

# Instantons in QFT and String Theory

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# Contents

- 1 What are instantons?
- 2 What does string theory tell us about them?
- 3 What is my project about?
- 4 Sources

# Motivation

Consider path-integral:

$$Z[J] = \int [dA] \exp \left( \frac{-1}{e^2} S[A] + \int d^4x \operatorname{Tr} (J^\mu A_\mu) \right)$$

Semi-classical approximation:

$$S[A] = S[A_{cl}] + \frac{1}{2} \int d^4x_1 d^4x_2 \delta A_\mu(x_1) \Delta(x_1, x_2)^{\mu\nu} \delta A_\nu(x_2) + \dots$$

Note that different contributions are weighted with  $\exp \left( -\frac{S_{cl}}{e^2} \right)$ .

# Finite action solutions

Let's look for *finite action* solutions to classical  $SU(N)$  Yang-Mills theory.

$$\mathcal{L} = \frac{1}{2} \text{Tr} F_{\mu\nu} F^{\mu\nu}$$

Finite action implies

$$F_{\mu\nu} \xrightarrow{r \rightarrow \infty} 0 \quad A_\mu \xrightarrow{r \rightarrow \infty} ig^{-1} \partial_\mu g,$$

where  $g(x) = e^{iT(x)} \in SU(N)$ .

# Homotopy

We are therefore dealing with maps from  $S^3$  to  $SU(N)$ .  
 Mathematicians classify these according to their *homotopy class*.  
 In our case:

$$\pi_3(SU(N)) \cong \mathbb{Z}$$

$\Rightarrow$  Solutions are labelled by single integer  $k$ . Two solutions of the same homotopy class can be smoothly deformed into each other.

The easiest example is  $SU(2)$ :

$$g^{(0)} = 1 \quad g^{(1)} = \frac{x_4 + ix_j \sigma^j}{r} \quad g^{(k)} = [g^{(1)}]^k$$

# The Bogomol'nyi bound

To obtain actual equations, we now exploit a trick. Let

$\star F_{\mu\nu} = \frac{1}{2}\epsilon_{\mu\nu\rho\sigma}F^{\rho\sigma}$ . Then  $\star F_{\mu\nu}^2 = F_{\mu\nu}^2$  and

$$\begin{aligned} S &= \frac{1}{2} \int d^4x F_{\mu\nu}^2 = \frac{1}{4} \int d^4x (F_{\mu\nu} \mp \star F_{\mu\nu})^2 \pm 2F_{\mu\nu} \star F^{\mu\nu} \\ &\geq \pm \frac{1}{2} \int d^4x \epsilon^{\mu\nu\rho\sigma} \partial_\nu \left( A_\rho F_{\mu\sigma} + \frac{2i}{3} A_\nu A_\rho A_\sigma \right) = 8\pi^2 |k| \end{aligned}$$

Hence  $S$  locally minimized  $\Leftrightarrow F_{\mu\nu} = \star F_{\mu\nu}$  for  $k > 0$ .

# The $SU(2)$ solution

The  $SU(2)$  solution for  $k = 1$  was found by Polyakov et. al (1975):

$$A_\mu(x) = \frac{\rho^2(x - X)^\nu}{(x - X)^2((x - X)^2 + \rho^2)} \bar{\eta}_{\mu\nu}^i (g \sigma^i g^{-1})$$

Note:

- Localized in 4 dimensions, hence the name
- 8 parameters called *collective coordinates*
- Singularity is only unphysical gauge artifact

# ADHM

To find an instanton solution one has to solve  $F_{\mu\nu} = \star F_{\mu\nu}$ . But differential equations are hard!

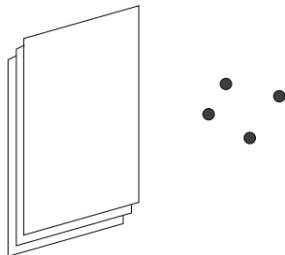
Fortunately, there exists a very elegant solution by Atiyah, Drinfeld, Hitchin and Manin (1978). Originally constructed in terms of twistor spaces, they reduced the differential equation to an algebraic one.

Their construction admits a string theory interpretation in terms of a certain brane set-up.



# Brane interpretation of ADHM I

- Stack  $N$  D $p$ -branes to obtain  $U(N)$  SYM on their worldvolume. It couples to bulk fields e.g. via  $\int_{D^p} d^{p+1}x C_{p-3} \wedge F \wedge F$ .
- Instanton gives  $\int_{D^{(p-4)}} d^{p-3}x C_{p-3} \Rightarrow$  identical to D $(p-4)$ -brane!
- In  $p+1$  dimensions, instanton is a  $(p-3)$ -dimensional object.



[stolen from D. Tong's lecture notes]

## Brane interpretation of ADHM II

One now considers the gauge theory on the  $k$   $D(p-4)$ -branes and looks for vacuum solutions.

It turns out that

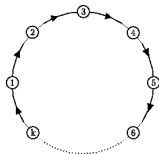
$$\mathcal{M}_{\text{Higgs}} \cong \mathcal{M}_{\text{instanton}} .$$

- ADHM works for all classical Lie groups.
- Instanton moduli space has Hyper-Kähler structure, the moment maps are obtained from gauge potential.
- For most QFT calculations one only needs the metric on the instanton moduli space.

# Applications & my work

What is my project about?

- Calculated the Hilbert series using plethystics ( Feng, Hanany, He (2007) ) for certain class of quiver gauge theories:



$$g_N(\{t_i\}, k) = \frac{1}{\prod_{i=1}^k (1 - t_k^N)} \prod_{j=1}^{N-1} \frac{1}{1 - \prod_{i=1}^k t_i^j}$$

- Apply these techniques to study instanton moduli space, see Hanany, Mekareeya et al (2010 & 2012).

# Sources

## Sources

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