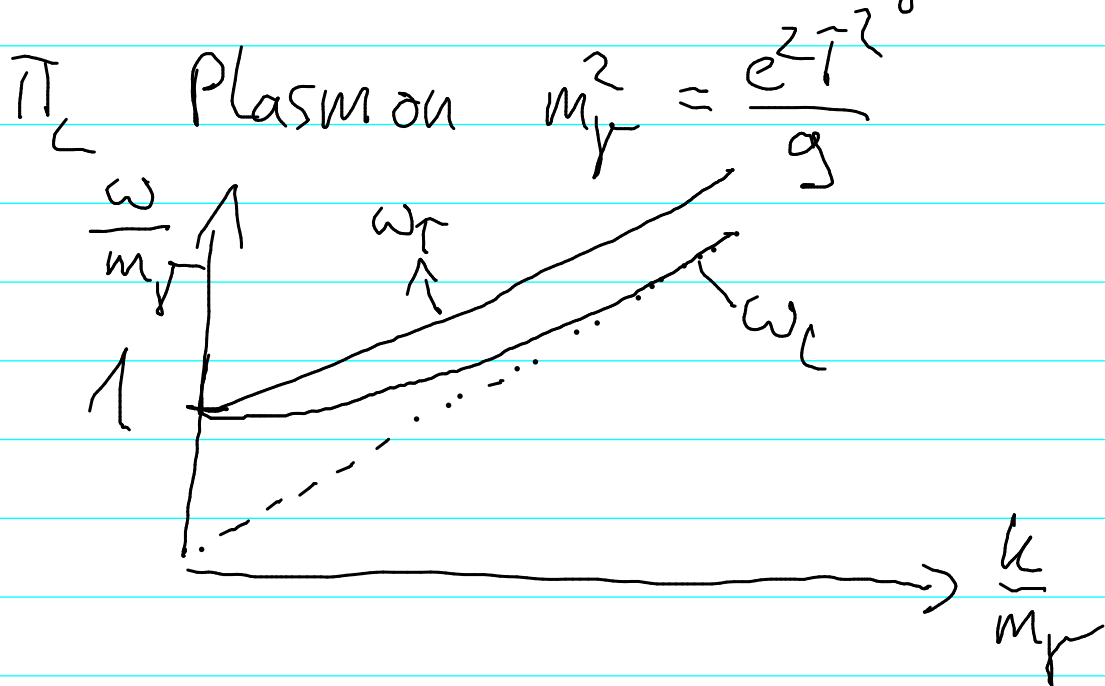
$$(\Delta^*)^{-1} = k^2 - m^2 - \pi \stackrel{!}{=} 0 \Rightarrow k_0 = \sqrt{k^2 + \underbrace{m^2 + \pi}_{m_{\text{eff}}^2}} \quad m_{\text{eff}}^2 = m^2 + \pi \approx \pi \text{ if } \pi \gg m^2$$

• Scalar field for $m \neq 0$: $m_{\text{th}}^2 = \Pi = g^2 T^2$

• Photon in QED: Propagator gets a physical longitudinal mode

$$i\mathcal{D} = i\mathcal{D} + i\mathcal{D} \overset{\sim T}{\bigcirc} i\mathcal{D} + i\mathcal{D} \overset{\sim T}{\bigcirc} \overset{\sim T}{\bigcirc} i\mathcal{D} + \dots$$

High T (HTL)-limit $k \sim T$



• Fermions ($m_f = 0$)

$$\overline{\psi} \not{\partial} \psi = \overline{\psi} \not{\partial} \psi + \overline{\psi} \not{\partial} \psi$$

$$S^* = - \frac{1}{2D_+} (\not{\gamma}_0 - \not{\gamma} \cdot \vec{k}) - \frac{1}{2D_-} (\not{\gamma}_0 + \not{\gamma} \cdot \vec{k})$$

ω_+

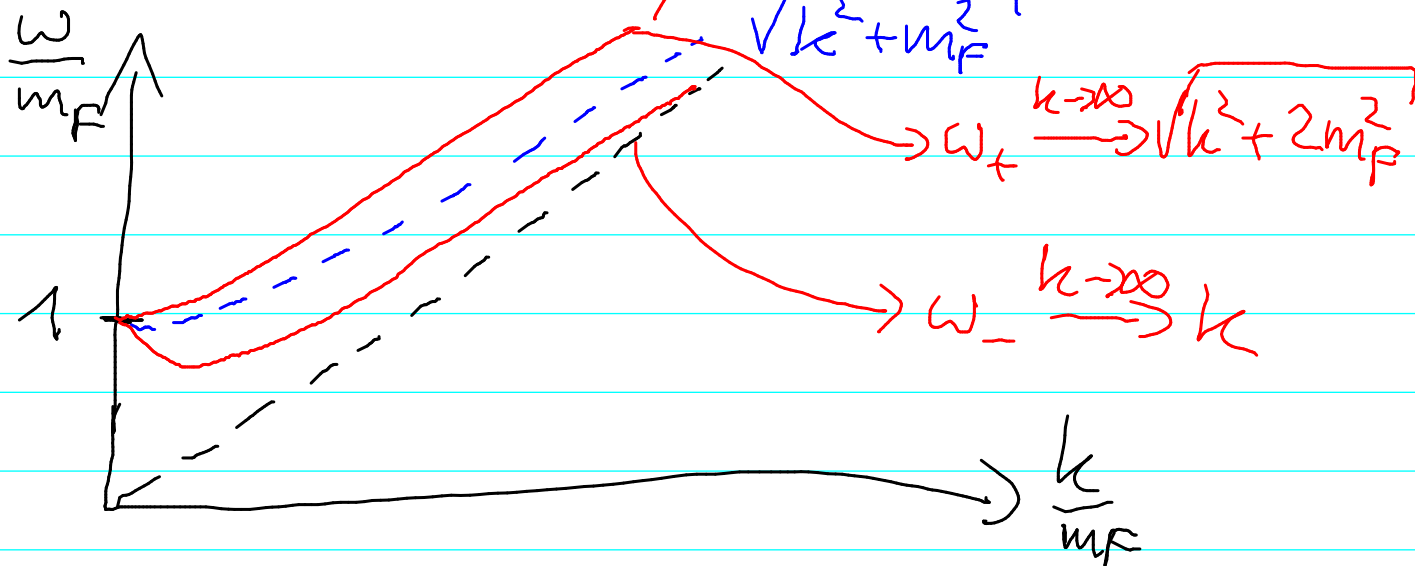
Helicity
Chirality $\equiv \chi = 1$

ω_-

$\chi = -1$

NEW @ $T > 0$!

$\frac{\omega}{m_f}$



$$m_f^2 = \frac{e^2 T^2}{8} \quad (\text{QCD})$$

S^* conserves chirality!

SUMMARY

- TFF: Average interactions
- $\langle O(A) \rangle \rightarrow \text{tr}(gA)$
- th. Propagator
- ITF \rightarrow Feynman rules
- Bath: Creation + absorption
- n_B, n_F occur naturally
- HFL res. gives improved pert. th.
- th. masses
- Modified dispersion relations
- New quasiparticles (Plasmon, Plasmino)