

Top Mass Measurement in the Fully Hadronic $t\bar{t}$ Channel With The ATLAS Detector

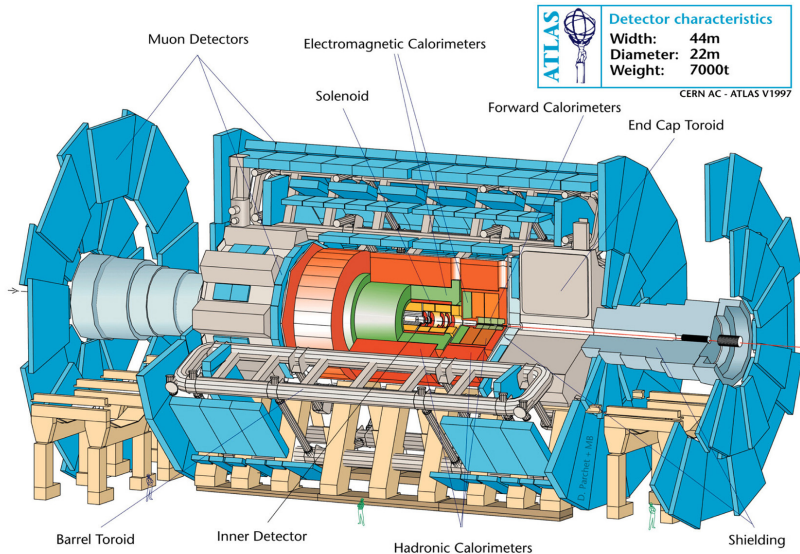
Stefanie Adomeit

LMU Munich

Young Scientists Workshop
Ringberg Castle, July 2012



The ATLAS Detector

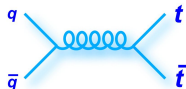
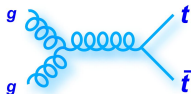
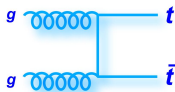


Introduction: The Fully Hadronic $t\bar{t}$ Decay Channel



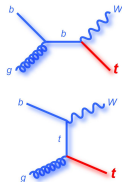
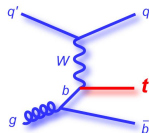
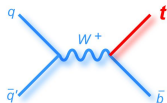
Top Quark Production At The LHC

top anti-top production:



- $\sigma_{t\bar{t}} \approx 164 \text{ pb}$ (@ 7TeV) \Rightarrow good statistics
- gluon-gluon fusion dominant production process

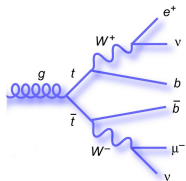
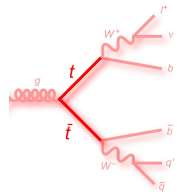
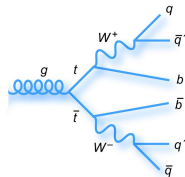
single top production:



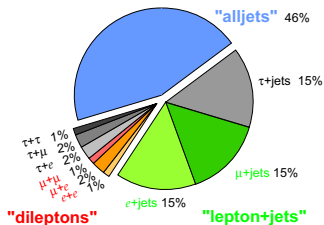
Top Anti-Top Decay Channels

- short lifetime \rightarrow no hadronization
- SM prediction: top exclusively decays into b-quark and W boson
- consecutive W-decay: leptonically (lepton + corresponding neutrino) or hadronically (quark anti-quark pair)

\Rightarrow 3 decay channels: allshadronic, single-lepton, dilepton

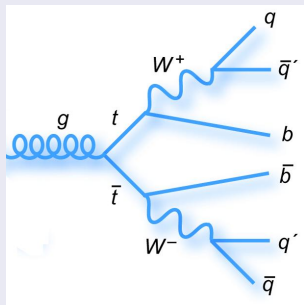


Top Pair Branching Fractions

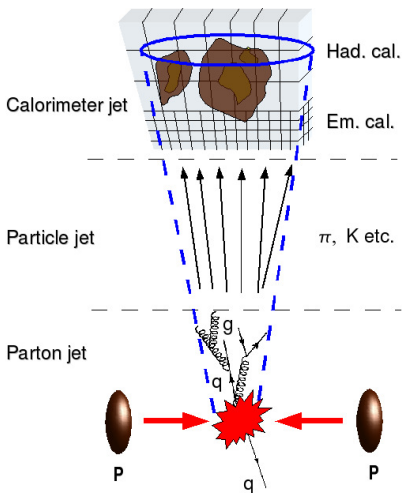


The Fully Hadronic $t\bar{t}$ Decay Channel

- 6 quarks from decaying $t\bar{t}$ pair
- ISR/FSR: additional partons



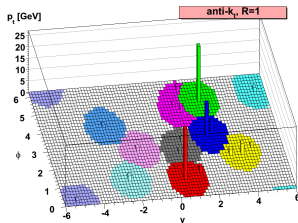
signature: ≥ 6 Jets



Jet Reconstruction

jets are reconstructed using the $\text{anti-}k_t$ clustering algorithm:

- based on $\min(p_i^{-2}, p_j^{-2})$ scaled distance measurement $\Delta_{ij}^2 = (y_i - y_j)^2 + (\phi_i - \phi_j)^2$
- combine hardest objects first until all objects are separated by $\Delta R_{ij} > R$ (ATLAS standard: $R=0.4/0.6$)

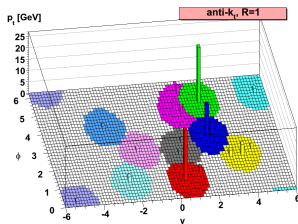


[Cacciari, Salam, Soyez, JHEP 0804:063,2008]

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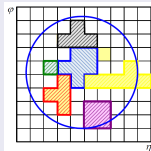


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CALORIMETER JETS

starting from $\text{topological clusters}$:

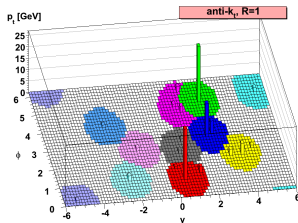
- group of calorimeter cells with a signal-to-noise ratio above a certain threshold
- contain a variable number of cells, no fixed shape



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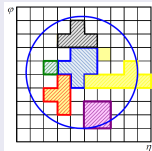


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CALORIMETER JETS

starting from **topological clusters**:

- group of calorimeter cells with a signal-to-noise ratio above a certain threshold
- contain a variable number of cells, no fixed shape



PARTICLE JETS

starting from **stable particles**

- particles with lifetime above 10 ps - excluding muons and neutrinos
- in Monte-Carlo samples

JES Calibration And JES Uncertainty

JES CALIBRATION:

correct kinematics of calorimeter jets to the values of the corresponding truth jet to compensate for

- non-compensating calorimeters
- inactive material
- out-of-cone effects

⇒ calibration constants (evaluated based on MC studies) to restore the JES

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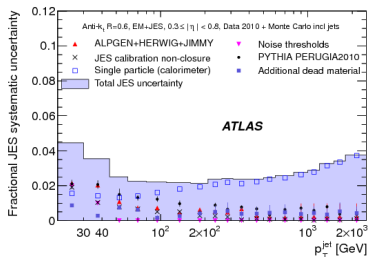
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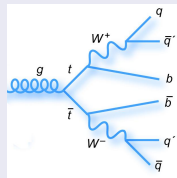
JES UNCERTAINTY = uncertainty on calibration constants:

- 2010 uncertainty: MC based studies (uncertainties on detector simulation, event modelling in Monte Carlo generators,...)
<2.5% for central jets ($p_T = 100$ GeV),
<9 (14)% in endcap (forward) region
- 2011 uncertainty: reduced uncertainty thanks to in-situ measurements



Measurement Of The Top Quark Mass In The Allhadronic $t\bar{t}$ Decay Channel

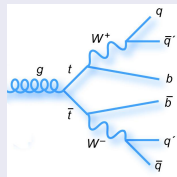
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top mass measurement in the allhadronic channel: challenges \Rightarrow analysis steps

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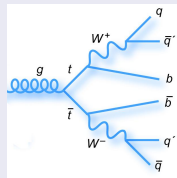


top mass measurement in the allhadronic channel: **challenges** \Rightarrow **analysis steps**

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 - \rightarrow background modelling
 - \rightarrow separation signal \leftrightarrow background

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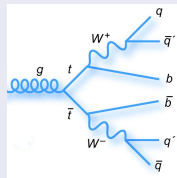


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- 2 large combinatorial background: 6 jets result in 90 possible ways of assigning jets to partons
 - \rightarrow use kinematic likelihood fit to reconstruct allhadronic top quarks
- 3 sensitive to systematic uncertainties e.g. jet energy scale
 - \rightarrow evaluate the top quark mass via 2D template method

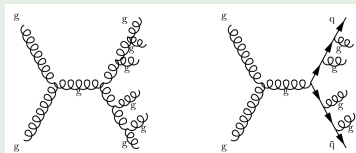
1. QCD Multijet Background



QCD-Multijet Background: Modelling

QCD Multijet Production:

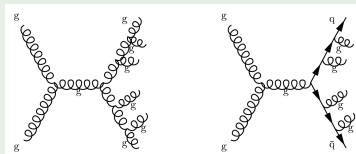
- parton pair (except $t\bar{t}$) from hard scatter
- additional partons from gluon radiation, gluon splitting



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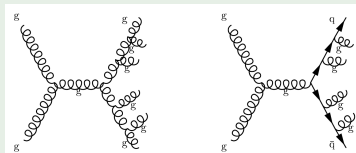


☹ Monte-Carlo based modelling of QCD multijet production:
final states with seven or more partons cannot be simulated accurately (production matrix elements not available)

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☺ Monte-Carlo based modelling of QCD multijet production:
final states with seven or more partons cannot be simulated accurately (production matrix elements not available)

☺ however, we have lots of data and the multijet production cross section is rather large...

→ model background based on a selected data sample dominated by QCD multijet production

Data Driven QCD-Multijet Background Modelling...

...an example: **EVENT MIXING**

Donour Events

select events with $n \geq 6$ jets and 0 leptons (\rightarrow expect $t\bar{t}$ contamination)

select 5th-nth jet with lowest
transverse momentum



Acceptor events

events with $n = 4$ jets and 0 leptons
 \rightarrow no significant $t\bar{t}$ contribution

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add jets from donour events to acceptor events if...

- 1 highest/lowest transverse momentum jets in donour and acceptor events are similar in p_t
- 2 there is no overlap between donour and acceptor event jets

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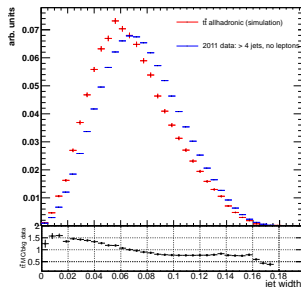
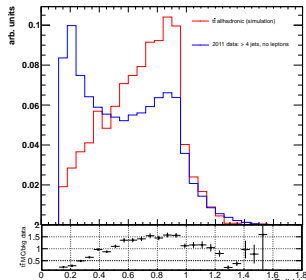
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\Rightarrow multijet events with $n \geq 6$ jets 😊

Separating $t\bar{t}$ From QCD Multijet Events

- **b-jet identification algorithms** (see next slide): require exactly 2 b-tagged jets
- **different b-production in signal and background**: large fraction of b-jets in multijet events due to gluon splitting → collinear emission
- **W-mass window**: check dijet invariant mass of light jets from W-decay
- **larger fraction of gluon jets in multijet events**: exploit the different properties of light quark and gluon initiated jets (jet width, number of jet constituents)
- ...



b-Jet Identification

long lifetime of b-flavoured hadrons

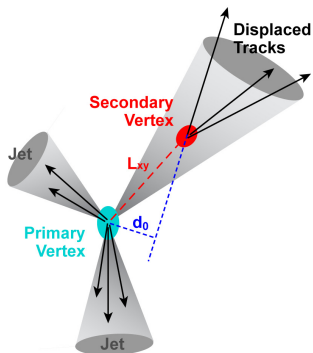
- significant **flight path length L** of b-hadrons
- measurable **secondary vertex**
- measurable **impact parameters** of tracks from b-hadron decay products

→ spatial taggers

semi-leptonic b-hadron decays

- tag soft leptons (from b-hadron decay) in the jets

→ soft lepton taggers



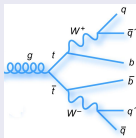
2. Finding Jet Parton Pairs - Kinematic Likelihood Fit



Reconstructing The Allhadronic Event Topology: Kinematic Likelihood Fit

MODEL: allhadronic $t\bar{t}$
event topology

- 4 light quark jets from W-decay
- 2 b-jets from $t\bar{t}$ decay

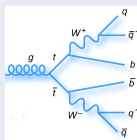


ASSUMPTIONS:

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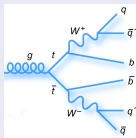
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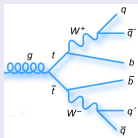
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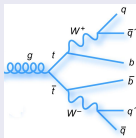
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→ use **transfer functions** which map the parton energy to the measured (calibrated) jet energy



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$$L = BW\{m(q_1, q_2)|m_W, \Gamma_W\} \cdot BW\{m(q_4, q_5)|m_W, \Gamma_W\} \cdot$$

$$BW\{m(q_1, q_2, q_3)|m_{top}, \Gamma_{top}\} \cdot BW\{m(q_4, q_5, q_6)|m_{top}, \Gamma_{top}\} \cdot \prod_{i=1}^6 W(\tilde{E}_i, E_i)$$

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⇒ include b-tagging information: veto permutations with **b-jets** on the position of **light quarks** from W decay

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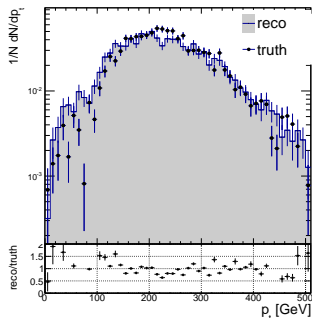
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Monte-Carlo based performance studies:

- right jet to parton assignment in 80% of all events when including b-tagging veto
- good agreement between reconstructed and true kinematical variables of top quarks

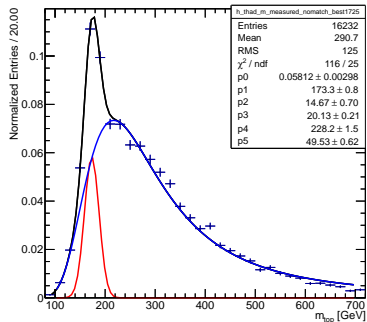


3. Measuring The Top Quark Mass - Building Templates



Measuring The Top Quark Mass: Building 1D Templates

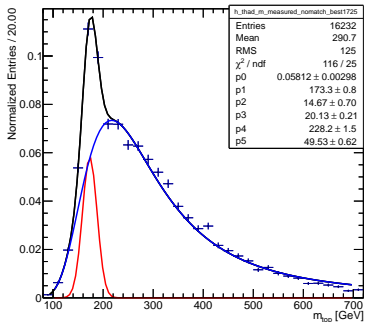
⇒ parameterize m_{top}^{reco} as a function of m_{top}



$$\text{Gaussian}(p_0, p_1, p_2) + \text{Landau}(p_3, p_4, p_5)$$

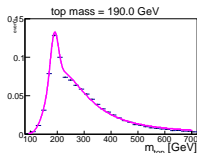
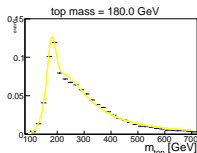
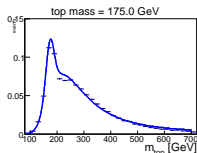
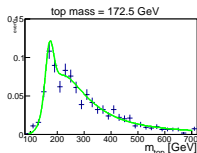
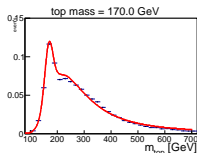
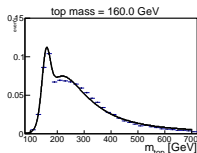
Measuring The Top Quark Mass: Building 1D Templates

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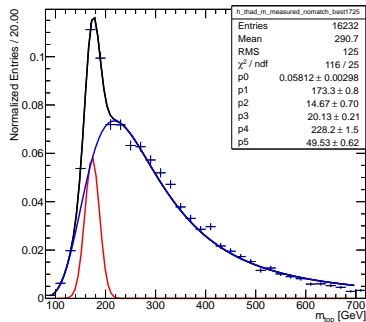
$$\text{Gaussian}(p_0, p_1, p_2) + \text{Landau}(p_3, p_4, p_5)$$

$$p_i = p_i(m_{top})$$



Measuring The Top Quark Mass And The Jet Energy Scale: Building 2D Templates

⇒ parameterize m_{top}^{reco}
as a function of m_{top} and ΔJES

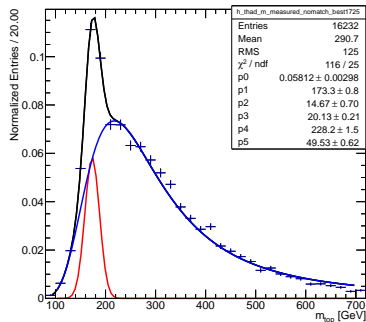


Gaussian(p_0, p_1, p_2) + Landau(p_3, p_4, p_5)

$$p_i = p_i(m_{top}, \Delta JES)$$

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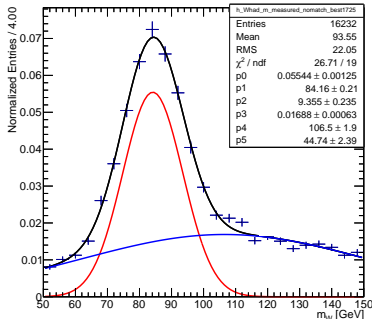
⇒ parameterize m_{top}^{reco}
as a function of m_{top} and ΔJES



Gaussian(p_0, p_1, p_2) + Landau(p_3, p_4, p_5)

$$p_i = p_i(m_{top}, \Delta JES)$$

⇒ parameterize m_W^{reco}
as a function of ΔJES

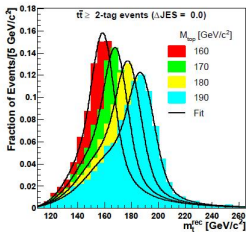


Gaussian(p_0, p_1, p_2) + Gaussian(p_3, p_4, p_5)

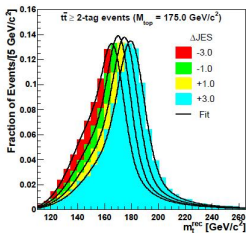
$$p_i = p_i(\Delta JES)$$

Building 2D Templates

Reconstructed Top Mass



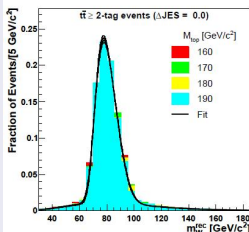
sensitive to
top mass



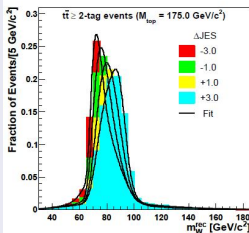
sensitive to
JES

⇒ parameterize $m_{\text{top}}^{\text{reco}}$ as a function of m_{top} and ΔJES

Reconstructed W Mass



insensitive to
top mass



sensitive to
JES

⇒ parameterize m_W^{reco} as a function of ΔJES

Measuring The Top Quark Mass Using Templates

- signal p.d.f.s: $P_s^{m_{top}^{reco}}(m_{top}^{reco} | m_{top}, \Delta JES)$, $P_s^{m_W^{reco}}(m_W^{reco} | \Delta JES)$ (previous slide)
- background p.d.f.s: $P_b^{m_{top}^{reco}}(m_{top}^{reco})$, $P_b^{m_W^{reco}}(m_W^{reco})$, independent of m_{top} and ΔJES (data driven estimate of QCD background)

⇒ find m_{top} and ΔJES which best reproduce the observed distribution of m_{top}^{reco} and m_W^{reco} in data

Measuring The Top Quark Mass Using Templates

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⇒ find m_{top} and ΔJES which best reproduce the observed distribution of m_{top}^{reco} and m_W^{reco} in data

recent results on the measurement of the top quark mass in the allhadronic channel:

- 2D: arXiv:1112.4891v2, submitted to Phys. Lett. B
→ $m_{top} = 172.5 \pm 1.4$ (stat) ± 1.0 (JES) ± 1.1 (syst) GeV/c²
- 1D: ATLAS-CONF-2012-030
→ $m_{top} = 174.9 \pm 2.1$ (stat) ± 3.8 (syst) GeV/c²

Summary

fully hadronic $t\bar{t}$ decay channel provides good statistics due to

- 1 relatively high $t\bar{t}$ production cross section
- 2 large allhadronic branching ratio

when measuring the allhadronic top mass one needs to get a handle on...

- the QCD background:
 - data driven modelling
 - finding efficient cuts in order to separate signal from background events
- the jet-to-parton assignment → kinematic likelihood fit
- the systematic uncertainties → in-situ measurement of light-JES via 2D template method