

Numerical evaluation of jet amplitudes at NLO

Valery Yundin

Niels Bohr International Academy & Discovery Center

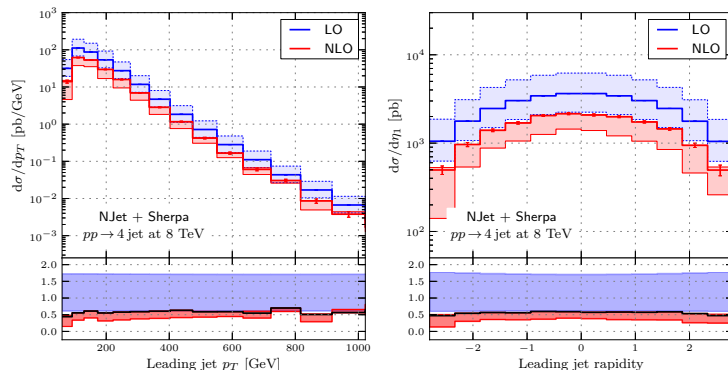
in collaboration with Simon Badger, Benedikt Biedermann and Peter Uwer

HP2, 6 September 2012, München

General solutions to virtual corrections

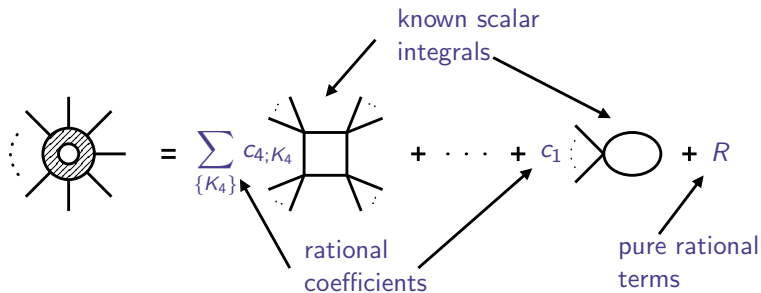
- HP2 {
- ▶ Helac-NLO [public] [arXiv:1110.1499]
 - ▶ GoSam [public] [arXiv:1111.2034]
 - ▶ BlackHat [private, W/Z+jets, multi-jets] [arXiv:0803.4180+...]
 - ▶ MadLoop(5) [private, SM+More] [arXiv:1103.0621]
 - ▶ Open Loops [private, QCD SM] [arXiv:1111.5206]
 - ▶ Rocket [private, W+jets, WW+jets, $t\bar{t}$ +jet] [arXiv:0805.2152+...]
 - ▶ MCFM [public, max. $2 \rightarrow 3$] [http://mcfm.fnal.gov]
 - ▶ Feynman based approaches:
VBFNLO [public], Denner et al., FeynCalc [public], Reina et al. ...
- New HP2
- ▶ NJet [public, multi-jets] [arXiv:1209.0100]

NJet (virtual) + Sherpa (real+integration) = full colour jet XS



NLO QCD corrections to 3 and 4 jets at 8 TeV [[arXiv:1209.0098](https://arxiv.org/abs/1209.0098)]

Structure of One-Loop Amplitudes



- ▶ Gauge theory amplitudes reduced to box topologies or simpler
[Passarino,Veltman;Melrose]
- ▶ Isolate logarithms with cuts and exploit on-shell simplifications
- ▶ General cutting principle:
[Bern,Dixon,Kosower]
 - ▶ apply δ -functions to left and right sides
 - ▶ generate and solve the linear system for the coefficients

NGLuon: public C++ library for evaluating multi-parton **primitive amplitudes** via unitarity (**now part of NJet**) [Comput.Phys.Commun. 182 (2011)]

Main features

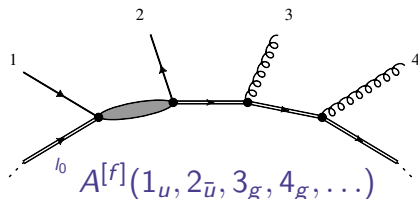
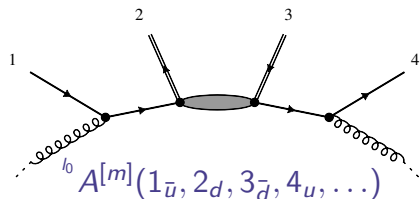
- ▶ Efficient tree amplitudes using Berends-Giele recursion.
- ▶ Rational terms from massive loop cuts.
- ▶ Extraction of integral coefficients via Fourier projections.
- ▶ Everything is in 4 dimensions (except loop integrals).

NJet: **public C++ library** for evaluating multi-parton **matrix elements** in massless QCD [<https://bitbucket.org/njet/njet>] [arXiv:1209.0100]

Additions since the first public version

- ▶ Multi-quark primitive amplitudes.
- ▶ NParton – interface to partonic primitive amplitudes.
- ▶ Full colour-summed amplitudes for up to 5 outgoing partons.
- ▶ Binoth Les Houches Accord interface for MC generators.

NParton computes arbitrary multi-fermion primitives.



All primitives are separated into two classes

- ▶ With **mixed** fermion and gluon loop content ($l_0 = \text{gluon}$)
- ▶ With internal **fermion loops** ($l_0 = \text{quark}$)

These two classes cover all partonic primitives in one loop QCD.

Colour decomposition of an L-loop amplitude:

$$\mathcal{A}_n^{(L)}(\{p_i\}) = \sum_c \underbrace{T_c(\{a_i\})}_{\text{colour basis}} \underbrace{A_{n;c}^{(L)}(p_1, \dots, p_n)}_{\text{partial amplitudes}}$$

Knowing partial amplitudes we can get matrix elements:

$$\sum_{\text{hel}} \sum_{\text{col}} \mathcal{A}_n^{(1)} \mathcal{A}_n^{(0)\dagger} = \sum_{\text{hel}} \sum_{cc'} A_{n;c}^{(1)} \cdot \mathcal{C}_{cc'} \cdot A_{n;c'}^{(0)\dagger}$$

Partial-Primitive decomposition for pure gluons:

- ▶ Tree level: Kleiss-Kuijf basis of $(n-2)!$ primitives
- ▶ One-loop: a basis of $(n-1)!/2$ primitives.

Partial-Primitive decomposition for multi-quark case:

No analytic formula. Reconstruct partials using diagram matching.

1. Generate all diagrams¹ D_i for a given n -parton amplitude $\mathcal{A}_n$

$$\mathcal{A}_n = \sum_{i=1}^{N_{\text{dia}}} D_i = \sum_{i=1}^{\hat{N}_{\text{dia}}} C_i K_i \quad C_i = \sum_c T_c F_{ci}$$

2. Write all possible **primitives** P_i as combinations of colour-stripped **diagrams** K_i using **matching matrix** M_{ij}

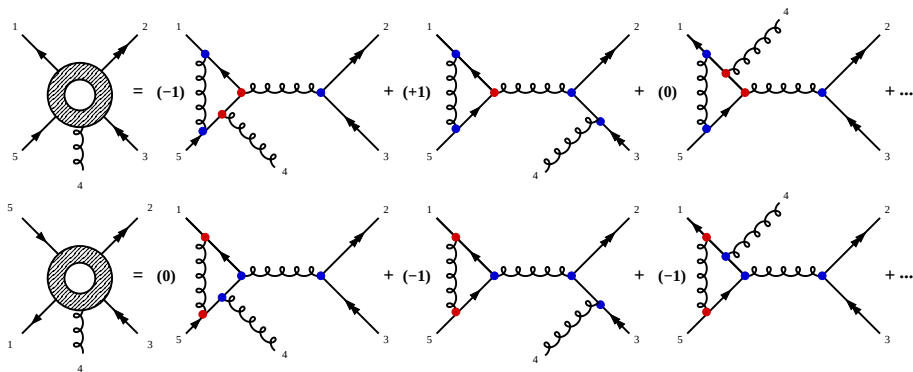
$$P_i = \sum_{j=1}^{\hat{N}_{\text{dia}}} M_{ij} K_j \quad i \in \{1, 2, \dots, N_{\text{pri}}\} \quad N_{\text{pri}} = N_{\text{pri}}^{[m]} + N_{\text{pri}}^{[f]}$$

$$N_{\text{pri}}^{[f]} = (n-1)! \quad N_{\text{pri}}^{[m]} = \begin{cases} (n-1)! & n_q = 0 \\ n_q(n-1)!/2 & n_q = 2, 4, \dots \end{cases}$$

¹only topologies are needed

Matching Matrix M_{ij}

$$P_i = \sum_{j=1}^{\hat{N}_{\text{dia}}} M_{ij} K_j \quad M_{ij} \in \{0, 1, -1\}$$



Matching of $A_5^{[m]}(1_{\bar{d}}, 2_u, 3_{\bar{u}}, 4_g, 5_d)$ and $A_5^{[m]}(1_{\bar{d}}, 4_g, 3_{\bar{u}}, 2_u, 5_d)$.

Each vertex is either **ordered** or **unordered** with respect to the colour ordered Feynman rules and the primitive in question.

$$\textcolor{red}{P}_i = \sum_{j=1}^{\hat{N}_{\text{dia}}} \textcolor{green}{M}_{ij} \textcolor{blue}{K}_j \quad M_{ij} \in \{0, 1, -1\}$$

Number of **independent** primitive amplitudes (denoted $\hat{P}_j$)

$$\hat{N}_{\text{pri}} = \mathbf{rank} \textcolor{green}{M}$$

$$\hat{N}_{\text{pri}} \leq (N_{\text{pri}} = N_{\text{rows}}) \quad \hat{N}_{\text{pri}} \leq (\hat{N}_{\text{dia}} = N_{\text{cols}})$$

Reduced row echelon form of $\hat{\mathbf{M}} = [\textcolor{green}{M} | -\mathbb{1}]$

- ▶ upper $\hat{N}_{\text{pri}}$ rows — **solution** of $\textcolor{blue}{K}_j$ in terms of $\textcolor{red}{P}_i$
- ▶ lower $N_{\text{pri}} - \hat{N}_{\text{pri}}$ rows — left null space of $\textcolor{green}{M}$ (relations)

$$\textcolor{blue}{K}_i = \sum_{j=1}^{\hat{N}_{\text{pri}}} B_{ij} \textcolor{red}{P}_j \quad \{\hat{P}_j\}_{\hat{N}_{\text{pri}}} \subset \{P_j\}_{N_{\text{pri}}}$$

Putting everything together

- ▶ **Colour factors** in terms of the colour “trace basis”
- ▶ **Kinematic factors** in terms of independent primitives

$$\begin{aligned} C_i &= \sum_c T_c F_{ci} & K_i &= \sum_{j=1}^{\hat{N}_{\text{pri}}} B_{ij} \hat{P}_j \\ A_n &= \sum_{i=1}^{\hat{N}_{\text{dia}}} C_i K_i = \sum_c T_c \sum_{j=1}^{\hat{N}_{\text{pri}}} \underbrace{\sum_{i=1}^{\hat{N}_{\text{dia}}} F_{ci} B_{ij}}_{Q_{cj}} \hat{P}_j \end{aligned}$$

We obtain **partial amplitudes** in terms of a basis of independent primitive amplitudes $\hat{P}_j$ for a given class of primitives

$$A_n = \sum_c T_c \sum_{j=1}^{\hat{N}_{\text{pri}}} Q_{cj} \hat{P}_j$$

Number of primitives in tree, mixed and fermion loop amplitudes

Process	$N_{\text{pri}}^{[0]}$	$N_{\text{pri}}^{[m]}$	$N_{\text{pri}}^{[f]}$	Process	$N_{\text{pri}}^{[0]}$	$N_{\text{pri}}^{[m]}$	$N_{\text{pri}}^{[f]}$
$4g$	2	3	3	$5g$	6	12	12
$\bar{u}u + 2g$	2	6	1	$\bar{u}u + 3g$	6	24	6
$\bar{u}udd$	1	4	1	$\bar{u}uddg$	3	16	3
Process	$N_{\text{pri}}^{[0]}$	$N_{\text{pri}}^{[m]}$	$N_{\text{pri}}^{[f]}$	Process	$N_{\text{pri}}^{[0]}$	$N_{\text{pri}}^{[m]}$	$N_{\text{pri}}^{[f]}$
$6g$	24	60	60	$7g$	120	360	360
$\bar{u}u + 4g$	24	120	33	$\bar{u}u + 5g$	120	720	230
$\bar{u}udd + 2g$	12	80	13	$\bar{u}udd + 3g$	60	480	75
$\bar{u}udd\bar{s}s$	4	32	4	$\bar{u}udd\bar{s}sg$	20	192	20

Example: fermion loop contribution to the six-quark amplitude

$$A_{6q}(1_{\bar{u}}, 2_u, 3_{\bar{d}}, 4_d, 5_{\bar{s}}, 6_s) = A_{6q}^{\text{mixed}} + n_f A_{6q}^{\text{q-loop}}$$

$$A_{6q}^{\text{q-loop}} = A^{[f]}(12\textcolor{blue}{3}\textcolor{green}{4}56) \cdot (T_3 + T_4 - T_1 N_c - T_6 / N_c) / N_c + A^{[f]}(12\textcolor{blue}{4}3\textcolor{green}{6}5) \cdot (T_5 - T_3) / N_c \\ + A^{[f]}(12\textcolor{blue}{5}6\textcolor{green}{3}4) \cdot (2T_5 - T_2 N_c - T_6 / N_c) / N_c + A^{[f]}(12\textcolor{blue}{6}5\textcolor{green}{3}4) \cdot (T_5 - T_4) / N_c$$

$$T_1 = \delta_{16} \delta_{32} \delta_{54}$$

$$T_2 = \delta_{14} \delta_{36} \delta_{52}$$

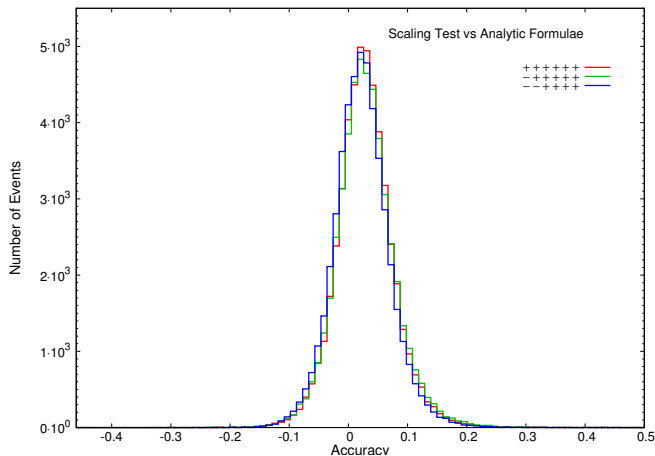
$$T_3 = \delta_{12} \delta_{36} \delta_{54}$$

$$T_4 = \delta_{14} \delta_{32} \delta_{56}$$

$$T_5 = \delta_{16} \delta_{34} \delta_{52}$$

$$T_6 = \delta_{12} \delta_{34} \delta_{56}$$

Scaling test to determine accuracy

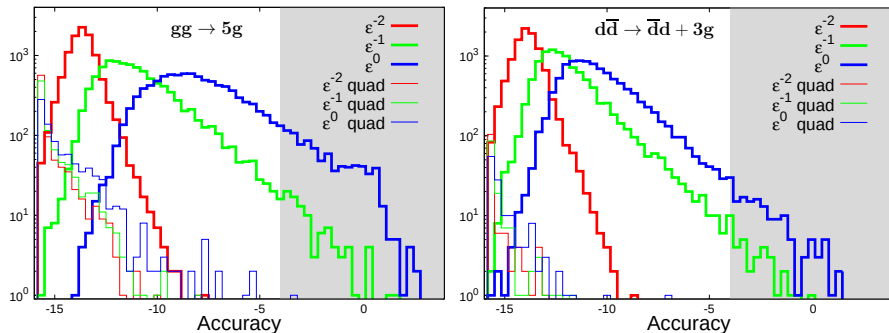


$$\log \left(\frac{A_{\text{NGluon}} - A_{\text{analytic}}}{A_{\text{analytic}}} \right) - \log \left(\frac{2(A_{\text{NGluon}} - A_{\text{NGluon}}^{\text{scaled}})}{A_{\text{NGluon}} + A_{\text{NGluon}}^{\text{scaled}}} \right)$$

Reliable, but essentially statistical.
A safety margin of 2 digits is advised.

Scaling test of 5 jet amplitudes

Left: 7 gluon squared amplitude. Right: 4 quark + 3 gluons.



Thick lines – double precision.

Thin lines – fixed with quadruple precision.

Full colour and helicity sum time per point [clang, Xeon 3.30 GHz].

process	$T_{sd}[s]$	T_4 dig. [s] (%)	process	$T_{sd}[s]$	T_4 dig. [s] (%)
$4g$	0.030	0.030 (0.00)	$5g$	0.22	0.22 (0.22)
$\bar{u}u+2g$	0.032	0.032 (0.00)	$\bar{u}u+3g$	0.34	0.35 (0.06)
$\bar{u}u\bar{d}d$	0.011	0.011 (0.00)	$\bar{u}u\bar{d}d+g$	0.11	0.11 (0.00)
$\bar{u}u\bar{u}u$	0.022	0.022 (0.00)	$\bar{u}u\bar{u}u+g$	0.22	0.22 (0.03)
process	$T_{sd}[s]$	T_4 dig. [s] (%)	process	$T_{sd}[s]$	T_4 dig. [s] (%)
$6g$	6.19	6.81 (1.37)	$7g$	171.3	276.7 (8.63)
$\bar{u}u+4g$	7.19	7.40 (0.38)	$\bar{u}u+5g$	195.1	241.2 (3.25)
$\bar{u}u\bar{d}d+2g$	2.05	2.06 (0.08)	$\bar{u}u\bar{d}d+3g$	45.7	48.8 (0.88)
$\bar{u}u\bar{u}u+2g$	4.08	4.15 (0.21)	$\bar{u}u\bar{u}u+3g$	92.5	101.5 (1.29)
$\bar{u}u\bar{d}d\bar{s}s$	0.38	0.38 (0.00)	$\bar{u}u\bar{d}d\bar{s}sg$	7.9	8.1 (0.23)
$\bar{u}u\bar{d}d\bar{d}d$	0.74	0.74 (0.00)	$\bar{u}u\bar{d}d\bar{d}dg$	15.8	16.2 (0.29)
$\bar{u}u\bar{u}u\bar{u}u$	2.16	2.17 (0.02)	$\bar{u}u\bar{u}u\bar{u}ug$	47.1	48.6 (0.41)

All times include two evaluations for the scaling test.

Both the phase-space and the amplitude are symmetric under permutations of the final state gluons:

$$\int dPS_n(g_1, \dots, g_n) = \int dPS_n(\sigma\{g_1, \dots, g_n\})$$

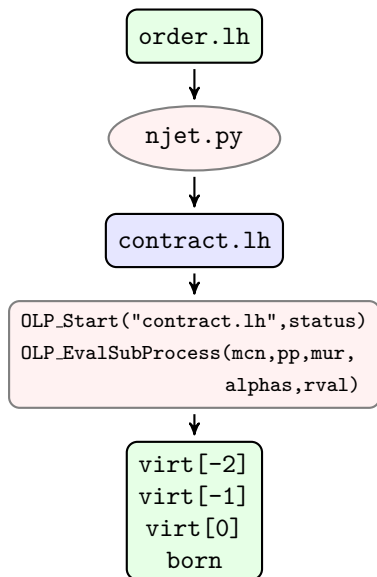
$$\mathcal{A}_{gg \rightarrow n(g)}(g, g, g_1, \dots, g_n) = \mathcal{A}_{gg \rightarrow n(g)}(g, g, \sigma\{g_1, \dots, g_n\})$$

We can replace symmetric colour sum with a **simpler object** which gives original sum after symmetrization (or PS integration)

$$\sigma_{gg \rightarrow n(g)}^V = \int dPS_n \left[A^{(0)\dagger} \cdot \mathcal{C}_{(n-2)! \times (n-1)!/2} \cdot A^{(1)} \right]$$

$$= (n-2)! \int dPS_n \left[A^{(0)\dagger} \cdot \mathcal{C}_{(n-2)! \times (n-1)}^{\text{dsym}} \cdot A^{(1), \text{dsym}} \right]$$

	$gg \rightarrow 3g$	$gg \rightarrow 4g$	$gg \rightarrow 5g$
Standard sum	0.22	6.19	171.31
De-symmetrized	0.07	0.50	2.76
Speedup	$\times 3$	$\times 12$	$\times 60$



Create an 'order' file

NJet takes an 'order' file
and returns a 'contract' file

Check that requested
options were accepted

Link with `libnjet.so`

Call 'Start' once to initialize

Call 'EvalSubProcess' to evaluate PS points

Use returned values to calculate XS

$1/eps^2$, $1/eps$, *finite*, *born*

order file

```
# OLE_order for 5jet production
```

```
MatrixElementSquareType CHsummed
CorrectionType QCD
IRregularisation CDR
AlphasPower 5
# process list
21 21 -> 21 21 21 21 21
1 -1 -> 21 21 21 21 21
1 -1 -> 21 -2 2 21 21
1 -1 -> 21 -1 1 21 21
1 -1 -> 21 -2 2 -3 3
1 -1 -> 21 -2 2 -2 2
1 -1 -> 21 -1 1 -1 1
```

contract file

```
# OLE_order for 5jet production
# Generated file. Do not edit by hand.
# Signed by NJet 3900867518.
# 12 1 1e-05 0.01 0 1 1 1 1 0 3 5
MatrixElementSquareType CHsummed | OK
CorrectionType QCD | OK
IRregularisation CDR | OK
AlphasPower 5 | OK
# process list
21 21 -> 21 21 21 21 21 | 1 1 # 70 120 4 64 0 (-2 -1 3 4 5 6 7)
1 -1 -> 21 21 21 21 21 | 1 2 # 71 120 4 9 0 (-1 -2 3 4 5 6 7)
1 -1 -> 21 -2 2 21 21 | 1 3 # 72 6 4 9 0 (-1 -2 4 5 3 6 7)
1 -1 -> 21 -1 1 21 21 | 1 4 # 73 6 4 9 0 (-1 -2 4 5 3 6 7)
1 -1 -> 21 -2 2 -3 3 | 1 5 # 74 1 4 9 0 (-1 -2 4 5 6 7 3)
1 -1 -> 21 -2 2 -2 2 | 1 6 # 75 4 4 9 0 (-1 -2 4 5 6 7 3)
1 -1 -> 21 -1 1 -1 1 | 1 7 # 76 4 4 9 0 (-1 -2 4 5 6 7 3)
```

Additional options

- ▶ NJetSwitchAcc — relative accuracy threshold
- ▶ NJetReturnAccuracy — return accuracy estimate
- ▶ NJetType — loop, tree or loops
- ▶ NJetNf — number of light flavours
- ▶ NJetNc — $SU(N_c)$ group parameter

Tools (linked together with BLHA interface)

- ▶ NJet — full colour virtual matrix elements
scalar integrals — QCDLoop/FF [Ellis,Zanderighi,van Oldenborgh]
extended precision — libqd [Hida,Li,Bailey]
- ▶ Sherpa/Amegic++ — trees, CS subtraction, PS integration [Gleisberg,Krauss,Kuhn,Soff,...]

Parameters

- ▶ $pp \rightarrow 3$ and 4 jets at 8 TeV
- ▶ MSTW2008 PDF set, $\alpha_s(M_Z)$ from PDFs
- ▶ $\mu_R = \mu_F = \hat{H}_T/2$, scale variations $\hat{H}_T/4$ and $\hat{H}_T$
- ▶ anti-kt $R = 0.4$, $p_T^{1st} > 80$ GeV, $p_T^{other} > 60$ GeV, $\eta < 2.8$

Total XS for 3 and 4 jet at 8 TeV

	3 jets	4 jets
σ_{LO}	$126.65(0.05)^{+66.56}_{-40.40} \text{ nb}$	$14.36(0.01)^{+10.38}_{-5.6} \text{ nb}$
σ_{NLO}	$72.57(0.16)^{+2.71}_{-28.08} \text{ nb}$	$8.15(0.09)^{+0.0}_{-3.24} \text{ nb}$
$\sigma_{8 \text{ TeV}}/\sigma_{7 \text{ TeV}}$	$\sim 35\%$	$\sim 46\%$

Reduced scale uncertainty

NLO scale variations are about 25% of the LO scale uncertainty for both 3 and 4 jet cross-sections

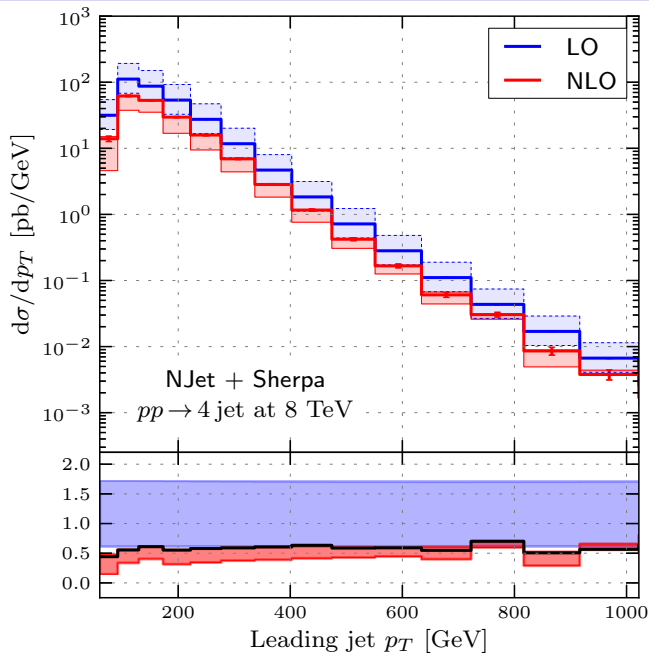
Cross-check at 7 TeV

Agreement with first 4 jet results by BlackHat collaboration

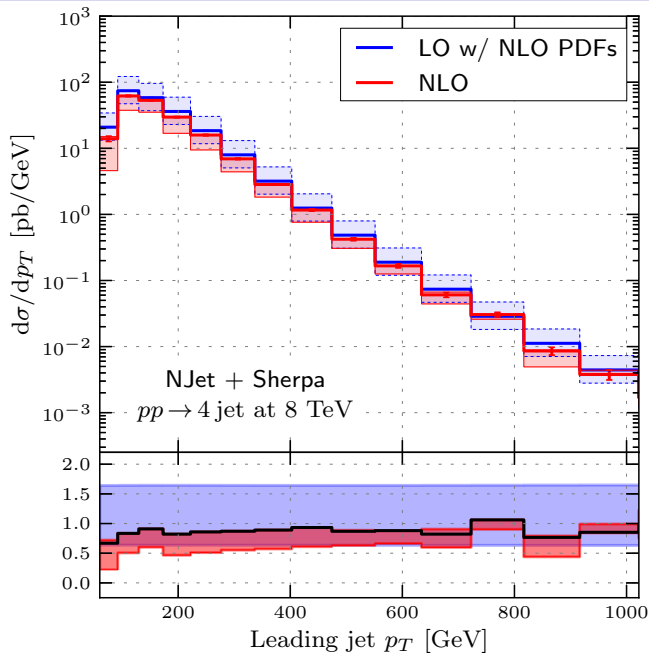
[Bern,Diana,Dixon,Febres Cordero,Hoeche,Kosower,Ita,Maitre,Ozeren]

[arXiv:1112.3940]

4 jets, leading jet p_T at NLO

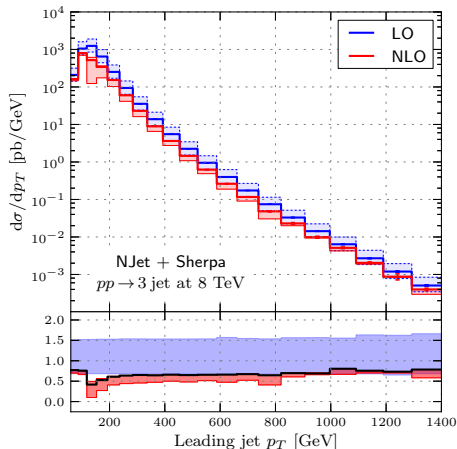


4 jets, leading jet p_T at LO with NLO PDFs and alphas

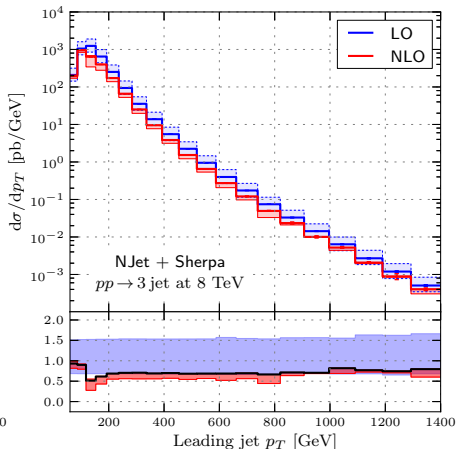


3 jets, leading jet pT at NLO, $\hat{H}_T$ vs H_T

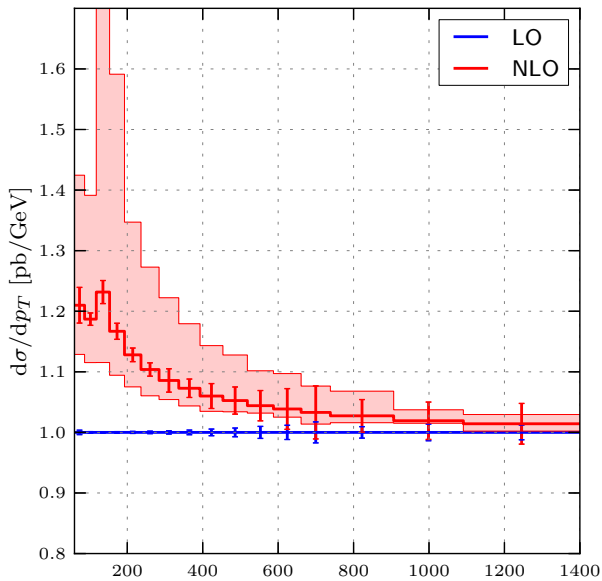
$$\frac{\hat{H}_T}{2} \\ 72.57(0.16)^{+2.71}_{-28.08} \text{ nb}$$



$$\frac{H_T}{2} \\ 85.92(0.17)^{+0.0}_{-16.72} \text{ nb}$$



3 jets, leading jet p_T at NLO, $d\sigma(H_T)/d\sigma(\hat{H}_T)$



Summary

- ▶ NJet: numerical evaluation of one-loop partonic amplitudes in massless QCD.
- ▶ General construction for primitive and partial amplitudes.
- ▶ Full colour results for ≤ 5 jets (+ fast 'desymmetrized' colour sums for gluons).
- ▶ Binoth Les Houches Accord interface.
- ▶ NJet+Sherpa: 3 and 4 jets NLO at 8 TeV.
- ▶ More applications feasible: 5 jets, parton shower matching. . .
- ▶ Publicly available from the NJet project page
<https://bitbucket.org/njet/njet>

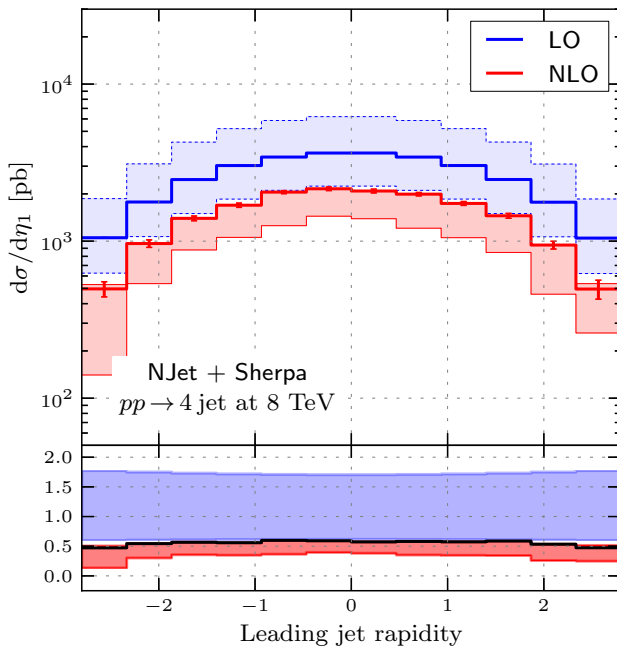
Bonus material

Desymmetrized 5 gluon amplitude

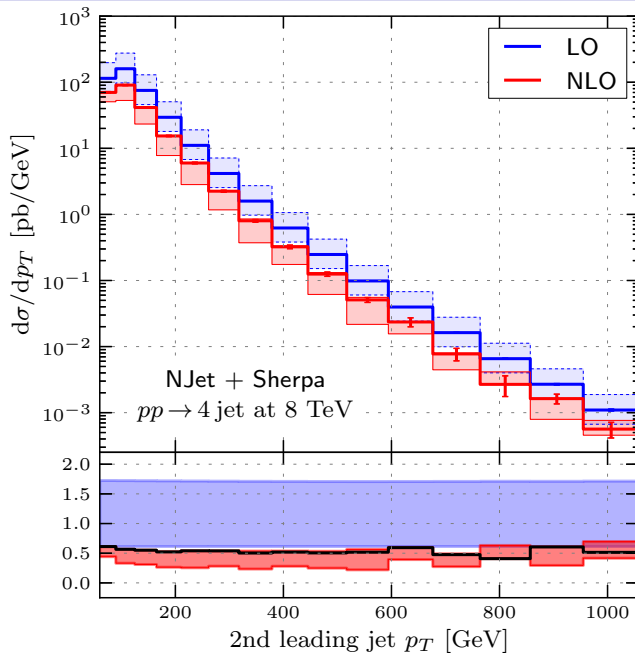
Desymmetrized sum contains 4 loop primitives of each kind (**mixed** and **q-loop**)

$$\begin{aligned}
 & \mathcal{A}_{5g}^{(0)\dagger} \otimes^{\text{dsym}} \mathcal{A}_{5g}^{(1)} = 6N_c(N_c^2 - 1) \times \\
 & \times \left\{ N_c \left[\begin{aligned}
 & \mathbf{A}^{[m]}(\mathbf{12345}) \left(12(A^{[0]}(12453) + A^{[0]}(12534) + A^{[0]}(12543)) + A^{[0]}(12345)N_c^2 \right)^\dagger \\
 & + \mathbf{A}^{[m]}(\mathbf{13245}) \left(12A^{[0]}(12534) - (A^{[0]}(12345) + A^{[0]}(12435) + A^{[0]}(12453))N_c^2 \right)^\dagger \\
 & + \mathbf{A}^{[m]}(\mathbf{13425}) \left(-12A^{[0]}(12354) + (A^{[0]}(12435) + A^{[0]}(12453) + A^{[0]}(12543))N_c^2 \right)^\dagger \\
 & + \mathbf{A}^{[m]}(\mathbf{13452}) \left(-12(A^{[0]}(12345) + A^{[0]}(12354) + A^{[0]}(12435)) - A^{[0]}(12543)N_c^2 \right)^\dagger \right] \\
 & - N_f \left[\begin{aligned}
 & \mathbf{A}^{[f]}(\mathbf{12345}) \left(2(A^{[0]}(12453) + A^{[0]}(12534) + A^{[0]}(12543)) + A^{[0]}(12345)N_c^2 \right)^\dagger \\
 & + \mathbf{A}^{[f]}(\mathbf{13245}) \left(2A^{[0]}(12534) - (A^{[0]}(12345) + A^{[0]}(12435) + A^{[0]}(12453))N_c^2 \right)^\dagger \\
 & + \mathbf{A}^{[f]}(\mathbf{13425}) \left(-2A^{[0]}(12354) + (A^{[0]}(12435) + A^{[0]}(12453) + A^{[0]}(12543))N_c^2 \right)^\dagger \\
 & + \mathbf{A}^{[f]}(\mathbf{13452}) \left(-2(A^{[0]}(12345) + A^{[0]}(12354) + A^{[0]}(12435)) - A^{[0]}(12543)N_c^2 \right)^\dagger \right] \right\}
 \end{aligned}
 \right.
 \end{aligned}$$

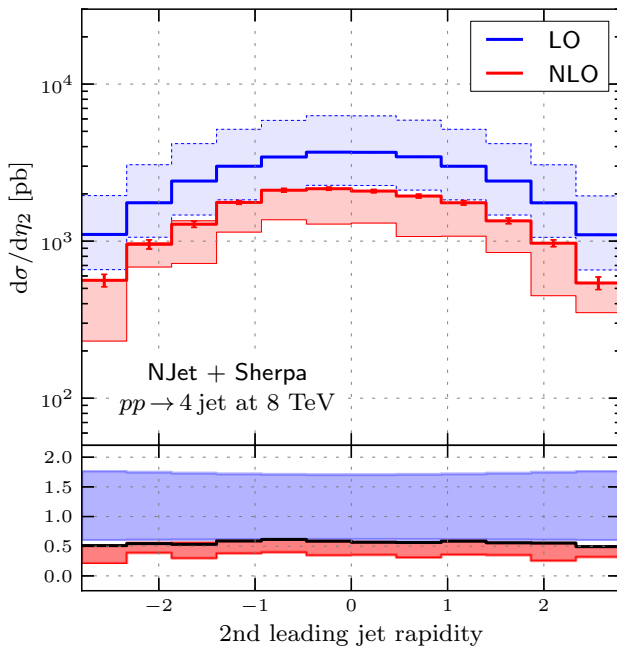
4 jets, leading jet rapidity at NLO



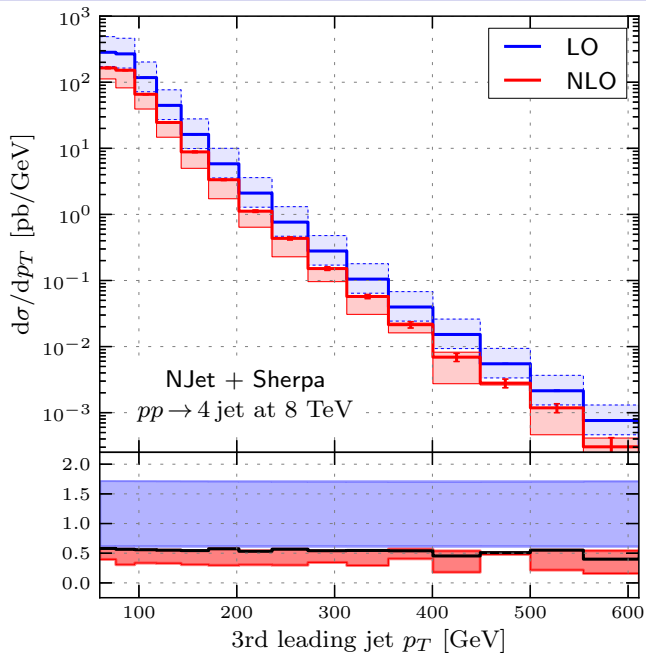
4 jets, 2nd jet p_T at NLO



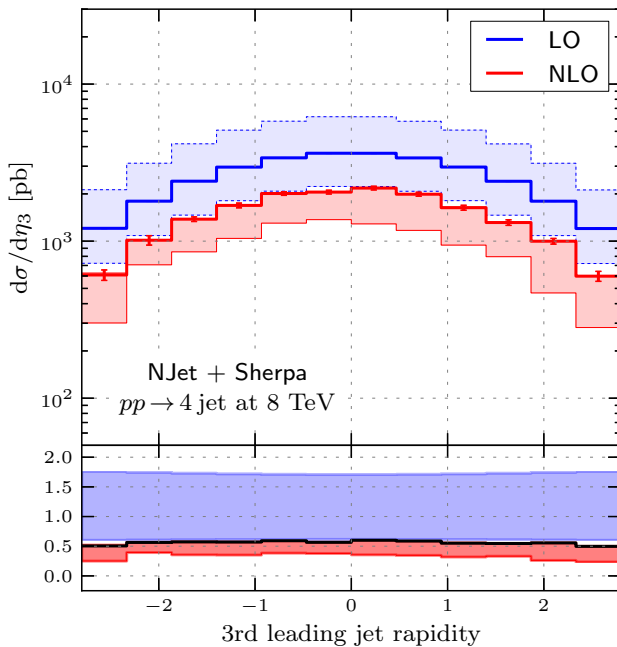
4 jets, 2nd jet rapidity at NLO



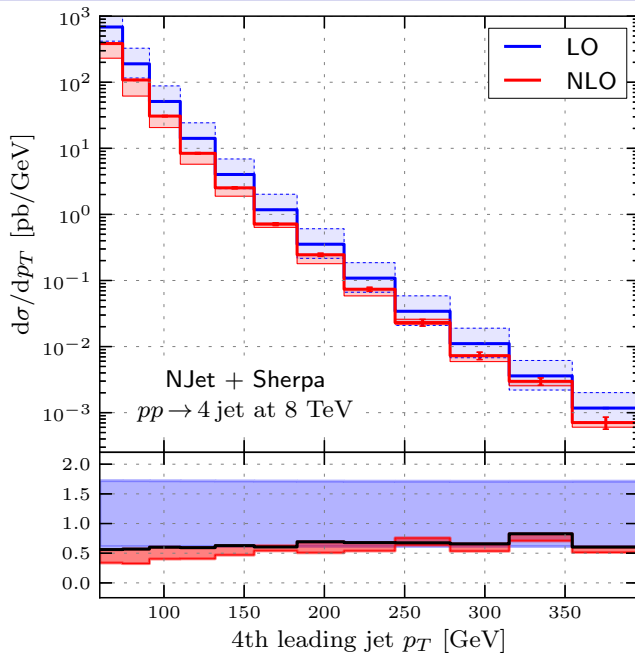
4 jets, 3d jet p_T at NLO



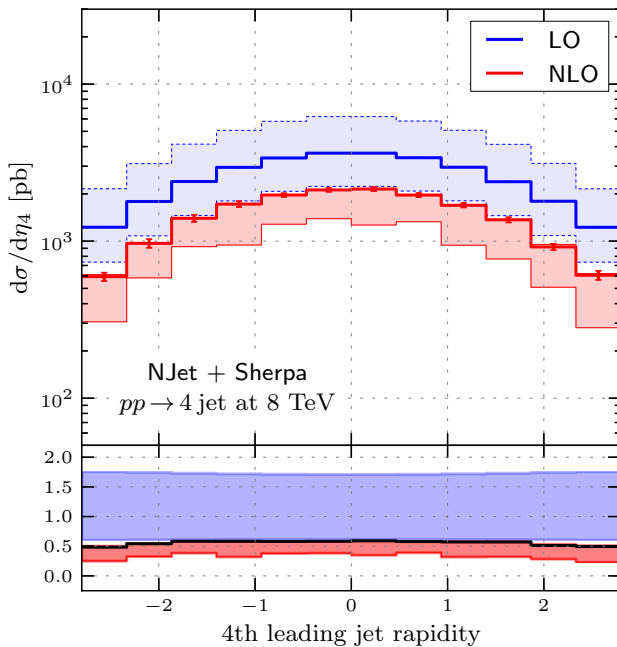
4 jets, 3rd jet rapidity at NLO



4 jets, 4th jet p_T at NLO

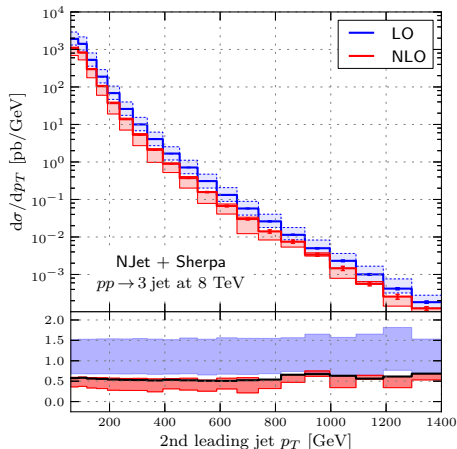


4 jets, 4th jet rapidity at NLO

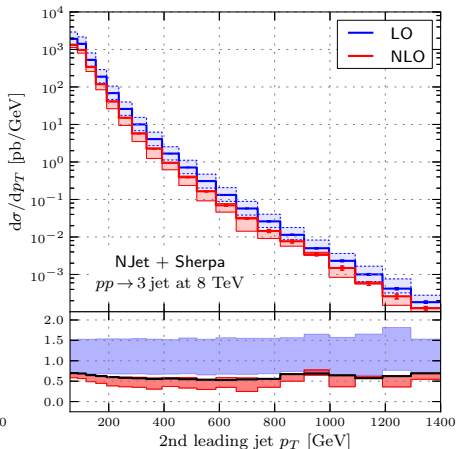


3 jets, 2nd jet pT at NLO, $\hat{H}_T$ vs H_T

$$\frac{\hat{H}_T}{2} \\ 72.57(0.16)^{+2.71}_{-28.08} \text{ nb}$$

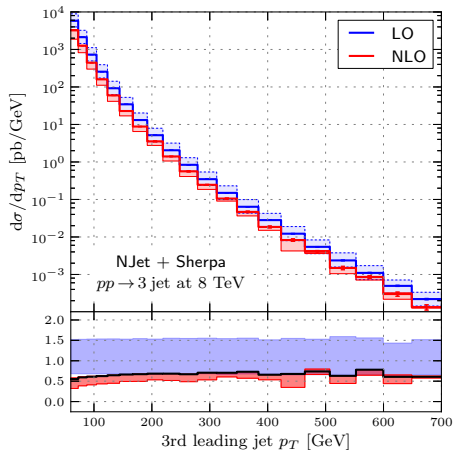


$$\frac{H_T}{2} \\ 85.92(0.17)^{+0.0}_{-16.72} \text{ nb}$$

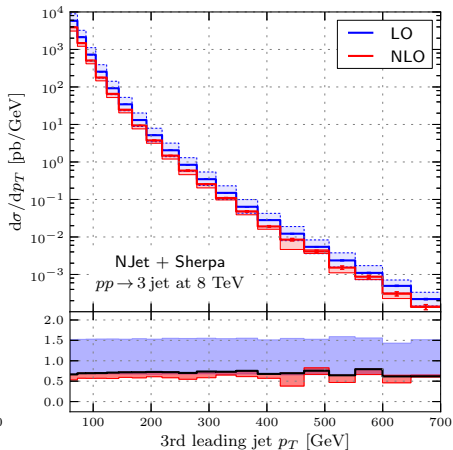


3 jets, 3rd jet pT at NLO, $\hat{H}_T$ vs H_T

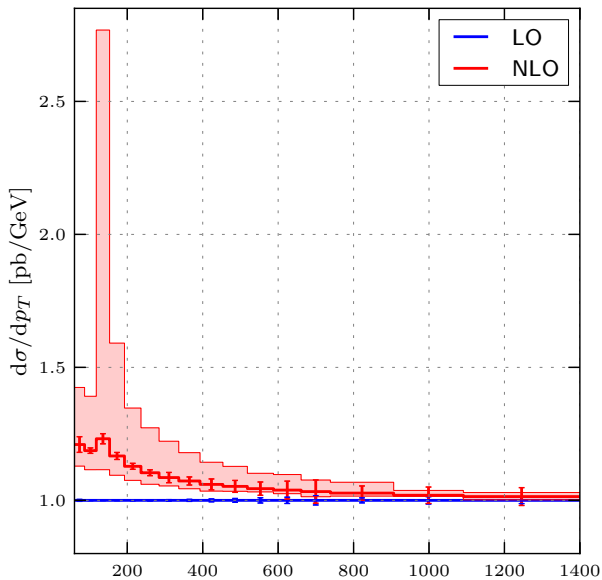
$$\frac{\hat{H}_T}{2}$$
$$72.57(0.16)^{+2.71}_{-28.08} \text{ nb}$$



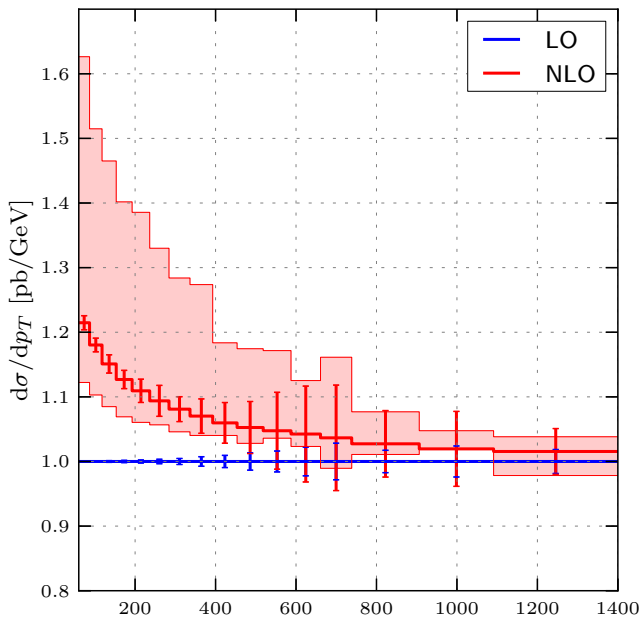
$$\frac{H_T}{2}$$
$$85.92(0.17)^{+0.0}_{-16.72} \text{ nb}$$



3 jets, leading jet p_T at NLO, $d\sigma(H_T)/d\sigma(\hat{H}_T)$



3 jets, 2nd jet pT at NLO, $d\sigma(H_T)/d\sigma(\hat{H}_T)$



3 jets, 3rd jet pT at NLO, $d\sigma(H_T)/d\sigma(\hat{H}_T)$

