

Black Holes Quantum N-Portrait

Gia Dvali

MPI Project Review

with: Cesar Gomez

Other directions with:

Pieter Lust, Sarah Folkerts,
Slava Mukhanov, Alex Richel,
Andra Franca, Daniel Flassing,
Konstantin Sluka, Nico Wintergerst,

(Semi) Classically,
Black hole physics is
mysterious.

We show that
Quantum-mechanically it is
the simplest (in certain
sense):

$\frac{1}{N}$ -coupled large- N
physics.

Black hole is a leaky
Bose-condensate of soft
gravitons:
(In Planck units)

Occupation number $= N$

Wave-length $= \sqrt{N}$

Interaction strength $= 1/N$

Mass $= \sqrt{N}$

Temperature $= 1/\sqrt{N}$

Entropy $= N$

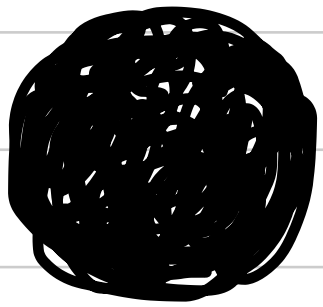
In $c=1$ units,

$$[G_N] = \ell / m$$

So we can define:

$$L_p^2 \equiv \hbar G_N \quad M_p \equiv \frac{\hbar}{L_p}$$

Both are quantum.

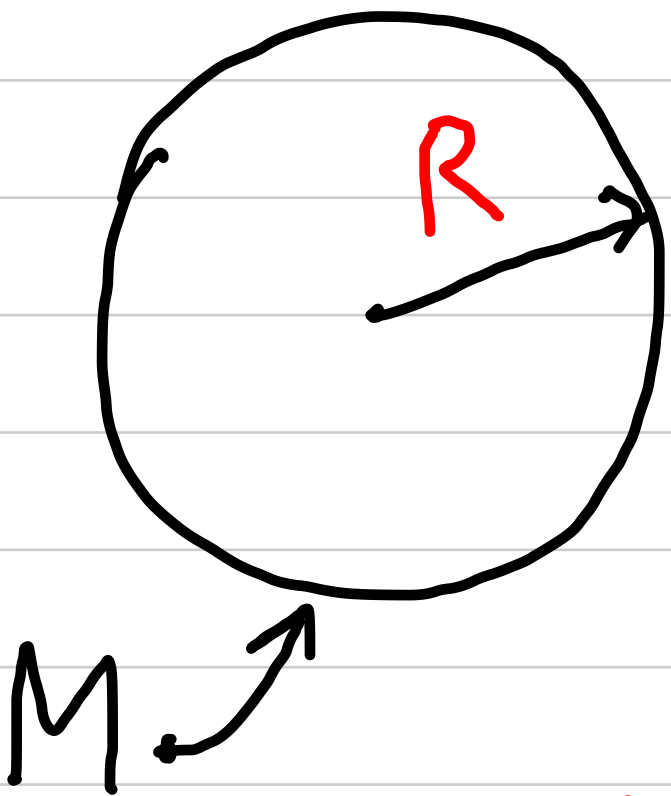


Source of mass M

Gravitational radius:

$$r_g \equiv M G_N = \frac{L_p^2}{\hbar M}$$

Classical Newtonian Geometry



$$\phi(r) = -\frac{\sqrt{g}}{r}$$

Gravitational energy

$$E_{\text{grav}} = M \frac{\sqrt{g}}{R}$$

But, quantum-mechanically
we should think about
many gravitons with:

$$\lambda = R \quad \text{and} \quad E = \hbar / R$$

Their occupation number is

$$N = \frac{E_{\text{grav}}}{E} = \frac{M r g}{\hbar}$$

or

$$N = \frac{v_g^2}{L_p^2}$$

$$= \frac{M^2}{M_p^2}$$

$$= \frac{L_p^2}{\hbar^2 M^{-2}}$$

For any source of size R
occupation number of
constituent weakly-coupled
particles is maximum

$$N_{\text{Source-max}} = \frac{MR}{\hbar}$$

and since $R \geq r_g$

for any fixed R ,

N_{Source} is maximized by

Black holes

For $R = v_g$, graviton Bose-condensate becomes self-sustained.

$$\lambda = \sqrt{N} L_p$$

These gravitons are very weakly interacting:

$$d_{gr} = \hbar G_N \lambda^{-2} = \frac{1}{N}$$

total mass


$$M = \sqrt{N} \frac{\hbar}{L_p}$$

Boundstate exist for any N , and it is leaky.

Hartree potential:

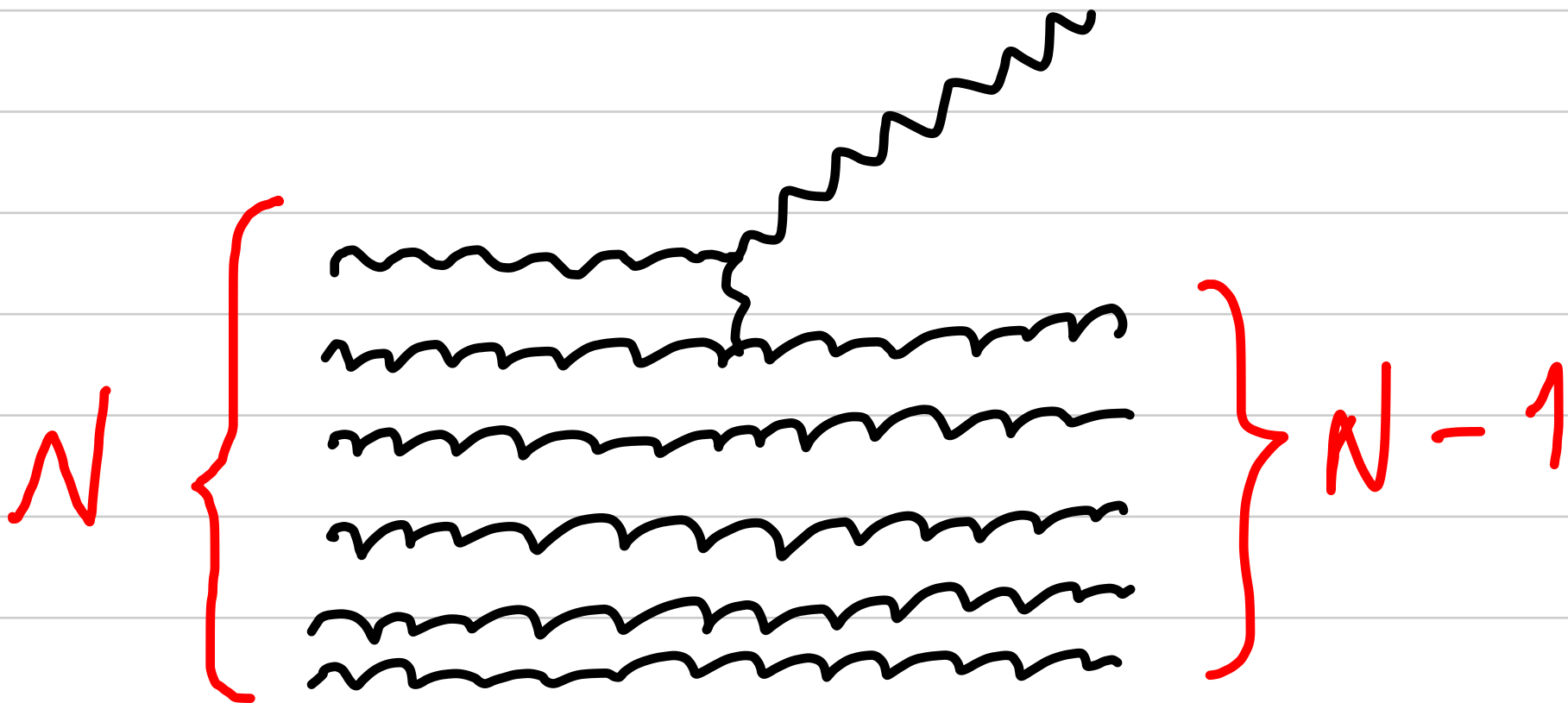
$$V = G_N N \frac{\hbar^2}{\lambda^2} \frac{1}{\lambda} = N \alpha_{gr} \frac{\hbar}{\lambda}$$

$$\alpha_{gr} = \frac{1}{N} \quad \lambda = \sqrt{N} L_P$$

$$V = \frac{\hbar}{\sqrt{N} L_P} = E_{\text{escape}}$$


Escape energy is just above!

Hawking radiation



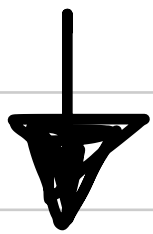
$$\Gamma = \frac{\hbar}{\Delta t} = \frac{\hbar}{\sqrt{N} L_P}$$

$$\frac{dN}{dt} = -\frac{1}{\sqrt{N}} L_P$$

$$\tau = N^{3/2} L_P$$

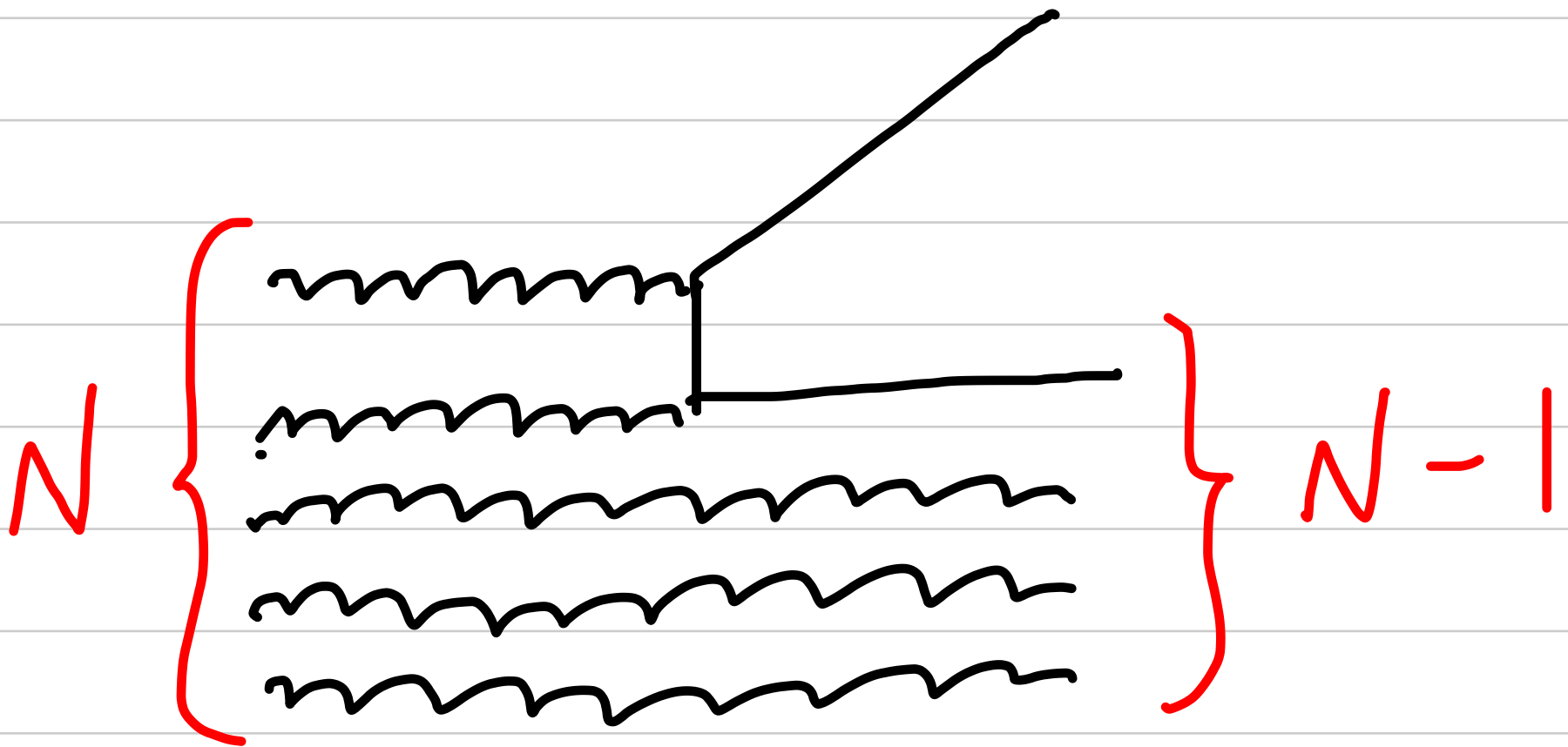
Bose-condensate is cold, but because of quantum depletion gravitons are radiated with a thermal spectrum of temperature

$$T = \frac{\hbar}{\sqrt{N} L P}$$



$$\frac{dM}{dt} = - \frac{T^2}{\hbar}$$

If there are extra species



$$\Gamma = \frac{\hbar}{\sqrt{N} L_P} N_{\text{species}}$$

Species bound:
Black holes cannot exist
for $N < N_{\text{species}}!$

Semi-classically this bound translates to the bound on \sqrt{g} :

$$N > N_{\text{species}}$$



$$\sqrt{g} > L_N \equiv \sqrt{N_{\text{species}}} L_P$$

Solitons (e.g. Monopole):

Bose-condensate of W-bosons

$$\alpha \equiv \hbar g^2 \ll 1$$

$$M = \frac{M_W}{\alpha}$$

$$N = \frac{L}{\alpha}$$

$$\lambda = \frac{\hbar}{M_W}$$

Defining: $L_* \equiv \frac{\hbar}{v}$

Higgs VEV

Soliton N-picture:

Mass $\rightarrow M = \sqrt{N_w} \frac{\hbar}{L_*}$

Strength $\rightarrow \alpha = \frac{1}{N_w}$

Number = N_w

wave-length = $\sqrt{N_w} L_*$

just like BH!

Now couple to gravity:

$$N = \frac{M^2}{M_P^2}$$

$$N_W = \frac{1}{\hbar g}$$

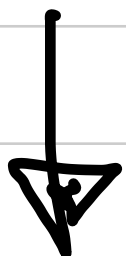
$N_W > N$ ← Soliton

$N_W = N$ ← Extremal BH

$N_W < N$ ← Black hole

Explains non-singularity
condition

$$Q_{\text{Magnetic}} \equiv \frac{l}{g} \leq \frac{M}{M_p} \sqrt{\hbar}$$



$$N_w < N$$

For $N_w > N$,

there are no BH!

Origin of Bekenstein

entropy:

$$N_f \sim N \text{ "flavors"}$$

Because of order

$\sim N$ subsets can form
"Unions".

At the number of
non-interacting flavors
is $\sim N$

BH wavefunction
is a Hartree wave-
function of
 N_f - flavors

$$\Psi_{BH} = \prod_j^{\sim N} \psi_j$$

$$N_{states} \sim \sum^N$$

Difference between
BH-s and solutions:

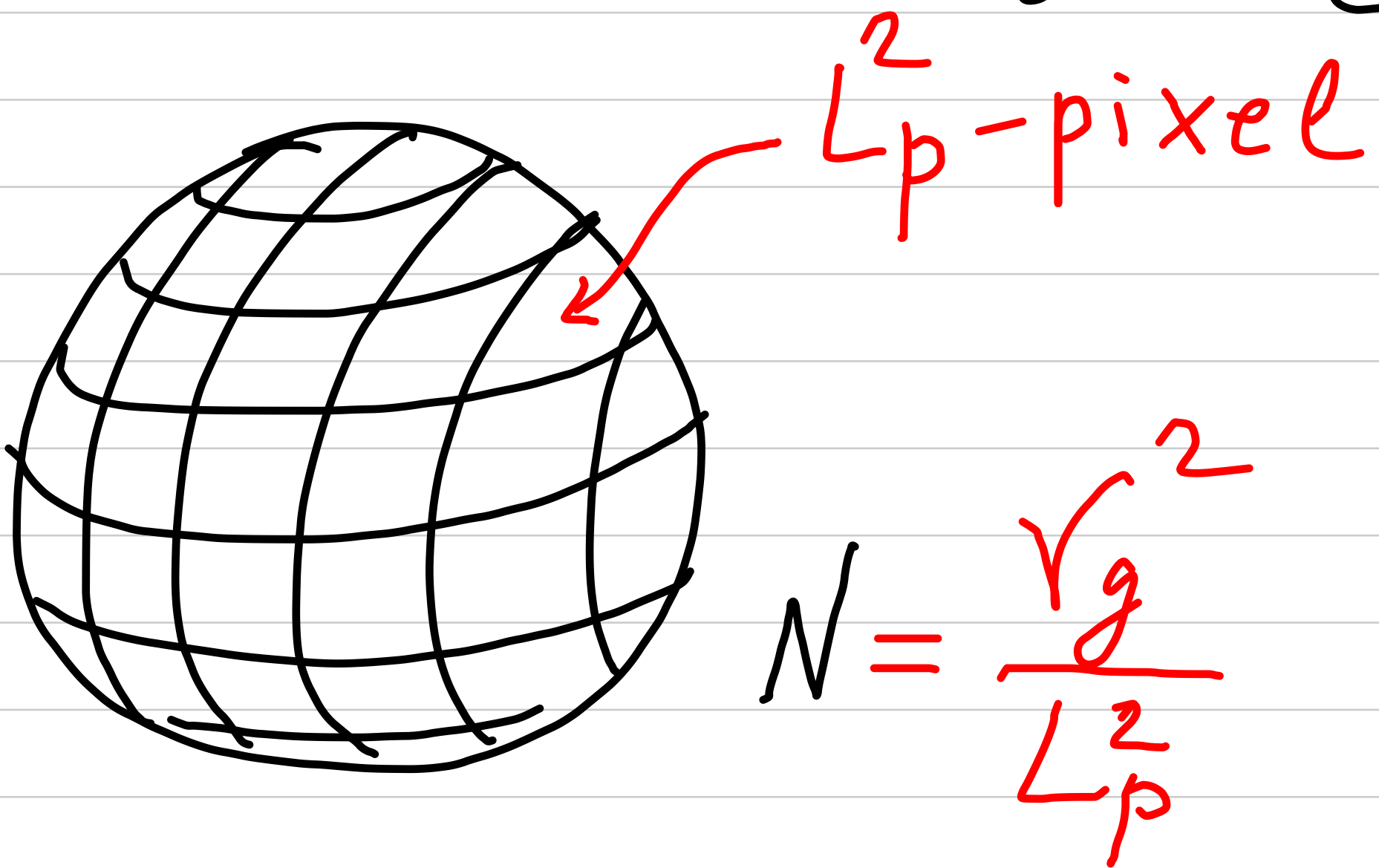
Solutions exist only
for fixed N_w .

No energy self-
sourcing!

Cannot be leaky!

We are learning that
overpacked systems
get oversimplified.

↖
Origin of holography



It is interesting that generalizing our idea to AdS/dS-geometry, we get the same N-portrait:

In D-dimensions:

$$N = \frac{R^{D-2}}{L_D^{D-2}}$$

$$\lambda = N^{\frac{1}{D-2}} L_D$$

$$\alpha_D = \frac{1}{N}$$

Notice, that N
coincides with the
central charge of
CFT

$$N_{\text{CFT}} = N = \left(\frac{R}{L_D} \right)^{D-2}$$

