

Electroweak effects in vector-boson pair production at the LHC

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[\[arXiv:1208.3147\]](#), [\[arXiv:1208.3404\]](#)

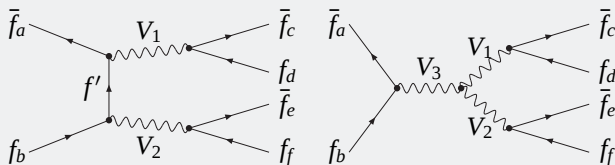
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Introduction – Weak-Boson Pair Production at the LHC

Vector-boson pair production: $pp \rightarrow WW/ZZ/WZ \rightarrow 4l$



- ZZ/WW production important **irreducible background** to inclusive SM Higgs-boson production
- Probe **non-abelian structure** of the Standard Model (SM) at high energies
- Search for **anomalous couplings**
- Backgrounds to **new-physics searches**, i.e. leptons + $\cancel{E}_T$ signatures
→ SUSY-particle pair production

Vector-Boson Pair Production – QCD Effects

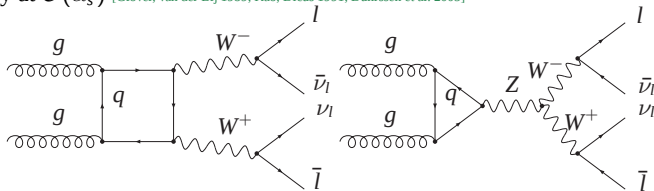
Extensive study of production of WW, WZ, ZZ, $W\gamma$, $Z\gamma$, $\gamma\gamma$ at NLO QCD [Campbell, Ellis,

Williams 2005; Campbell, Ellis 1999, . . .]

- Results matched with **parton showers** $\oplus$ combined with **soft gluon resummation**

[Nason, Ridolfi 2006; Frixione, Webber 2006]

- On-shell **leptonic decays** of the vector bosons taken into account (narrow-width approximation) retaining all spin information
- anomalous couplings included
- Corrections dominated by the $q\bar{q}$ channels
 - Significant contributions of the channels $gg \rightarrow V_1 V_2 \sim 10\%$ to LO, although formally at $\mathcal{O}(\alpha_s^2)$ [Glover, van der Bij 1989; Kao, Dicus 1991; Duhrssen et al. 2005]



- Even larger corrections of 30% if event selection for Higgs searches is applied**

[Binoth et al. 2006]

High-energy (Sudakov) limit

$$s, |t|, |u| \gg M_V^2, \quad V = W, Z$$

→ **bosons have to be produced at large p_T**

- EW corrections at high energies dominated by **universal large logarithms**

$$\propto \alpha^L \ln^{2L}(M_V/\sqrt{s}) \quad (\text{LL}),$$

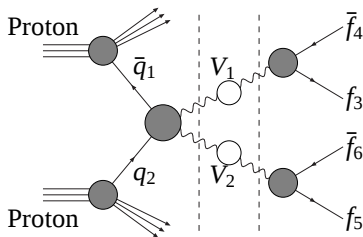
$$\propto \alpha^L \ln^{2L-1}(M_V/\sqrt{s}) \quad (\text{NLL}), \dots$$

at the L -loop level

- Corrections of $\sim -50\%$ at $\sqrt{s} = 2 \text{ TeV}$ (W-pair production)
- Change of sign going from LL to NLL (to NNLL ...)
→ **substantial cancellations possible!**

EW Corrections to V-Pair Production – Overview

- $\mathcal{O}(\alpha)$ high-energy approximation known for WW/WZ/ZZ production, vector bosons treated in [pole-approximation](#) [Accomando et al. 2002–2006]
→ **final-state leptons phenomenologically accessible**



- [NNLL effects](#) at two loops published for the WW channel [Kühn, Metzler, Uccirati, Penin 2011]
- **We have calculated the full one-loop corrections to on-shell V-pair production at the LHC** [Bierweiler, TK, Kühn, Uccirati 2012]

Motivation – Why the full Calculation?

- **First LHC results on total cross sections** for $pp \rightarrow WW/WZ/ZZ$ have been published

- **ATLAS:** $pp \rightarrow W^+W^-$ at 7 TeV, 1.02 fb^{-1} [arXiv:1203.6232 [hep-ex]],
 $pp \rightarrow W^+W^-$ at 8 TeV, 5.8 fb^{-1} [ATLAS-CONF-2012-090]
- **CMS:** $pp \rightarrow WW/WZ/ZZ$ at 7 TeV, 1.1 fb^{-1} [CMS PAS EWK-11-010],
 $pp \rightarrow WW(ZZ)$ at 8 TeV, $3.5(5.3) \text{ fb}^{-1}$ [CMS PAS SMP-12-013, CMS PAS SMP-12-014]

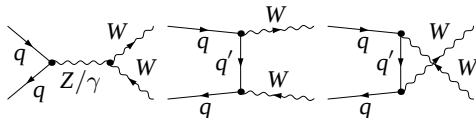
experimental error at the level of 15%; WW dominated by systematics, ZZ dominated by statistics

⇒ **Quantify EW effects in this SM benchmark process**

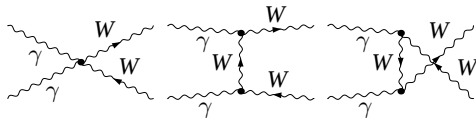
- **Focus of LHC physics:** V-pair production at high energies, $\hat{s} \gg M_W^2$:
 - V-pair production dominated by **small scattering angles**, i.e. small $\hat{t}, \hat{u}$
 - Potentially sizable effects due to anomalous couplings
 - **Until now, no predictions for EW effects in this particular kinematic regime!**
- Electroweak corrections **potentially large**

Lowest-Order Contributions to $pp \rightarrow W^-W^+ + X$

- **Partonic LO contributions at $\mathcal{O}(\alpha^2)$**



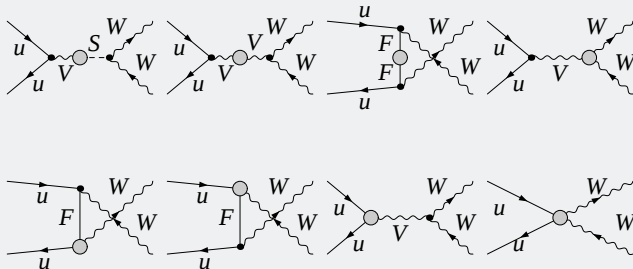
- **Photon-induced contributions at $\mathcal{O}(\alpha^2)$**



- Adopt MRST2004QED PDF set [\[Martin et al. 2005\]](#)
- Potentially large contribution at high invariant masses!
- **Caveat:** large theoretical uncertainties!

Virtual EW Corrections to $pp \rightarrow W^-W^+ + X$

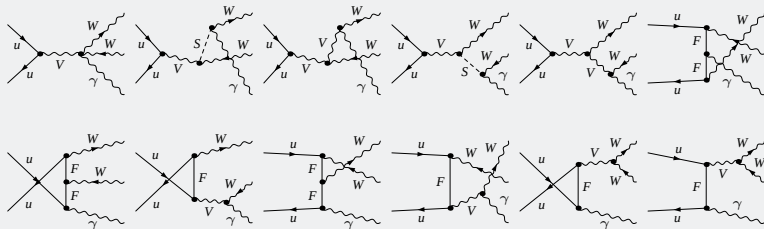
One-loop contributions at $\mathcal{O}(\alpha^3)$



- On-shell renormalization of SM parameters
- We use the Fermi scheme to calculate the loop corrections.
→ universal corrections to Δr absorbed in effective LO coupling
- $V_{ij}^{\text{CKM}} = \delta_{ij}$ within the loops → no renormalization of V_{ij}^{CKM}

Real EW Corrections – Infrared Singularities

Real photon radiation at $\mathcal{O}(\alpha^3)$ (generic diagrams): $q\bar{q} \rightarrow W^-W^+ + \gamma$



- **Soft singularities** due to soft photons
- **Initial-state collinear singularities** due to collinear photon radiation off initial-state quarks $\rightarrow$ renormalization of PDFs
- Introduce small **quark mass m_q** and infinitesimal **photon mass λ** to regularize divergences $\rightarrow$ results exhibit unphysical $\ln m_q$ and $\ln \lambda$ terms

Apply phase-space slicing for numerically-stable evaluation of phase-space integral

Numerical Results (I) – Total Cross Sections

LHC at 8 TeV:

no cuts	$\sigma_{\text{LO}}^{q\bar{q}'} \text{ (pbarn)}$	$\delta_{\text{EW}} \text{ (%)}$	$\delta_{\gamma\gamma} \text{ (%)}$
W^+W^-	35	-0.4	1.5
W^+Z	9	-1.4	—
W^-Z	5	-1.3	—
ZZ	6	-4.1	—

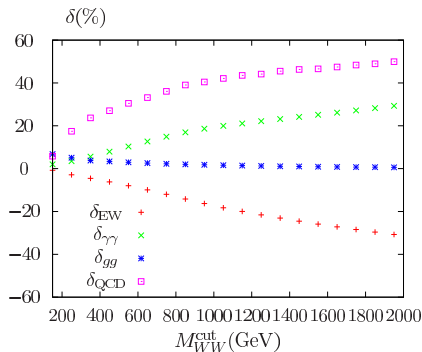
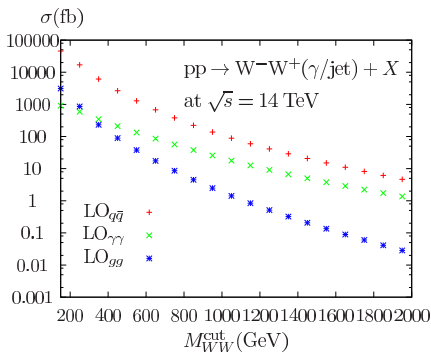
LHC at 14 TeV:

no cuts	$\sigma_{\text{LO}}^{q\bar{q}'} \text{ (pbarn)}$	$\delta_{\text{EW}} \text{ (%)}$	$\delta_{\gamma\gamma} \text{ (%)}$
W^+W^-	75	-0.5	1.6
W^+Z	19	-1.4	—
W^-Z	12	-1.4	—
ZZ	12	-4.2	—

- $\sigma = \sigma_{\text{LO}} \times (1 + \delta)$
- **Setup:** MSTW2008LO PDFs [\[Martin et al. 2009\]](#), $\mu_R = \mu_F = (m_{T,1} + m_{T,2})/2$

Numerical Results (II) – $pp \rightarrow W^-W^+(\gamma)$

Default cuts: $p_{T,W^\pm} > 15$ GeV, $y_{W^\pm} < 2.5$, **small scattering angles**



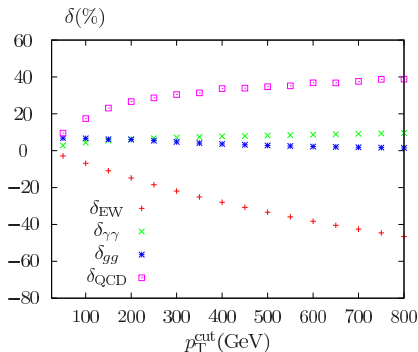
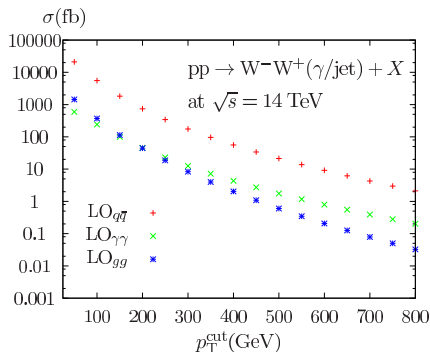
- assume $\int \mathcal{L} dt = 200 \text{ fb}^{-1}$
 $\Rightarrow$ 1000 WW events with $M_{WW} > 2$ TeV

- rapidly increasing admixture of $\gamma\gamma$

- **large admixture** of $\gamma\gamma$ (up to 30%)
- **sizable negative EW corrections** (-30%) at high energies,
comparable to QCD corrections

Numerical Results (III) – $pp \rightarrow W^-W^+(\gamma)$

Default cuts: $p_{T,W^\pm} > 15$ GeV, $|y_{W^\pm}| < 2.5$, **large scattering angles**



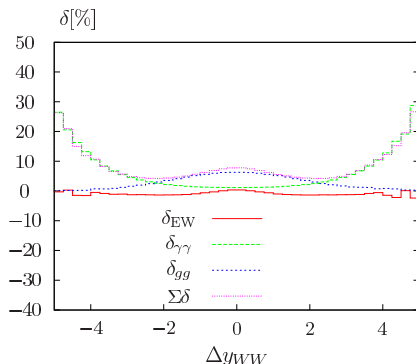
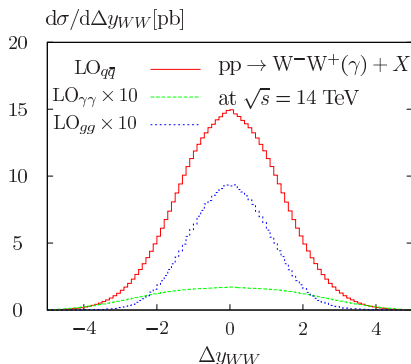
- assume $\int \mathcal{L} dt = 200 \text{ fb}^{-1}$
 $\Rightarrow$ 1000 WW events with $p_T > 500$ GeV
- decreasing admixture of gg, increasing admixture of $\gamma\gamma$
- large admixture of $\gamma\gamma$ (10%!)
- large negative EW corrections (-45%), comparable to QCD corrections

No compensation between $\gamma\gamma \rightarrow WW$ and weak corrections!
 $\Rightarrow$ Different angular distributions!

- $\sigma(\gamma\gamma \rightarrow WW) \rightarrow \frac{8\pi\alpha^2}{M_W^2}$ for large $\hat{s}$
 $\Rightarrow$ **strong enhancement** in forward & backward directions
- **weak corrections:**
negative Sudakov logs for large $\hat{s}$ and $\hat{t}$
 $\Rightarrow$ negative corrections for large scattering angles
- $gg \rightarrow WW$ **small, isotropic**
- implications for $d\sigma/d\Delta y_{WW}$ with $\Delta y_{WW} = y_{W^+} - y_{W^-}$
(for fixed M_{WW} this corresponds to the angular distribution in the W-W rest frame!)

Numerical Results (IV) – $pp \rightarrow W^-W^+$

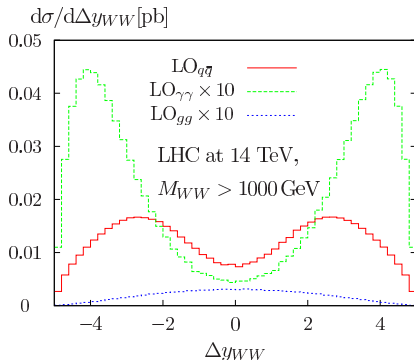
Default cuts: $p_{T,W^\pm} > 15 \text{ GeV}$, $y_{W^\pm} < 2.5$



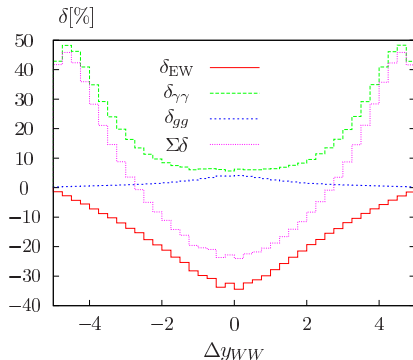
- WW production dominated by events near threshold, isotropic production at small Δy_{WW}
- 5% increase of cross section by gg channel
- EW corrections at the percent level
- Sizable contributions from $\gamma\gamma$ at large $|\Delta y_{WW}|$

Numerical Results (V) – $pp \rightarrow W^-W^+(\gamma)$

High-energy cuts: $p_{T,W^\pm} > 15$ GeV, $y_{W^\pm} < 2.5$, $M_{WW} > 1$ TeV



- WW production dominated by small scattering angles
- drastic forward-backward peaking of $\gamma\gamma \rightarrow WW$



- drastic distortion of angular distribution
- $\Sigma\delta$ varies from -30% and $+45\%$ for $M_{WW} > 1$ TeV!

“Real Radiation” of massive vector bosons

Low energies: Phase-space and perturbative suppression of $pp \rightarrow WW + (W/Z)$
 $\Rightarrow$ contribution below 1%

High energies: Logarithmic enhancement of additional soft/collinear W- or Z-boson radiation

$\Rightarrow$ Investigation of $WW + W/Z$ production as background to W pairs at large p_T , M_{WW}

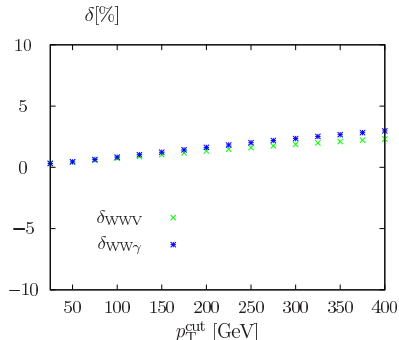
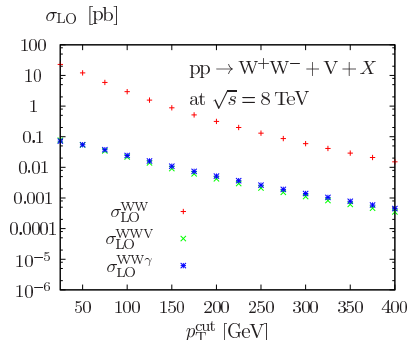
- invisible decay of $Z \rightarrow \nu\bar{\nu}$
- collinear emission
- ...

Simplified approach: (details depend on experimental analysis)

- 1 Include $pp \rightarrow W^-W^+Z$ with totally inclusive Z
- 2 Include $pp \rightarrow W^-W^+W^\pm$; treat $W^\pm$ with lowest p_T totally inclusively

Numerical Results (VI) – “Real Radiation” of W/Z

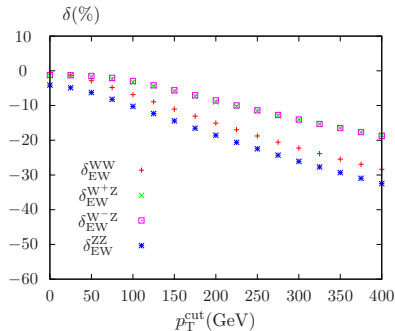
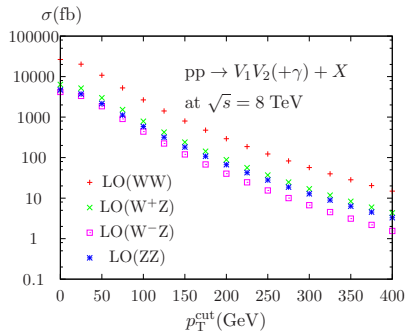
LHC at 8 TeV, default cuts: $p_{T,W^\pm} > 15 \text{ GeV}$, $|y_{W^\pm}| < 2.5$



- Corrections due to hard photons ($p_{T,\gamma} > 15 \text{ GeV}$, $|y_\gamma| < 2.5$) below 5 %
- Contributions of massive-boson radiation below 5 %

Numerical Results (VII) – WW, W⁺Z and ZZ production

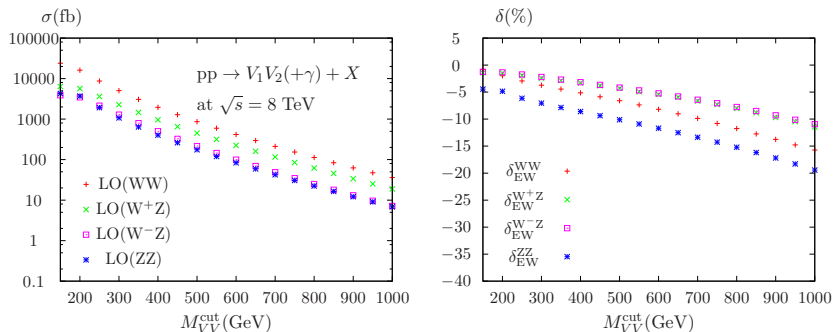
LHC at 8 TeV, default cuts: $p_{T,V} > 15$ GeV, $|y_V| < 2.5$, **large scattering angles**



- moderate EW corrections to WZ production
- large EW corrections to ZZ production

Numerical Results (VIII) – WW, W⁺Z and ZZ production

LHC at 8 TeV, default cuts: $p_{T,V} > 15$ GeV, $|y_V| < 2.5$, **small scattering angles**



- moderate EW corrections to WZ production
- large EW corrections to ZZ production
- **EW corrections important at high energies and small scattering angles**

- **Proper understanding of V-pair production processes crucial at the LHC:**
 - Understand **SM at highest energies**
 - Understand **backgrounds to Higgs- and BSM-physics searches**
- **We have computed the full EW corrections to V-pair production at the LHC**
 - Small corrections to the total cross section
 - **But:** **Sizable negative corrections** at high energies
 - **W-pair production:** significant contributions of photon-induced channels at large scattering angles
 - Real-radiation contributions of W/Z small; further investigation needed?
- **Future work:**
 - **Leptonic decays** of the vector bosons should be included.
 - Discuss combination of EW and QCD effects
 - Discuss interplay of **anomalous couplings and EW effects**

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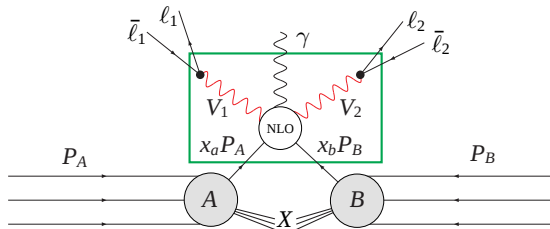
Thank You!

Reminder: Calculation of Hadronic Cross Sections

Schematic illustration for

$$pp \rightarrow V_1 V_2 (+\gamma) + X$$

$$\rightarrow \ell_1 \ell_2 \bar{\ell}_1 \bar{\ell}_2 (+\gamma) + X$$



Hadronic cross sections

$$d\sigma_{AB}(p_A, p_B) = \sum_{a,b} \int_0^1 dx_a \int_0^1 dx_b f_{a/A}(x_a, \mu_F) f_{b/B}(x_b, \mu_F) d\hat{\sigma}_{ab}^{\text{NLO}}(p_a, p_b, \mu_F, \mu_R) \\ \times \mathcal{F}^{(4\ell+\gamma)}(\{\mathcal{O}_{\text{FS}}\}), \quad p_{\{a,b\}}^\mu = x_{\{a,b\}} P_{\{A,B\}}^\mu$$

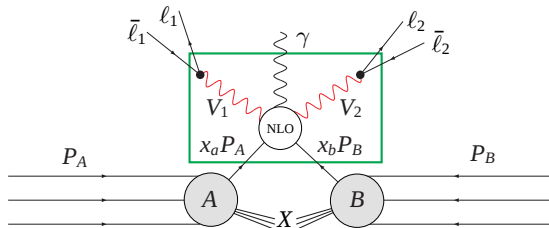
- Dependence on μ_R , μ_F reduced by inclusion of higher perturbative orders
- $\mathcal{F}^{(4\ell+\gamma)}$ incorporates definition of observables + phase-space cuts

Reminder: Calculation of Hadronic Cross Sections

Schematic illustration for

$$pp \rightarrow V_1 V_2 (+\gamma) + X$$

$$\rightarrow \ell_1 \ell_2 \bar{\ell}_1 \bar{\ell}_2 (+\gamma) + X$$



Hadronic cross sections

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NLO partonic cross section:

$$\hat{\sigma}_{ab}^{\text{NLO}} = \hat{\sigma}_{ab}^{\text{LO}} + \hat{\sigma}_{ab}^{\text{virt}} + \hat{\sigma}_{ab}^{\text{real}}$$

Double-Pole Approximation (DPA)

- **Lowest order:** Amplitude given as a product of **on-shell (OS) production amplitude** $\otimes$ **on-shell decay amplitude** $\otimes$ **Breit-Wigner**:

$$\mathcal{M}_{\text{Born,DPA}}^{\bar{q}_1 q_2 \rightarrow V_1 V_2 \rightarrow 4f} = \frac{1}{k_1^2 - M_1^2 + iM_1\Gamma_1} \frac{1}{k_2^2 - M_2^2 + iM_2\Gamma_2} \\ \times \sum_{\lambda_1, \lambda_2} \mathcal{M}_{\text{Born}}^{\bar{q}_1 q_2 \rightarrow V_{1,\lambda_1} V_{2,\lambda_2}} \mathcal{M}_{\text{Born}}^{V_{1,\lambda_1} \rightarrow f_3 \bar{f}_4} \mathcal{M}_{\text{Born}}^{V_{2,\lambda_2} \rightarrow f_5 \bar{f}_6}$$

- Use **OS-projected momenta** $\hat{k}$ [Denner, Dittmaier, Roth, Wackerath 2000] in the OS matrix elements:

$$\hat{k}_{1,0} = \frac{1}{2} \sqrt{\hat{s}}, \quad \hat{\mathbf{k}}_1 = \frac{\mathbf{k}_1}{|\mathbf{k}_1|} \beta_W \frac{\sqrt{\hat{s}}}{2}, \quad \dots$$

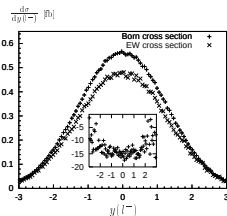
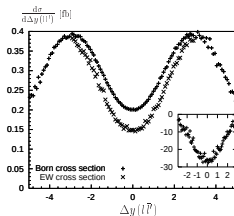
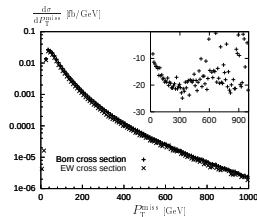
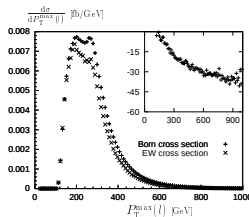
- **NLO:** EW corrections consist of **factorizable** and **non-factorizable** contributions, e.g.

$$\mathcal{M}_{\text{fact}} = \frac{R(k_1, k_2, \theta)}{(k_1^2 - M_1^2 + iM_1\Gamma_1) (k_2^2 - M_2^2 + iM_2\Gamma_2)}$$

Caution: Gauge invariance!

EW corrections to $pp \rightarrow W^+W^- \rightarrow \nu_e e^+ \mu^- \bar{\nu}_\mu$ (DPA)

- Standard LHC event selection cuts applied to final-state leptons and missing transverse momentum; **additionally** $M_{e^+\mu^-} > 500 \text{ GeV}$ required
- Large negative corrections at large transverse momenta
- Substantial negative corrections to inclusive observables
- Error due to DPA about 10% in the relative corrections
- EW corrections significantly larger than experimental error throughout the whole energy range (for $L \sim 30 \text{ fb}^{-1}$)**



[Accomando, Denner, Kaiser: arXiv:0409247 [hep-ph]]

- Simple LL ansatz for $f_{\gamma/p}(x, Q_0^2)$

$$f_{\gamma/p}(x, Q_0^2) = \frac{\alpha}{2\pi} \left[\frac{4}{9} \ln \left(\frac{Q_0^2}{m_u^2} \right) f_{u/p,v}(x, Q_0^2) + \frac{1}{9} \ln \left(\frac{Q_0^2}{m_d^2} \right) f_{d/p,v}(x, Q_0^2) \right] \otimes \frac{1 + (1-x)^2}{x}$$

- Running of $f_{q/p}(x, Q^2)$ at $\mathcal{O}(\alpha)$ affected by photon PDFs!

$$\frac{\partial f_{q/p}(x, \mu^2)}{\partial \ln \mu^2} = \frac{\alpha}{2\pi} \int_x^1 \frac{dy}{y} [P_{qq}(y) Q_q^2 f_{q/p}(x/y, \mu^2) + P_{q\gamma}(y) Q_q^2 f_{\gamma/p}(x/y, \mu^2)]$$

- Momentum conservation

$$\int_0^1 dx \, x \left[\sum_q f_{q/p}(x, \mu^2) + f_{g/p}(x, \mu^2) + f_{\gamma/p}(x, \mu^2) \right] = 1$$

⇒ QED effects on $f_{q/p}(x, \mu^2)$ small!

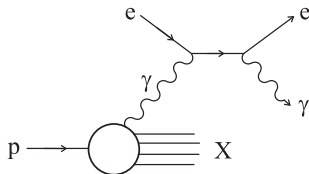
⇒ Still large conceptual uncertainties in $f_{\gamma,0}$

Measure Photon PDFs?

Consider the DIS process

$$ep \rightarrow e\gamma + X$$

with high- p_T back-to-back e, γ in the final state



$$\sigma(ep \rightarrow e\gamma + X) = \int dx^\gamma f_{\gamma/p}(x^\gamma, \mu^2) \hat{\sigma}(e\gamma \rightarrow e\gamma),$$

related to Compton scattering

- $x^\gamma = \frac{E_T^\gamma E_e \exp(\eta^\gamma)}{2E_p E_e - E_T^\gamma E_e \exp(-\eta^\gamma)}$
- $f_{\gamma/p}(x^\gamma, \mu^2)$ could be in principle extracted from HERA data!

EW Input Schemes – Definition of α

- $\alpha(0)$: On-shell definition in the Thomson-limit (zero momentum transfer)

$$\bar{u}(p)\Gamma_{\mu}^{Ae\bar{e}}(p,p)u(p)\big|_{p^2=m_e^2} = e(0)\bar{u}(p)\gamma_{\mu}u(p), \alpha(0) = e(0)^2/4\pi$$

- $\alpha(M_Z)$ obtained via renormalization-group running from 0 to weak scale M_Z

$$\alpha(M_Z) = \frac{\alpha(0)}{1 - \Delta\alpha(M_Z)}, \quad \Delta\alpha(M_Z) = \Pi_{f\neq t}^{AA}(0) - \text{Re} \Pi_{f\neq t}^{AA}(M_Z^2)$$

- $\alpha_{G_{\mu}}$ defined through the Fermi constant related to the muon lifetime

$$\alpha_{G_{\mu}} = \frac{\sqrt{2}G_{\mu}M_W^2s_W^2}{\pi} = \frac{\alpha(0)}{1 - \Delta r}$$

Δr includes corrections to muon lifetime not contained in QED-improved Fermi model

- **light-fermion mass logs contained in $\Pi_{f\neq t}^{AA}(0)$ resummed in effective couplings $\alpha(M_Z)$ and $\alpha_{G_{\mu}}$**

Two-cut-off phase-space slicing

- Definition of bremsstrahlung phase space:

$$\sigma_{\text{real}} = \int \text{dPS}(W^- W^+ \gamma) |\mathcal{M}^\gamma|^2$$

- Phase-space decomposition:

$$\sigma_{\text{real}} = \sigma_{\text{hard}} + \sigma_{\text{soft}} + \sigma_{\text{coll}}$$

- **Soft limit:** $E_\gamma < \Delta E \ll M_W$

$$\sigma_{\text{soft}}(\Delta E) = -\sigma_{\text{LO}} \left[\frac{e^2}{(2\pi)^3} \int_{|\mathbf{k}_\gamma| < \Delta E} \frac{d^3\mathbf{k}_\gamma}{2\sqrt{\mathbf{k}_\gamma^2 + \lambda^2}} \sum_{ij} \frac{\pm Q_i Q_j (p_i p_j)}{(p_i k_\gamma)(p_j k_\gamma)} \right]$$

- **Collinear limit:** $\theta_{q\gamma} < \Delta\theta \ll 1$, $E_\gamma > \Delta E$

$$\sigma_{\text{coll},q}(\Delta E, \Delta\theta) = \frac{\alpha Q_q^2}{2\pi} \int_0^{1-2\Delta E/\sqrt{s}} dz \frac{(1+z^2)}{1-z} \left(\ln \frac{\hat{s}(\Delta\theta)^2}{4m_q^2} - \frac{2z}{1+z^2} \right) \sigma_{\text{LO}}(z\hat{s})$$

- **Hard bremsstrahlung:** $\theta_{q\gamma} > \Delta\theta$, $E_\gamma > \Delta E$;
numerical evaluation of $\sigma_{\text{hard}}(\Delta E, \Delta\theta)$ without regulators
- Numerical result independent of $\ln \Delta E$ and $\ln \Delta\theta$

$\ln m_q$ and $\ln \lambda$ terms cancel in the sum $\sigma_{\text{virt}} + \sigma_{\text{soft}} + \sigma_{\text{coll}}$ in infrared-safe observables

Virtual corrections computed in the [FeynArts/FormCalc/LoopTools](#)(FF) framework [(FA): Küblbeck, Böhm, Denner 1990; (FC,LT): Hahn, Pérez Victoria 1999; Hahn 2001; (FF): van Oldenborgh, Vermaseren 1990]

1 FeynArts-3.5:

- Automatic generation of diagrams
- Calculation of amplitudes

2 FormCalc-6.1:

- Algebraical simplification of amplitudes, introduction of tensor coefficients
- Analytical calculation of squared amplitudes
- Spin-, colour- and polarization sums
- Generation of Fortran code

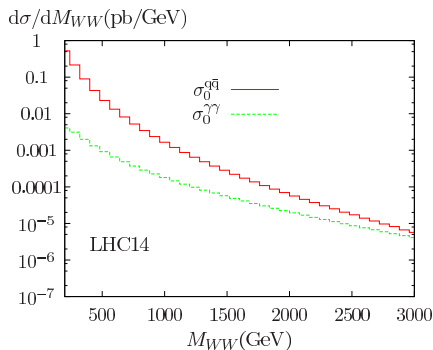
3 LoopTools-2.5:

- Numerical Passarino–Veltman reduction within Fortran
- Numerically-stable evaluation of scalar integrals

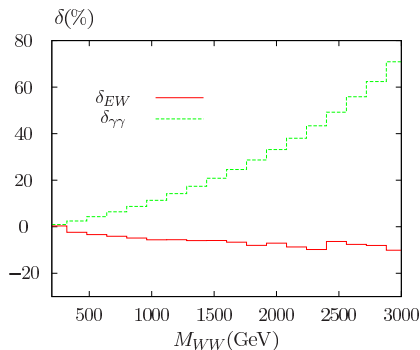
Bremsstrahlung amplitudes computed with [FeynArts/FeynCalc](#) $\oplus$ [Madgraph](#) [Alwall et al.], numerical phase-space integration within Fortran using the Vegas algorithm

pp $\rightarrow$ W⁻W⁺(γ) – Numerical Results

No cuts



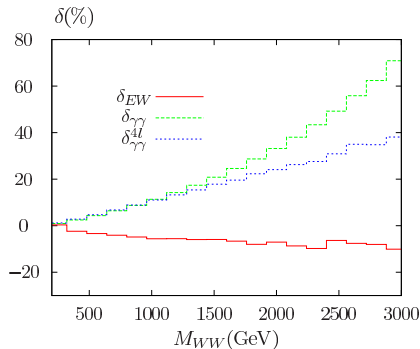
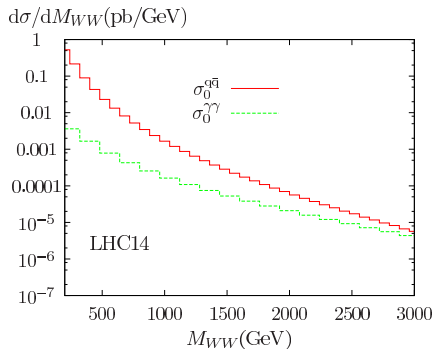
- LO cross section dominated by $q\bar{q}$ contributions
- Rapid decrease of cross section for increasing invariant masses



- EW corrections small even for large values of M_{WW}
 - Large contributions (+80%!) from $\gamma\gamma \rightarrow WW$ at high invariant masses
- ⇒ Leptonic decays?

pp $\rightarrow$ W⁻W⁺(γ) – Numerical Results

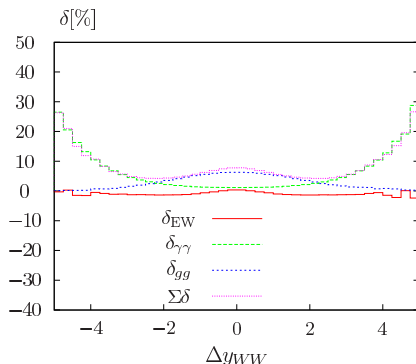
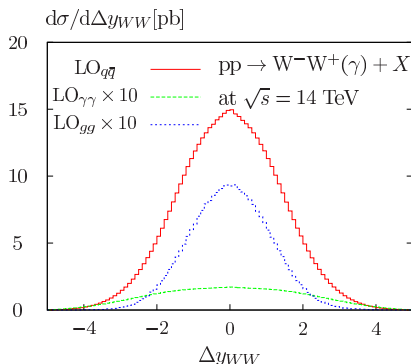
LHC acceptance cuts



- LO cross section dominated by $q\bar{q}$ contributions
- Rapid decrease of cross section for increasing invariant masses
- Employ LHC cuts on decay products:
 $p_{T,l} > 20$ GeV, $|y_l| < 3$, $p_{T,\text{miss}} > 25$ GeV
⇒ relative effect of $\gamma\gamma \rightarrow WW$ reduced by factor 2 at large M_{WW}

EW Corrections to $pp \rightarrow W^-W^+ - \text{Numerical Results}$

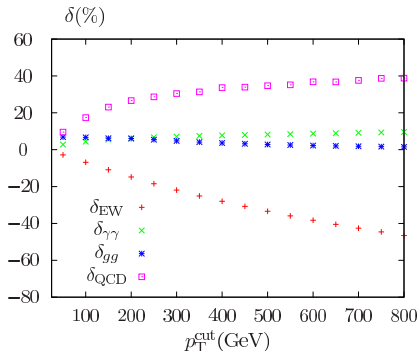
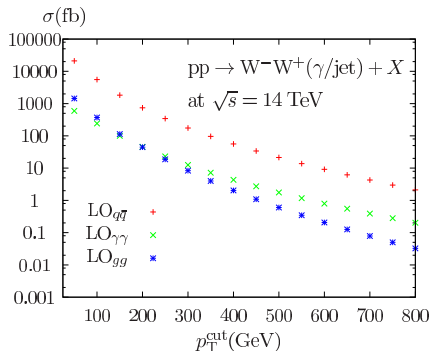
Default cuts: $p_{T,W^\pm} > 15 \text{ GeV}$, $y_{W^\pm} < 2.5$



- WW production dominated by events near threshold, isotropic production at small Δy_{WW}
- 5% increase of cross section by gg channel
- EW corrections at the percent level
- Sizable contributions from $\gamma\gamma$ at large $|\Delta y_{WW}|$

pp $\rightarrow$ W⁻W⁺(γ) – Numerical Results

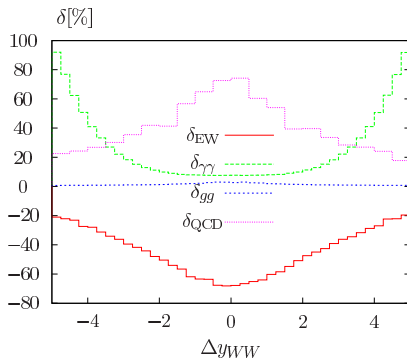
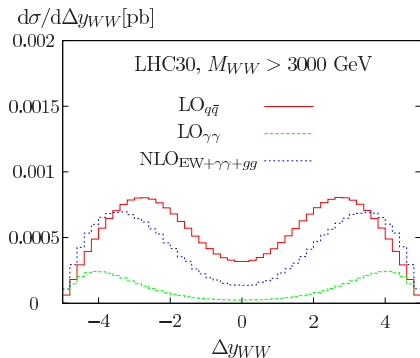
Default cuts: $p_{T,W^\pm} > 15$ GeV, $|y_{W^\pm}| < 2.5$



- assume $\int \mathcal{L} dt = 200 \text{ fb}^{-1}$
 $\Rightarrow$ 1000 WW events with $p_T > 500$ GeV
- decreasing admixture of gg, increasing admixture of $\gamma\gamma$
- large admixture of $\gamma\gamma$ (10%!)
- large negative EW corrections (-45%), comparable to QCD corrections

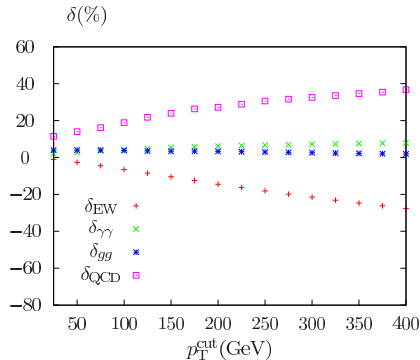
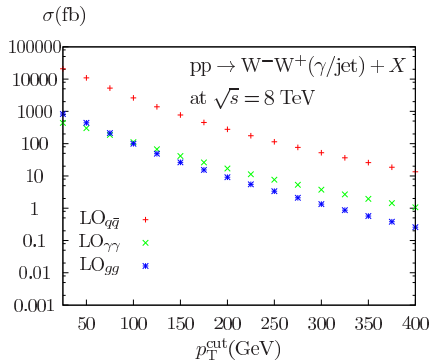
EW Corrections to $pp \rightarrow W^-W^+$ – Numerical Results

Very-high-energy cuts: $p_{T,W^\pm} > 15$ GeV, $y_{W^\pm} < 2.5$, $M_{WW} > 3$ TeV

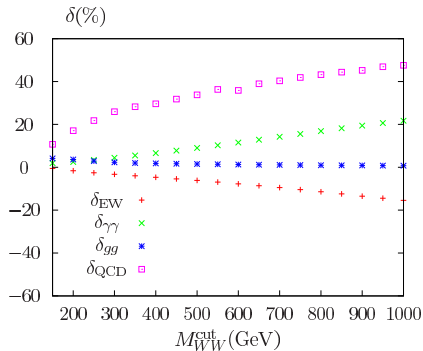
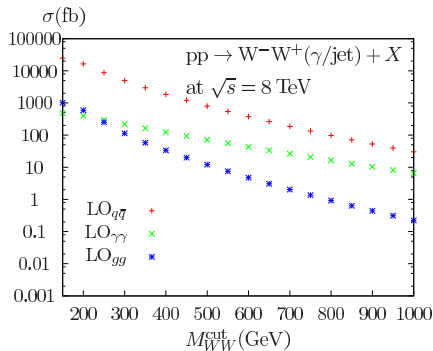


- NLO EW as important as QCD
- extreme distortion due to $\gamma\gamma$ (**caveat:** high uncertainty in photon PDFs)

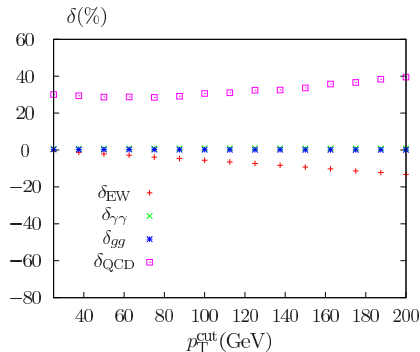
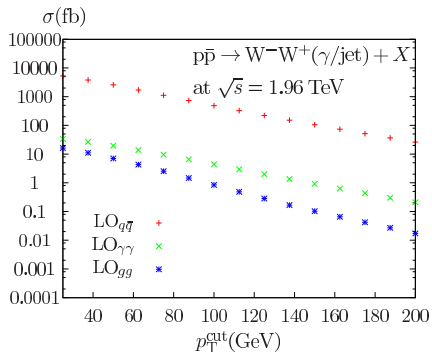
Transverse-momentum distribution at the LHC8



Invariant-mass distribution at the LHC8



Transverse-momentum distribution at the Tevatron



Invariant-mass distribution at the Tevatron

