

Higher order QCD corrections for associated VH production at hadron colliders

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In collaboration with: M. Grazzini & F. Tramontano

Outline

- 1 Associated VH production at hadron colliders
- 2 q_T -subtraction formalism at NNLO
- 3 Associated VH production at NNLO: numerical results
- 4 Conclusions

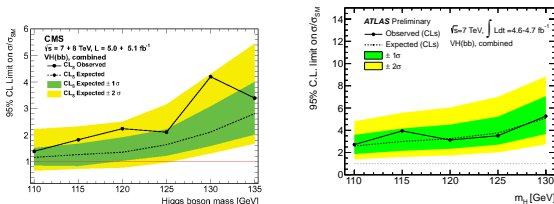


Motivations

Associated vector boson Higgs (VH) production (with $H \rightarrow b\bar{b}$ and $V \rightarrow l_1 l_2$ decay) is an important mechanism for discovery and study the properties of the Higgs boson.

- At the LHC it is important channel through boosted analysis with jet reconstruction and decomposition techniques [Butterworth et al. ('08)].
- At the Tevatron is the main search channel in the low Higgs mass region.

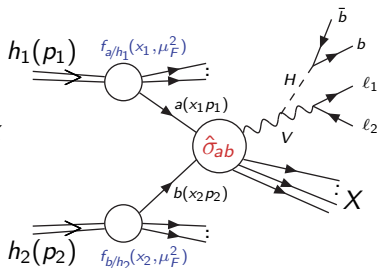
To get closer to SM VH sensitivity with the LHC 2012 data, precise theoretical predictions needed \implies computation of higher-order QCD corrections.



Associated VH production

$$h_1(p_1) + h_2(p_2) \rightarrow V + H + X \rightarrow l_1 l_2 + b\bar{b} + X$$

$$\text{where } V = Z^0, W^\pm \quad \text{and} \quad l_1 l_2 = e^+ e^-, \nu \bar{\nu}$$



According to the QCD factorization theorem:

$$d\sigma(p_1, p_2) = \sum_{a,b} \int_0^1 dx_1 \int_0^1 dx_2 f_{a/h_1}(x_1, \mu_F^2) f_{b/h_2}(x_2, \mu_F^2) d\hat{\sigma}_{ab}(x_1 p_1, x_2 p_2; \mu_F^2).$$

$$d\hat{\sigma}_{ab}(\hat{p}_1, \hat{p}_2; \mu_F^2) = d\hat{\sigma}_{ab}^{(0)}(\hat{p}_1, \hat{p}_2; \mu_F^2) + \alpha_S(\mu_R^2) d\hat{\sigma}_{ab}^{(1)}(\hat{p}_1, \hat{p}_2; \mu_F^2) \\ + \alpha_S^2(\mu_R^2) d\hat{\sigma}_{ab}^{(2)}(\hat{p}_1, \hat{p}_2; \mu_F^2, \mu_R^2) + \mathcal{O}(\alpha_S^3).$$

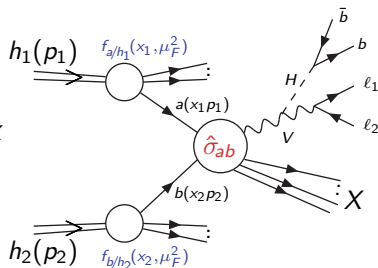
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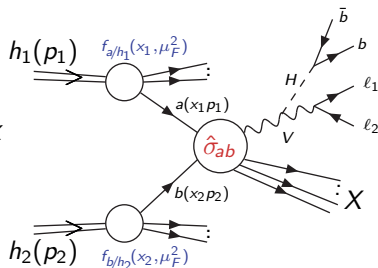
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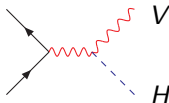
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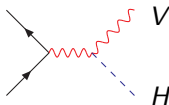
Associated VH production: total cross section



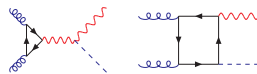
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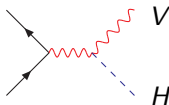
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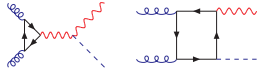
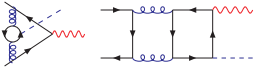


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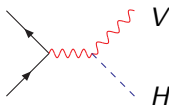
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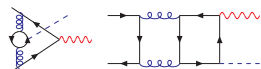
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- Experiments have finite acceptance, in particular VH experimental analyses performed in extreme kinematical regimes (e.g. boosted analysis with jet veto): **important to provide exclusive theoretical predictions.**
- At NLO general algorithms (e.g. Dipole formalism [Catani, Seymour('98)]) allow (relative) straightforward fully-exclusive calculations.
- At NNLO in hadronic collisions only few fully exclusive calculations exist:
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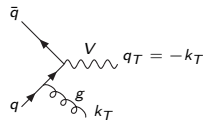
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V is one or more **colourless** particles (vector bosons, leptons, photons, Higgs bosons, ...) [Catani, Grazzini('07)].

- **Key point I:** at LO the q_T of the V is exactly zero.

$$d\sigma_{(N)NLO}^V|_{q_T \neq 0} = d\sigma_{(N)LO}^{V+jets},$$



for $q_T \neq 0$ the NNLO IR divergences cancelled with the NLO subtraction method.

- The only remaining NNLO singularities are associated with the $q_T \rightarrow 0$ limit.
- **Key point II:** treat the NNLO singularities at $q_T = 0$ by an additional subtraction using the universality of logarithmically-enhanced contributions from q_T resummation formalism [Catani, de Florian, Grazzini('00)].

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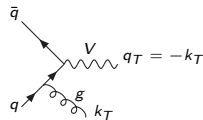
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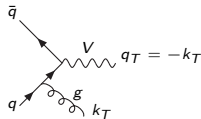
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- The choice of the counter-term has some arbitrariness but it must behave $d\sigma^{CT} \xrightarrow{q_T \rightarrow 0} d\sigma_{LO}^V \otimes \Sigma(q_T/M) dq_T^2$ where $\Sigma(q_T/M)$ is universal.
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Associated WH production in NNLO QCD

G.F., Grazzini, Tramontano arXiv:1107.1164

- A NLO calculation for $h_1 h_2 \rightarrow V + X$ requires:
 - $d\sigma_{LO}^{V+\text{jets}}$ (and $d\sigma_{LO}^V$).
 - $\mathcal{H}^{V(1)}$ [de Florian, Grazzini('01)]: contains the finite-part of the one-loop amplitude $c\bar{c} \rightarrow V$.
 - $d\sigma_{LO}^{CT}$: depends by the (universal) q_T -resummation coeff. A_1 and B_1 .
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- WH production at NNLO within q_T -subtraction:
 - $d\sigma_{NLO}^{WH+\text{jets}}$.
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Fully-exclusive NNLO calculation, implemented in the parton-level Monte

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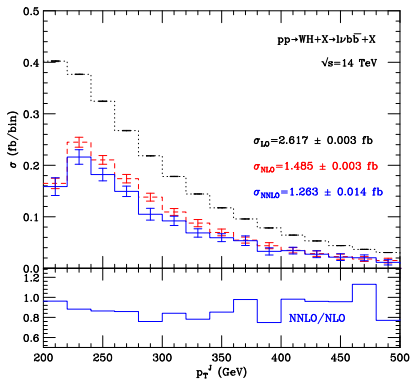
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Numerical results at the LHC and the Tevatron





$pp \rightarrow WH(\rightarrow l\nu b\bar{b})$

p_T spectra of the fat jet at the LHC@14TeV for $m_H = 120\text{GeV}$ at LO (dots), NLO (dashes) and NNLO (solid).

- Selection strategy of [Butterworth et al. ('08)]: search a large- p_T Higgs boson through a collimated $b\bar{b}$ pair decay.

Cuts:

Leptons: $p_T^l > 30\text{GeV}$, $|\eta^l| < 2.5$,

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Jets: Cambridge/Aachen algorithm with $R=1.2$.

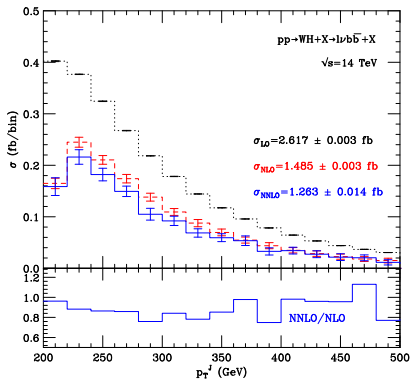
Fat jet (contain the $b\bar{b}$) $p_T^J > 200\text{GeV}$,

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Jet veto: No other jets with $p_T > 20\text{GeV}$ and $|\eta| < 5$.

- Large negative higher-order corrections: NLO (NNLO) effects -52%/-36% (-6%/-19%), depending on the scale choice (factor two around $\mu_F = \mu_R = m_W + m_H$).
- Jet veto strongly affect the higher order corrections \Rightarrow stability of fixed order calculation challenged.



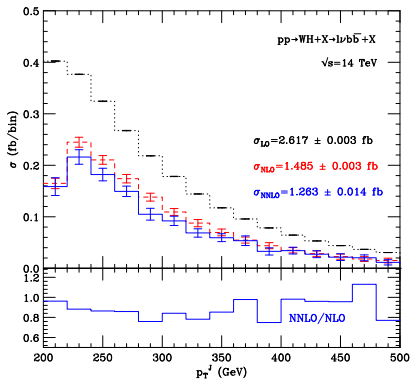


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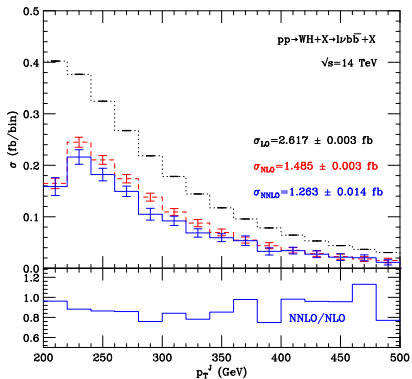


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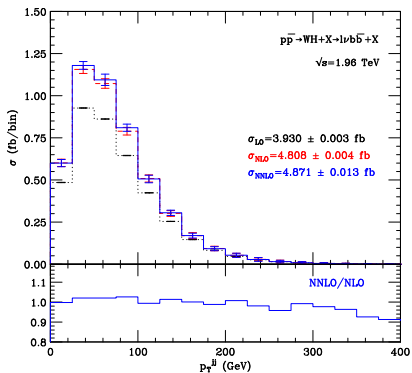


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$p\bar{p} \rightarrow WH(\rightarrow l\nu b\bar{b})$

p_T spectra of the dijet system at the Tevatron for $m_H = 120\text{GeV}$ at LO (dots), NLO (dashes) and NNLO (solid).

- Cuts:

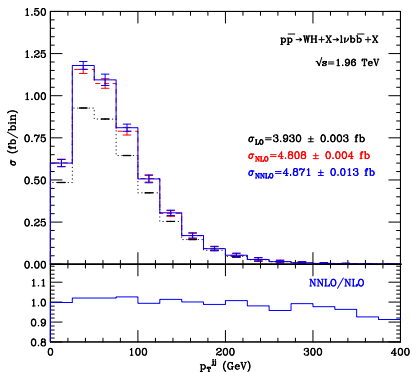
Leptons: $p_T^l > 20\text{GeV}$, $|\eta^l| < 2$, $p_T^{\text{miss}} > 20\text{GeV}$.

Jets: k_T algorithm with $R=0.4$.

Exactly two jets (with $p_T > 20\text{GeV}$ and $|\eta| < 2$) at least one of them has to be a b jet (with $|\eta| < 1$).

- Higher-order corrections: NLO (NNLO) effects from +13% to +30% (from -1% to +4%) depending on the scale choice (factor two around $\mu_F = \mu_R = m_W + m_H$). The scale dependence is at the level of about $\pm 1\%$ both at NLO and NNLO.
- The shape of the distribution is stable against perturbative corrections. Perturbative expansion under good control.



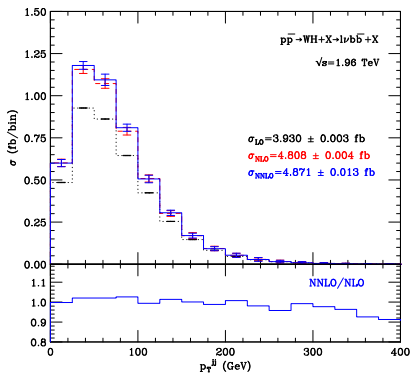


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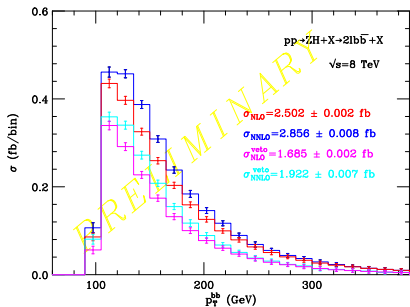
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NEW: associated ZH production at NNLO:

G.F., Grazzini, Tramontano (in preparation)



$pp \rightarrow ZH(\rightarrow 2l b\bar{b})$

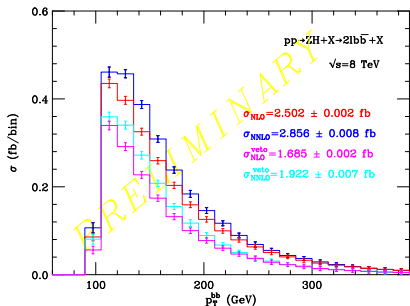
p_T spectra of the $b\bar{b}$ system at the LHC for $m_H = 125\text{GeV}$ at NLO and NNLO with and without the jet veto.

- $gg \rightarrow HZ$ top-loop $\sim g^2 \lambda_t^2 \alpha_S^2$ (non DY-like) corrections included.
- Cuts (we follow CMS analysis):
Leptons: $p_T^l > 20\text{GeV}$, $|\eta^l| < 2.5$,
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Jets: anti- k_T algorithm with $R=0.5$.
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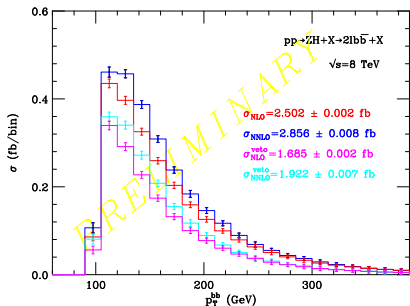
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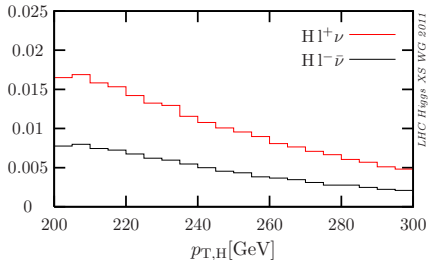
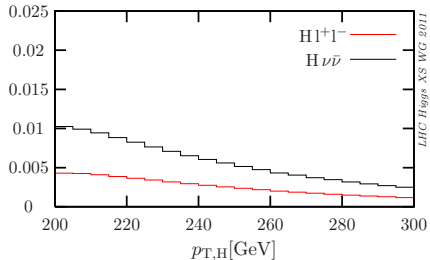


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$d\sigma/dp_{T,H}[\text{fb}/\text{GeV}]$

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Yellow Report II:arXiv:1201.3084

Distributions in $p_{T,H}$ for $pp \rightarrow WH \rightarrow l\nu H$ (NNLO QCD + NLO EW) and for $pp \rightarrow ZH \rightarrow ll\nu\nu H$ (NLO QCD + NLO EW) at $\sqrt{s} = 7 \text{ TeV}$.

Boosted setup: $|\eta_l| < 2.5$, $p_{T,l} > 20 \text{ GeV}$, $p_{T,\nu} > 25 \text{ GeV}$, $p_{T,H} > 200 \text{ GeV}$, $p_{T,W/Z} > 190 \text{ GeV}$.

We produced similar results at $\sqrt{s} = 8 \text{ TeV}$.



Conclusions

- Calculation of **NNLO QCD** corrections to **VH production** in hadron collision using the **q_T -subtraction** formalism, included in a **fully-exclusive** parton-level Monte Carlo code.
- NNLO corrections can be important:
large and negative: $\sim -20\%$ for the WH fat-jet analysis at the LHC@14 TeV when a jet veto is applied;
 sizeable for the CMS analysis at the LHC;
 moderate for the WH Tevatron analysis.
- **Outlook/Work in progress:**
 Public release of the parton-level numerical code.
 Inclusion of the higher-order QCD correction to $H \rightarrow b\bar{b}$ decay.
 Extension to the ZH production.
 Inclusion of $H \rightarrow WW/ZZ \rightarrow 2l2\nu/4l$ decay.
 Comparison with parton-shower Monte Carlo predictions.
 Study of the NNLO uncertainty band: **first reliable estimate** of perturbative uncertainty.

