

Electroweak Corrections to squark-antisquark pair production at the LHC

Edoardo Mirabella

Max-Planck-Institut
für Physik
(Werner-Heisenberg-Institut)



In collaboration with W. Hollik.

Outline

- Motivations
- Closer look into (diagonal) squark pair production
 - State of the art
 - squark pair production at parton level
 - IR divergencies
- Numerical Analysis
- Conclusions

Electroweak Corrections to squark-antisquark pair production at the LHC

- LHC ... “The” event of 2008

- ↳ Proton proton collider at 14 TeV.

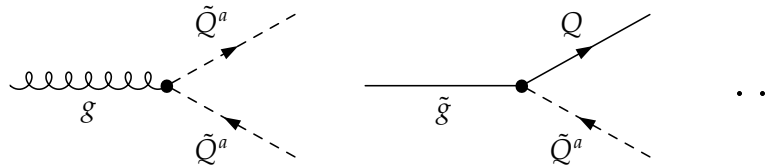
- ↳ Multi purpose collider ... Among the major goals SUSY searches

Electroweak Corrections to squark-antisquark pair production at the LHC

● Squarks ... why?

↪ SUSY partners of quarks.

↪ Colored particles



↪ ... detectable (Same sign di-lepton signal, \cancel{E}_T , jet multiplicity)

Electroweak Corrections to squark-antisquark pair production at the LHC

● EW corrections ... why?

↪ Formally higher order with respect to QCD corrections (already available) but ...

↪ ... potentially important for distributions

● $PP \rightarrow t\bar{t}$ in SM: EW contribution to p_T distribution from -5 to -15 %
[Kühn *et al* '06]

● $PP \rightarrow W j$ in SM: EW corrections on p_T distribution amount up to 40%.
[Kühn *et al* '07, Hollik *et al* '07]

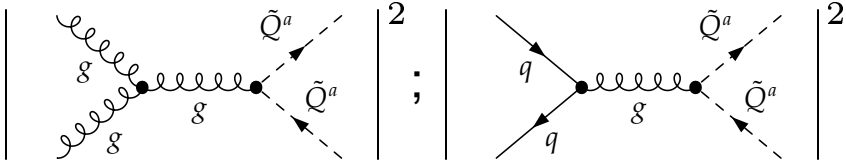
Diagonal squark PP, status

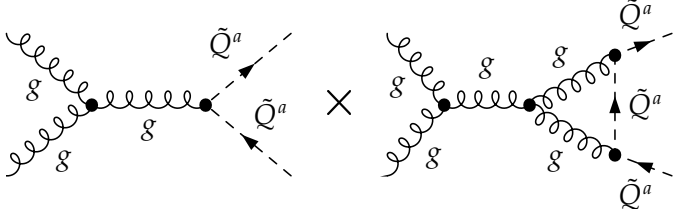
$\bullet \mathcal{O}(\alpha_s^2):$
[Baer and Tata '85]

$\bullet \mathcal{O}(\alpha_s^3):$
[Beenakker et al '96, '98]

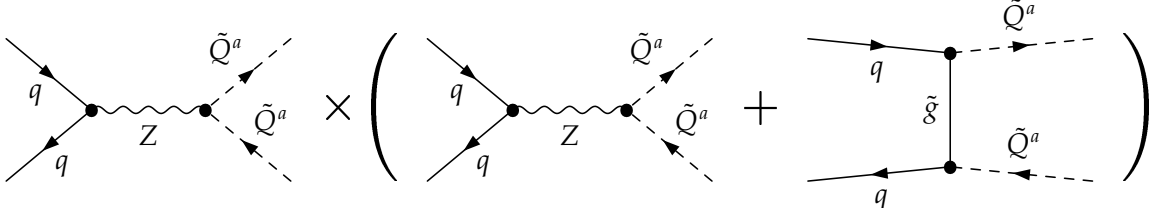
\hookrightarrow big positive corrections to the total cross section (from 5 to 90 %)

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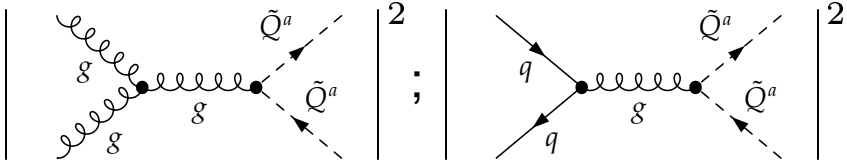
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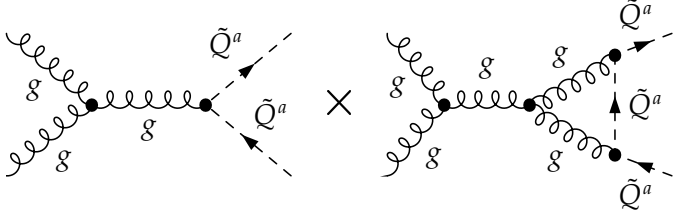
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$\bullet \mathcal{O}(\alpha_s \alpha + \alpha^2):$

[Bornhauser *et al* '07]

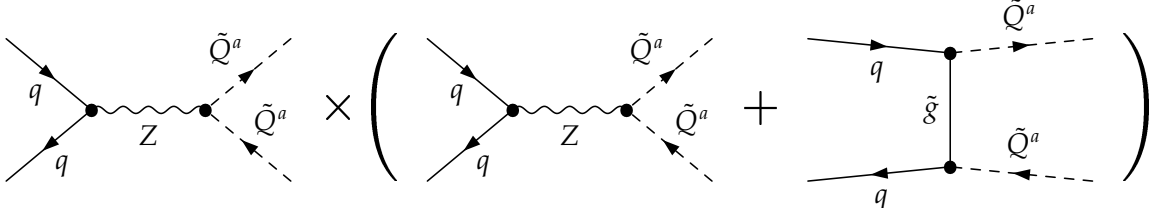
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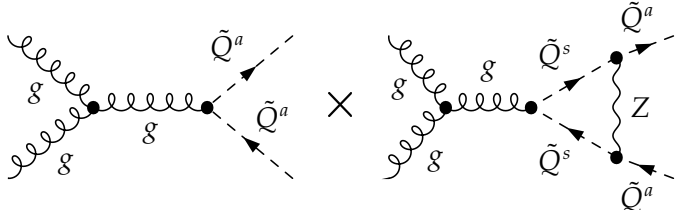
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$\bullet \mathcal{O}(\alpha_s \alpha + \alpha^2):$

[Bornhauser et al '07]

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$\bullet \mathcal{O}(\alpha_s^2 \alpha):$ stop case only

[Hollik et al '07]

↪ 10 – 20 % level in p_T & Invariant Mass distributions

$\bullet \mathcal{O}(\alpha_s^2 \alpha)$ corrections for arbitrary squarks are still missing ...

... And the generalization from the stop case is not trivial!

Diagonal squark PP at the parton level

Stop case, $\mathcal{O}(\alpha_s^2 \alpha)$

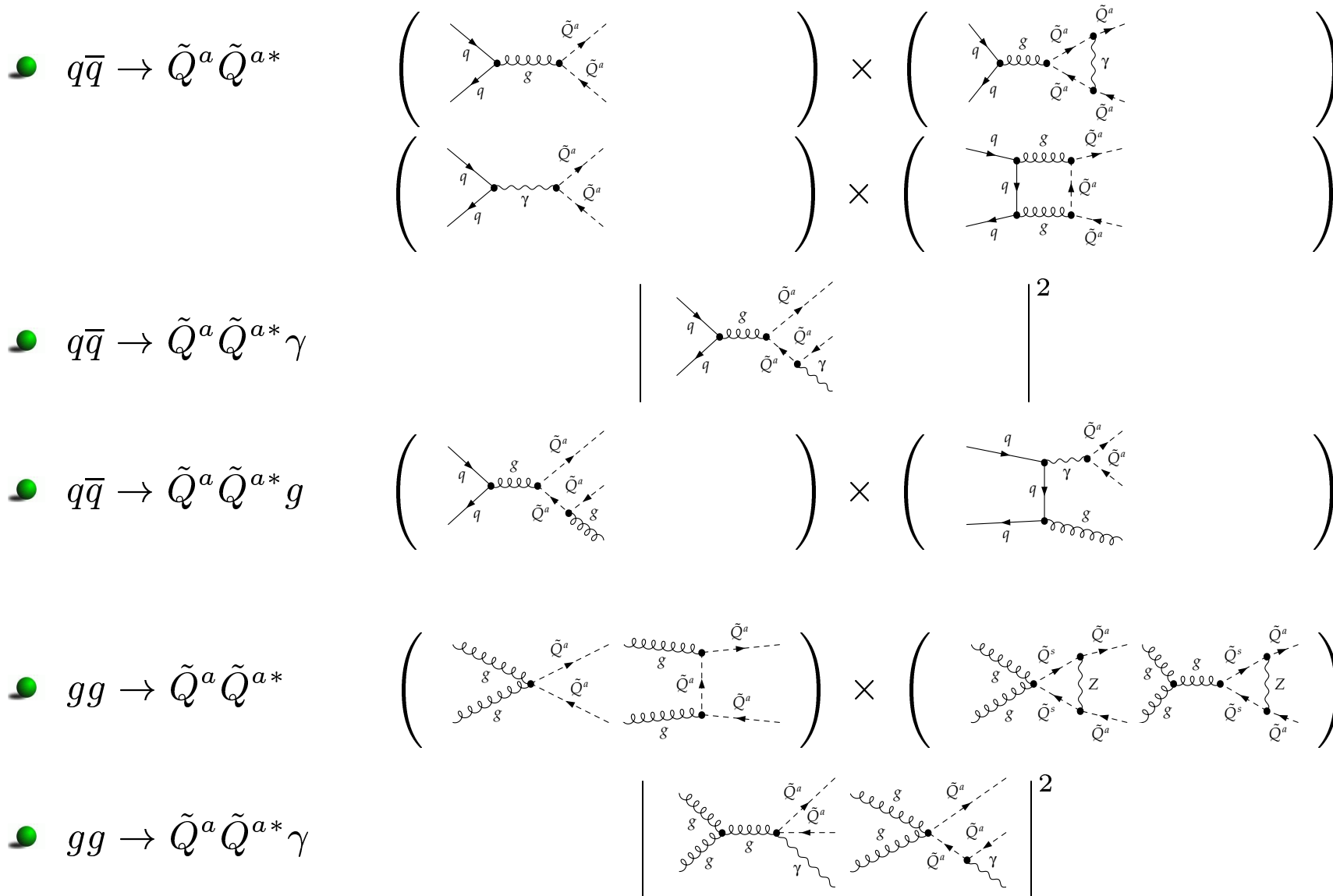
$\bullet \quad q\bar{q} \rightarrow \tilde{Q}^a \tilde{Q}^{a*}$
 $\left(\begin{array}{c} \text{[Diagram 1]} \\ \text{[Diagram 2]} \end{array} \right) \times \left(\begin{array}{c} \text{[Diagram 3]} \\ \text{[Diagram 4]} \end{array} \right)$

$\bullet \quad q\bar{q} \rightarrow \tilde{Q}^a \tilde{Q}^{a*} \gamma$
 $\left| \begin{array}{c} \text{[Diagram 5]} \end{array} \right| \begin{array}{c} 2 \\ | \end{array}$

$\bullet \quad q\bar{q} \rightarrow \tilde{Q}^a \tilde{Q}^{a*} g$
 $\left(\begin{array}{c} \text{[Diagram 6]} \end{array} \right) \times \left(\begin{array}{c} \text{[Diagram 7]} \end{array} \right)$

Diagonal squark PP at the parton level

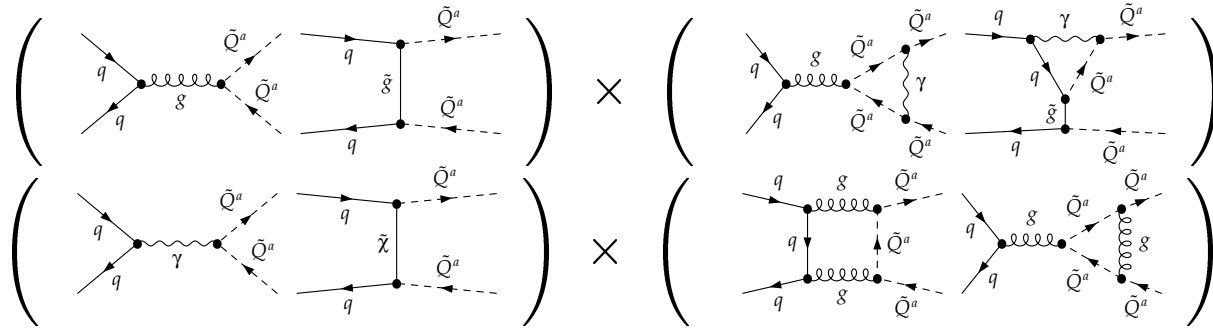
Stop case, $\mathcal{O}(\alpha_s^2 \alpha)$



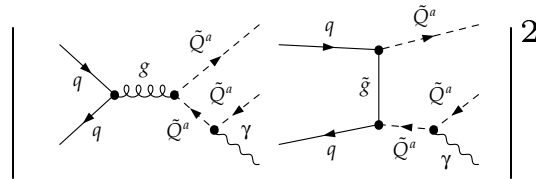
Diagonal squark PP at the parton level

General case, $\mathcal{O}(\alpha_s^2 \alpha)$

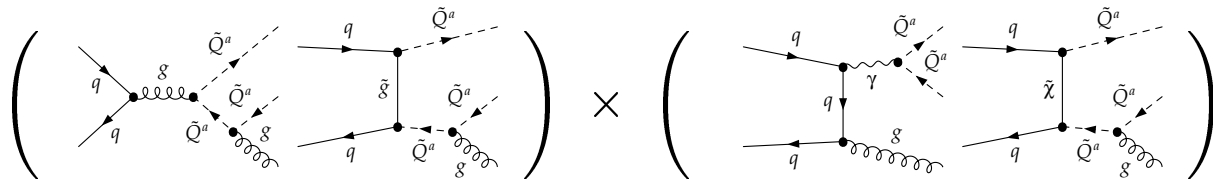
$q\bar{q} \rightarrow \tilde{Q}^a \tilde{Q}^{a*}$



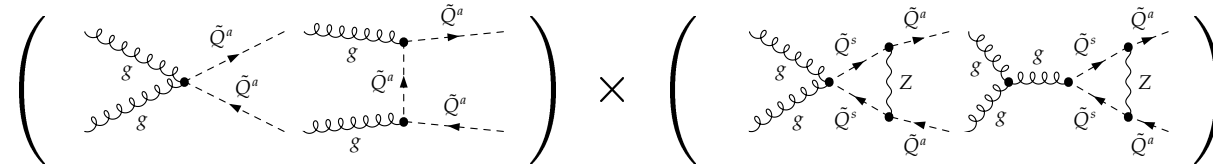
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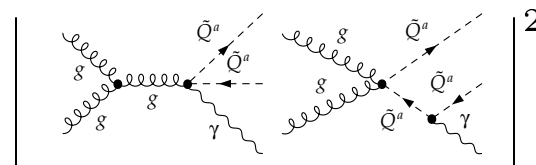
$q\bar{q} \rightarrow \tilde{Q}^a \tilde{Q}^{a*} g$



$gg \rightarrow \tilde{Q}^a \tilde{Q}^{a*}$



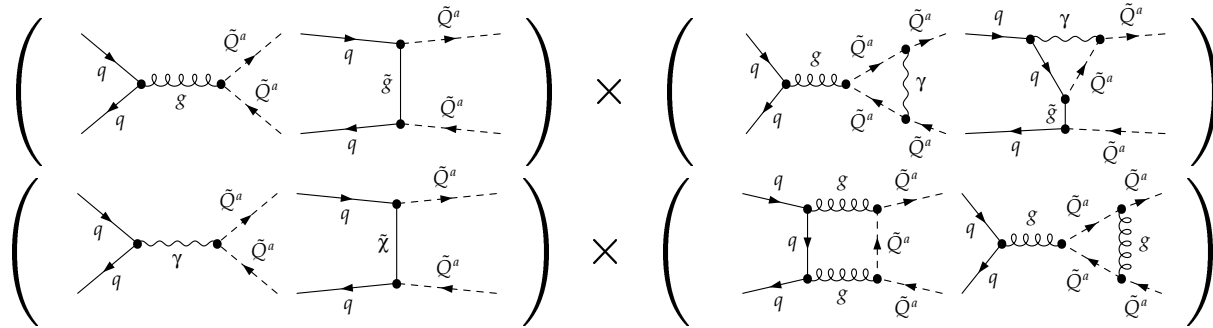
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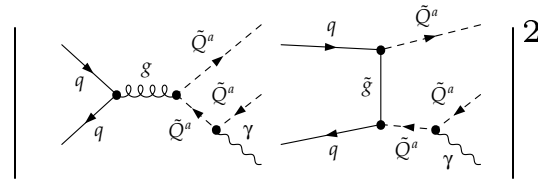
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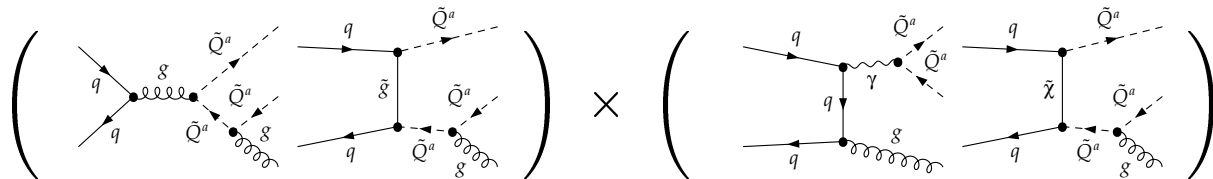
● $q\bar{q} \rightarrow \tilde{Q}^a \tilde{Q}^{a*}$



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● $q\bar{q} \rightarrow \tilde{Q}^a \tilde{Q}^{a*} g$

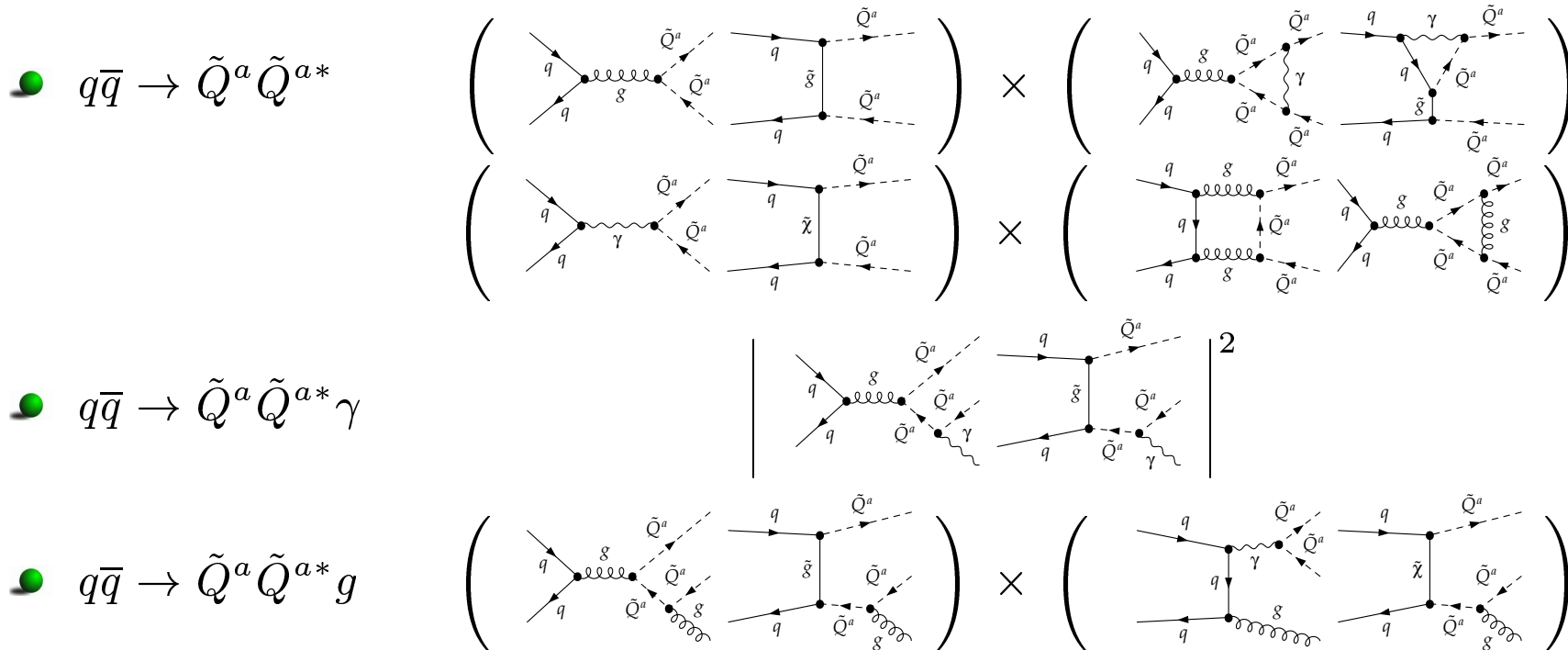


($q = Q$) \implies (tree level QCD (EW) with \tilde{g} ($\tilde{\chi}^0$) in the T-channel)
 ((q, Q) SU(2) doublet) \implies (tree level EW with $\tilde{\chi}^+$ in the T-channel)

Numerically Speaking: New interferences can be important (from $\mathcal{O}(\alpha_s \alpha)$ experience)

Diagonal squark PP at the parton level

General case, $\mathcal{O}(\alpha_s^2 \alpha)$

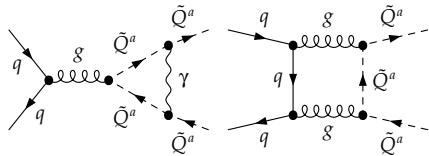


$(q = Q) \implies (\text{tree level QCD (EW) with } \tilde{g} (\tilde{\chi}^0) \text{ in the T-channel})$
 $((q, Q) \text{ SU}(2) \text{ doublet}) \implies (\text{tree level EW with } \tilde{\chi}^+ \text{ in the T-channel})$

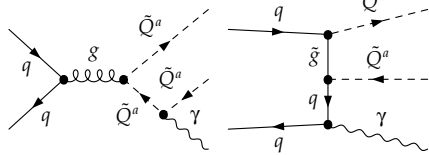
Technically speaking:

- Full one loop QCD corrections needed
- New UV divergencies $\implies \delta\alpha_s$ & new CTT's enter the game
- Richer IR structure (more on this later)

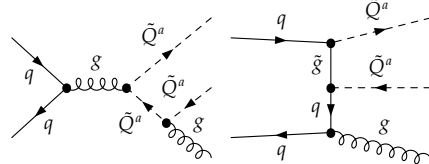
IR & Collinear Divergences



... \rightsquigarrow IR & Collinear singularities to $q\bar{q} \rightarrow \tilde{Q}^a \tilde{Q}^{a*}$

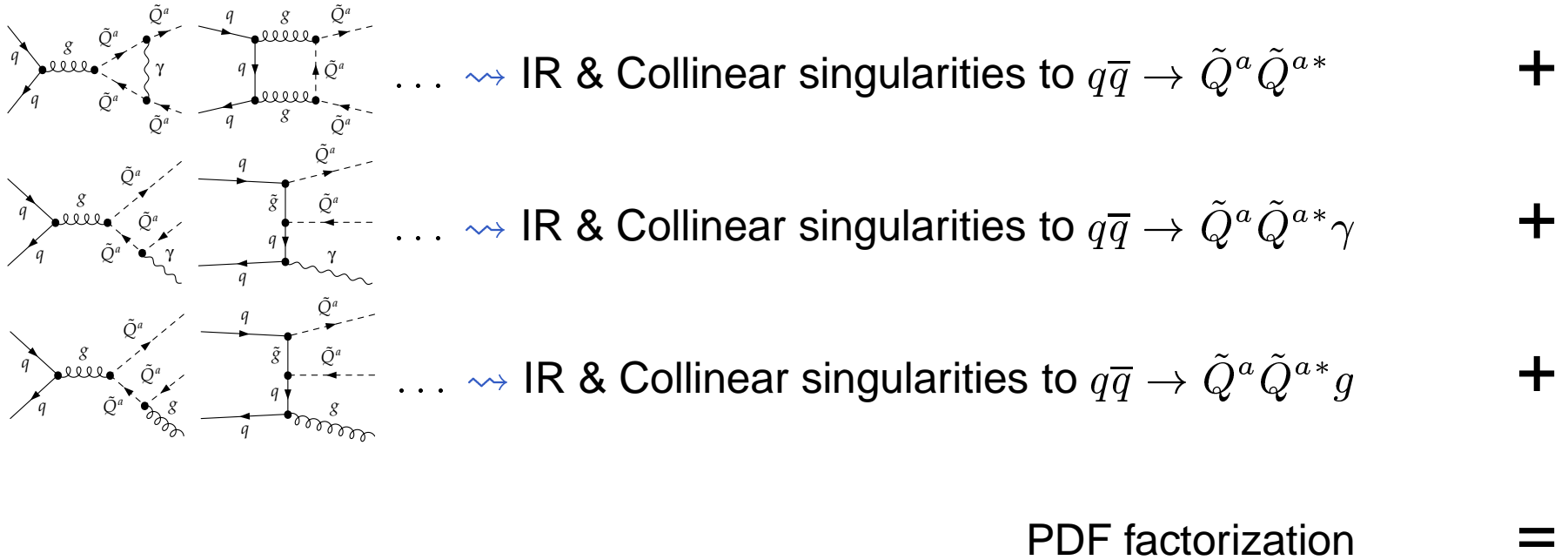


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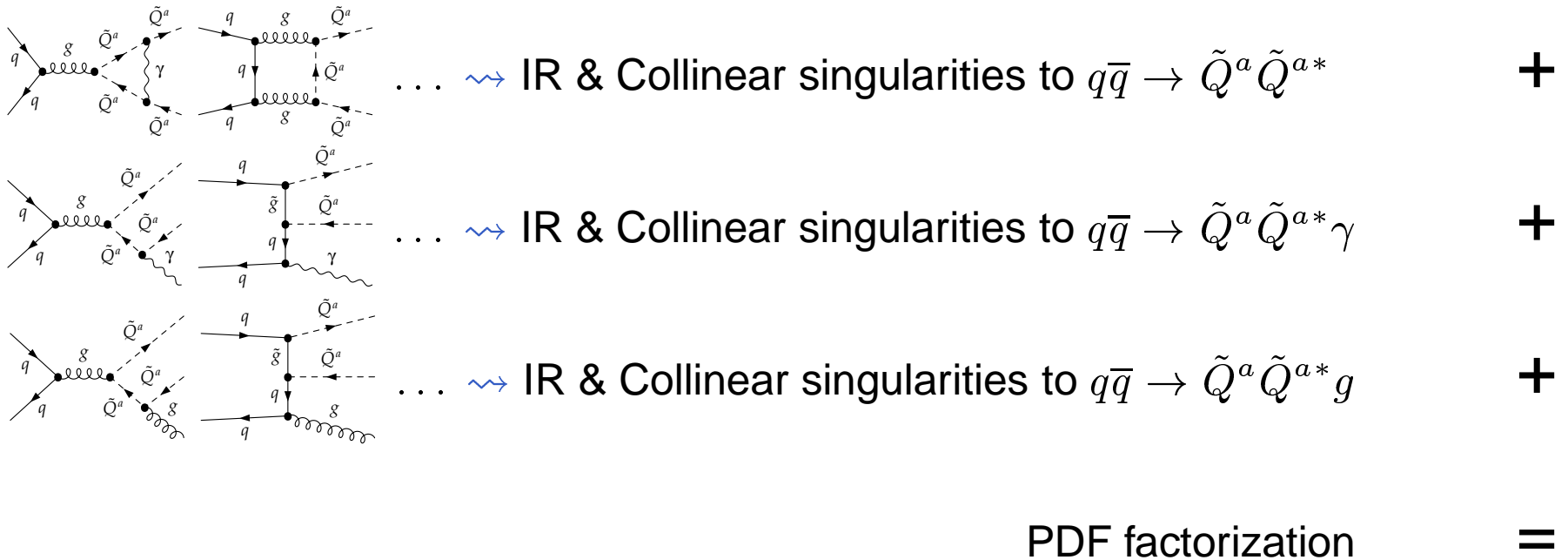
IR & Collinear Divergences



[KLN theorem]

IR & Collinear Safe

IR & Collinear Divergences



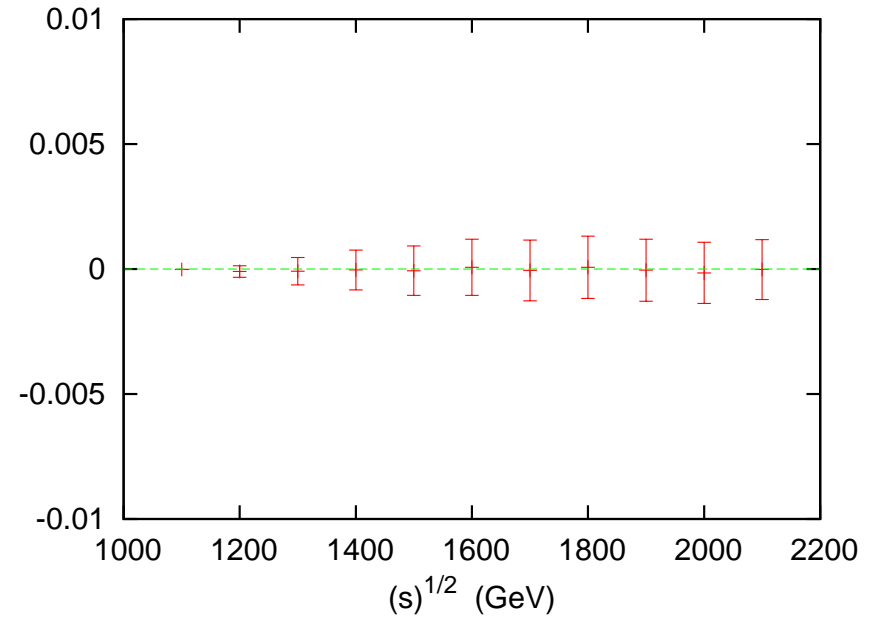
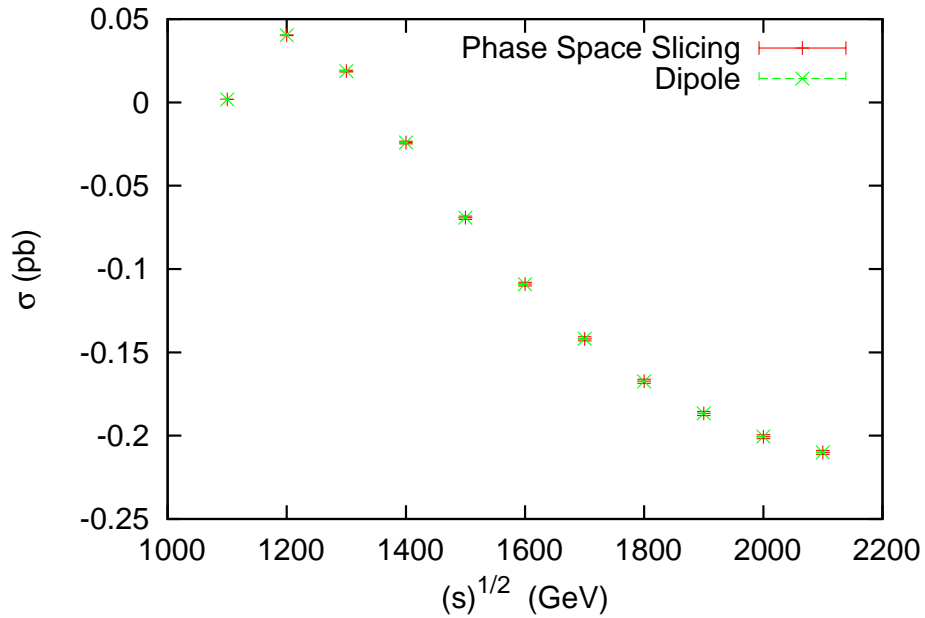
[KLN theorem]

IR & Collinear Safe

- Mass regularization in both cases (g singularities essentially Abelian)
- Two methods for γ (g) phase space integration:
 - Phase Space Slicing
 - Dipole Formalism
- Colour Algebra introducing color-charge operators in color space

IR & Collinear Divergences

cross check, $u\bar{u} \rightarrow \tilde{u}^a \tilde{u}^{a*} g$

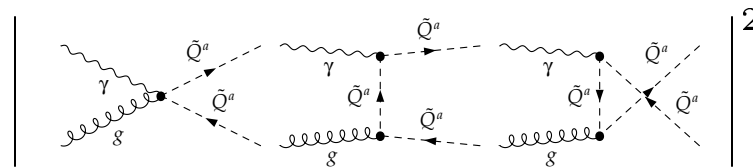


$\Delta = (\text{Dipole} - \text{Slicing}), \text{ point SPS1a}'$

Numerical Results

General Informations

- The $\mathcal{O}(\alpha_s \alpha + \alpha^2)$ interference contributions have been included
↳ their importance already known [[Bornhauser et al](#)]
- The photon induced process $g\gamma \rightarrow \tilde{Q}^a \tilde{Q}^{a*}$:



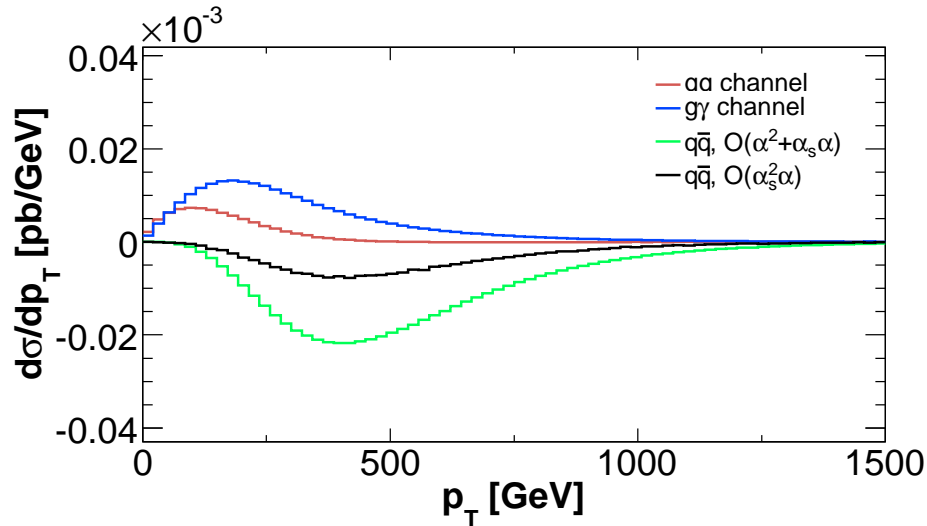
has been included

↳ in the stop case important [[Hollik et al](#)]

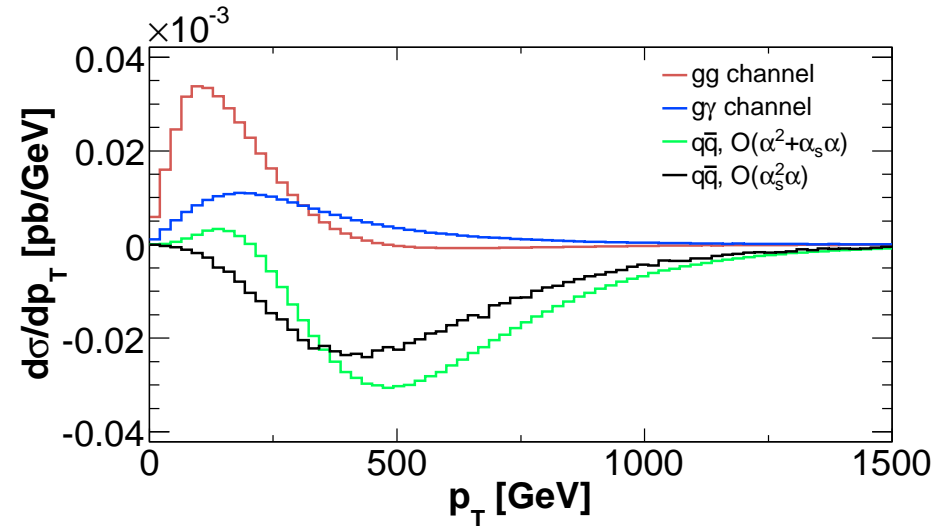
Numerical Results

Dependence on the Chirality (Transverse Momentum distribution), $\left[\delta = \frac{\mathcal{O}^{\text{NLO}} - \mathcal{O}^{\text{LO}}}{\mathcal{O}^{\text{NLO}}} \right]$

\tilde{u}^R , point SPS1a'



\tilde{u}^L , point SPS1a'

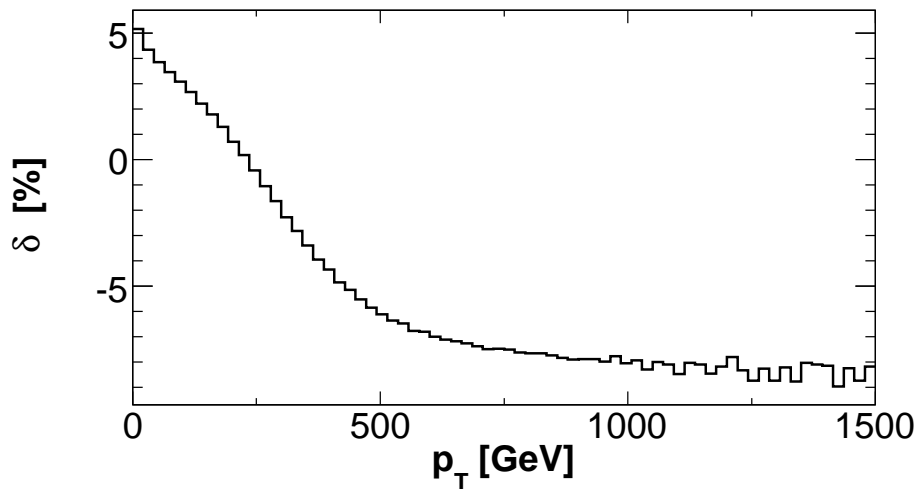


- $g\gamma$ channel chirality-independent, the others more important in the \tilde{u}^L case
- $\mathcal{O}(\alpha_s^2 \alpha)$ & $\mathcal{O}(\alpha_s \alpha + \alpha^2)$ comparable (at least in the left handed case)
- $q\bar{q}$ @ $\mathcal{O}(\alpha_s \alpha^2)$ have different sign in the low p_T region in the two cases
 $\hookrightarrow (d\bar{d} \rightarrow \tilde{u}^L \tilde{u}^{L*} > 0, d\bar{d} \rightarrow \tilde{u}^R \tilde{u}^{R*} = 0)$

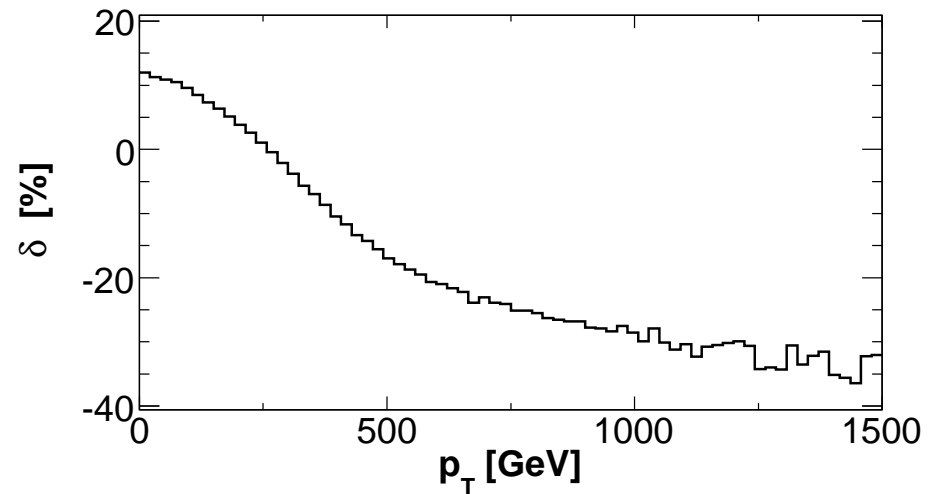
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\tilde{u}^R , point SPS1a'



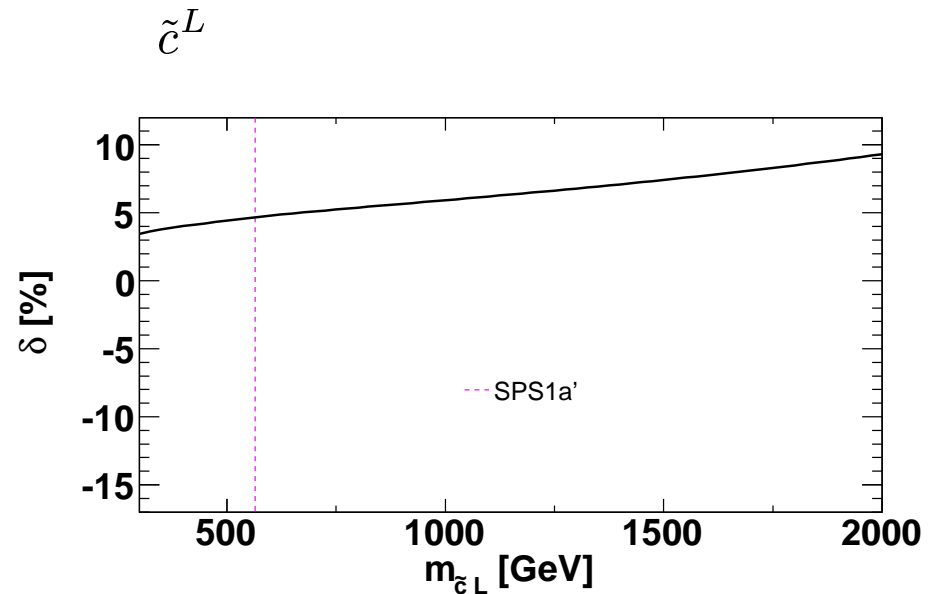
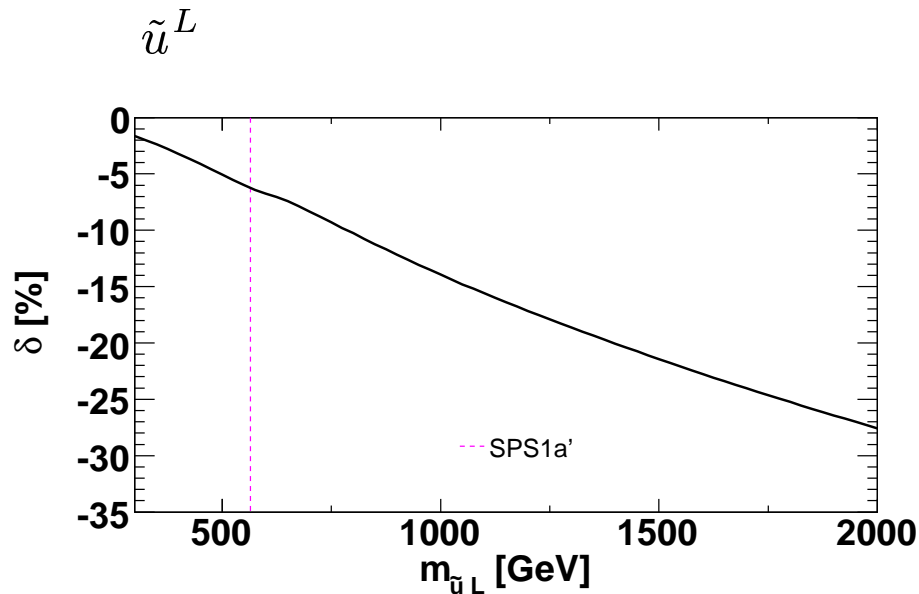
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- $q\bar{q}$ @ $\mathcal{O}(\alpha_s\alpha^2)$ have different sign in the low p_T region in the two cases
 $\rightarrow (d\bar{d} \rightarrow \tilde{u}^L\tilde{u}^{L*} > 0, d\bar{d} \rightarrow \tilde{u}^R\tilde{u}^{R*} = 0)$
- % contributions of NLO EW corrections are important in the left handed case.

Numerical Results

Dependence on the Flavour (Total Cross Section V_s squark mass), $\left[\delta = \frac{\sigma^{\text{NLO}} - \sigma^{\text{LO}}}{\sigma^{\text{NLO}}} \right]$

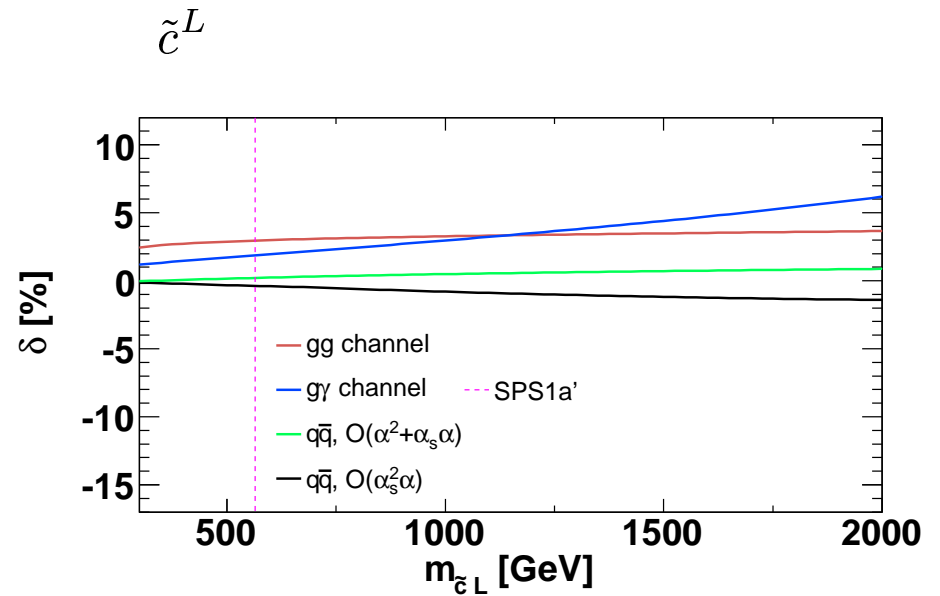
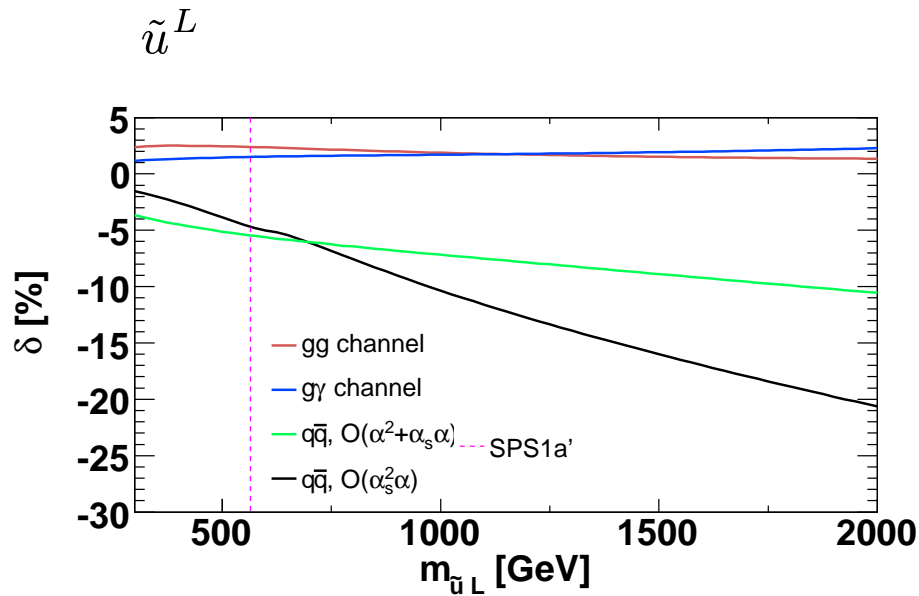


- NLO EW corrections differ in the two cases

\tilde{u} :	Negative, enhanced if $m_{\tilde{u}}$ big
\tilde{c} :	Positive, below 10 %

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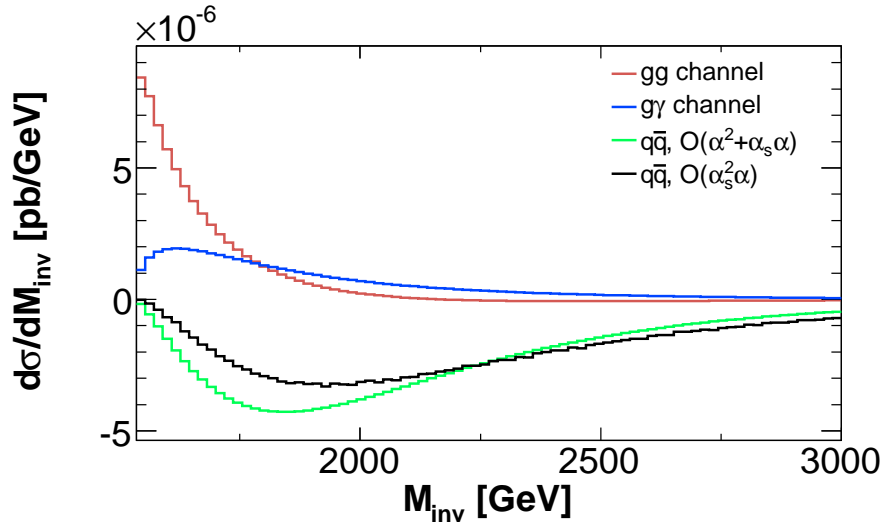


- NLO EW corrections differ in the two cases $\left\{ \begin{array}{l} \tilde{u} : \text{Negative, enhanced if } m_{\tilde{u}} \text{ big} \\ \tilde{c} : \text{Positive, below } 10 \% \end{array} \right.$
- The leading channels are different in the two cases (PDF suppression for \tilde{c})
- \tilde{u} case: $q\bar{q}$ @ $\mathcal{O}(\alpha_s^2 \alpha)$ leading at high $m_{\tilde{u},L}$ values
 - ↪ Sudakov enhancement + enhancement of $q\bar{q}$ channel @ LO.
 - ↪ $q\bar{q}$ channel @ $\mathcal{O}(\alpha_s \alpha + \alpha^2)$ suppressed at high $m_{\tilde{u},L}$ values

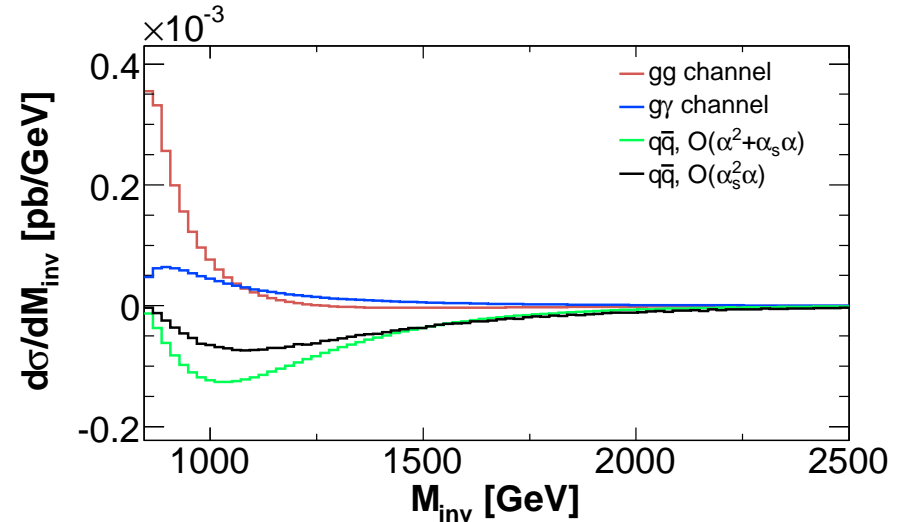
Numerical Results

Dependence on the Benchmark Point (Invariant Mass distribution for \tilde{u}^L prod.)

point SU1, $m_{\tilde{u},L} = 766$ GeV



point SU4, $m_{\tilde{u},L} = 420$ GeV



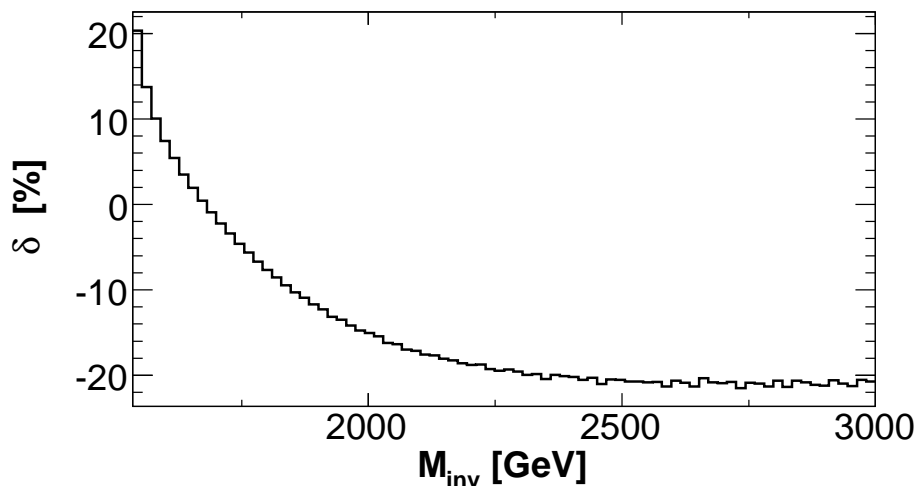
• The size of the contributions decrease as the squark mass increases

• Corrections dominated by $\begin{cases} gg, & g\gamma \\ q\bar{q} \end{cases}$ channels in the $\begin{cases} \text{low} \\ \text{high} \end{cases} M_{\text{inv}}$ region

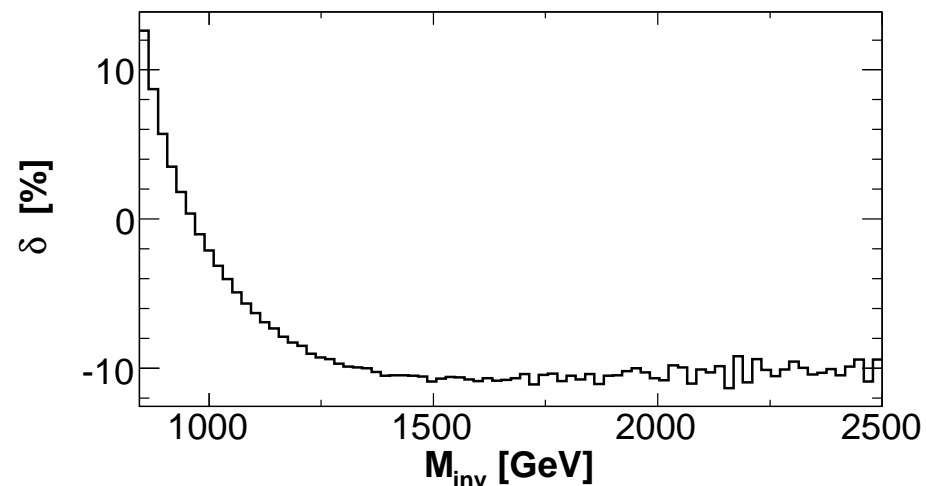
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- Corrections dominated by $\begin{cases} gg, g\gamma \\ q\bar{q} \end{cases}$ channels in the $\begin{cases} \text{low} \\ \text{high} \end{cases} M_{\text{inv}}$ region
- NLO EW corrections are sizable, their importance increase as the squark mass increases

Conclusions & Outlook

- (diagonal) squark-antisquark pair production is an important process in the hunting for SUSY
- $\mathcal{O}(\alpha_s^2 \alpha)$ corrections to this process are available for arbitrary squark-antisquark pair
- They strongly depend on flavour, chirality and mass of the final states
- They can give rise to sizable corrections not only to the distributions but also to the total cross section.

road map:

- Compute the $\mathcal{O}(\alpha_s^2 \alpha)$ corrections to other processes of production of SUSY colored particle:
 - gluino pair production → completed
 - associated production of a squark and a gluino → w.i.p.
 - (non diagonal) squark pair production → on the wish list
- Link these corrections to PROSPINO