Multi-Loop Integrand Reduction Techniques

Simon Badger (NBIA & Discovery Center)

7^{th} September 2012

HP2 2012, Munich, Germany

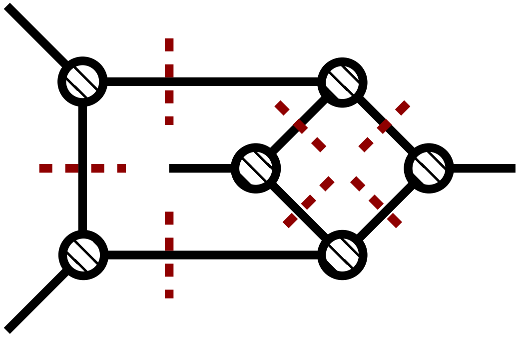
Outline

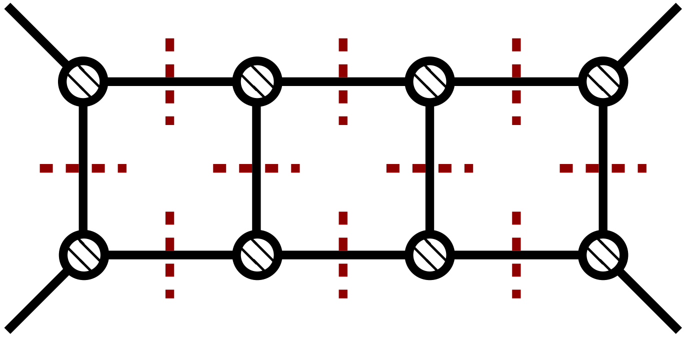
- Integrand reduction and generalized unitarity: going beyond one-loop
- Multi-loop integral coefficients via computational algebraic geometry
- Two-loop hepta-cuts : planar and non-planar

  [SB, Frellesvig, Zhang arXiv:1202.2019, JHEP 1204:055 (2012)]

- Three-loop maximal cuts : triple box

  [SB, Frellesvig, Zhang arXiv:1207:2976, JHEP 1208:065 (2012)]





Background

- One-loop techniques: [Ossola,Papadopoulos,Pittau (2006)]
  [Ellis,Giele,Kunszt,Melnikov (2007-2008)]
  [Bern,Dixon,Dunbar,Kosower (1994)][Britto,Cachazo,Feng (2004)]
- Recent progress in extensions to two-loops:
  - OPP reduction at two-loops [Mastrolia, Ossola arXiv:1107.6041]
    [Mastrolia, Mirabella, Ossola, Peraro arXiv:1205.7087]
    [Kleiss, Malamos, Papadopoulos, Verheyen arXiv:1206.4180]
  - Maximal cuts via contour integration [Kosower, Larsen arXiv:1108.1180]
    [Larsen arXiv:1205.0297], [Larsen, Caron-Huot arXiv:1205.0801]
    [Johansson, Kosower, Larsen arXiv:1208.1754]

Background

- Feynman diagrams and integration-by-parts reduction

  current state-of-the-art for QCD corrections

  - $2 \to 2$ scattering amplitudes:
    - massless QCD [Anastasiou, Glover, Tejeda-Yeomans, Oleari (2000-2002)]

      [Bern, Dixon, Kosower (2000)][Bern, De-Frietas ,Dixon (2002)]
    - $pp \to W + j / e^+e^- \to 3j$ [Garland, Gerhmann, Glover, Koukoutsakis, Remiddi (2002)]
    - $pp \to H + 1j$ [Gerhmann, Jaquier, Glover, Koukoutsakis (2011)]
  - Full NNLO predicitions for $2 \to 2$ processes
    - $e^+e^- \to 3j$ [Gehrmann-De Ridder, Gehrmann, Glover, Heinrich (2007)]
    - $q\bar{q} \to t\bar{t}$ [Bernreuther, Czakon, Mitov (2012)]
    - Towards $pp \to t\bar{t}$ and $pp \to 2$ jets [talks from Abelof, Gerhmann-De Ridder]
- On-shell methods for higher multiplicity at two loops?

Background

- Maximal cut and techniques and leading singularity methods well established in super-symmetric theories

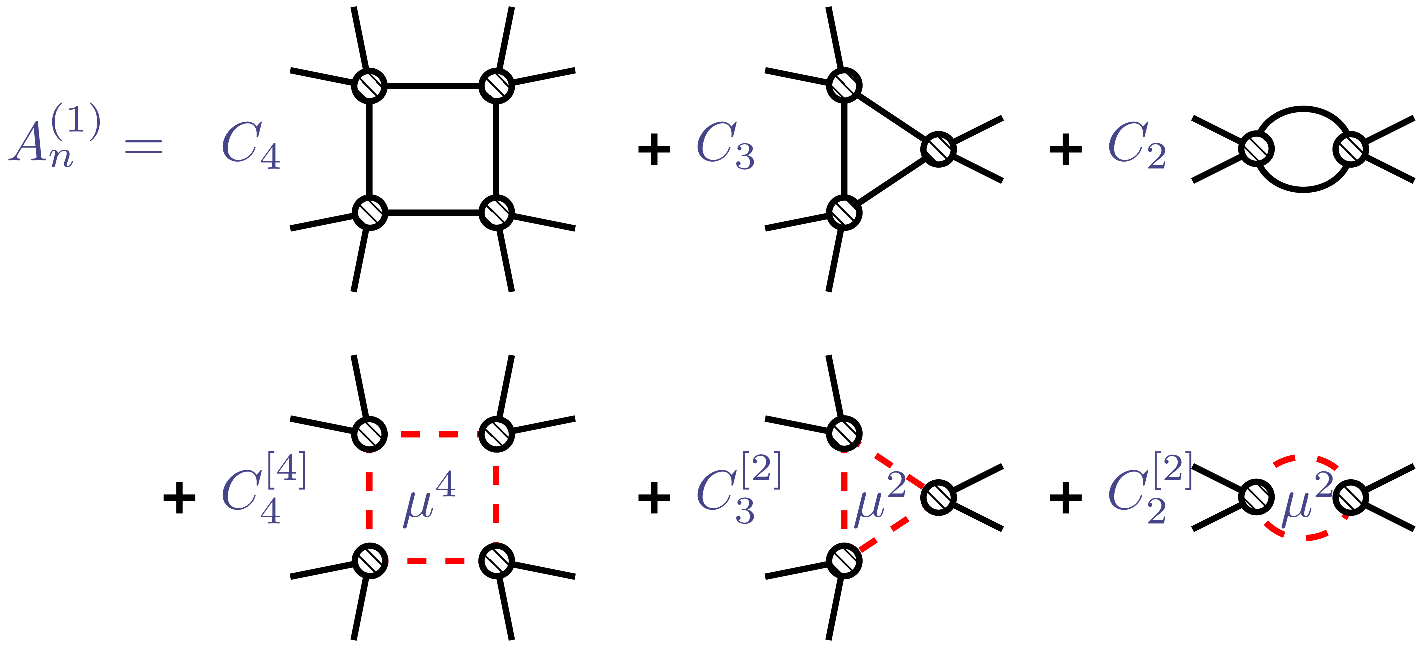
	$2 \to 2$	$2 \to 3$	$2 \to 4$
2-loop	$\mathcal{N} \leq 4$	$\mathcal{N} = 4$	$\mathcal{N} = 4$
3-loop	$\mathcal{N} = 4$	$\mathcal{N} = 4$	
4-loop	$\mathcal{N} = 4$		
5-loop	$\mathcal{N} = 4$		

Bern, Dixon, Kosower, Carrasco, Johansson, Cachzao, Buchbinder, Vergu, Spradlin, Volovich, Wen, Roiban, Drummond, Henn, Korchemsky, Sokatchev, Plefka, Alday, Schuster

- Additional symmetries make amplitudes simpler, e.g. dual conformal symmetry
- Would nice if some of this applied to QCD. . .

One-Loop Overview

- Scalar integral ≤ 4-point functions form a basis with rational coefficients

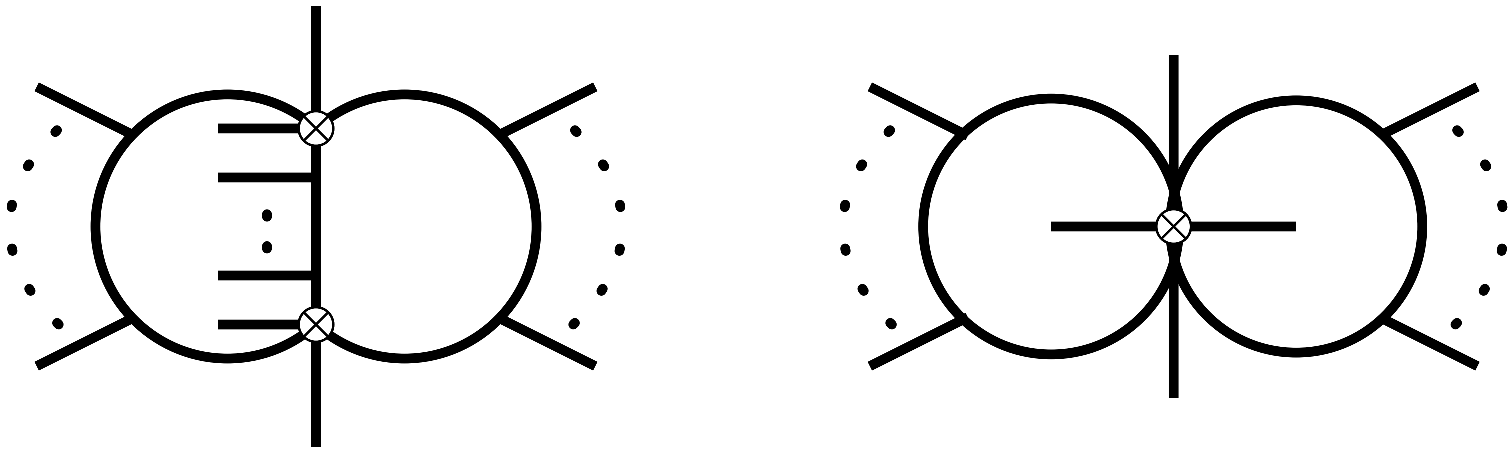


- Integrand representation (OPP) : $\Delta_4(k \cdot \omega) = C_4 + \tilde{C}_4 k \cdot \omega$

Two-Loop Integral Bases

[Gluza, Kosower, Kajda arXiv:1009.0472]

- IBP relations without doubled propagators to obtain a unitarity compatible basis
- Algebraic geometry methods used to solve IBP systems
- Alternative algorithm proposed by Schabinger [Schabinger arXiv:1111.4220]



- No longer just scalar integrals, also tensor integrals in basis

A Two-Loop Integrand Basis

- Integrand is polynomial in irreducible scalar products (ISPs)
  spanned by indep. ext. moms. : $\{p_1, \ldots, p_k\}$ and spurious vecs. : $\{\omega_1, \ldots, \omega_j\}$.
- Gram matrix gives (non-linear) constraints on the polynomial form.

$$G\begin{pmatrix} v_1 \ldots v_n \\ r_1 \ldots r_n \end{pmatrix}, G_{ij} = v_i \cdot r_j$$

- Important to identify spurious terms which integrate to zero.

$$A_n^{(2)} = \int\int \frac{d^D k_1}{(4\pi)^{D/2}} \frac{d^D k_2}{(4\pi)^{D/2}} \sum_{p=3}^{11} \sum_{T_p \in \text{topologies}} \frac{\Delta_{p,T_p}(\{k_i \cdot p_j, k_i \cdot \omega_j\}, \epsilon)}{\prod_{i=1}^{p} l_i(k_1, k_2)}$$

A Two-Loop Integrand Basis

- Integrand is polynomial in irreducible scalar products (ISPs)
  spanned by indep. ext. moms. : $\{p_1, \ldots, p_k\}$ and spurious vecs. : $\{\omega_1, \ldots, \omega_j\}$.
- Gram matrix gives (non-linear) constraints on the polynomial form.

$$G \begin{pmatrix} v_1 \ldots v_n \\ r_1 \ldots r_n \end{pmatrix}, G_{ij} = v_i \cdot r_j$$

- Important to identify spurious terms which integrate to zero.

$$A_n^{(2)} = \int\int \frac{d^D k_1}{(4\pi)^{D/2}} \frac{d^D k_2}{(4\pi)^{D/2}} \sum_{p=3}^{11} \sum_{T_p \in \text{topologies}} \frac{\Delta_{p,T_p}(\{k_i \cdot p_j, k_i \cdot \omega_j\}, \epsilon)}{\prod_{i=1}^{p} l_i(k_1, k_2)}$$

sum from 3 to 11 propagators (8 in 4-D)

A Two-Loop Integrand Basis

- Integrand is polynomial in irreducible scalar products (ISPs)
  spanned by indep. ext. moms. : $\{p_1, \ldots, p_k\}$ and spurious vecs. : $\{\omega_1, \ldots, \omega_j\}$.
- Gram matrix gives (non-linear) constraints on the polynomial form.

$$G \begin{pmatrix} v_1 \ldots v_n \\ r_1 \ldots r_n \end{pmatrix}, G_{ij} = v_i \cdot r_j$$

- Important to identify spurious terms which integrate to zero.

$$A_n^{(2)} = \int\int \frac{d^D k_1}{(4\pi)^{D/2}} \frac{d^D k_2}{(4\pi)^{D/2}} \sum_{p=3}^{11} \sum_{T_p \in \text{topologies}} \frac{\Delta_{p,T_p}(\{k_i \cdot p_j, k_i \cdot \omega_j\}, \epsilon)}{\prod_{i=1}^{p} l_i(k_1, k_2)}$$

sum over all topologies e.g. planar and non-planar

A Two-Loop Integrand Basis

- Integrand is polynomial in irreducible scalar products (ISPs)
  spanned by indep. ext. moms. : $\{p_1, \ldots, p_k\}$ and spurious vecs. : $\{\omega_1, \ldots, \omega_j\}$.
- Gram matrix gives (non-linear) constraints on the polynomial form.

$$G\begin{pmatrix} v_1 \ldots v_n \\ r_1 \ldots r_n \end{pmatrix}, G_{ij} = v_i \cdot r_j$$

- Important to identify spurious terms which integrate to zero.

$$A_n^{(2)} = \int\int \frac{d^D k_1}{(4\pi)^{D/2}} \frac{d^D k_2}{(4\pi)^{D/2}} \sum_{p=3}^{11} \sum_{T_p \in \text{topologies}} \frac{\Delta_{p,T_p}(\{k_i \cdot p_j, k_i \cdot \omega_j\}, \epsilon)}{\prod_{i=1}^{p} l_i(k_1, k_2)}$$

Integrand parametrised as coeff $\times$ ISP

A Two-Loop Integrand Basis

- Integrand is polynomial in irreducible scalar products (ISPs)
  spanned by indep. ext. moms. : $\{p_1, \ldots, p_k\}$ and spurious vecs. : $\{\omega_1, \ldots, \omega_j\}$.
- Gram matrix gives (non-linear) constraints on the polynomial form.

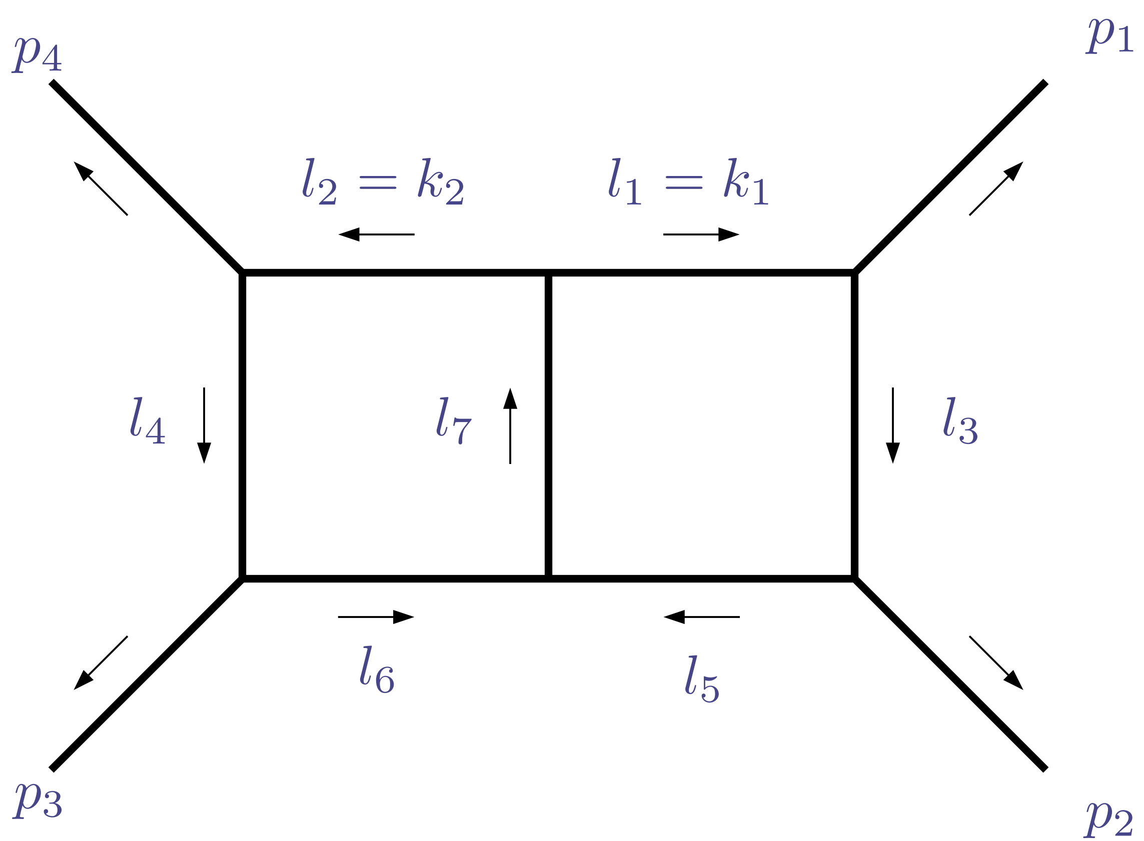
$$G \begin{pmatrix} v_1 \ldots v_n \\ r_1 \ldots r_n \end{pmatrix}, G_{ij} = v_i \cdot r_j$$

- Important to identify spurious terms which integrate to zero.

$$A_n^{(2)} = \int\int \frac{d^D k_1}{(4\pi)^{D/2}} \frac{d^D k_2}{(4\pi)^{D/2}} \sum_{p=3}^{11} \sum_{T_p \in \text{topologies}} \frac{\Delta_{p,T_p}(\{k_i \cdot p_j, k_i \cdot \omega_j\}, \epsilon)}{\prod_{i=1}^{p} l_i(k_1, k_2)}$$

Set of propagators for given topology

Example : Planar Double Box



RSPs :

$2k_1 \cdot p_1 = -(l_1 - p_1)^2 + l_1^2$

$2k_1 \cdot p_1 = -l_3^2 + l_1^2$

- ISPs $= \{k_1 \cdot p_4, k_2 \cdot p_1, k_1 \cdot \omega, k_2 \cdot \omega\}$
- $\Delta = c_0 + c_1 k_1 \cdot p_4 + \cdots + c_{16} k_1 \cdot \omega + \ldots$
- Everything will be simplifed to 4 dimensions from now on. . .

Generalized Unitarity Cuts

- On a solution to $\{l_i^2 = 0\}$ the integrand factorizes onto a product of tree-level amplitudes
- Algorithm to fit a generic integrand:
  - Parametrize the full set of on-shell solutions, $l_i^{(s)}(\tau_1, \ldots, \tau_p)$
  - Identify the ISPs on this solution:

$$k_i \cdot p_j = f_{ij}(\tau_1, \ldots, \tau_p)$$

  - Construct and solve the resulting linear system:

$$\Delta^{(s)}(\tau_1, \ldots, \tau_p) = \sum d_a \tau_1 \ldots \tau_p$$

$$\boxed{\mathbf{M} \cdot \vec{c} = \vec{d}}$$

Integrand Reduction

- Take a top-down approach to fitting each Δ
- Subtract previously determined poles, e.g.

$$\Delta_{6;\text{tri}|\text{box}} = \prod_{i=1}^{5} A_i^{(0)} - \frac{\Delta_{7;\text{box}|\text{box}}}{(k_1 - p_1)^2} = \sum_{i,j} d_{ij}\tau_i\tau_j$$

- Fitting can be done numerically or analytically
- Total number of topologies is still very large. . . .
- Towards automation:
  Solving the non-linear integrand constraints using algebraic geometry

  [Zhang arXiv:1205.5705]
- Public Mathematica code `BasisDet` [http://www.nbi.dk/~zhang/BasisDet.html]

An Algorithm for the Integrand Basis

- $B = \{v_1, v_2, v_3, v_4\}$, $[G_4]_{ij} = v_i \cdot v_j$, $P = \{l_1^2, \ldots, l_p^2\}$
- Gram matrix $[G_4]_{ij} = v_i \cdot v_j$. to re-write scalar products:

$$a \cdot b = (a \cdot v_1 \;\; a \cdot v_2 \;\; a \cdot v_3 \;\; a \cdot v_4) \;\; G_4^{-1} \begin{pmatrix} b \cdot v_1 \\ b \cdot v_2 \\ b \cdot v_3 \\ b \cdot v_4 \end{pmatrix} \tag{1}$$

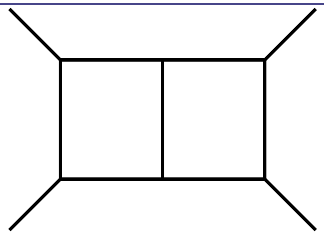
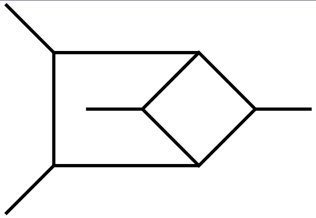
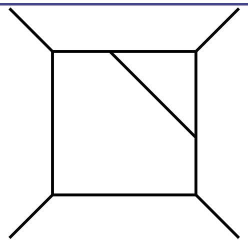
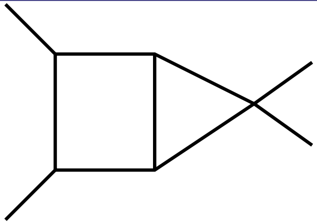
- Re-write P using (1) $\Rightarrow$ set of equations for the scalar products.
- $\{P_i = 0\}$ has linear parts (RSPs) non-linear parts : ISP constraints $= I$
- Construct general ISP polynomial using renormalization constraints $= R$
- Remove I from R $(R/I) \Rightarrow$ Integrand Basis $= \Delta(ISPs)$.
  - Carried out using Gröbner bases and polynomial division

[see also Mastrolia's talk]

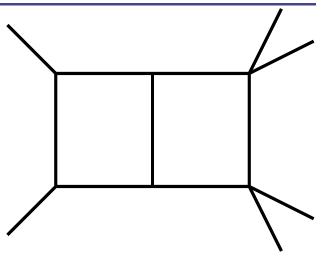
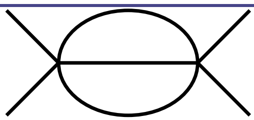
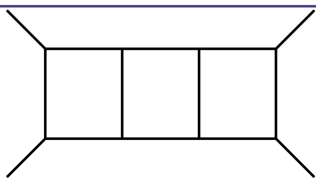
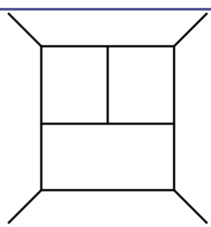
Solving the On-Shell Contraints

- *primary decomposition* to identify all on-shell solution branches
  - Prime factorization for polynomials [Lasker-Noether theorem (1905,1921)]
- Decompose $Z(I) \sim \{I = 0\}$ into a finite number of irreducible components
  - e.g. consider $I = \{x^2 - y^2\}$
    $I = \{x + y\} \cup \{x - y\} \Rightarrow Z(I) = \{x + y = 0\} \cup \{x - y = 0\}$
- Prime decomposition "factorizes" all solutions to identify different branches.
- Available in the public `Macaulay2` program [http://www.math.uiuc.edu/Macaulay2/]

A Few Examples

Topology	ISPs (non-spurious+spurious)	$\lvert\Delta\rvert$(non-sp.+sp.)	#branches(dimension)
	2+2	$32(16+16)$	$6(1)$
	2+2	$38(19+19)$	$8(1)$
	2+1	$20(10+10)$	$2(2)$
	1+4	$69(18+51)$	$4(2)$

A Few Examples

Topology	ISPs (non-spurious+spurious)	$\lvert\Delta\rvert$(non-sp.+sp.)	#branches(dimension)
	2+2	$32(16+16)$	$4(1)$
	2+6	$42(12+30)$	$1(5)$
	4+3	$398(199+199)$	$14(2)$
	5+3	$584(292+292)$	$12(2)+4(3)$

Further reduction to MI's via IBPs

- The integrand representation contains hundreds of integrals
- From this form we can apply further identities from conventional IBPs

[Tkachov, Chetyrkin (1980)]

- Doubled propagators can be allowed at this stage
- Public codes :

[AIR, Anastasiou, Lazopoulos (2004)]

[FIRE, Smirnov ,Smirnov (2008)]

[Reduze2, Studerus, von Manteuffel (2009-2011)]

$$A_n^{(2)} = \int\int \frac{d^4k_1}{(4\pi)^2}\frac{d^4k_2}{(4\pi)^2}\sum_{p=3}^{8}\sum_{T_p\in\text{topologies}}\frac{\Delta_{p,T_p}(\{k_i\cdot p_j, k_i\cdot\omega_j\})}{\prod_{i=1}^{p} l_i(k_1,k_2)}$$

Further reduction to MI's via IBPs

- The integrand representation contains hundreds of integrals
- From this form we can apply further identities from conventional IBPs

[Tkachov, Chetyrkin (1980)]

- Doubled propagators can be allowed at this stage
- Public codes :

[AIR, Anastasiou, Lazopoulos (2004)]

[FIRE, Smirnov ,Smirnov (2008)]

[Reduze2, Studerus, von Manteuffel (2009-2011)]

$$A_n^{(2)} = \int\int \frac{d^4k_1}{(4\pi)^2}\frac{d^4k_2}{(4\pi)^2}\sum_{p=3}^{8}\sum_{T_p\in\text{topologies}}\frac{c_{\vec{T_p}}\cdot\vec{B}_{T_p}}{\prod_{i=1}^{p} l_i(k_1,k_2)}$$

Further reduction to MI's via IBPs

- The integrand representation contains hundreds of integrals
- From this form we can apply further identities from conventional IBPs

[Tkachov, Chetyrkin (1980)]

- Doubled propagators can be allowed at this stage
- Public codes :

[AIR, Anastasiou, Lazopoulos (2004)]

[FIRE, Smirnov ,Smirnov (2008)]

[Reduze2, Studerus, von Manteuffel (2009-2011)]

$$A_n^{(2)} = \int\int \frac{d^4k_1}{(4\pi)^2}\frac{d^4k_2}{(4\pi)^2}\frac{\vec{C}\cdot\vec{B}}{\prod_{i=1}^n l_i(k_1,k_2)}$$

Further reduction to MI's via IBPs

- The integrand representation contains hundreds of integrals
- From this form we can apply further identities from conventional IBPs

[Tkachov, Chetyrkin (1980)]

- Doubled propagators can be allowed at this stage
- Public codes :

[AIR, Anastasiou, Lazopoulos (2004)]

[FIRE, Smirnov ,Smirnov (2008)]

[Reduze2, Studerus, von Manteuffel (2009-2011)]

$$A_n^{(2)} = \int\int \frac{d^4k_1}{(4\pi)^2}\frac{d^4k_2}{(4\pi)^2}\frac{\vec{C}\cdot\vec{B}}{\prod_{i=1}^{n} l_i(k_1,k_2)}$$

$$A_n^{(2)} = \vec{C}\cdot M_{IBP}\cdot\int\int \frac{d^4k_1}{(4\pi)^2}\frac{d^4k_2}{(4\pi)^2}\frac{\vec{B}\prime}{\prod_{i=1}^{n} l_i(k_1,k_2)}$$

Further reduction to MI's via IBPs

- The integrand representation contains hundreds of integrals
- From this form we can apply further identities from conventional IBPs

[Tkachov, Chetyrkin (1980)]

- Doubled propagators can be allowed at this stage
- Public codes :

[AIR, Anastasiou, Lazopoulos (2004)]

[FIRE, Smirnov ,Smirnov (2008)]

[Reduze2, Studerus, von Manteuffel (2009-2011)]

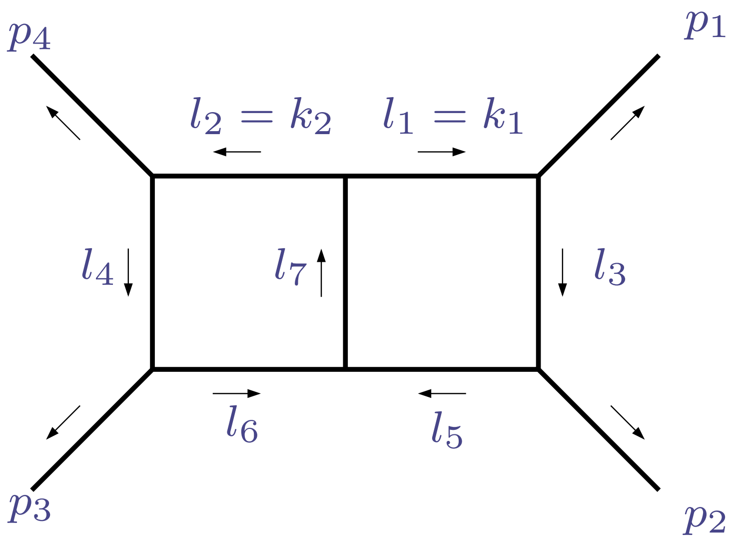
$$A_n^{(2)} = \int\int \frac{d^4k_1}{(4\pi)^2}\frac{d^4k_2}{(4\pi)^2}\frac{\vec{C}\cdot\vec{B}}{\prod_{i=1}^n l_i(k_1,k_2)}$$

solution to system of IBPs :

$$\int\int \vec{B} = M_{IBP}\cdot\int\int \vec{B}'$$

$$A_n^{(2)} = \vec{C}\cdot M_{IBP}\cdot\int\int \frac{d^4k_1}{(4\pi)^2}\frac{d^4k_2}{(4\pi)^2}\frac{\vec{B}'}{\prod_{i=1}^n l_i(k_1,k_2)}$$

Two-Loop Hepta-cuts

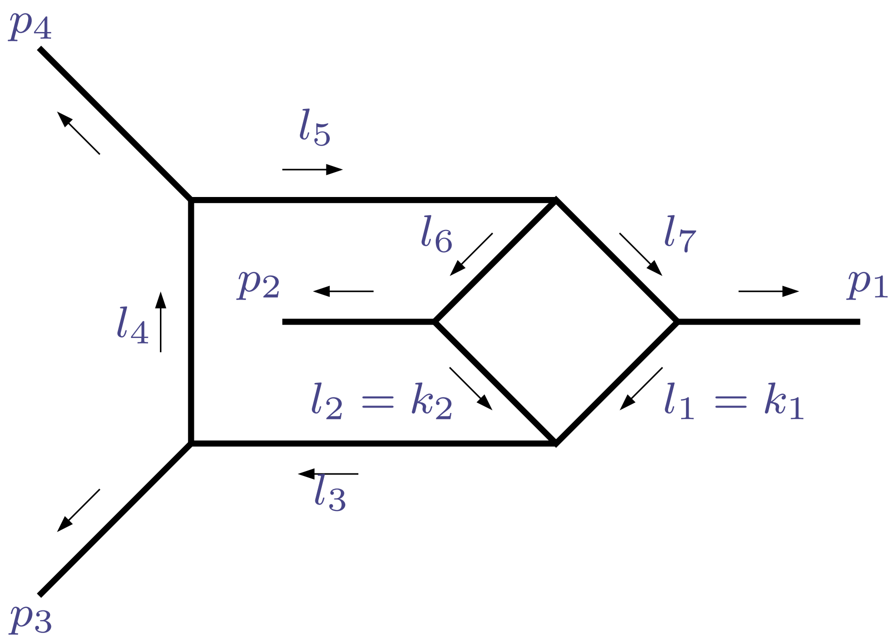


- Solve 38×32 system (4 ISPs) + IBPs
  $\Rightarrow$ 2 MIs $I_7[1], I_7[k_1 \cdot p_4]$
- IBPs with FIRE [Smirnov,Smirnov]
- 4-gluon scattering in Yang-Mills with adjoint fermions and scalars
- SUSY limit : $nf \to \mathcal{N}, ns \to \mathcal{N} - 1$

$$
\begin{aligned}
C_1^{-++-} &= -\frac{s^2}{4t^2} A^{(0)} \big(2s\left(10s^2 + 11ts + 2t^2\right)\left(1 - n_f + n_s\right) \\
&\quad + t\left(\left(4 - n_f\right)\left(3 - n_s\right) s\left(2s + t\right) - \left(4 - n_f\right) s\left(4s + t\right) + 4t^2\right)\big) \\
C_2^{-++-} &= \frac{3s^3}{2t^3} A^{(0)} \big(\left(20s^2 + 22ts + 4t^2\right)\left(1 - n_f + n_s\right) \\
&\quad + t\left(\left(4 - n_f\right)\left(3 - n_s\right)\left(2s + t\right) - \left(4 - n_f\right)\left(4s + t\right)\right)\big)
\end{aligned}
$$

[cross checked with Bern, De-Freitas, Dixon (2002)]

Two-Loop Hepta-cuts : Non-planar



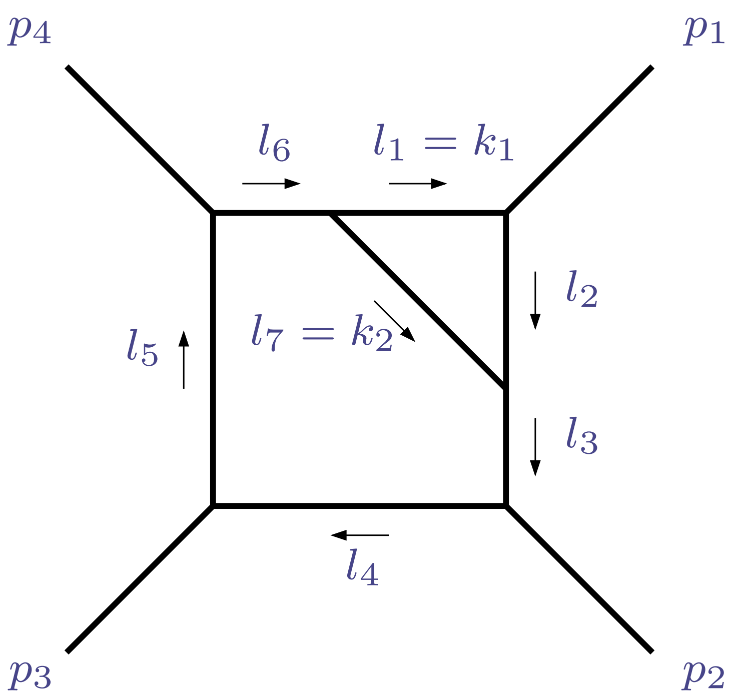
- Solve 48×38 system (4 ISPs) + IBPs
  $\Rightarrow$ 2 MIs $I_7[1], I_7[k_1 \cdot p_4]$
- IBPs with FIRE [Smirnov,Smirnov]

$$C_1^{-++-} = \frac{s^2}{4t^2} A^{(0)} \left(2su^2(1 - n_f + n_s) - t^2\left((4 - n_f)u + 4t\right)\right)$$

$$C_2^{-++-} = \frac{s^2}{2t^3} A^{(0)} \left(3u + t\right) \left((4 - n_f)t^2 - 2su(1 - n_f + n_s)\right)$$

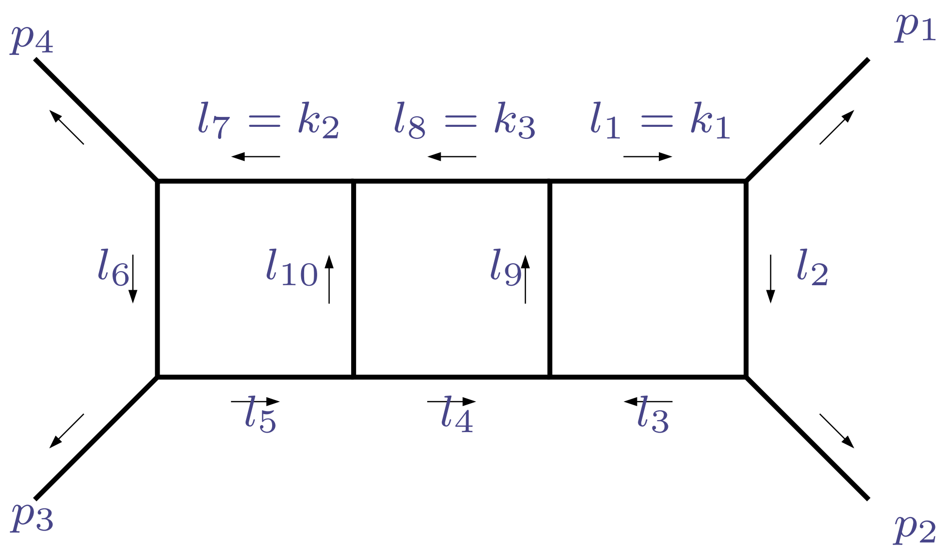
[cross checked with Bern, De-Freitas, Dixon (2002)]

Two-Loop Hepta-cuts : pentagon|triangle



- Solve 20×20 system (3 ISPs)
- No 6-propagator MIs after IBPs
- Important contribution for subtractions

The Three-Loop Triple Box



- Solve 622×398 system (7 ISPs) : standard getting slow. . .

- Option 1 : split spurious and non-spurious systems and solve with linear algebra
- Option 2 : branch-by-branch fitting and algebraic geometry
  - Fit each branch over a reduced set of ISP monomials
  - Merge branches using intersection of ideals
- Test in 4-gluon scattering
- 34 flavour configuration inside the loop $n_f, n_s, n_f n_s, n_f^2 n_s, \ldots$

The Three-Loop Triple Box II

- IBPs with Reduze2 $\Rightarrow$ 3 ten-propagator MIs [Studerus,von Manteuffel]

$$I_{10}[1], I_{10}[k_1 \cdot p_4], I_{10}[k_3 \cdot p_4]$$

- New analytic results for non-super-symmetric gauge theory

$$\hat{C}_1^{--++}(s,t) = -1$$
$$\hat{C}_2^{--++}(s,t) = 0$$
$$\hat{C}_3^{--++}(s,t) = 0$$

The Three-Loop Triple Box II

- IBPs with Reduze2 $\Rightarrow$ 3 ten-propagator MIs [Studerus,von Manteuffel]

$$I_{10}[1], I_{10}[k_1 \cdot p_4], I_{10}[k_3 \cdot p_4]$$

- New analytic results for non-super-symmetric gauge theory

$$\begin{aligned}
\hat{C}_1^{-++-}(s,t) &= \\
&- 1 - (4-n_f)(3-n_s)\frac{s(t+2s)}{2t^2} + (4-n_f)\frac{s(t+4s)}{2t^2} \\
&- (1+n_s-n_f)\frac{s}{t^3}\left(2t^2+11st+10s^2\right) \\
\hat{C}_2^{-++-}(s,t) &= \frac{2}{t}\left(1-\hat{C}_1^{-++-}\right) \\
\hat{C}_3^{-++-}(s,t) &= 0,
\end{aligned}$$

The Three-Loop Triple Box II

- IBPs with Reduze2 $\Rightarrow$ 3 ten-propagator MIs [Studerus,von Manteuffel]

$$I_{10}[1], I_{10}[k_1 \cdot p_4], I_{10}[k_3 \cdot p_4]$$

- New analytic results for non-super-symmetric gauge theory

$$\begin{aligned}
\hat{C}_1^{-+-+}(s,t) &= -1 + (4-n_f)\frac{st}{u^2} - 2(1+n_s-n_f)\frac{s^2t^2}{u^4} \\
&+ \big(2(1-2n_s)+n_f\big)(4-n_f)\frac{s^2t(2t-s)}{4u^4} \\
&- \big(n_f(3-n_s)^2 - 2(4-n_f)^2\big)\frac{st(t^2-4st+s^2)}{8u^4} \\
\hat{C}_2^{-+-+}(s,t) &= -(4-n_f)\frac{s}{u^2} + 2(1+n_s-n_f)\frac{s^2t}{u^4} \\
&- \big(2(1-2n_s)+n_f\big)(4-n_f)\frac{s^2(2t-s)}{u^4} \\
&+ \big(n_f(3-n_s)^2 - 2(4-n_f)^2\big)\frac{s(t^2-4st+s^2)}{2u^4}
\end{aligned}$$

The Three-Loop Triple Box II

- IBPs with Reduze2 $\Rightarrow$ 3 ten-propagator MIs [Studerus,von Manteuffel]

$$I_{10}[1], I_{10}[k_1 \cdot p_4], I_{10}[k_3 \cdot p_4]$$

- New analytic results for non-super-symmetric gauge theory

$$\begin{aligned}\hat{C}_3^{-+-+}(s,t) =& \\ &+ \big(2(1-2n_s) + n_f\big)(4-n_f)\frac{3s^2(2t-s)}{2u^4} \\ &- \big(n_f(3-n_s)^2 - 2(4-n_f)^2\big)\frac{3s(t^2-4st+s^2)}{4u^4}.\end{aligned}$$

The Three-Loop Triple Box II

- IBPs with Reduze2 $\Rightarrow$ 3 ten-propagator MIs [Studerus,von Manteuffel]

$$I_{10}[1], I_{10}[k_1 \cdot p_4], I_{10}[k_3 \cdot p_4]$$

- New analytic results for non-super-symmetric gauge theory

$$\begin{aligned}\hat{C}_3^{-+-+}(s,t) = \\ &+ \big(3(\mathcal{N}-2)\big)(\mathcal{N}-4)\frac{3s^2(2t-s)}{2u^4} \\ &- \big((\mathcal{N}-2)(\mathcal{N}-4)^2\big)\frac{3s(t^2-4st+s^2)}{4u^4}.\end{aligned}$$

- 3rd MI vanishes in $\mathcal{N}=2$ SUSY

Outlook

- A few small steps towards automated multi-loop amplitudes
- Computational algebraic geometry for integrand reduction
  - efficent solutions to constraint equations
  - generalizes easily to D-dimensional systems
- We didn't address the evaluation of the Master Integrals
- IBPs with many scales are hard:
  - massive amplitudes, higher multiplicity