

Next-to-leading order predictions for production and decay of squarks

in collaboration with Wolfgang Hollik and Davide Pagani



MAX-PLANCK-GESELLSCHAFT



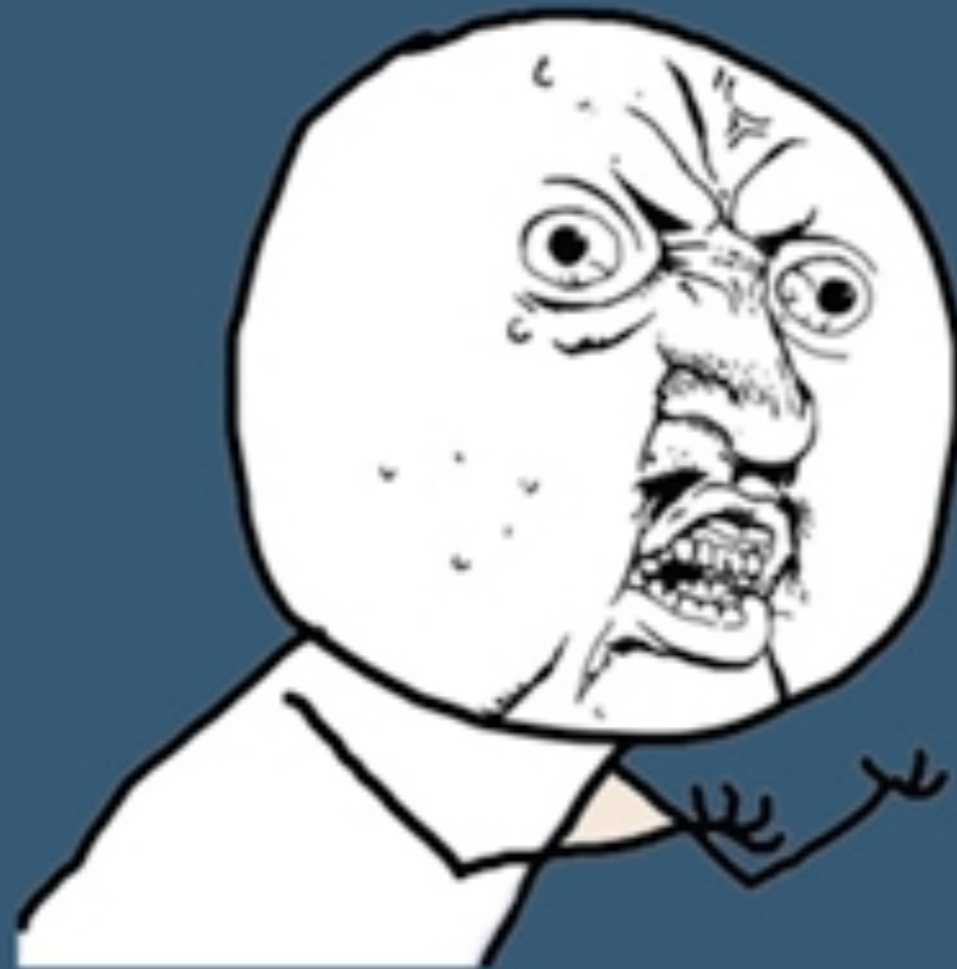
Max-Planck-Institut für Physik
(Werner-Heisenberg-Institut)

Jonas M. Lindert

Max Planck Institut für Physik, München

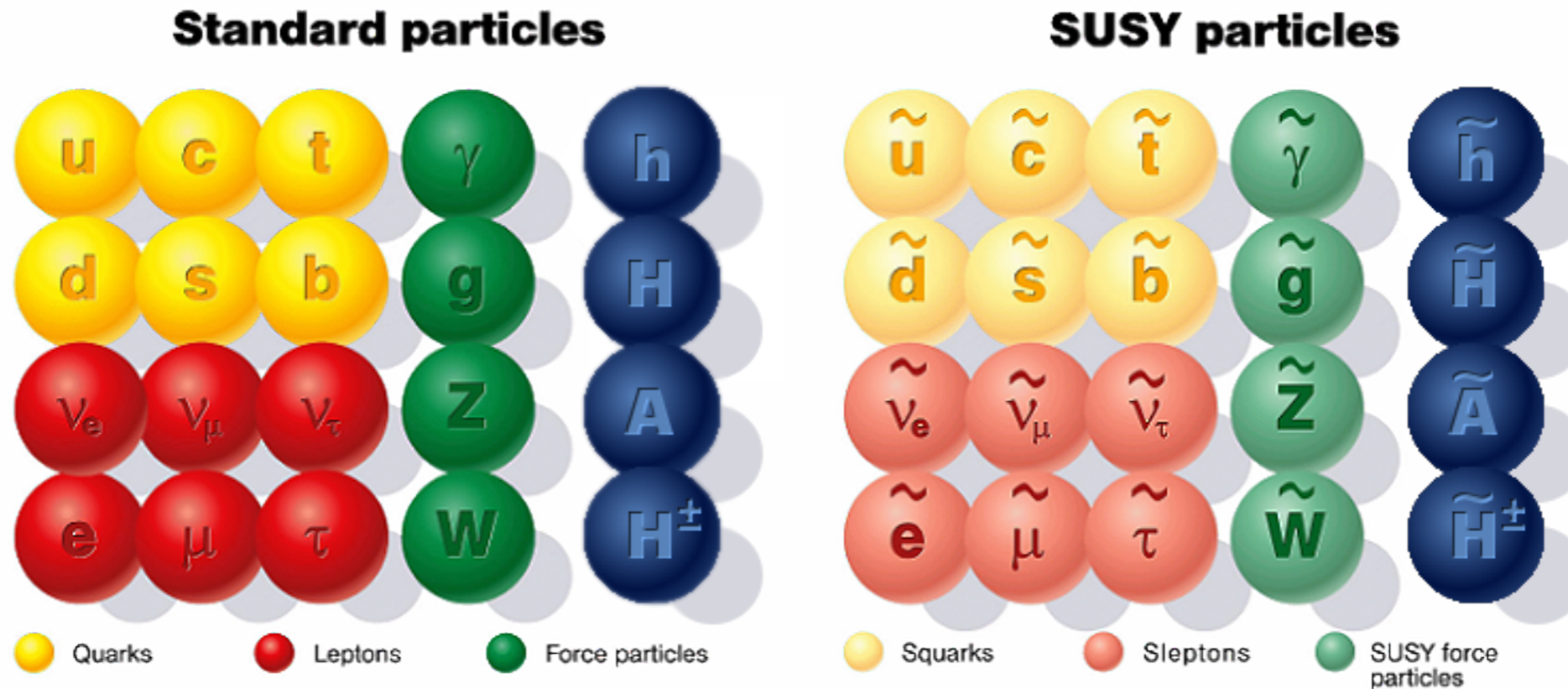
Particle Physics School Munich Colloquium
MPP, Munich, 11 January 2013

LHC



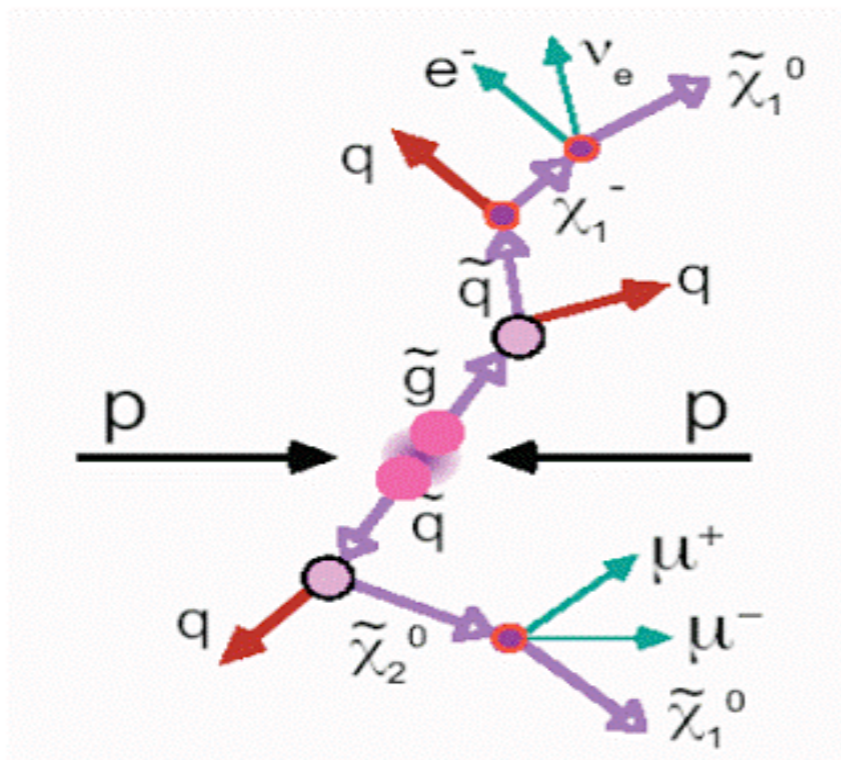
Y U NO SEE SUSY?

THE MSSM AT THE LHC



Often the lightest neutralino assumed to be the (stable) LSP \rightarrow missing Energy

e.g.



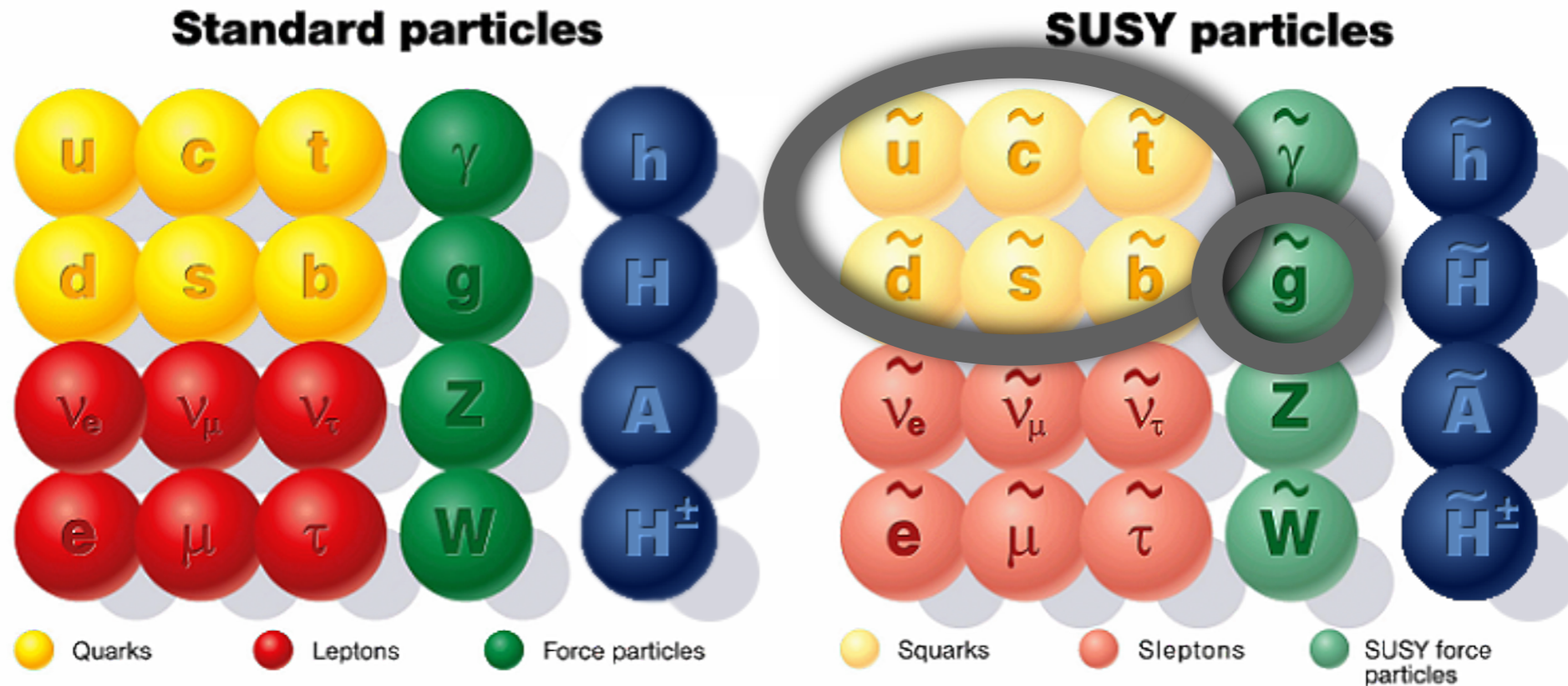
but also:

- more or less jets
- more or less leptons
- b-tagged jets

in general

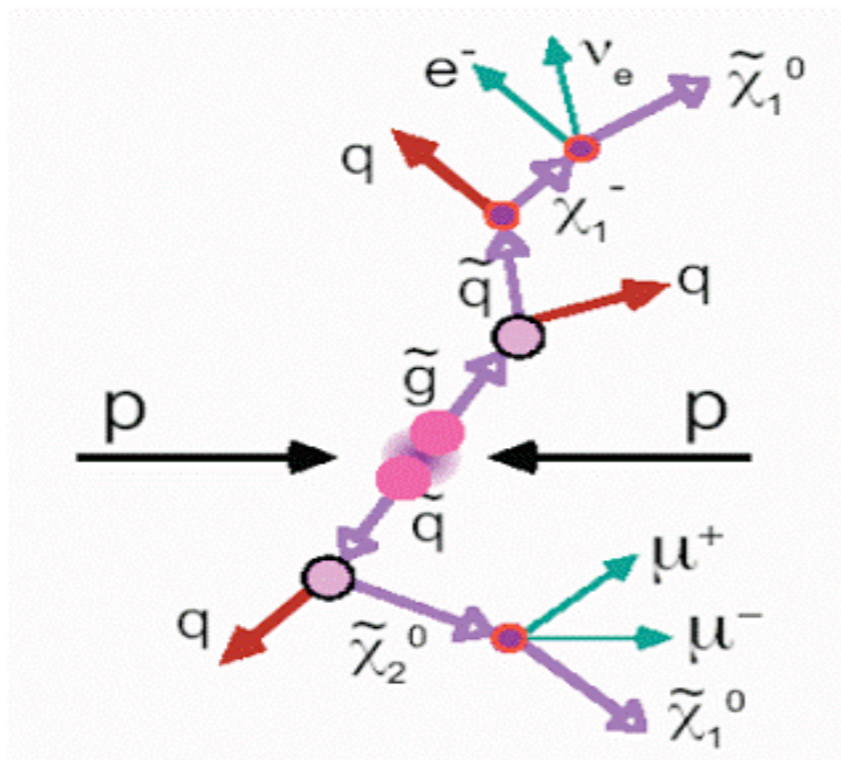
- many particles in the final state

THE MSSM AT THE LHC



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but also:

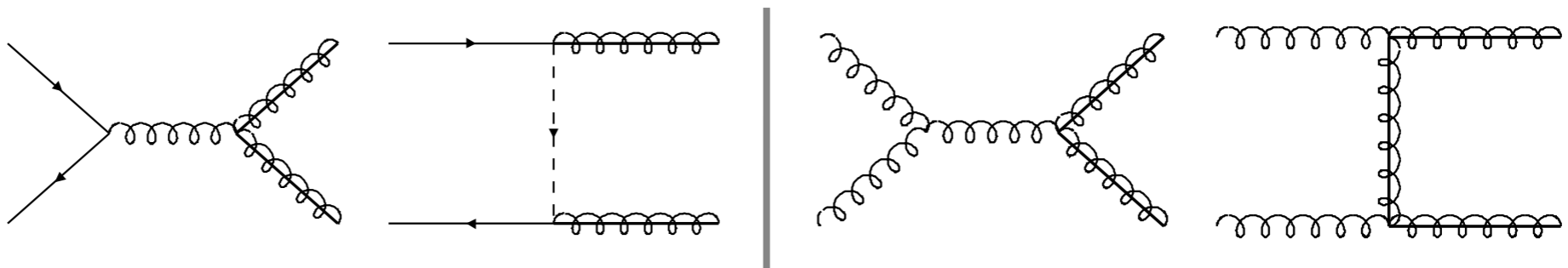
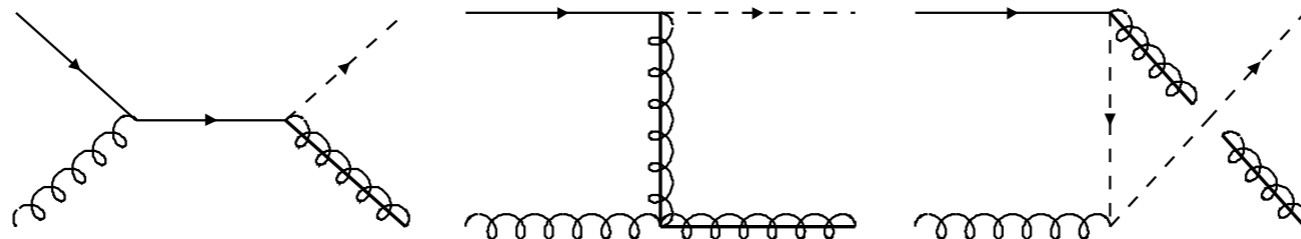
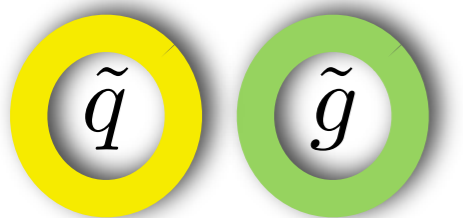
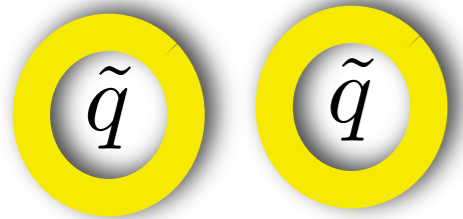
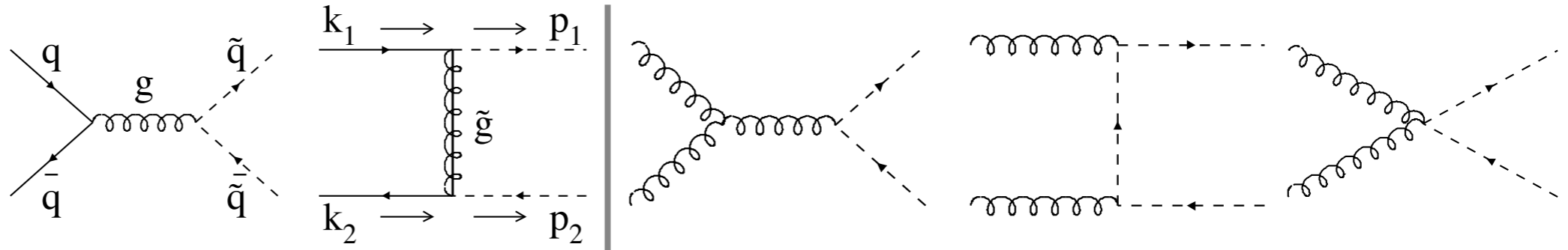
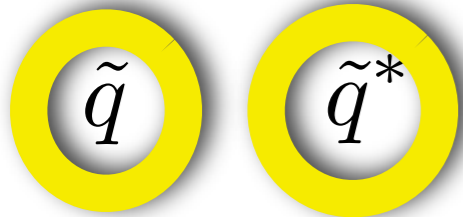
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in general

- many particles in the final state

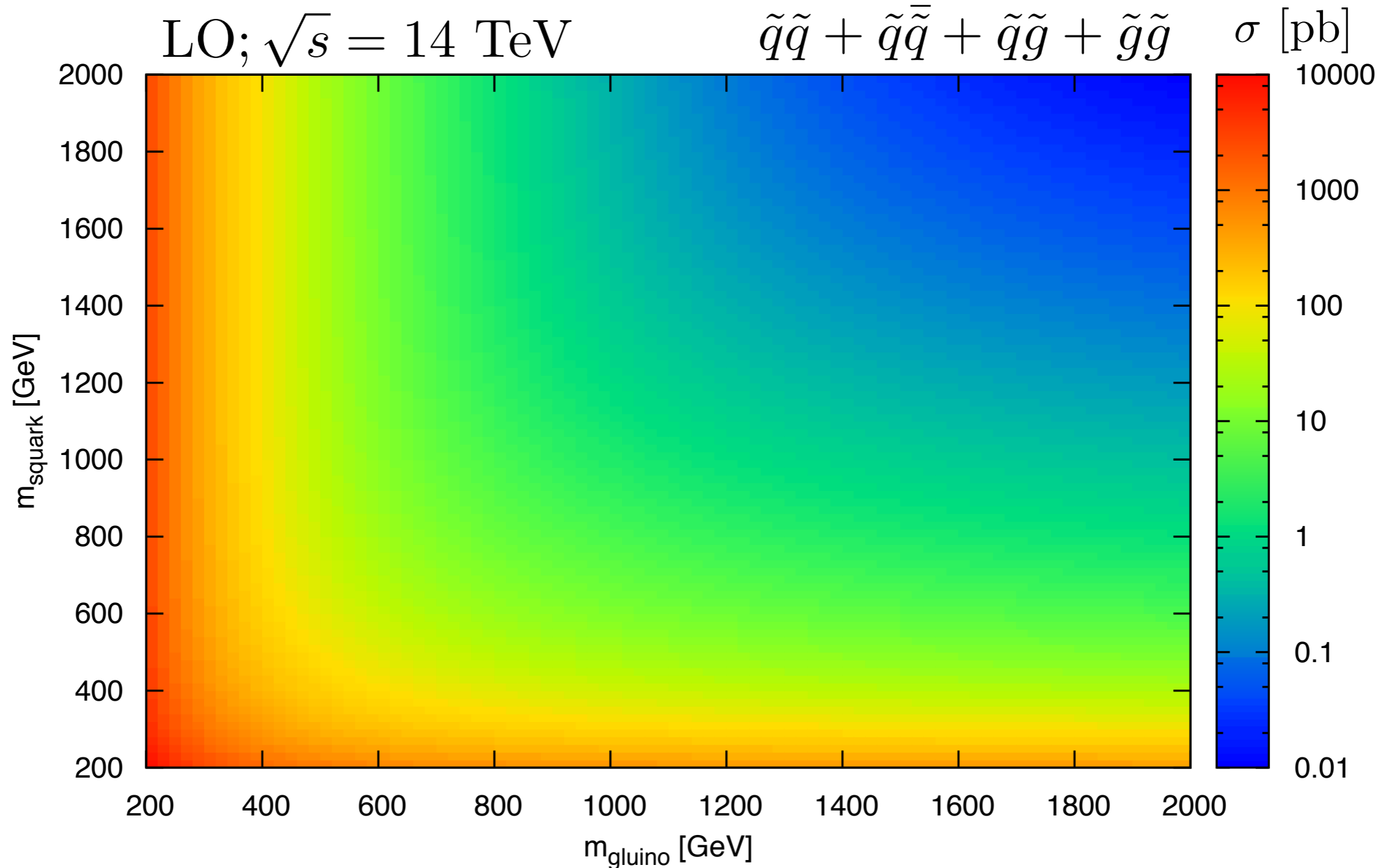
SQUARKS AND GLUINOS AT THE LHC

PRODUCTION



SQUARKS AND GLUINOS AT THE LHC

PRODUCTION

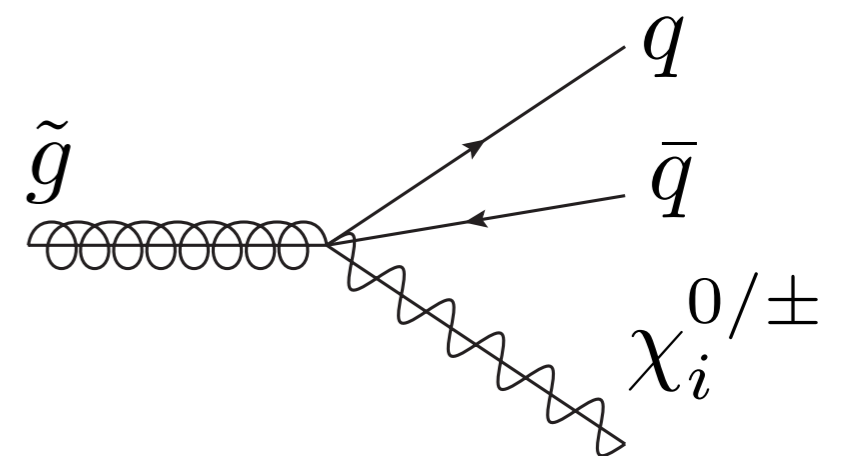
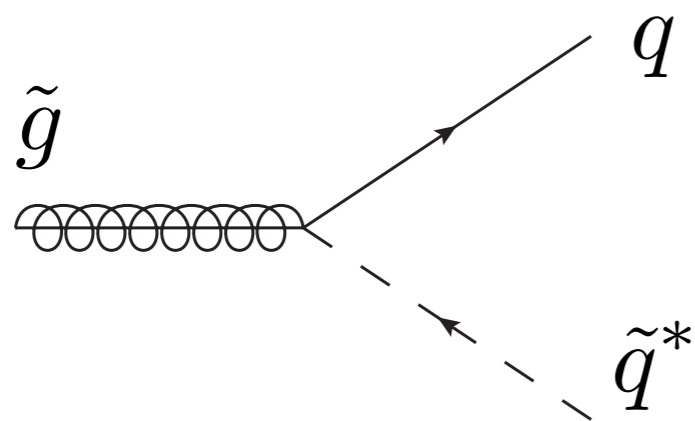
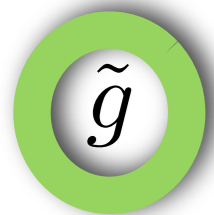
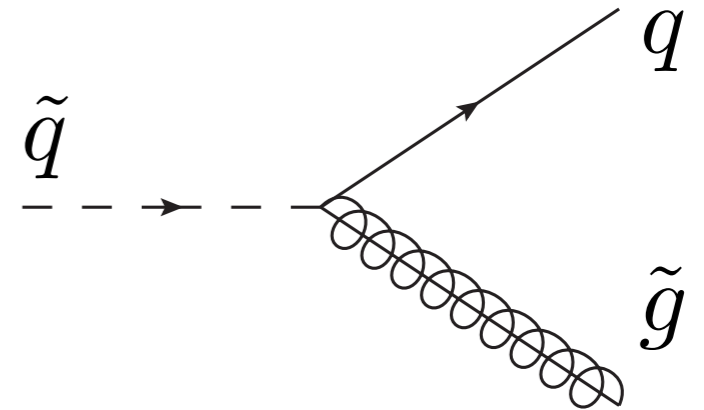
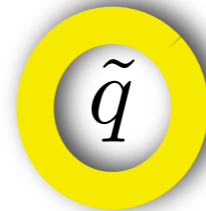
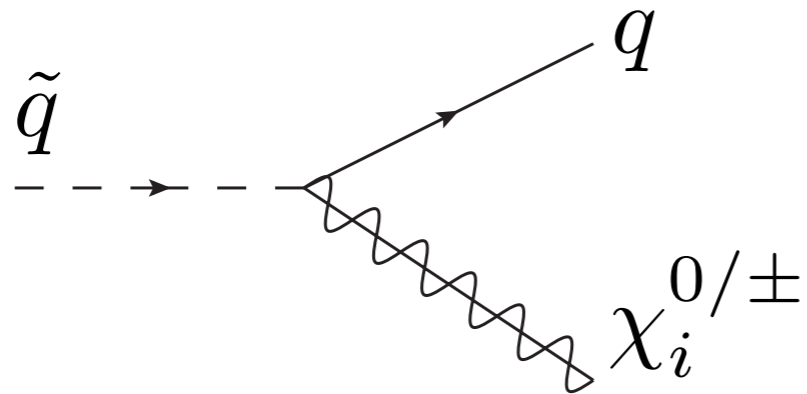
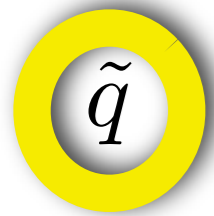


SQUARKS AND GLUINOS AT THE LHC

DECAY

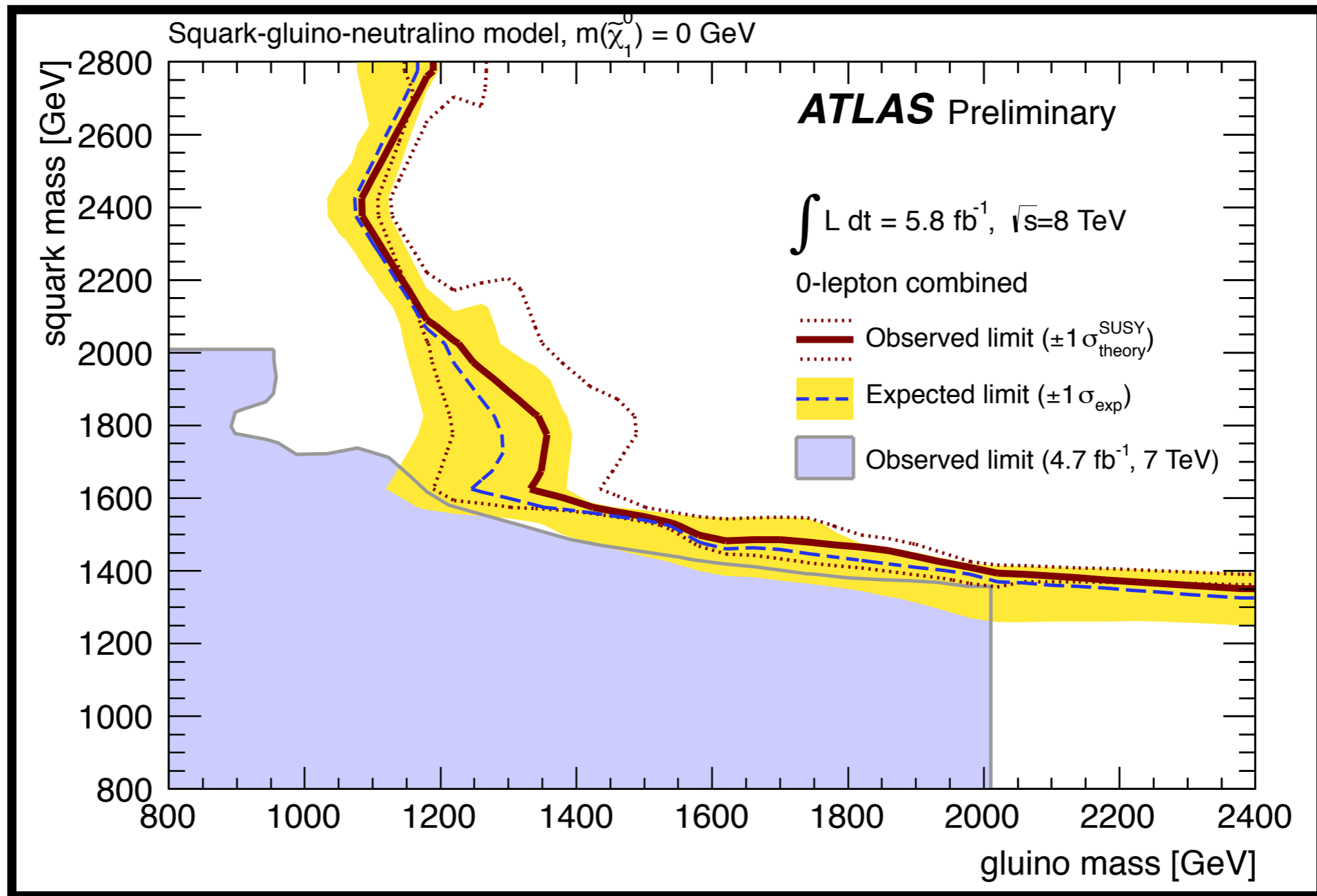
$$m_{\tilde{g}} > m_{\tilde{q}}$$

$$m_{\tilde{q}} > m_{\tilde{g}}$$



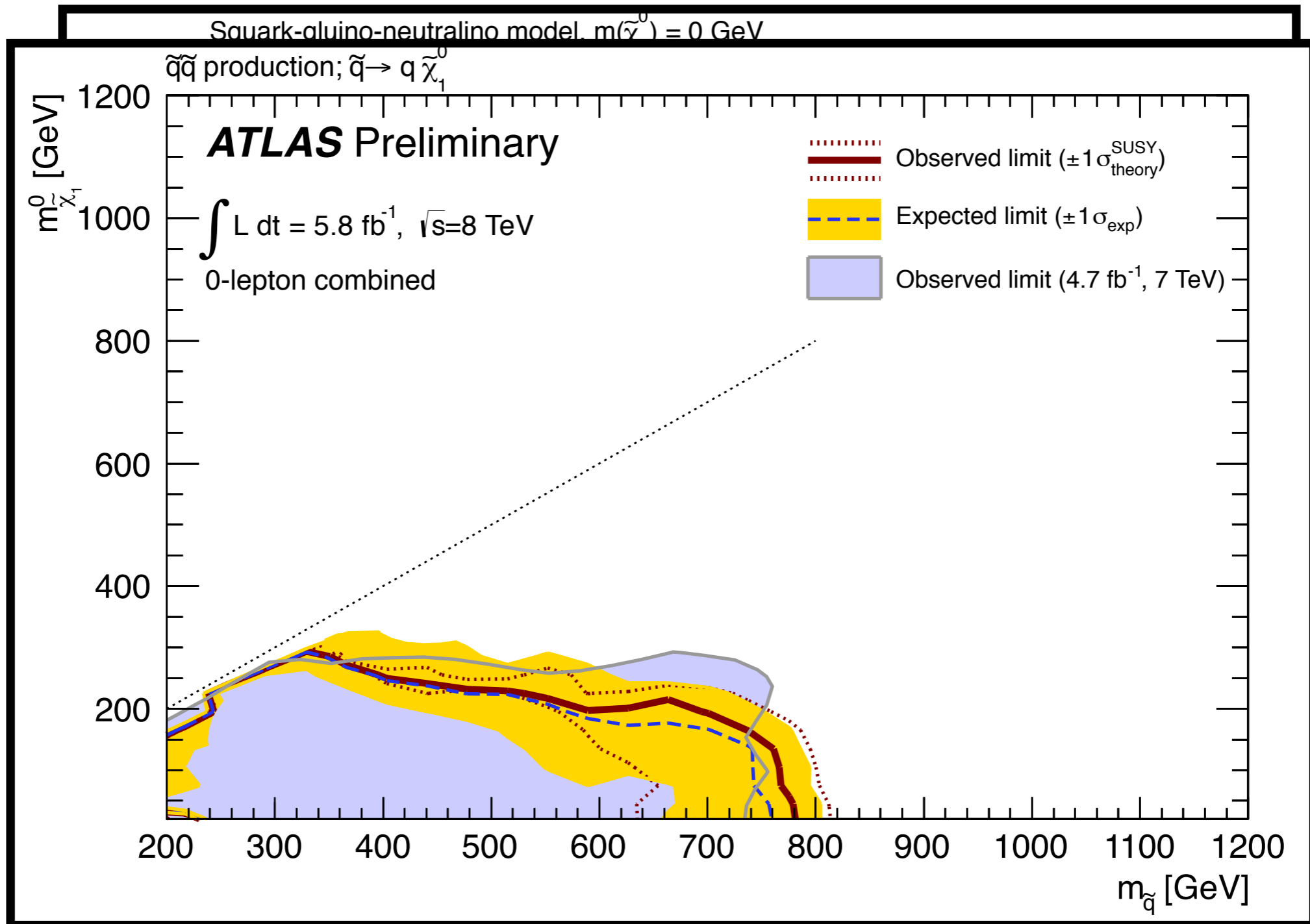
SQUARKS AND GLUINOS AT THE LHC

EXCLUSION LIMITS

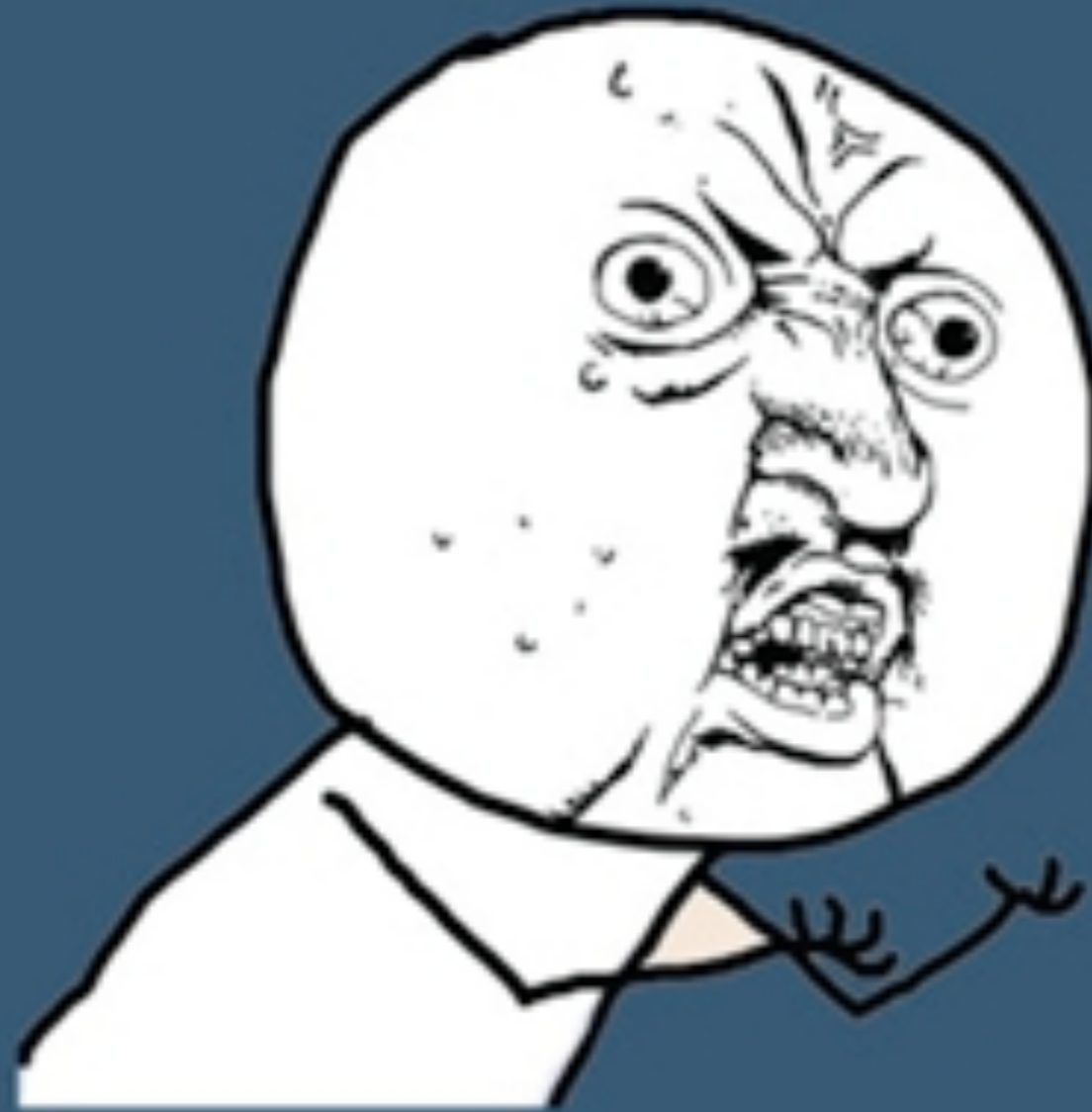


SQUARKS AND GLUINOS AT THE LHC

EXCLUSION LIMITS



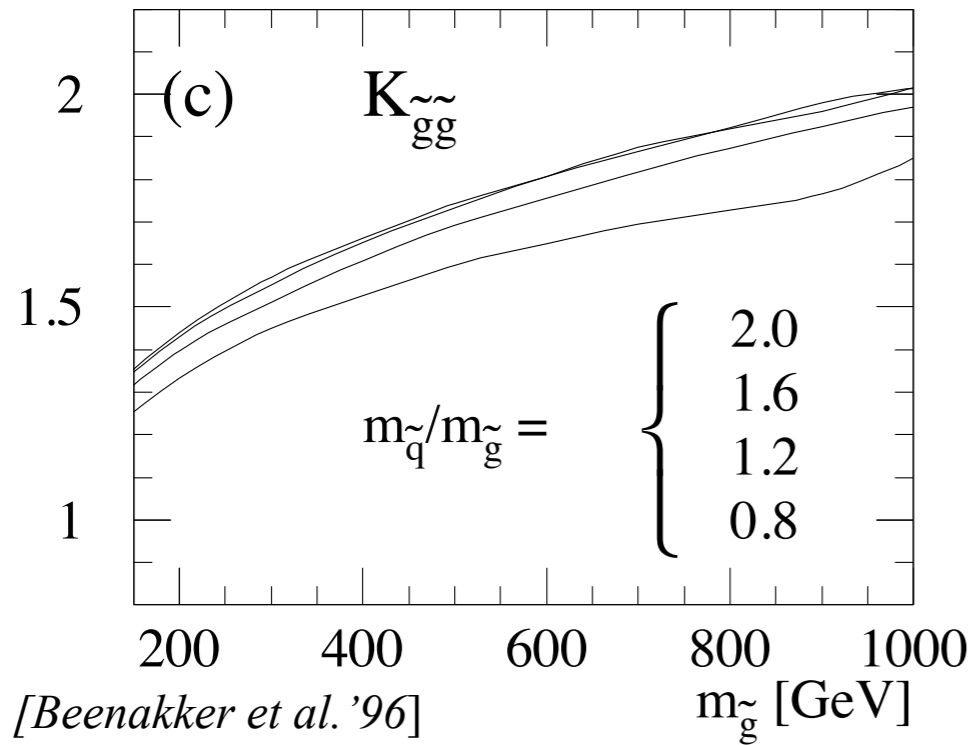
LHC



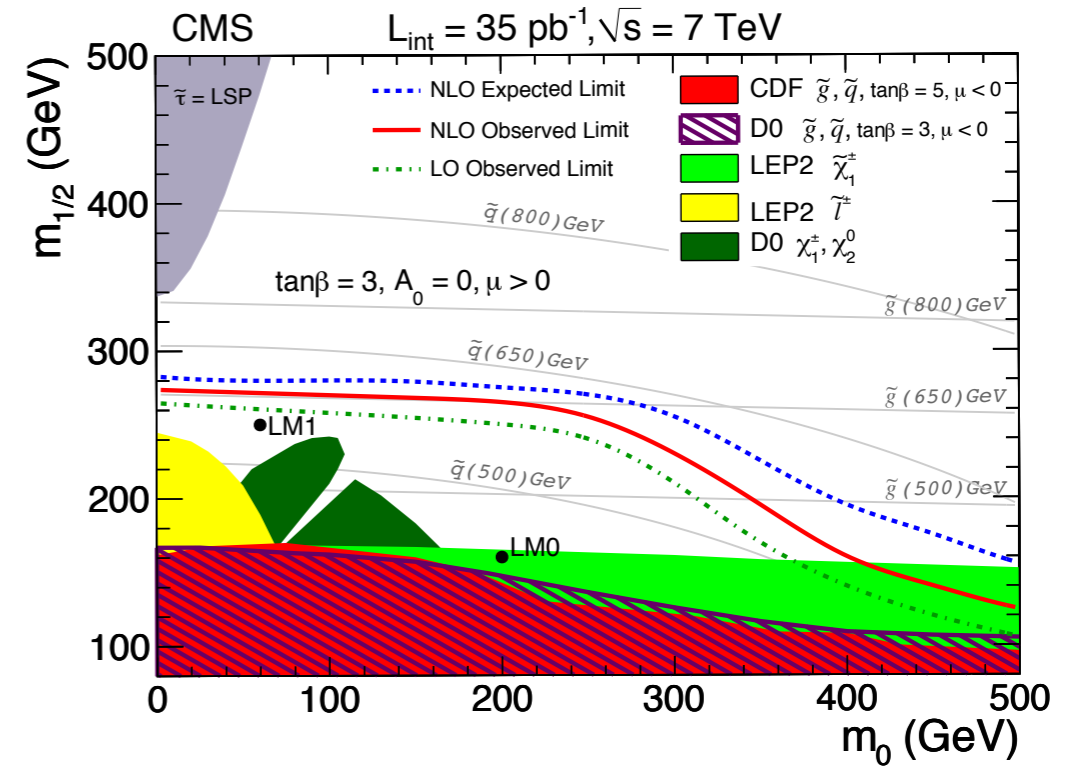
Y U NO SEE SUSY?

Why Higher Orders?

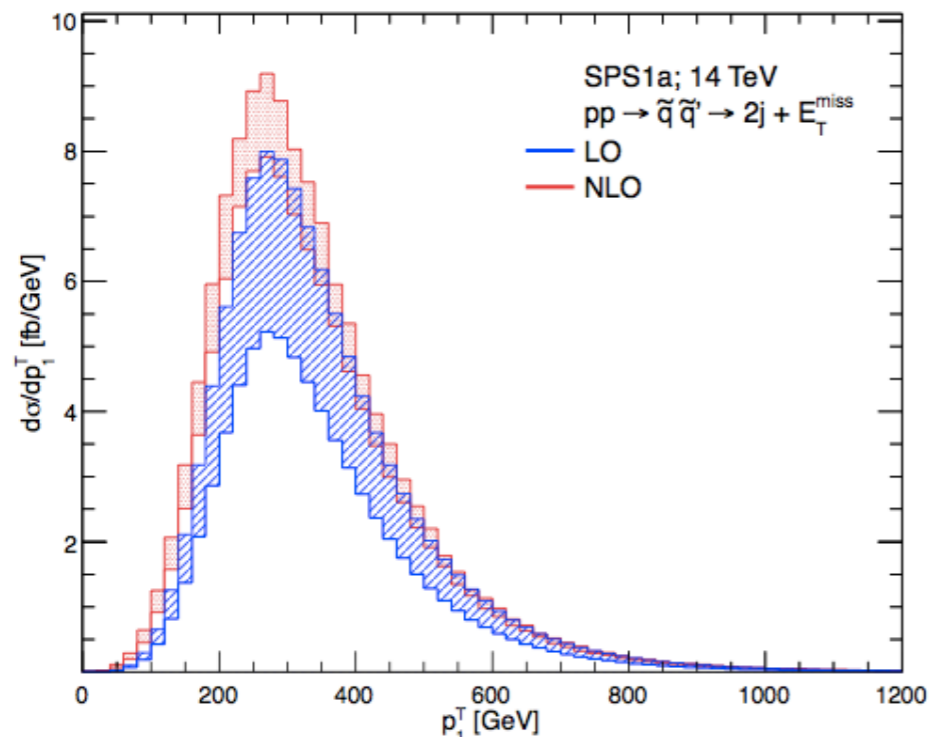
Corrections can be large!



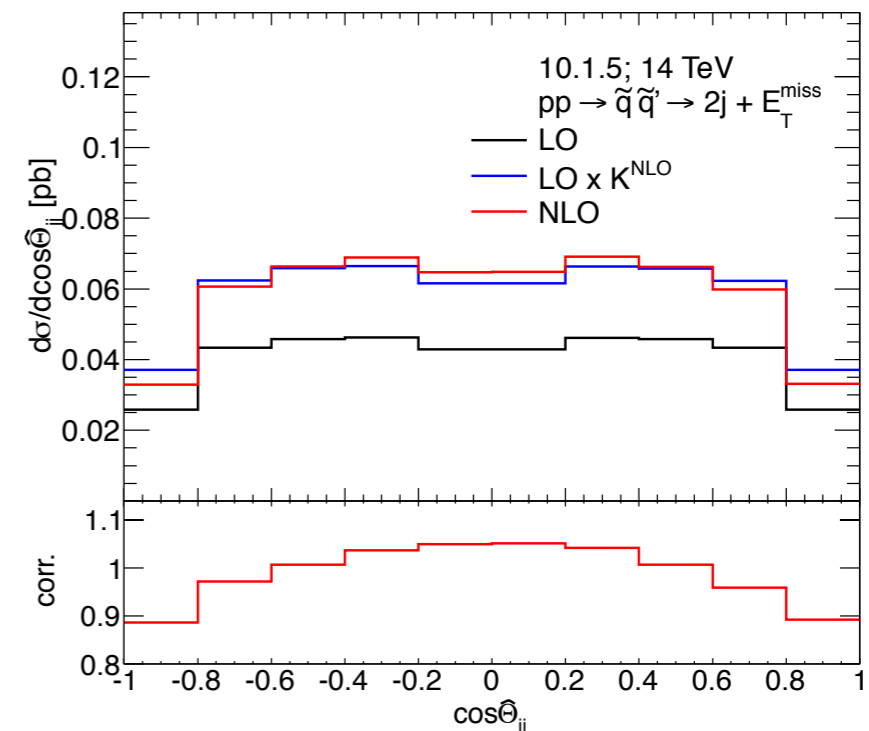
Accurate exclusion limits.



Study & reduce theoretical uncertainties.



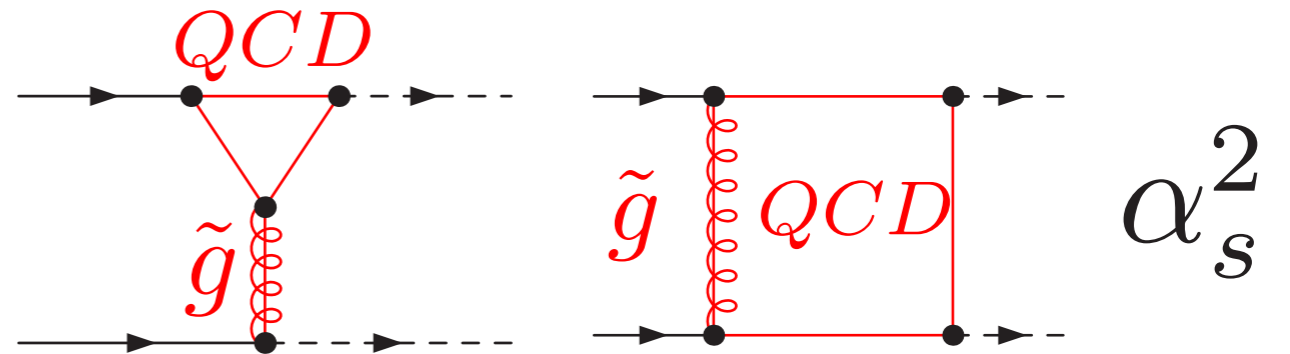
Necessary for parameter determination.



Higher Order Corrections to Production

NLO QCD/SUSY-QCD

Beenakker et al. '96, Goncalves-Netto et al. '12



α_s^2

LO EW

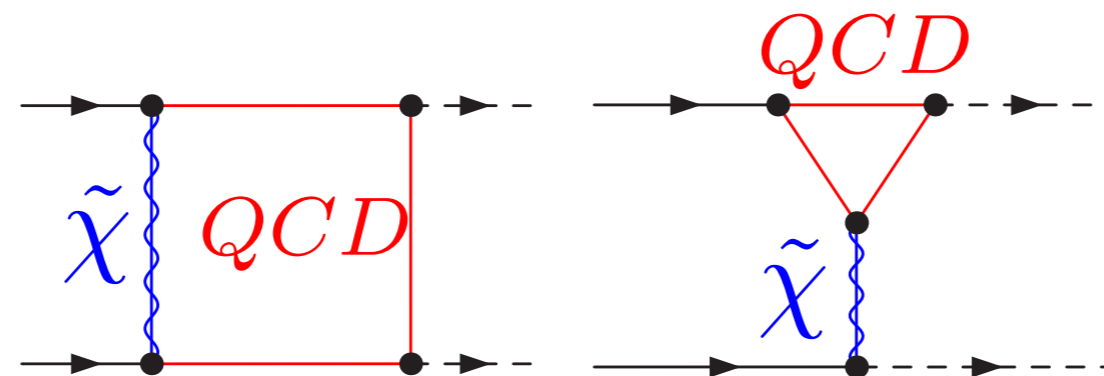
Bornhauser et al. '07



α

NLO EW

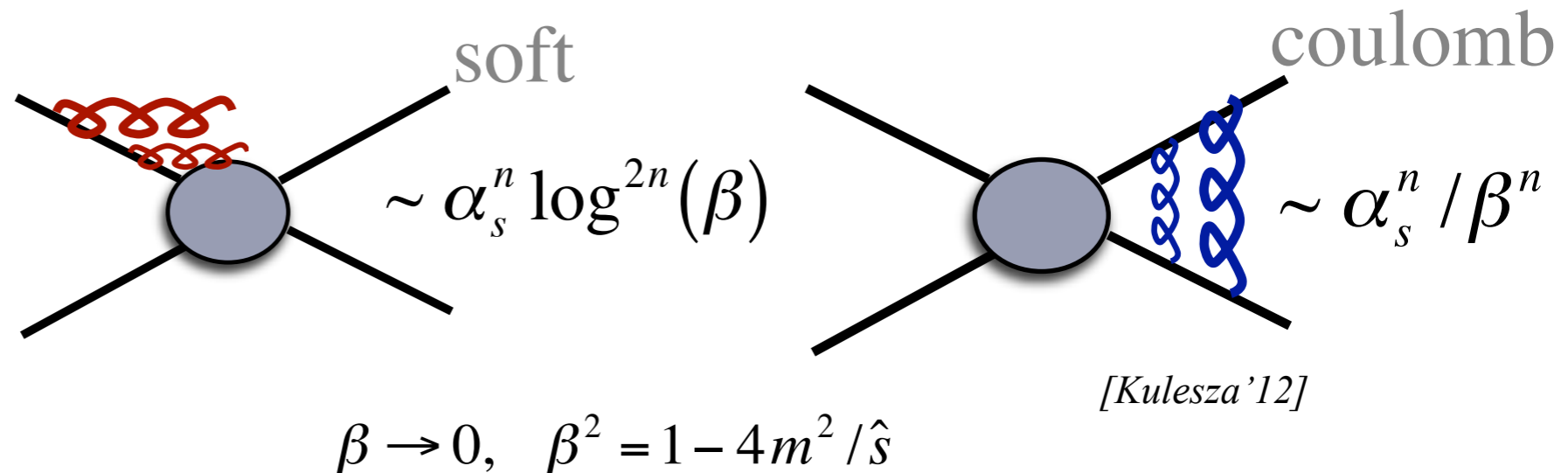
Germer, Hollik, Mirabella, Trenkel '10, ...



$\alpha_s \alpha$

Beyond NLO (resummed)

Beenakker et al. '09, Falgari, Schwinn, Wever '12, ...

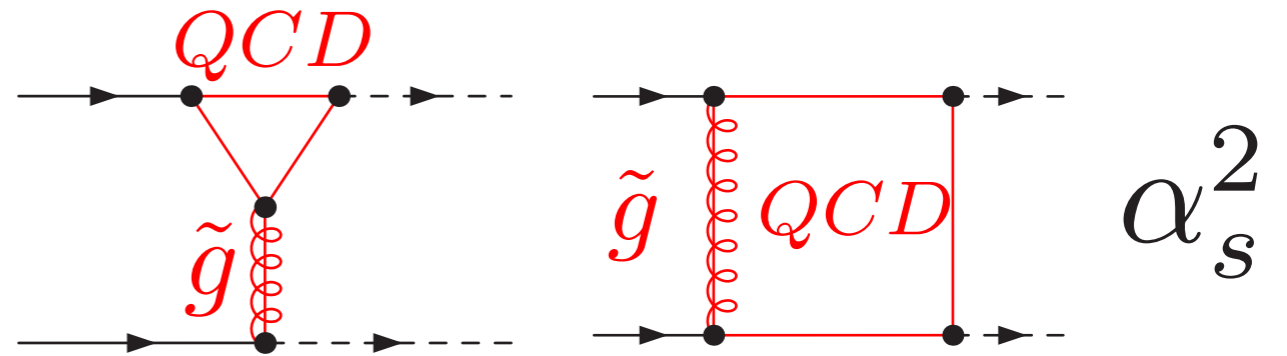


[Kulesza '12]

Higher Order Corrections to Production

NLO QCD/SUSY-QCD

Beenakker et al. '96, Goncalves-Netto et al. '12



α_s^2

LO EW

Born, 1907



α

mostly corrections to inclusive cross sections
or unphysical states!

N²LO

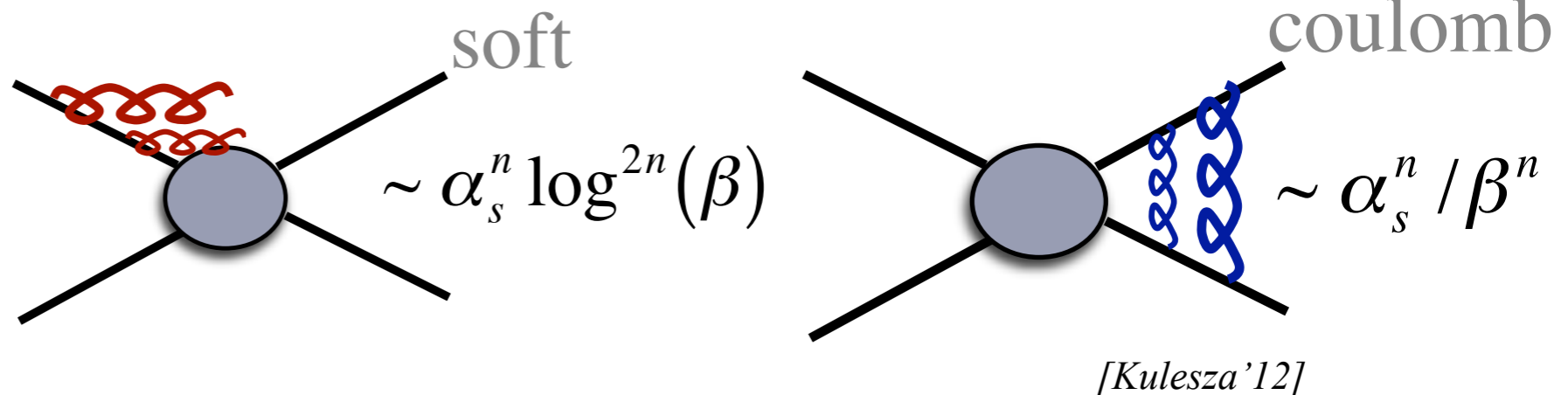
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$\alpha_s \alpha$

Beyond NLO (resummed)

Beenakker et al. '09, Falgari, Schwinn, Wever '12, ...

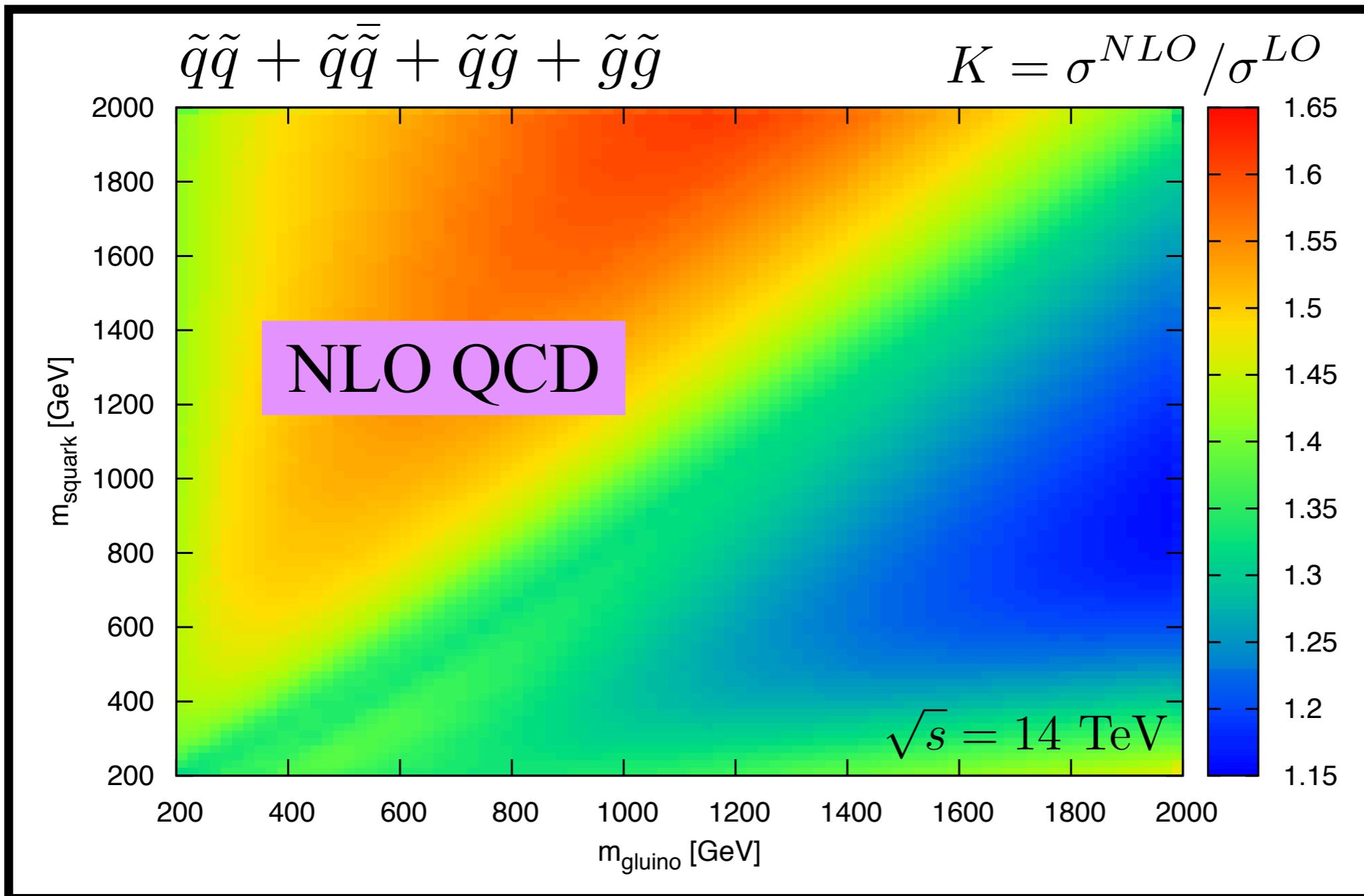


$\beta \rightarrow 0, \quad \beta^2 = 1 - 4m^2 / \hat{s}$

[Kulesza '12]

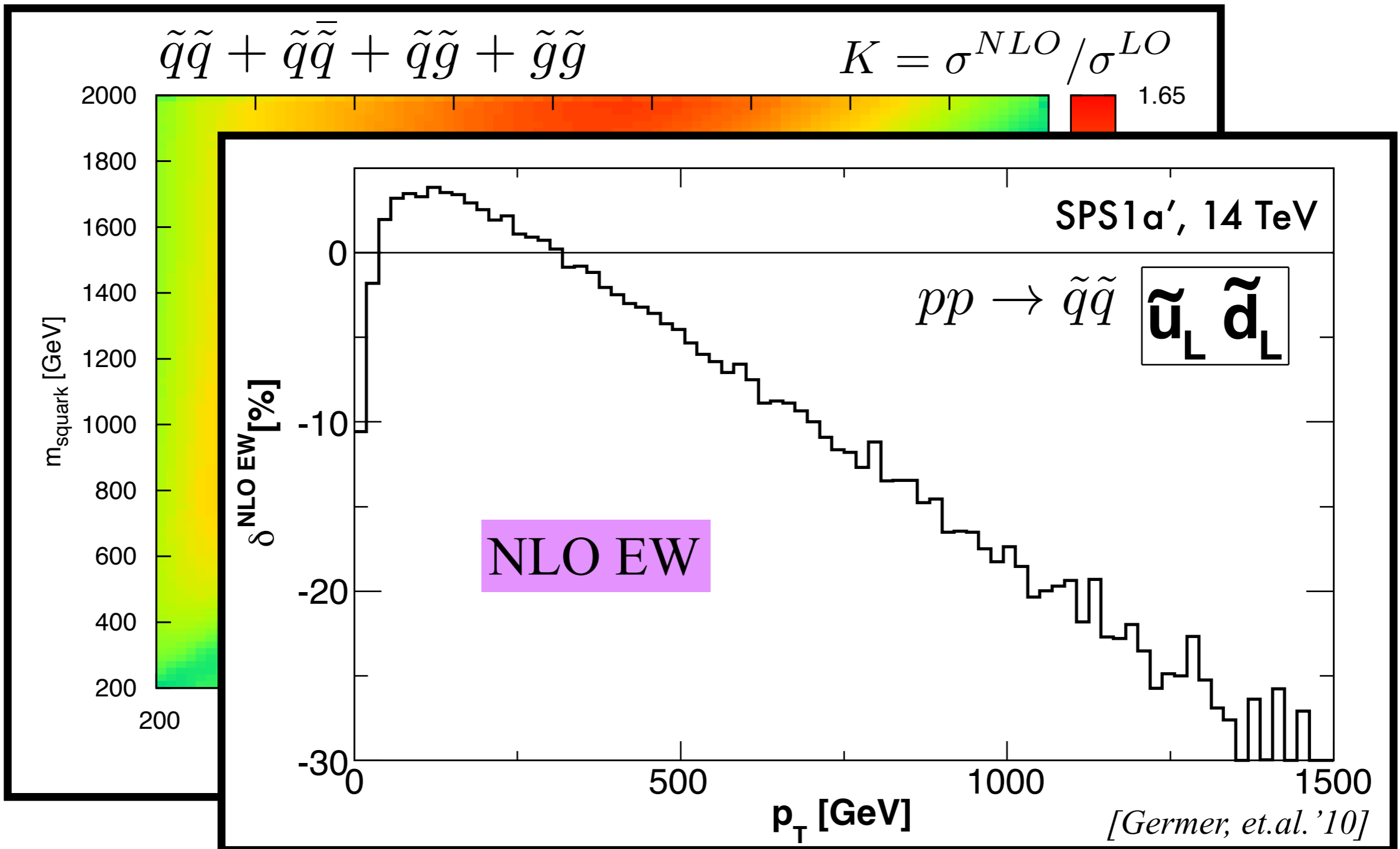
Higher Order Corrections to Production

some examples



Higher Order Corrections to Production

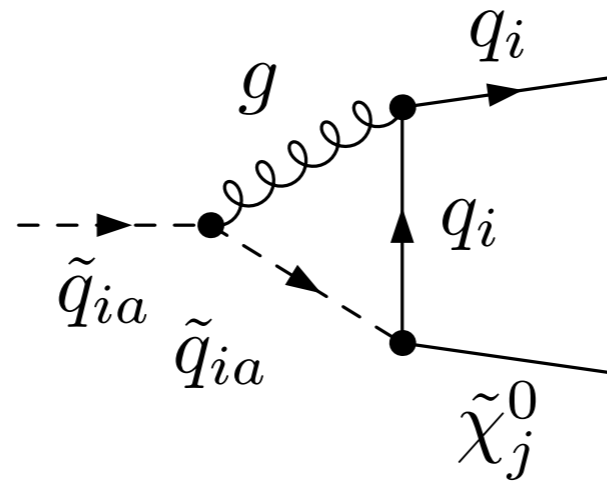
some examples



Higher Order Corrections to Decay

NLO QCD/ SUSY-QCD

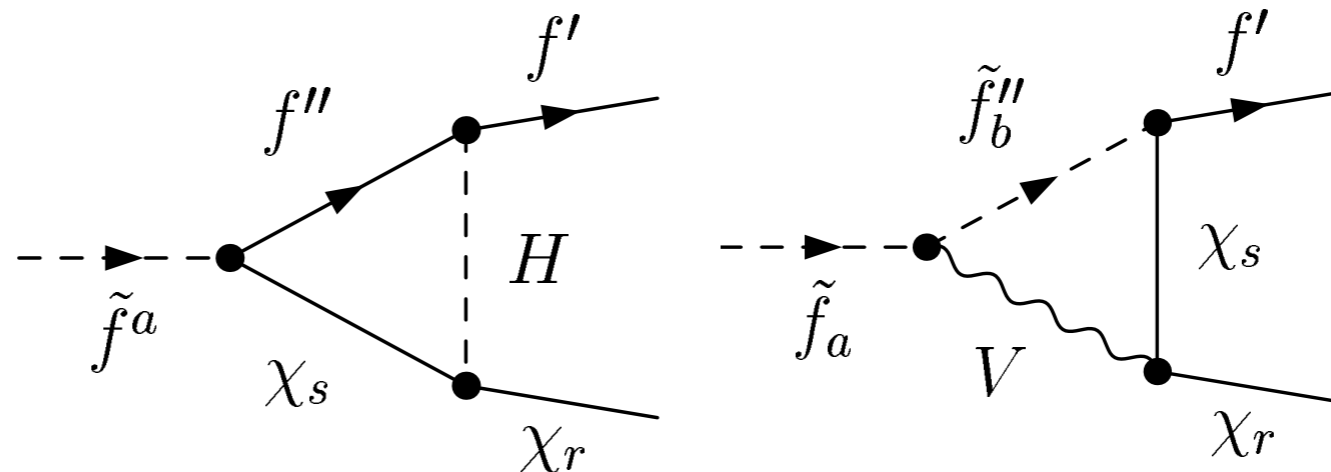
Djouadi, Hollik, Junger '96



α_S

NLO EW

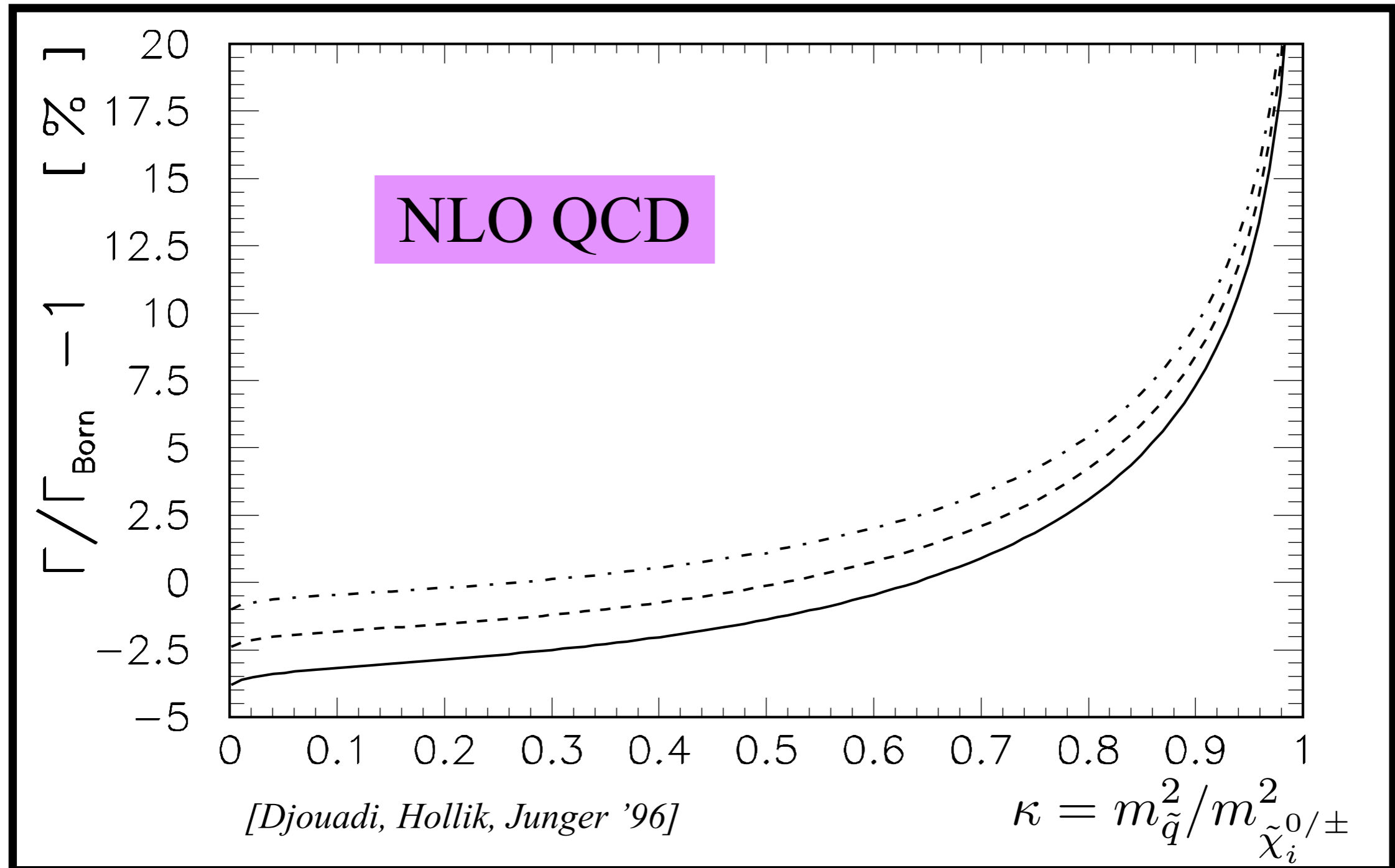
Guasch, Hollik, Sola '02

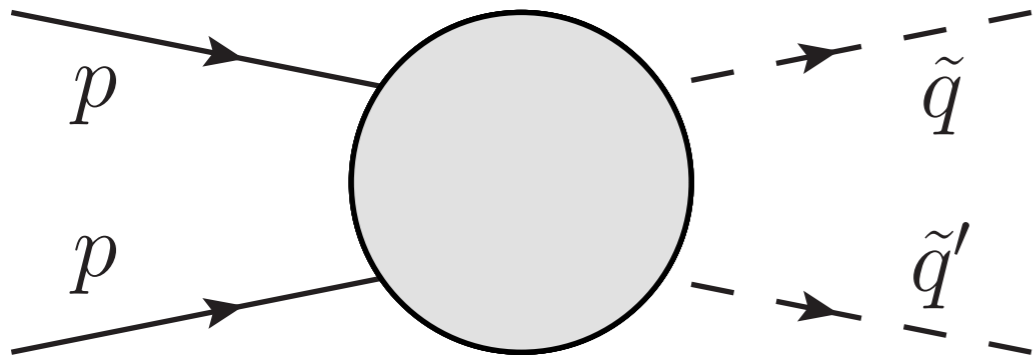


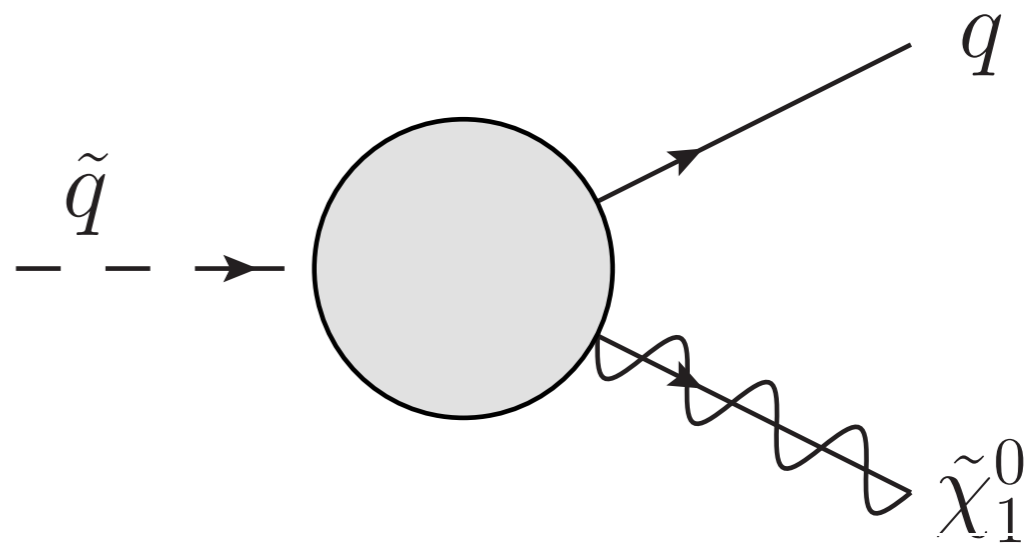
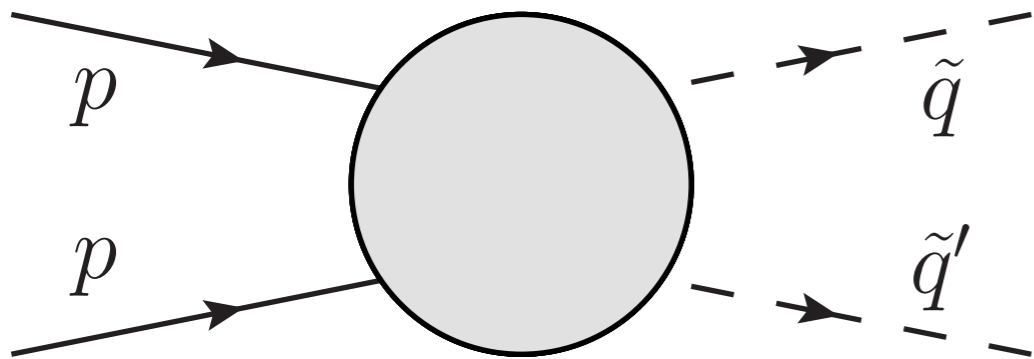
α

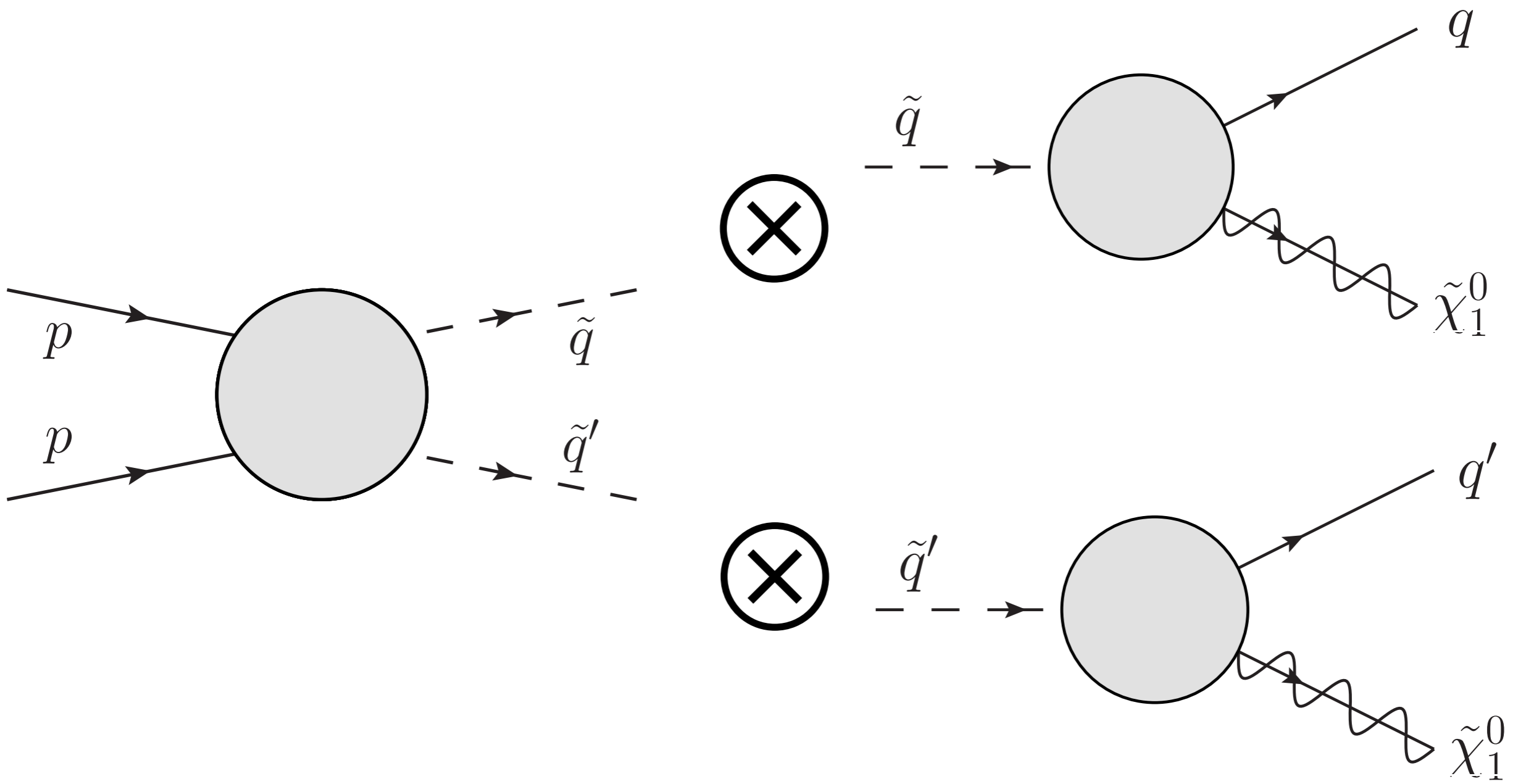
Higher Order Corrections to Decay

one example

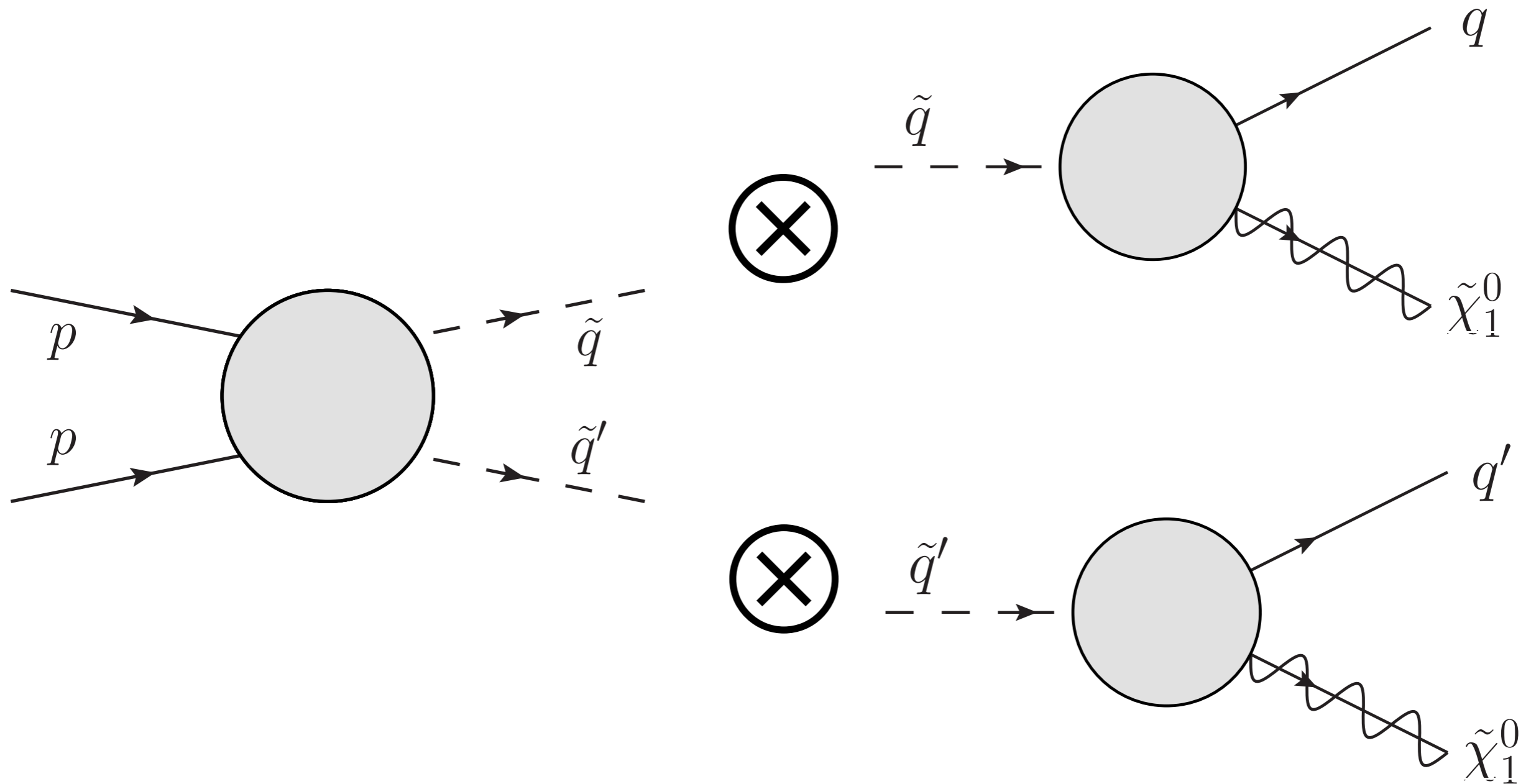








Combining Production and Decay at NLO



We study the experimental signature

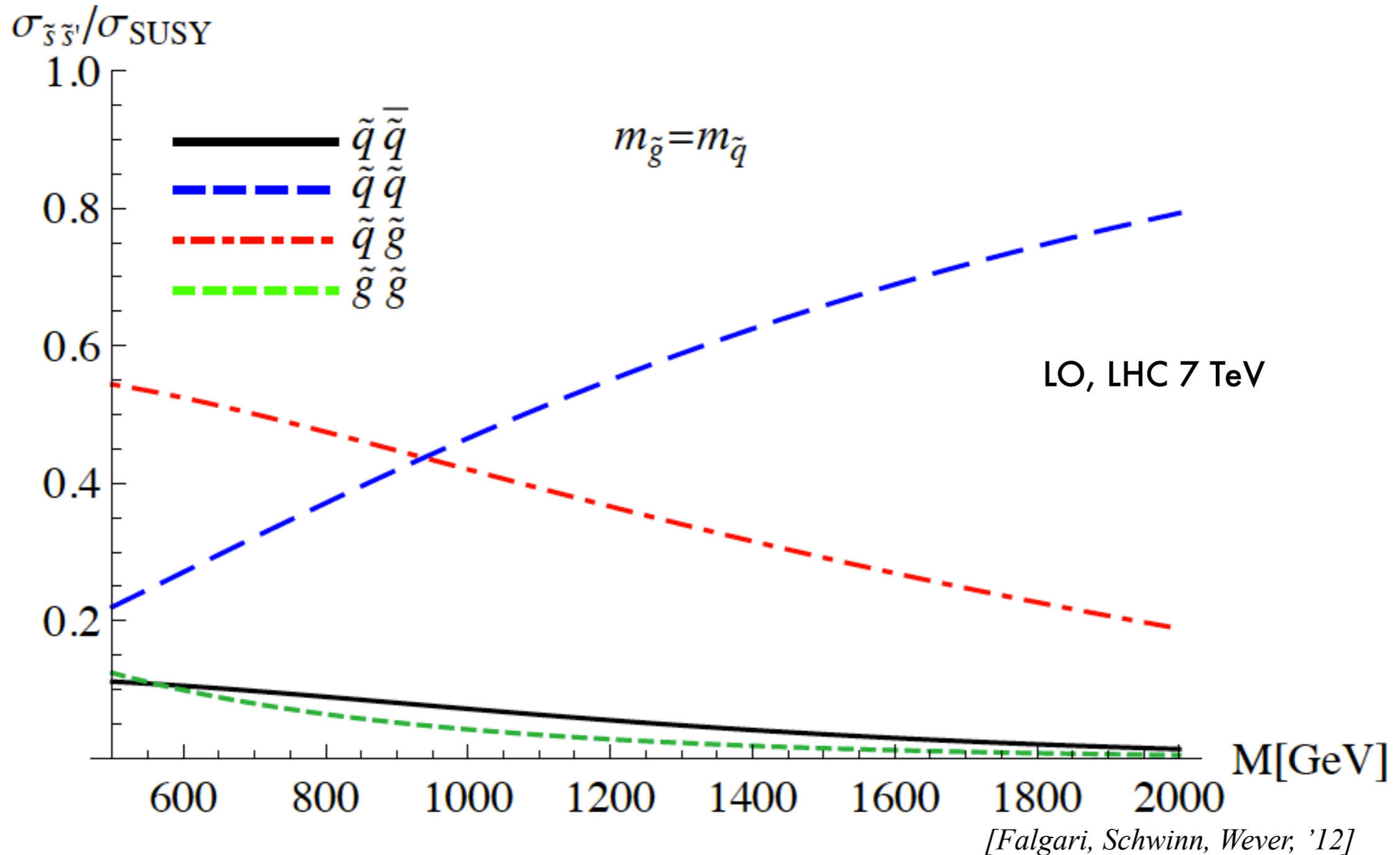
via squark-squark production and direct decay into the lightest neutralino.

$$2j + \cancel{E}_T (+X)$$

$$pp \rightarrow \tilde{q}\tilde{q}' \rightarrow qq' \tilde{\chi}_1^0 \tilde{\chi}_1^0 (+X)$$

Combining Production and Decay at NLO

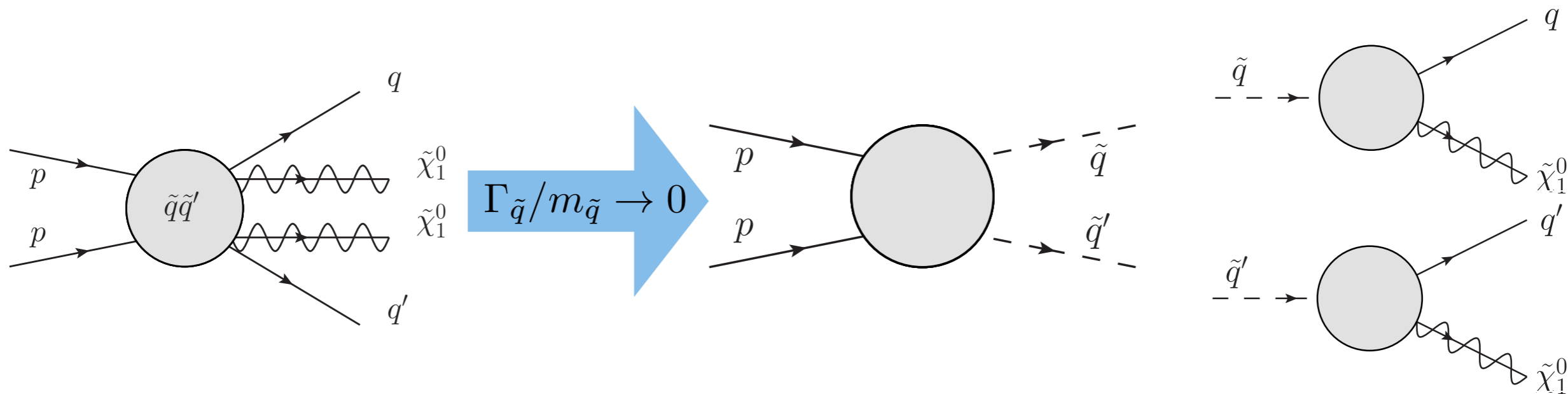
Why squark-squark channel?



LO in NWA

$$qq' \rightarrow \tilde{q}\tilde{q}' \rightarrow q\tilde{\chi}_1^0 q'\tilde{\chi}_1^0$$

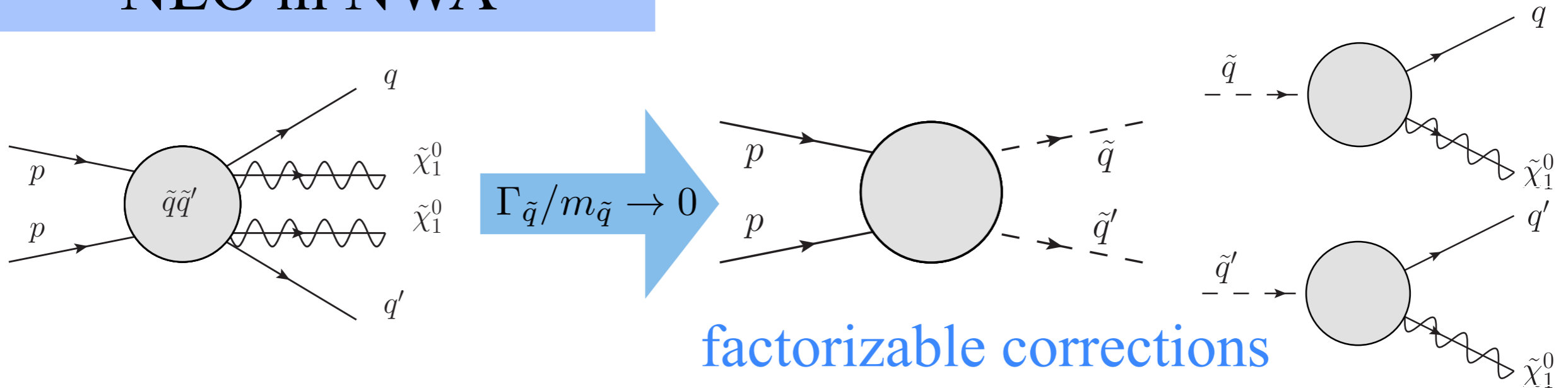
$$\Gamma_{\tilde{q}}/m_{\tilde{q}} \rightarrow 0 \quad \hat{\sigma}_{\text{NWA}}^{(0)} = \hat{\sigma}^{(0)}(qq' \rightarrow \tilde{q}\tilde{q}') \times BR^{(0)}(\tilde{q} \rightarrow q\tilde{\chi}_1^0) \times BR^{(0)}(\tilde{q}' \rightarrow q'\tilde{\chi}_1^0)$$



Hadronic differential LO cross section in NWA

$$d\sigma_{\text{NWA}}^{(0)}(pp \rightarrow \tilde{q}\tilde{q}' \rightarrow q\tilde{\chi}_1^0 q'\tilde{\chi}_1^0 (+X)) = \frac{1}{\Gamma_{\tilde{q}}^{(0)} \Gamma_{\tilde{q}'}^{(0)}} \left[d\sigma_{pp \rightarrow \tilde{q}\tilde{q}'}^{(0)} d\Gamma_{\tilde{q} \rightarrow q\tilde{\chi}_1^0}^{(0)} d\Gamma_{\tilde{q}' \rightarrow q'\tilde{\chi}_1^0}^{(0)} \right]$$

NLO in NWA



Formal expansion in α_s :

Born

$$\begin{aligned}
 d\sigma_{\text{NWA}}^{(0+1)}(pp \rightarrow \tilde{q}\tilde{q}' \rightarrow q\tilde{\chi}_1^0 q'\tilde{\chi}_1^0 (+X)) &= \frac{1}{\Gamma_{\tilde{q}}^{(0)} \Gamma_{\tilde{q}'}^{(0)}} \left[d\sigma_{pp \rightarrow \tilde{q}\tilde{q}'}^{(0)} d\Gamma_{\tilde{q} \rightarrow q\tilde{\chi}_1^0}^{(0)} d\Gamma_{\tilde{q}' \rightarrow q'\tilde{\chi}_1^0}^{(0)} \left(1 - \frac{\Gamma_{\tilde{q}}^{(1)}}{\Gamma_{\tilde{q}}^{(0)}} - \frac{\Gamma_{\tilde{q}'}^{(1)}}{\Gamma_{\tilde{q}'}^{(0)}} \right) \right. \\
 &\quad \left. + d\sigma_{pp \rightarrow \tilde{q}\tilde{q}'}^{(0)} d\Gamma_{\tilde{q} \rightarrow q\tilde{\chi}_1^0}^{(1)} d\Gamma_{\tilde{q}' \rightarrow q'\tilde{\chi}_1^0}^{(0)} + d\sigma_{pp \rightarrow \tilde{q}\tilde{q}'}^{(0)} d\Gamma_{\tilde{q} \rightarrow q\tilde{\chi}_1^0}^{(0)} d\Gamma_{\tilde{q}' \rightarrow q'\tilde{\chi}_1^0}^{(1)} \right. \\
 &\quad \left. + d\sigma_{pp \rightarrow \tilde{q}\tilde{q}'}^{(1)} d\Gamma_{\tilde{q} \rightarrow q\tilde{\chi}_1^0}^{(0)} d\Gamma_{\tilde{q}' \rightarrow q'\tilde{\chi}_1^0}^{(0)} \right]
 \end{aligned}$$

NLO decay

NLO production

“master formula”

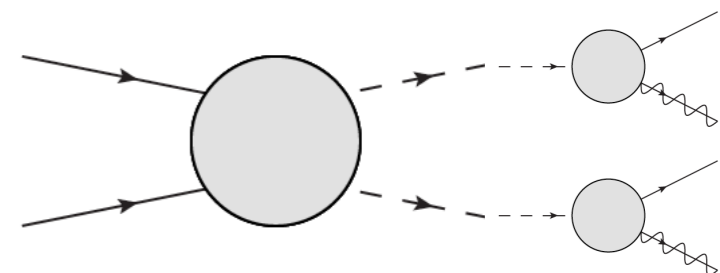
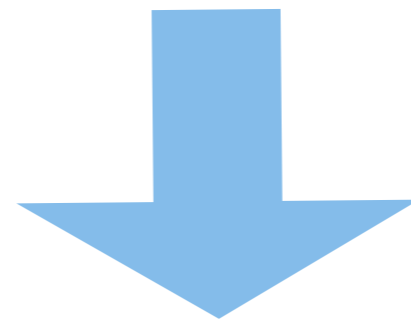
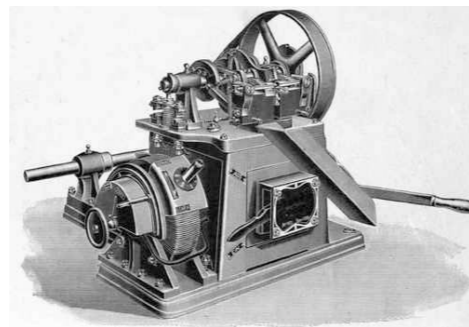
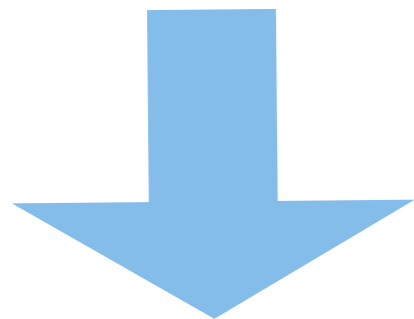
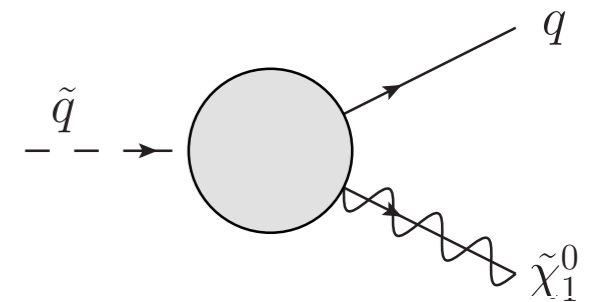
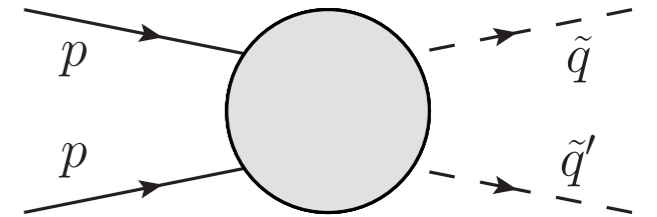
COMBINATION

For all different combinations of light flavours and chiralities, weighted events for squark-squark production are produced in the LAB frame.

Weighted decay events are generated in the respective squark rest-frame.

boost of decay events + “master formula”

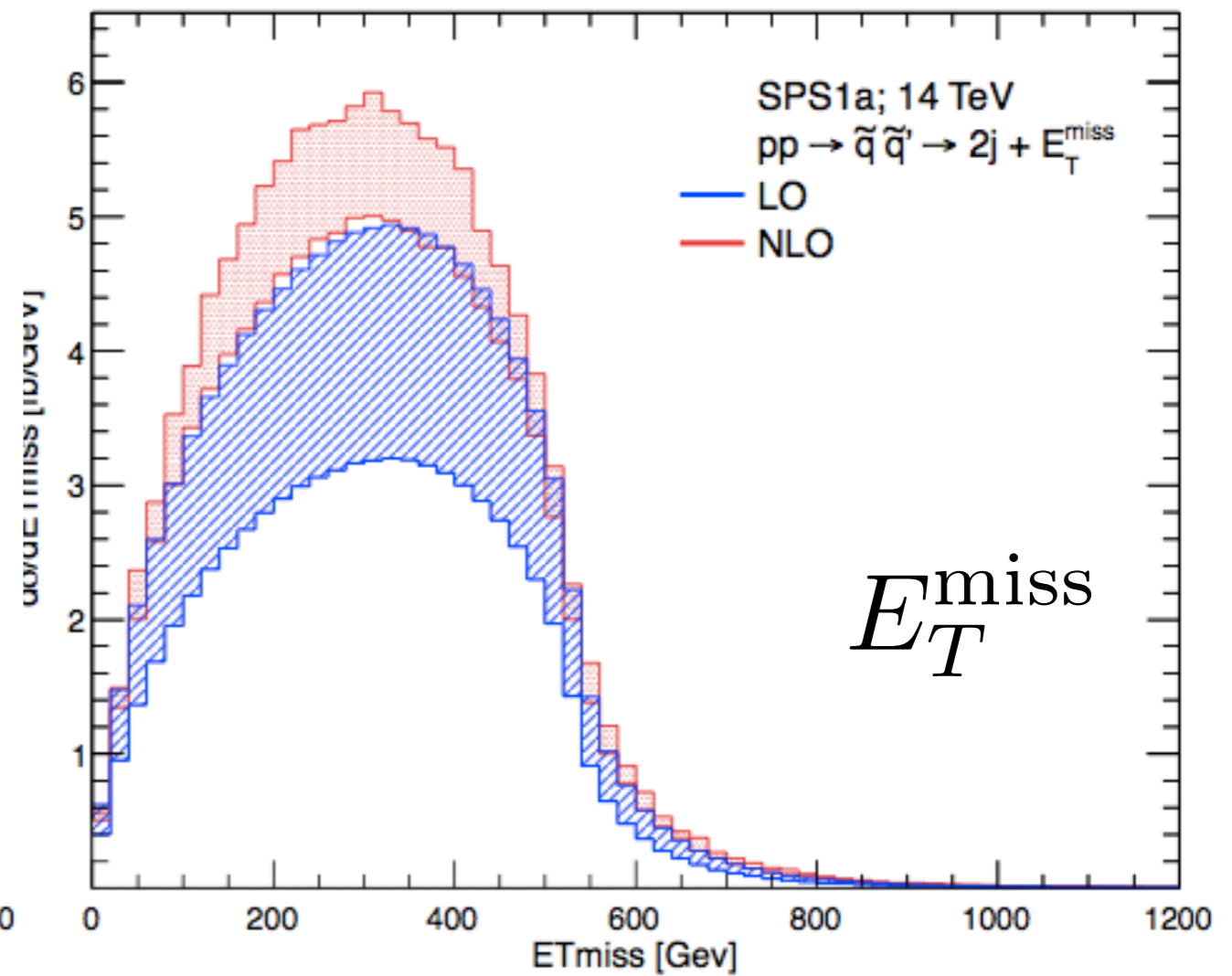
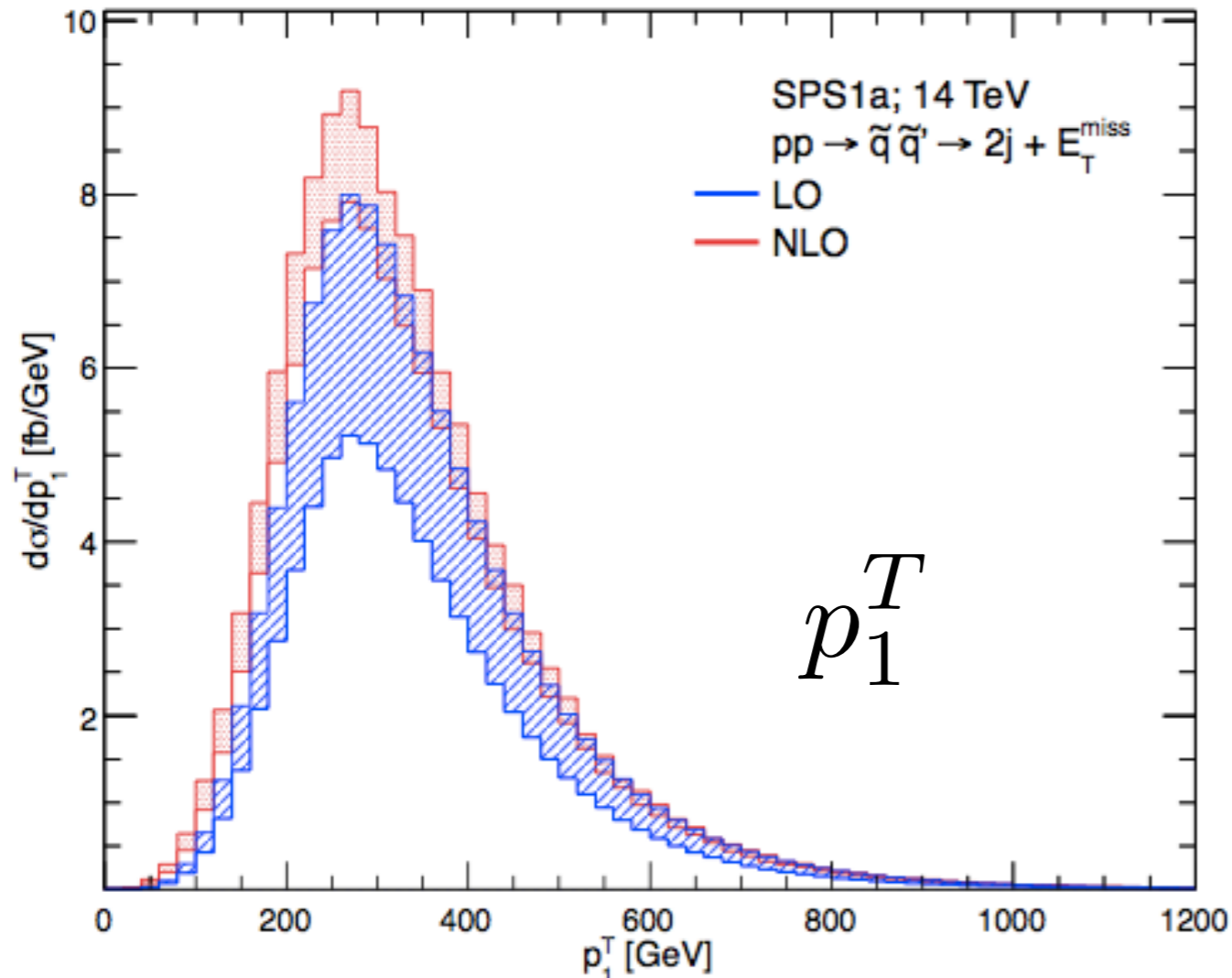
Fully differential prediction of factorizable NLO contributions in NWA.



NUMERICAL RESULTS

For SPS1a (14 TeV): Scale variation: $\mu_f = \mu_r = (m/2, m, 2m)$, m : average \tilde{q} mass

SPS1a	\tilde{u}_L	\tilde{u}_R	\tilde{d}_L	\tilde{d}_R	\tilde{g}	$\tilde{\chi}_1^0$	(PDFs: CTEQ6.6 both for LO and NLO)
mass (GeV)	563.6	546.7	569.0	546.6	608.5	97.0	



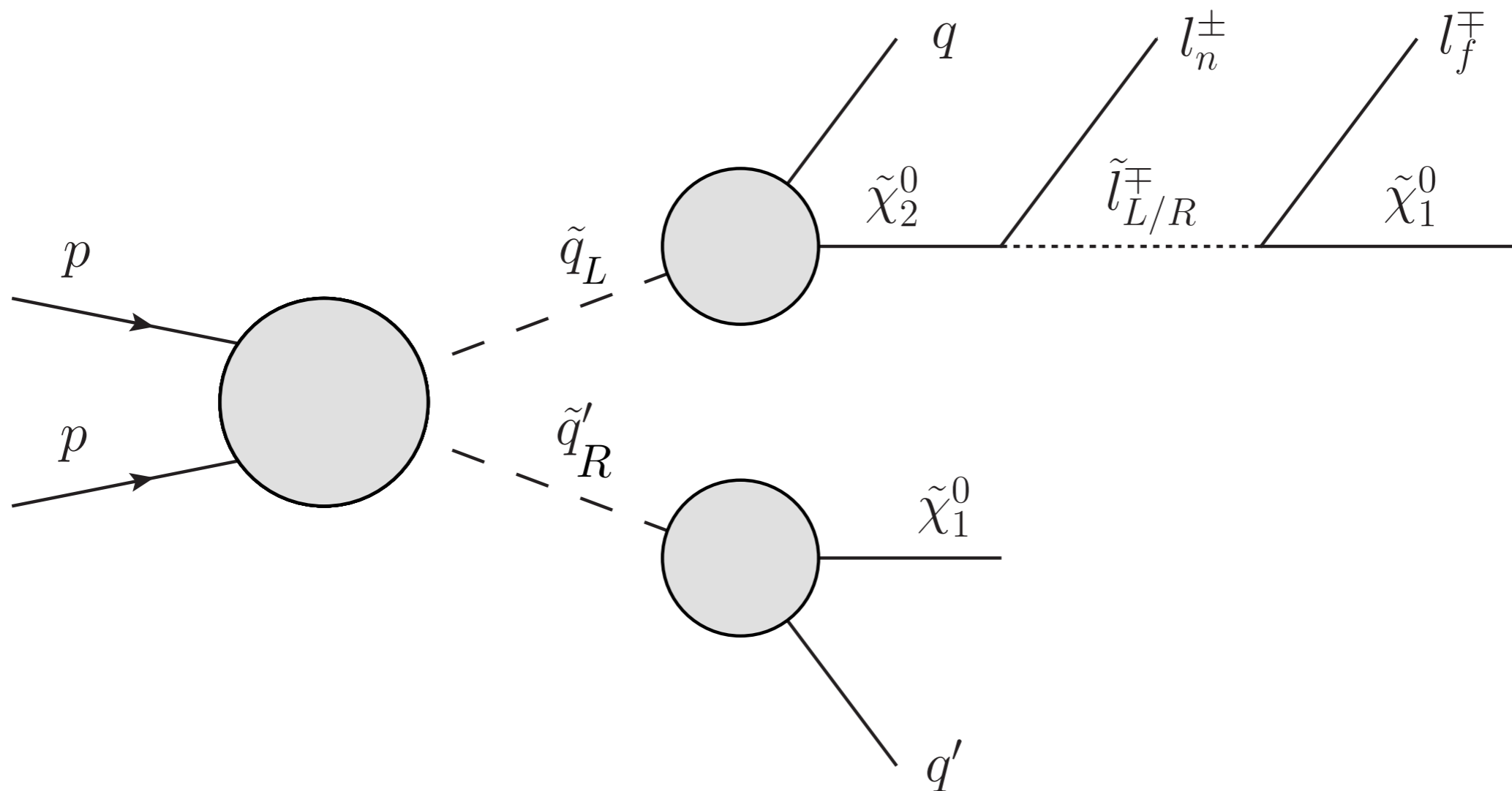
Combining Production and a Decay Chain at NLO

We study the experimental signature

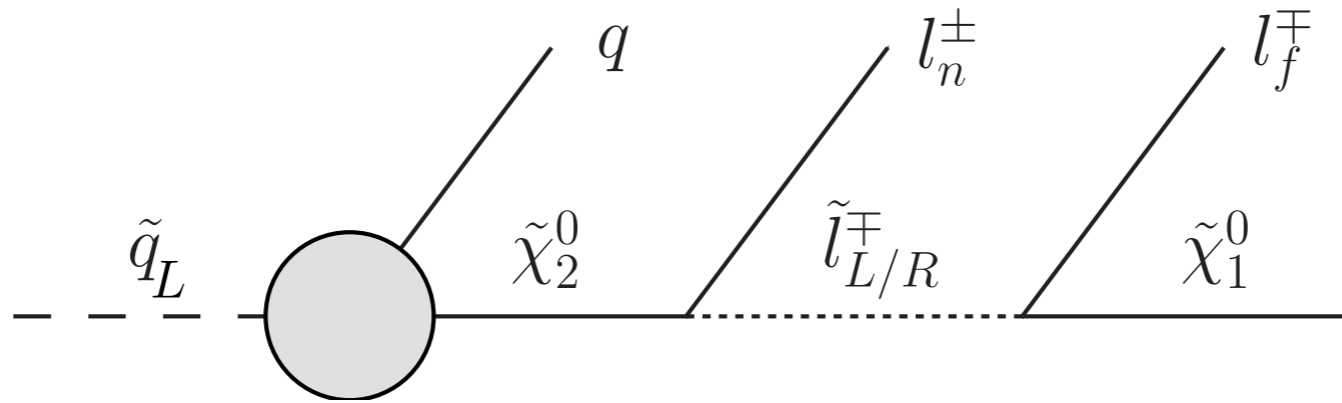
$$2j + 2l + \cancel{E}_T (+X)$$

via squark-squark production and an attached EW decay chain.

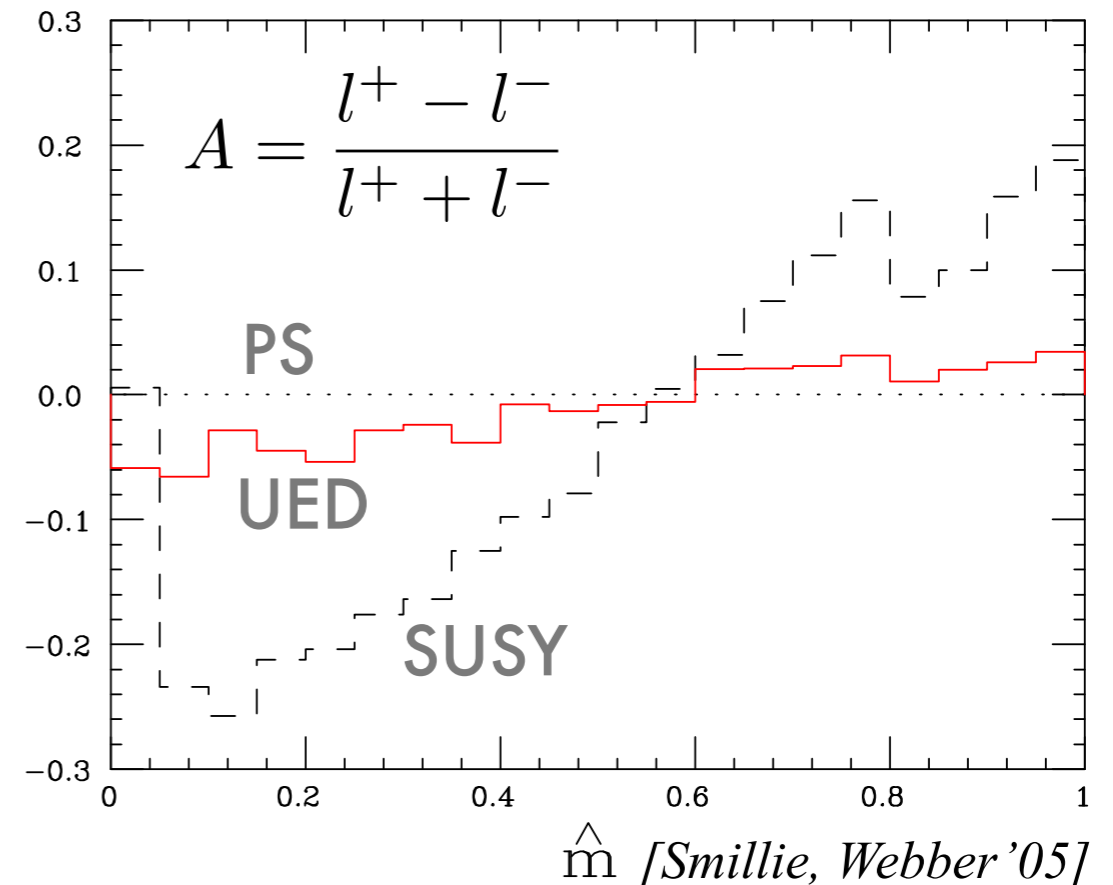
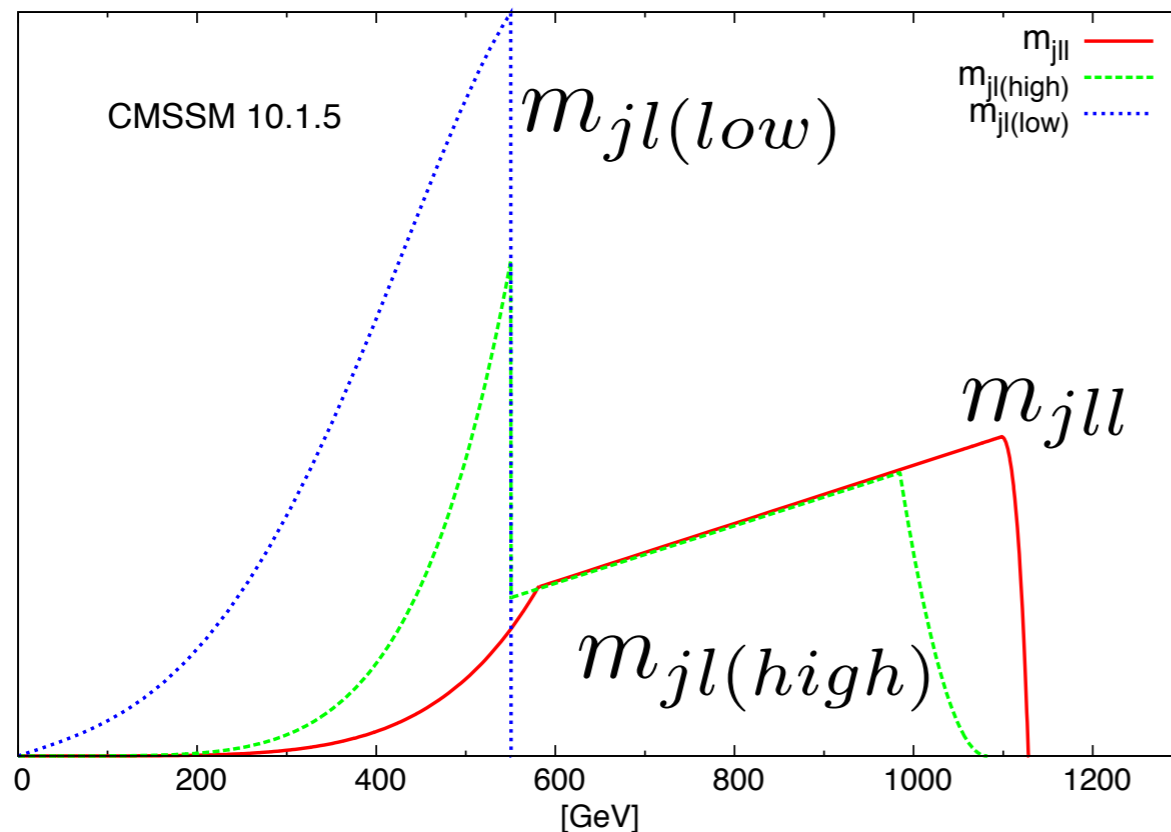
$$pp \rightarrow \tilde{q}_L \tilde{q}'_R \rightarrow q \tilde{\chi}_1^0 q' l^+ l^- \tilde{\chi}_1^0 (+X)$$



The “golden” decay chain



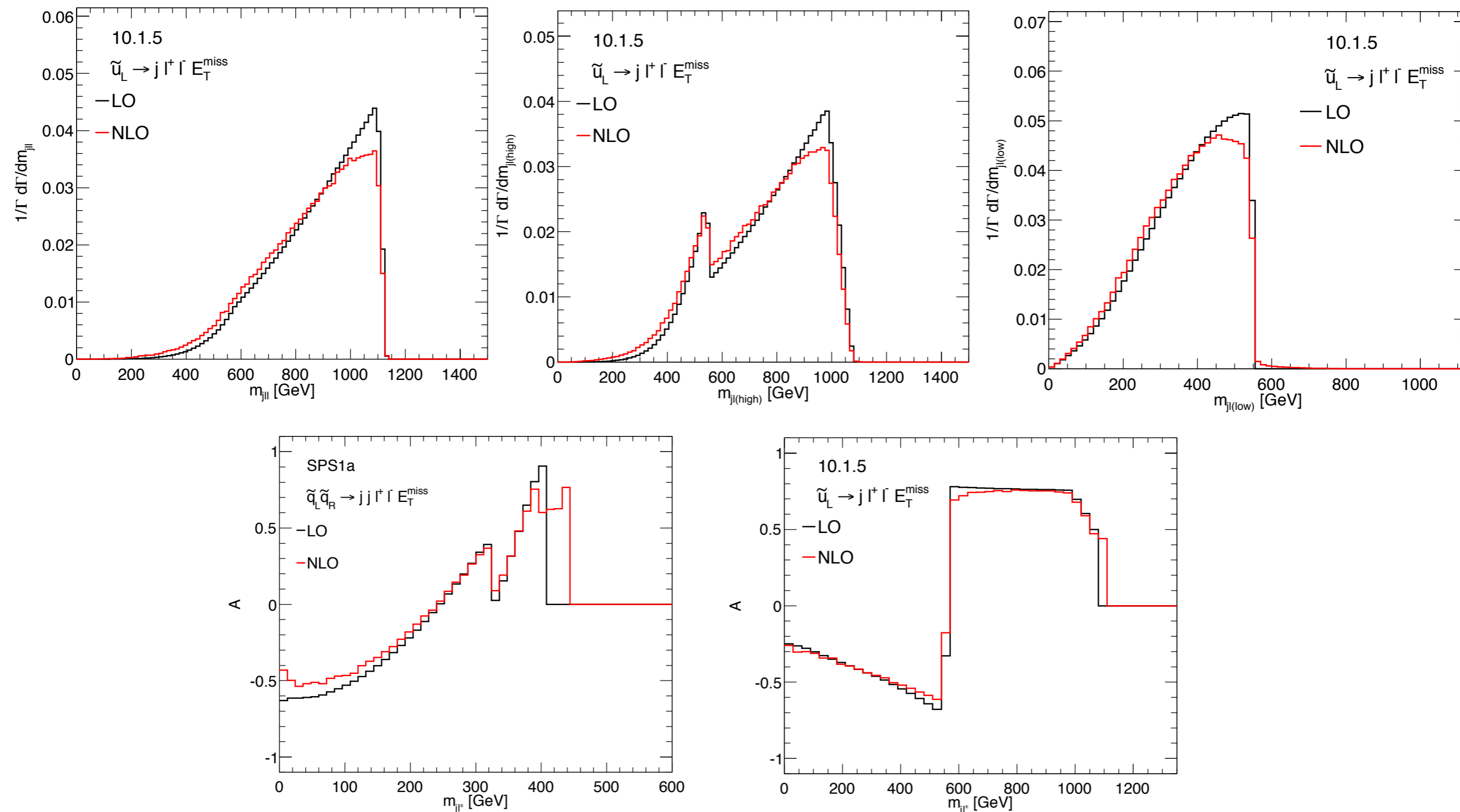
- Search for SUSY in “jets + OSSF leptons” channel
- Possible to measure **masses** of intermediate sparticles from invariant mass distribution endpoints and shapes ($m_{jll}, m_{jl(high)}, m_{jl(low)}, \dots$).
- Possible to measure **spin** of sparticles via charge asymmetries.



Comparison between NLO and LO corrections purely in the **shapes** of distributions.

10.1.5	\tilde{u}_L	\tilde{u}_R	\tilde{d}_L	\tilde{d}_R	\tilde{g}	$\tilde{\chi}_1^0$
mass (GeV)	1437.7	1382.3	1439.7	1376.9	1568.6	291.3

LHC 14 TeV



CONCLUSION

Knowledge of higher-order corrections to squark/gluino processes are important for precise description of physical observables and thus for setting **accurate limits** and even more for **parameter determination**.

We provide a fully differential calculation of factorizable NLO QCD corrections in NWA for squark-squark production and different decays.

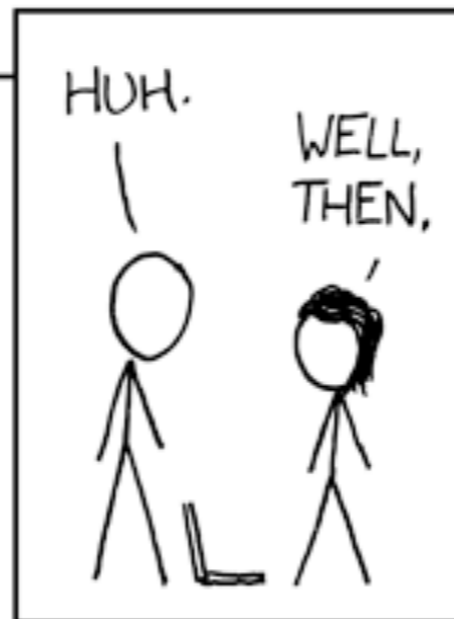
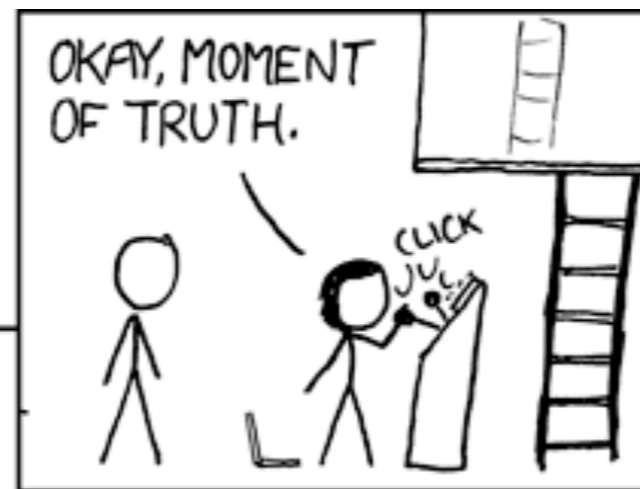
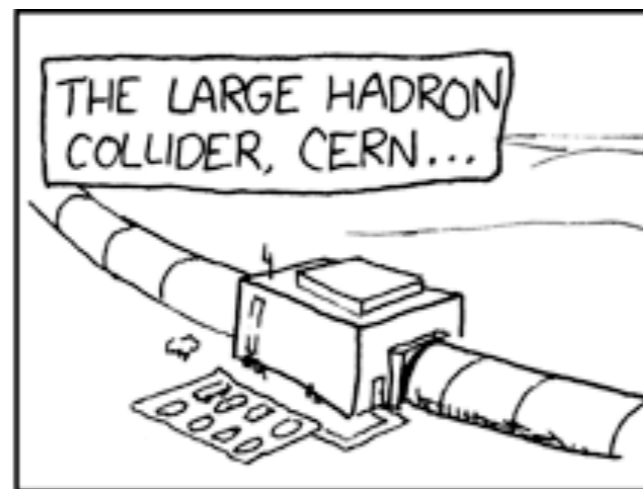
OUTLOOK

Study of further experimental signatures (monojets) under way.

Fully differential NLO QCD predictions of combined production and decay for all squark/gluino channels are desirable (matched to a NLO PS).

Discovery of SUSY in the next run of the LHC.

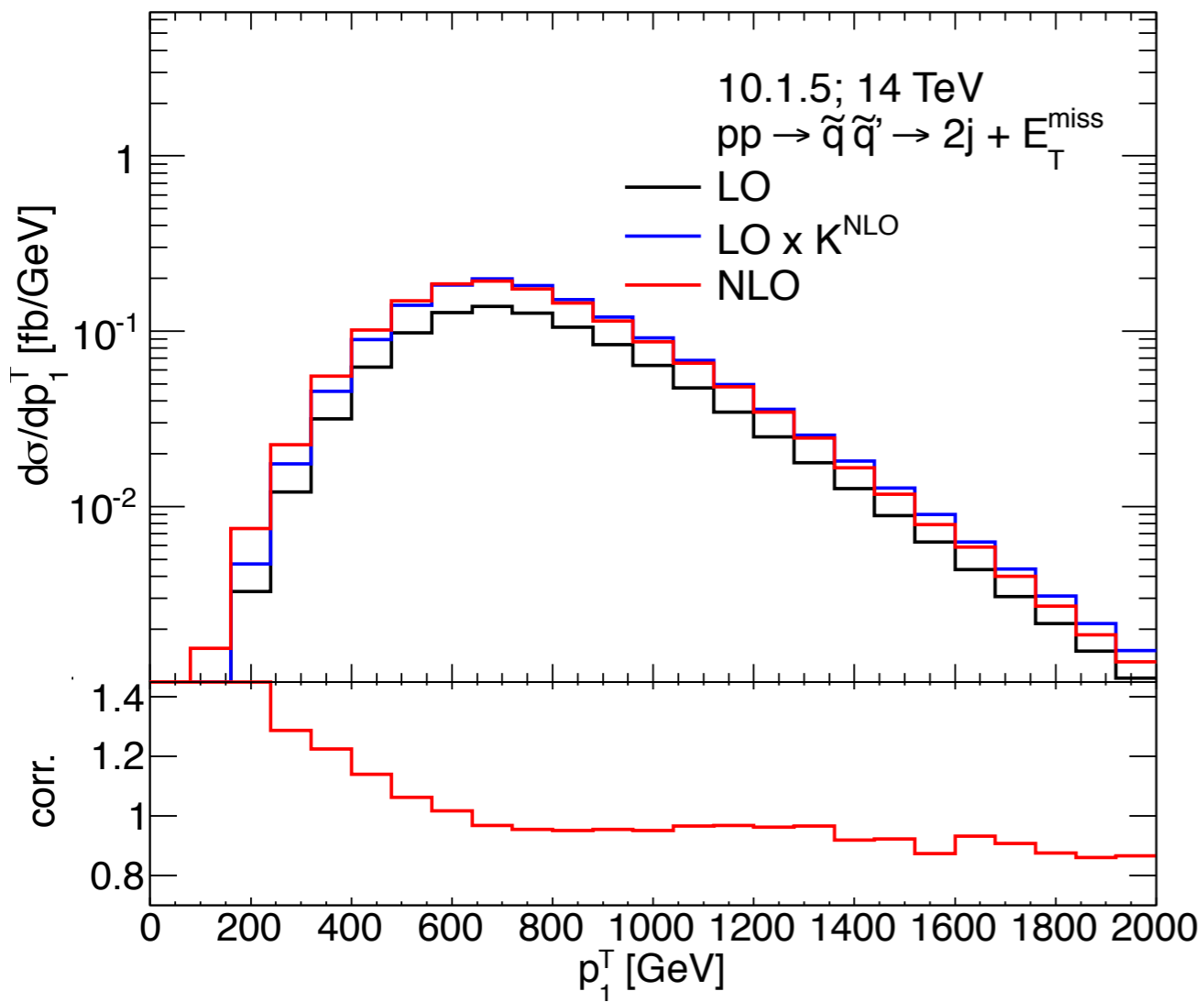
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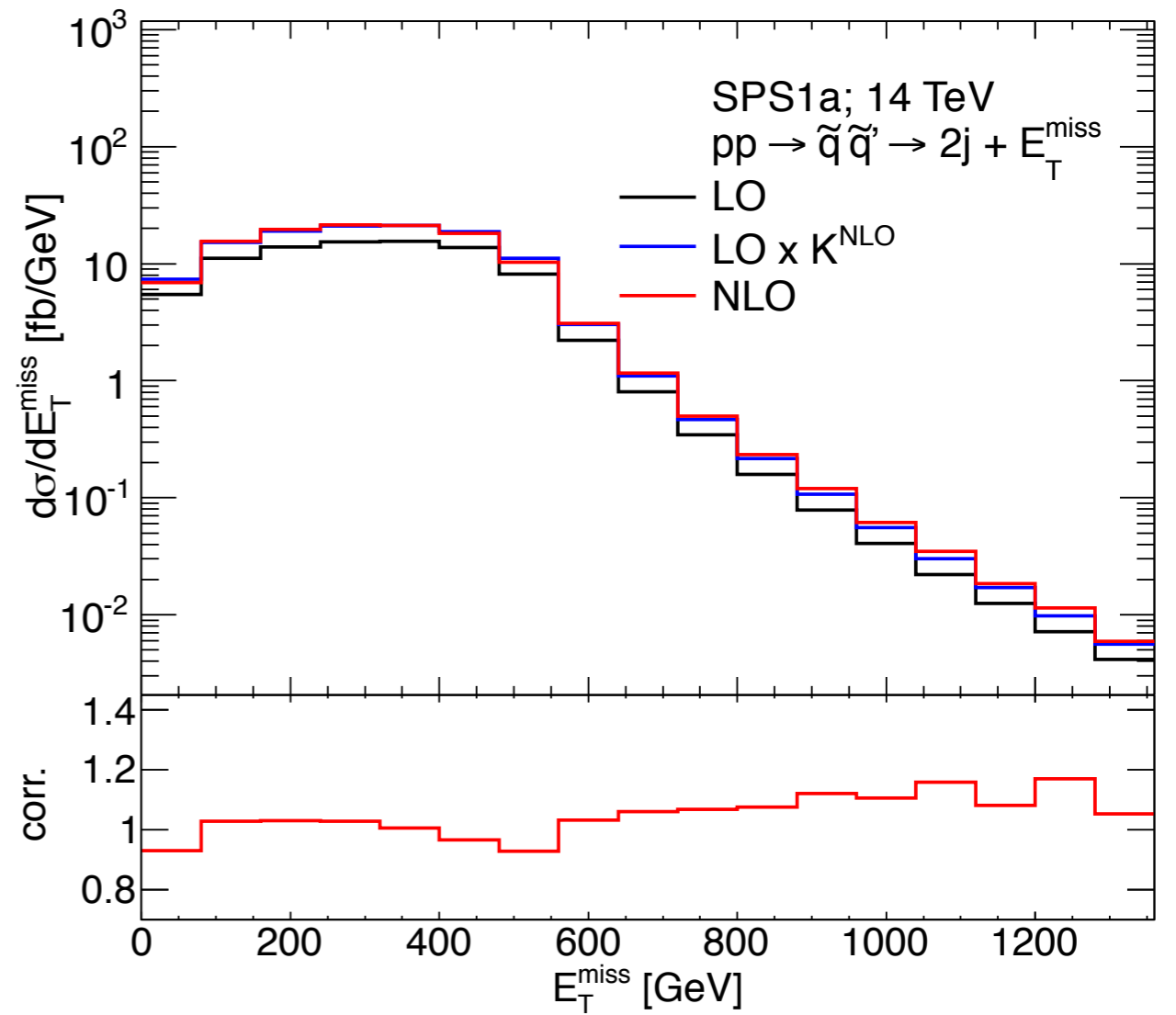
CMSSM 10.1.5 (14 TeV)

Comparison between NLO and LO rescaled by global K-factor:
corrections purely in the **shapes**

10.1.5	\tilde{u}_L	\tilde{u}_R	\tilde{d}_L	\tilde{d}_R	\tilde{g}	$\tilde{\chi}_1^0$
mass (GeV)	1437.7	1382.3	1439.7	1376.9	1568.6	291.3



p_1^T

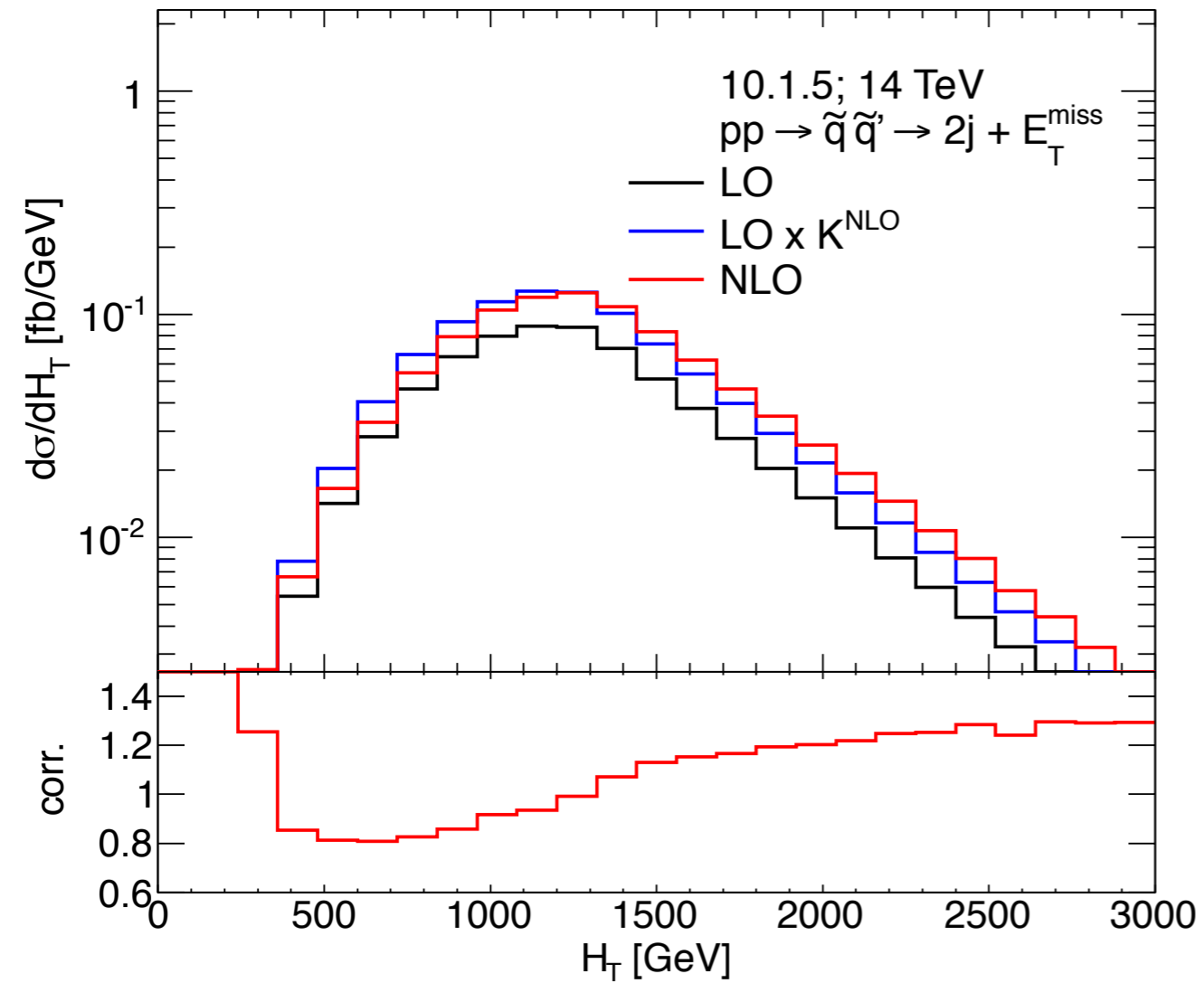


E_T^{miss}

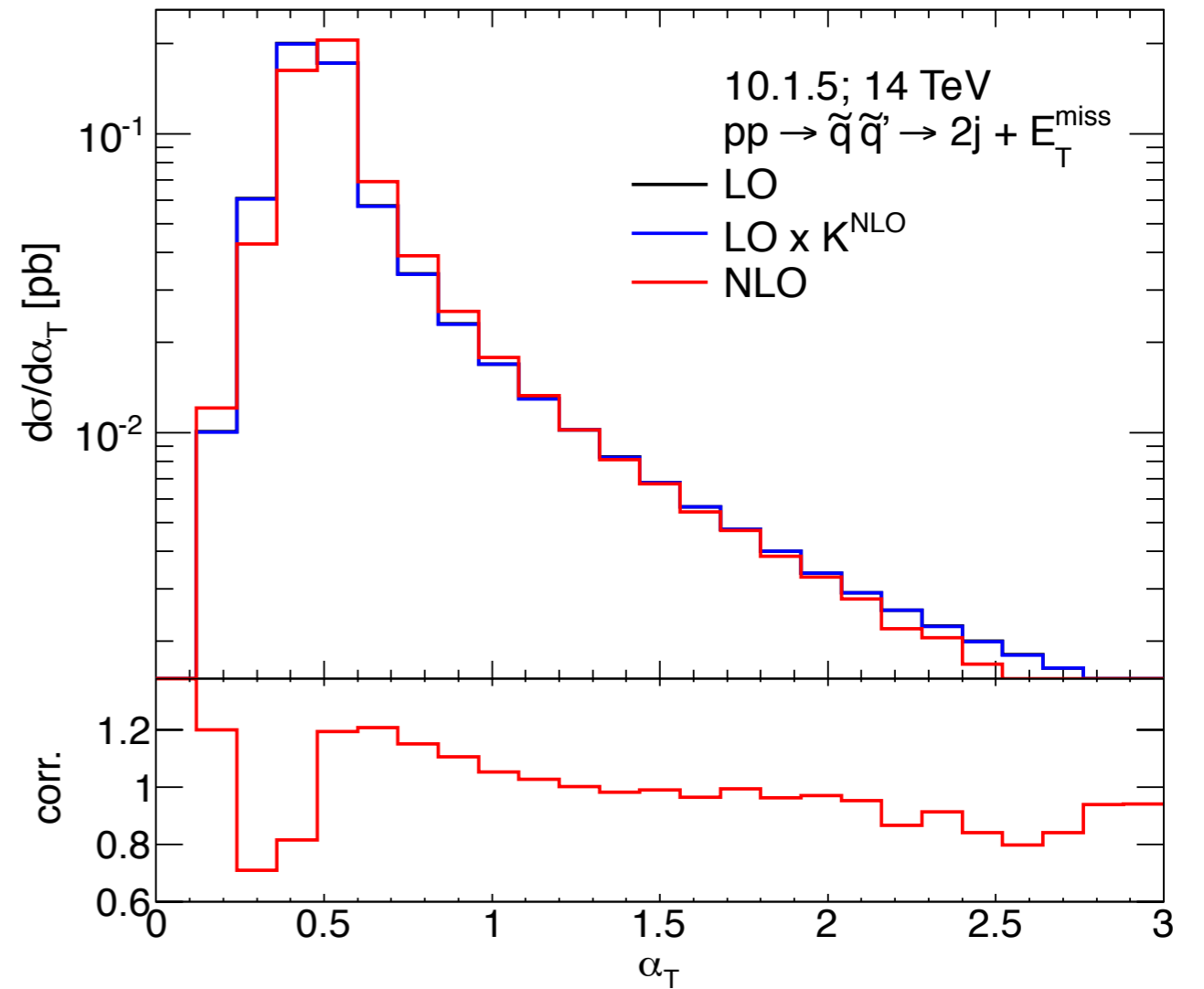
CMSSM 10.1.5 (14 TeV)

Comparison between NLO and LO rescaled by global K-factor:
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10.1.5	\tilde{u}_L	\tilde{u}_R	\tilde{d}_L	\tilde{d}_R	\tilde{g}	$\tilde{\chi}_1^0$
mass (GeV)	1437.7	1382.3	1439.7	1376.9	1568.6	291.3



$$H_T = \sum_{i=1,2(,3)} p_i^T$$

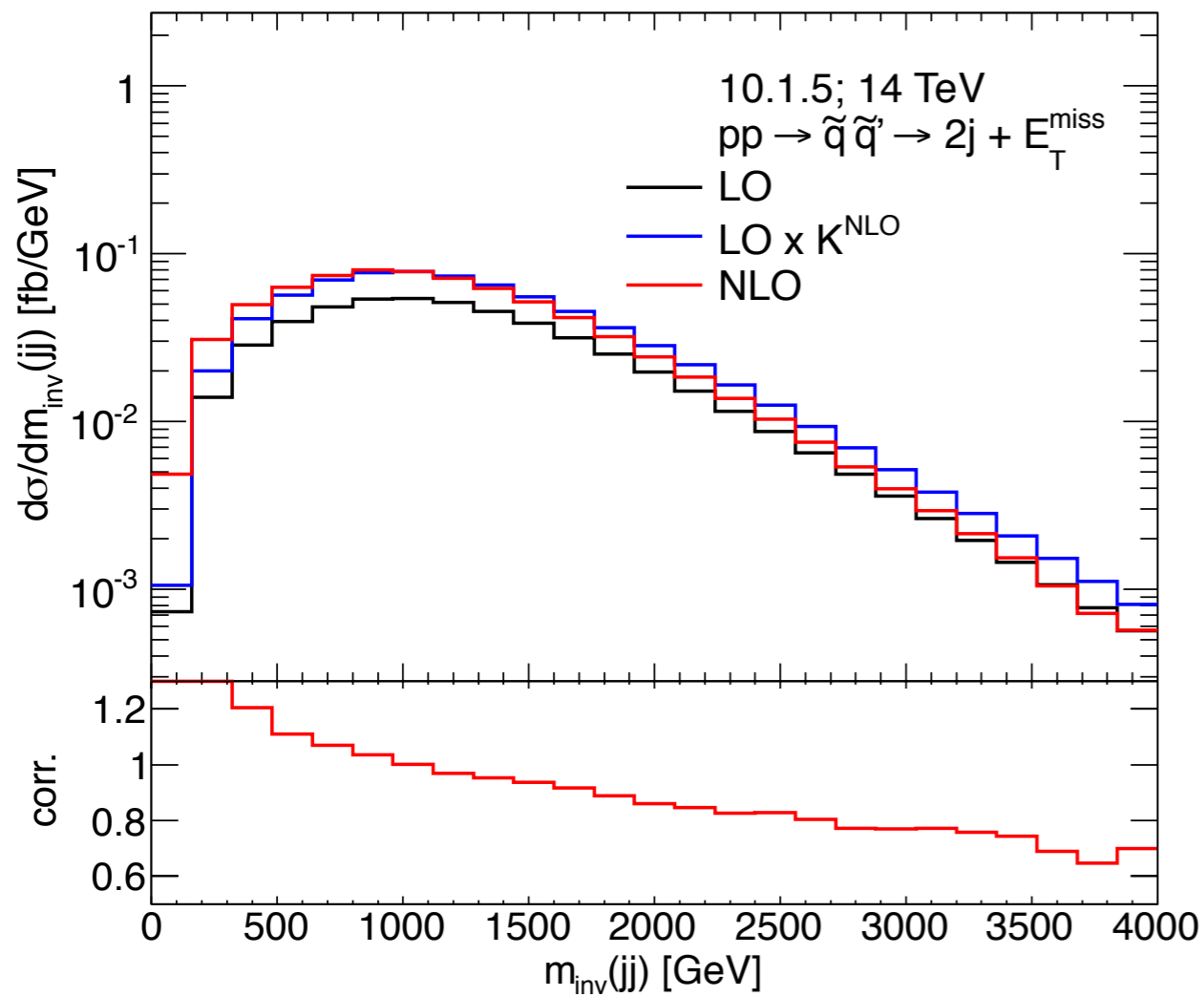


$$\alpha_T = E_T^{j2} / M_T$$

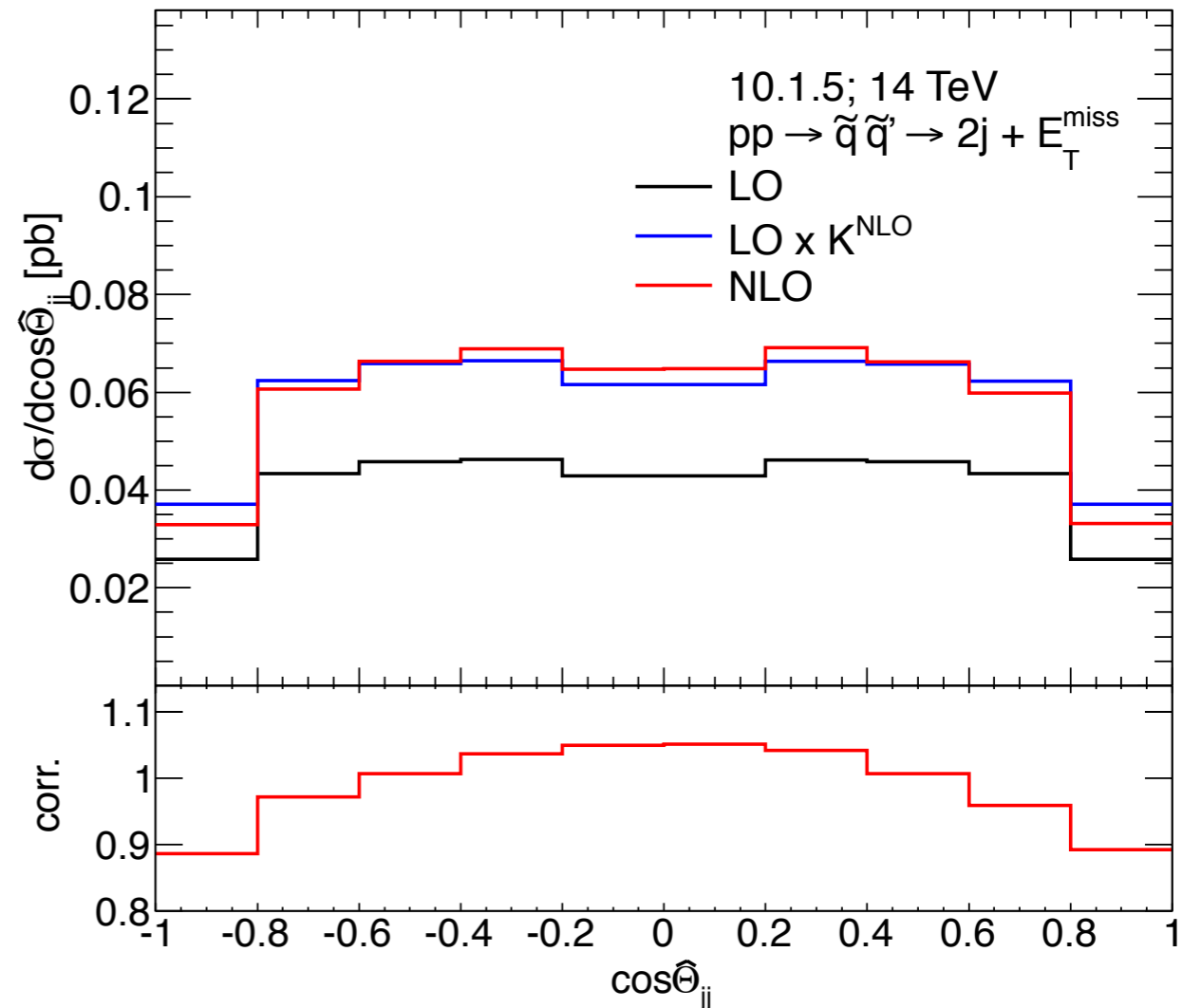
CMSSM 10.1.5 (14 TeV)

Comparison between NLO and LO rescaled by global K-factor:
corrections purely in the **shapes**

10.1.5	\tilde{u}_L	\tilde{u}_R	\tilde{d}_L	\tilde{d}_R	\tilde{g}	$\tilde{\chi}_1^0$
mass (GeV)	1437.7	1382.3	1439.7	1376.9	1568.6	291.3



$$m_{\text{inv}}(jj)$$

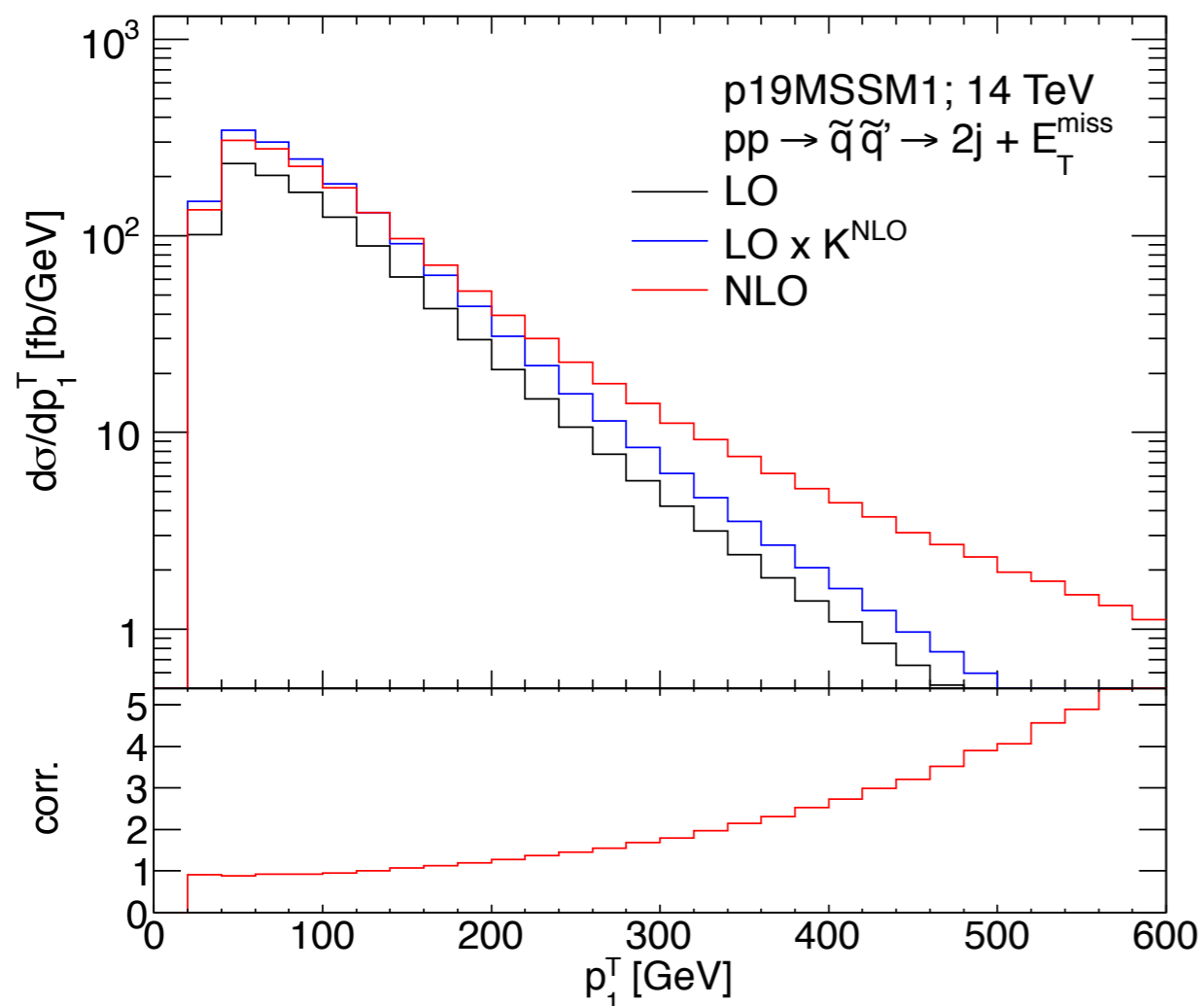


$$\cos \hat{\Theta} = \tanh \left(\frac{\Delta \eta_{jj}}{2} \right)$$

p19MSSM1A (14 TeV)

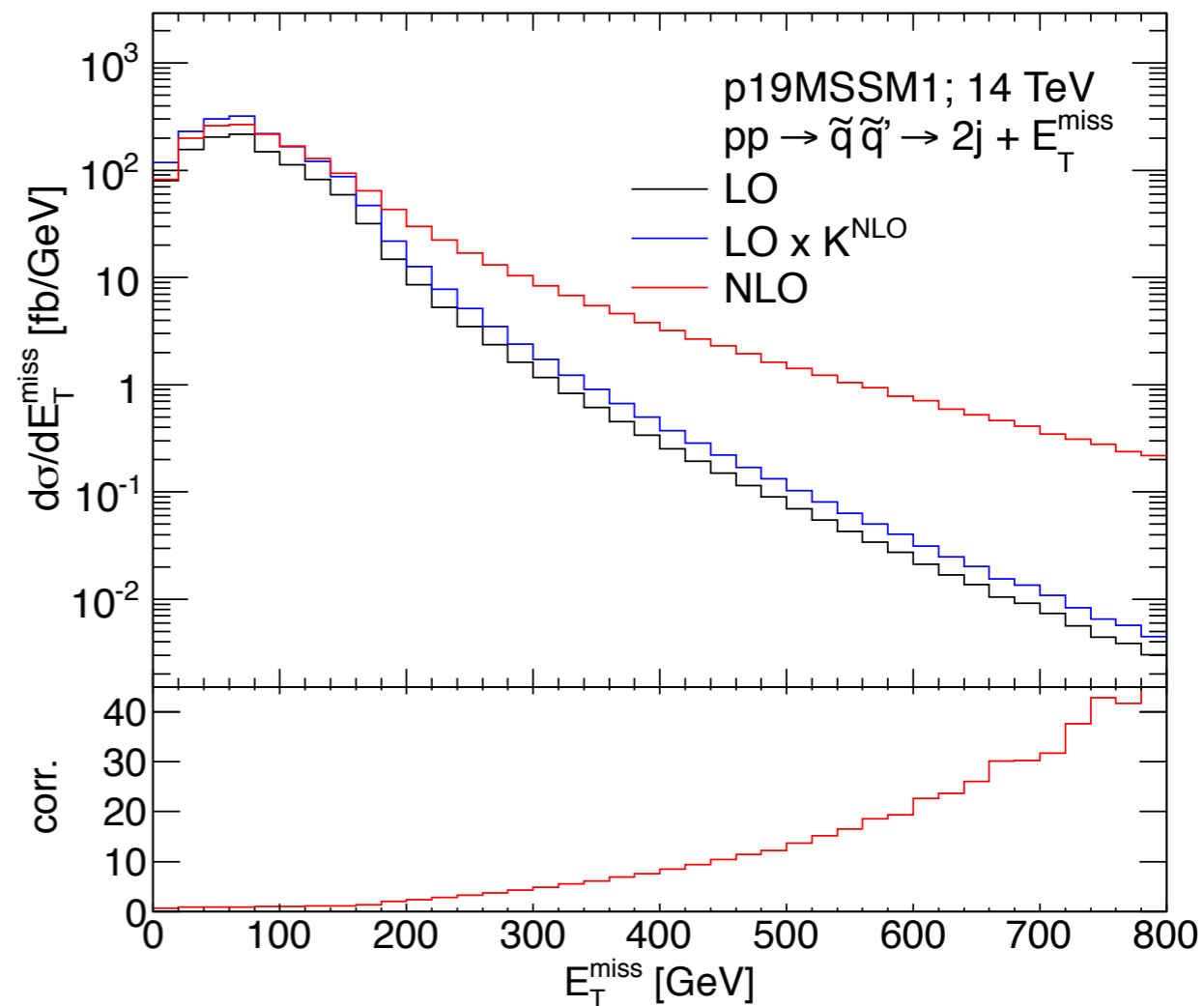
Comparison between NLO and LO rescaled by global K-factor:
corrections purely in the **shapes**

p19MSSM1	\tilde{u}_L	\tilde{u}_R	\tilde{d}_L	\tilde{d}_R	\tilde{g}	$\tilde{\chi}_1^0$
mass (GeV)	339.6	394.8	348.3	392.7	414.7	299.1



p_1^T

see also: Plehn, Rainwater, Skands '07;
 Alwall, de Visscher, Maltoni '08



E_T^{miss}

Effect on cut-and-count searches performed by **ATLAS**.

2j - signal region:

$$p_{j_1}^T > 130 \text{ GeV}, p_{j_2}^T > 40 \text{ GeV}, |\eta_{j_{1/2}}| < 2.8, \Delta\phi(j_{1/2}, \vec{E}_T) > 0.4.$$

$$m_{\text{eff}} > 1 \text{ TeV}, \cancel{E}_T/m_{\text{eff}} > 0.3,$$

benchmarkpoint	Energy [TeV]	$N_{\text{ATLAS}}^{(0)}$	$N_{\text{ATLAS}}^{(0+1)}$	$K_{N_{\text{ATLAS}}}$	$K_{pp \rightarrow \tilde{q}\tilde{q}'}$
SPS1a	7	0.066 pb	0.083 pb	1.26	1.37
	8	0.097 pb	0.121 pb	1.25	1.35
	14	0.347 pb	0.424 pb	1.22	1.28
10.1.5	7	0.313 fb	0.503 fb	1.61	1.57
	8	0.861 fb	1.344 fb	1.56	1.52
	14	13.82 fb	19.77 fb	1.43	1.40
p19MSSM1	7	0.140 fb	20.76 fb	~ 150	1.40
	8	0.339 fb	37.96 fb	~ 110	1.39
	14	0.0044 pb	0.264 pb	~ 60	1.34

differential
combined NLO

flat K-factor
of just production

Effect on cut-and-count searches performed by **CMS**.

α_T - signal region:

$$p_{j_{1/2}}^T > 100 \text{ GeV}, |\eta_{j_1}| < 2.5, |\eta_{j_2}| < 3.0,$$

$$H_T > 350 \text{ GeV}, \cancel{H}_T/\cancel{E}_T < 1.25, \alpha_T > 0.55,$$

benchmarkpoint	Energy [TeV]	$N_{\text{CMS}}^{(0)}$	$N_{\text{CMS}}^{(0+1)}$	$K_{N_{\text{CMS}}}$	$K_{pp \rightarrow \tilde{q}\tilde{q}'}$
SPS1a	7	0.112 pb	0.141 pb	1.26	1.37
	8	0.157 pb	0.197 pb	1.25	1.35
	14	0.488 pb	0.614 pb	1.26	1.28
10.1.5	7	0.201 pb	0.261 pb	1.30	1.57
	8	0.542 fb	0.674 fb	1.24	1.52
	14	8.129 fb	8.884 fb	1.09	1.40
p19MSSM1	7	10^{-6} pb	0.095 pb	$\mathcal{O}(10^4)$	1.40
	8	10^{-6} pb	0.151 pb	$\mathcal{O}(10^4)$	1.39
	14	$2 \cdot 10^{-5}$ pb	0.687 pb	$\mathcal{O}(10^4)$	1.34

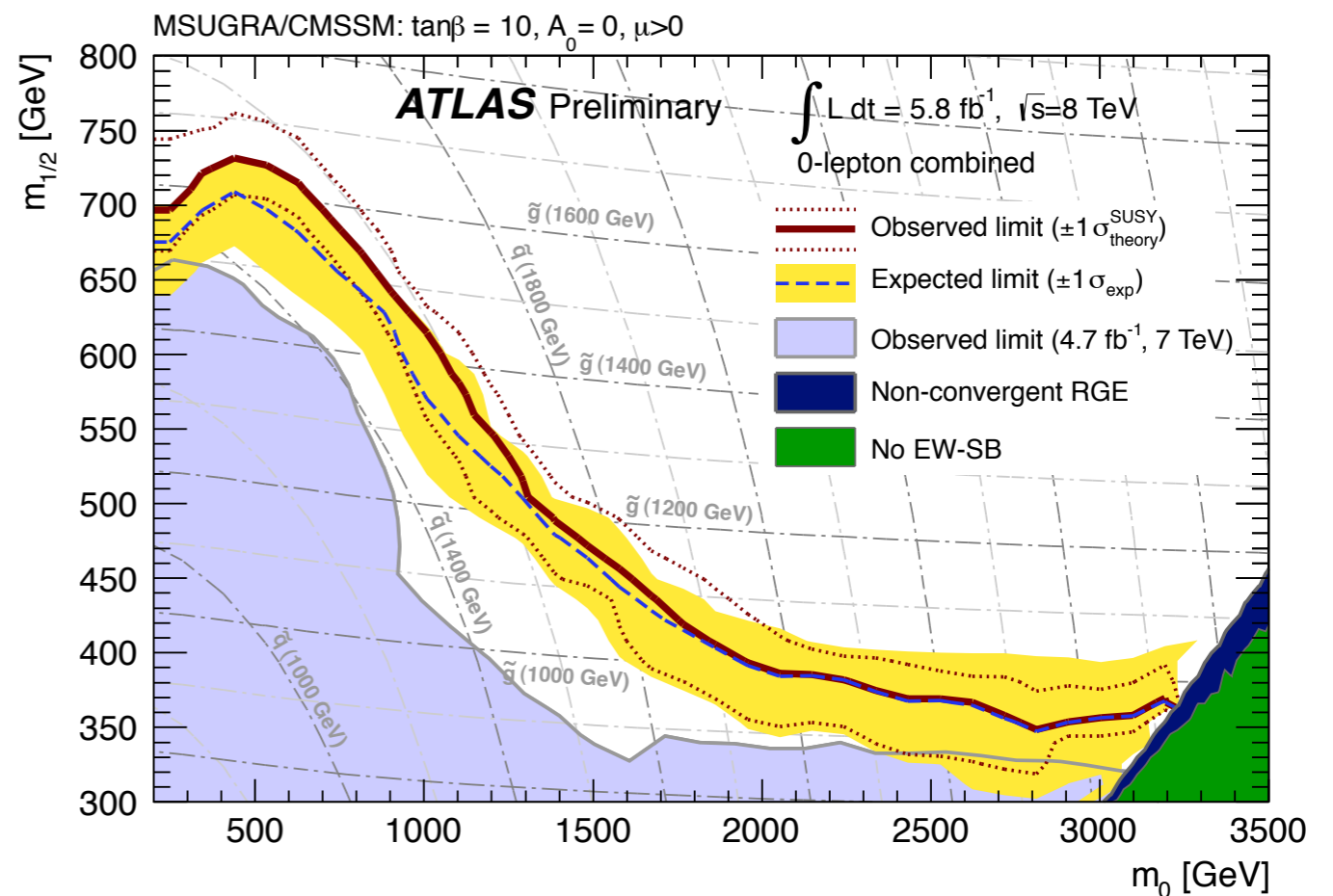
differential
combined NLO

flat K-factor
of just production

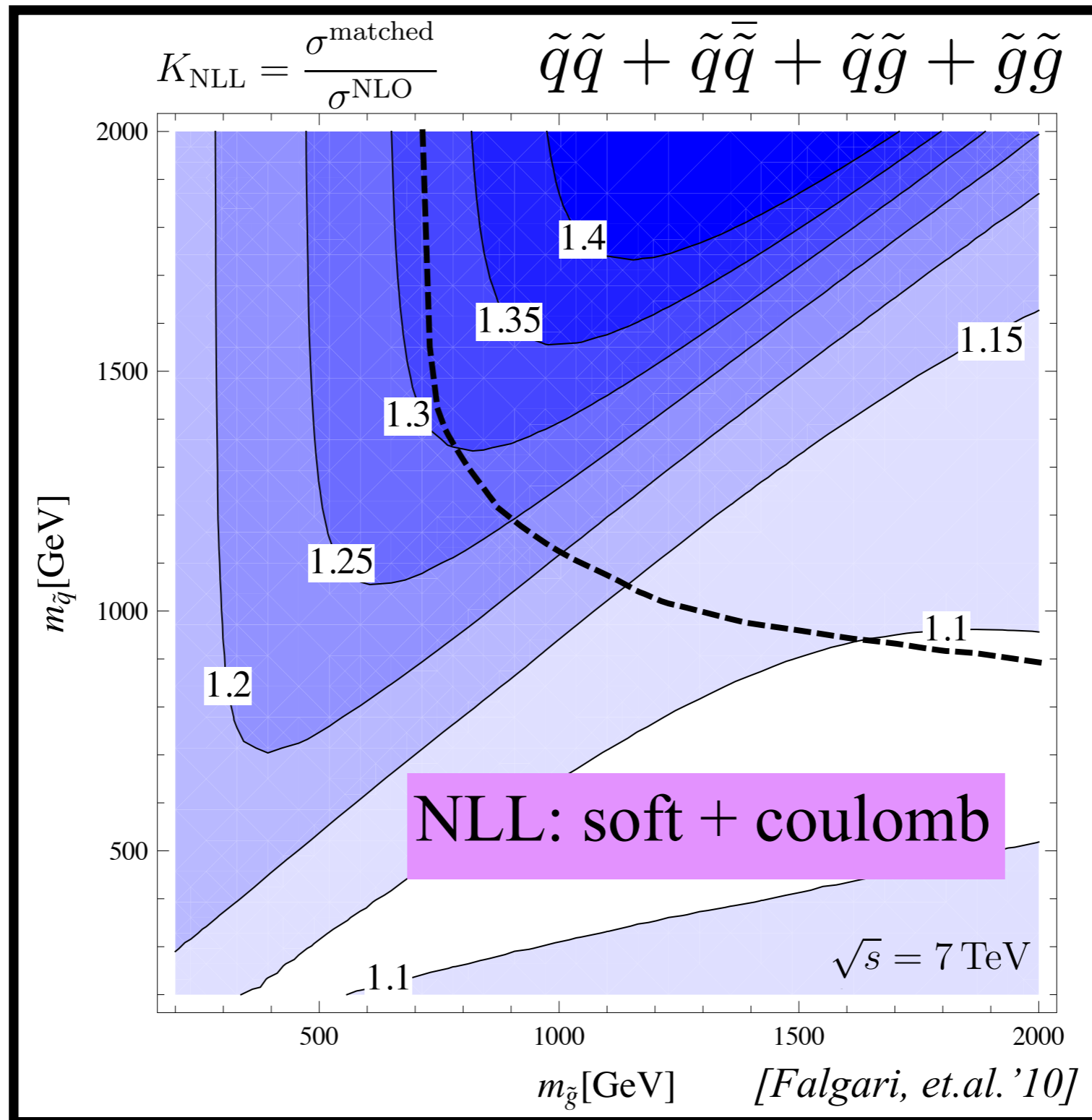
SQUARKS AND GLUINOS AT THE LHC

Requirement	Channel				
	A 2-jets	B 3-jets	C 4-jets	D 5-jets	E 6-jets
$E_T^{\text{miss}} [\text{GeV}] >$	160				
$p_T(j_1) [\text{GeV}] >$	130				
$p_T(j_2) [\text{GeV}] >$	60				
$p_T(j_3) [\text{GeV}] >$	-	60	60	60	60
$p_T(j_4) [\text{GeV}] >$	-	-	60	60	60
$p_T(j_5) [\text{GeV}] >$	-	-	-	60	60
$p_T(j_6) [\text{GeV}] >$	-	-	-	-	60
$\Delta\phi(\text{jet}, \mathbf{E}_T^{\text{miss}})_{\text{min}} [\text{rad}] >$	0.4 ($i = \{1, 2, (3)\}$)		0.4 ($i = \{1, 2, 3\}$), 0.2 ($p_T > 40 \text{ GeV jets}$)		
$E_T^{\text{miss}} / m_{\text{eff}}(Nj) >$	0.3/0.4/0.4 (2j)	0.25/0.3/- (3j)	0.25/0.3/0.3 (4j)	0.15 (5j)	0.15/0.25/0.3 (6j)
$m_{\text{eff}}(\text{incl.}) [\text{GeV}] >$	1900/1300/1000	1900/1300/-	1900/1300/1000	1700/-/-	1400/1300/1000

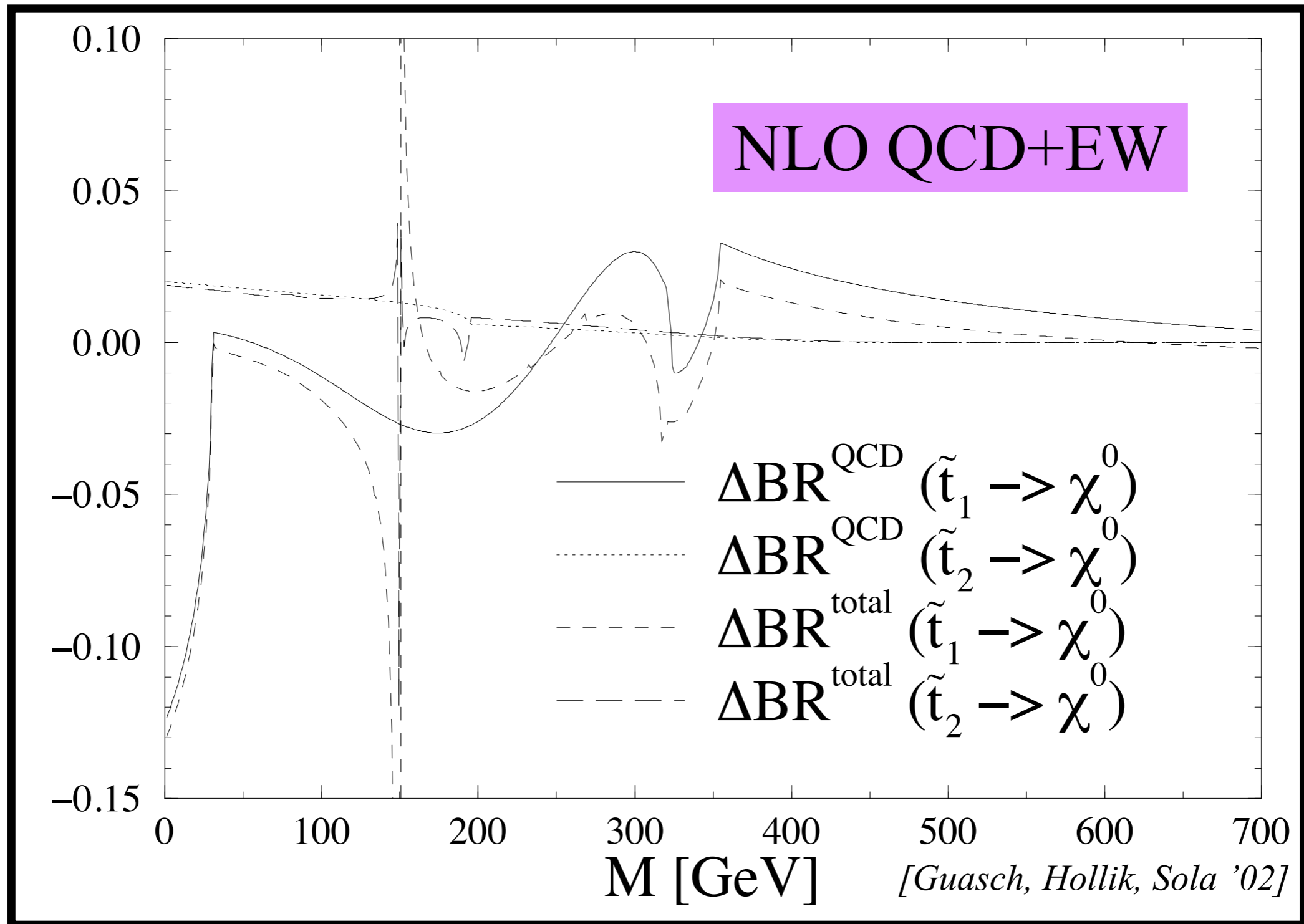
ATLAS search regions

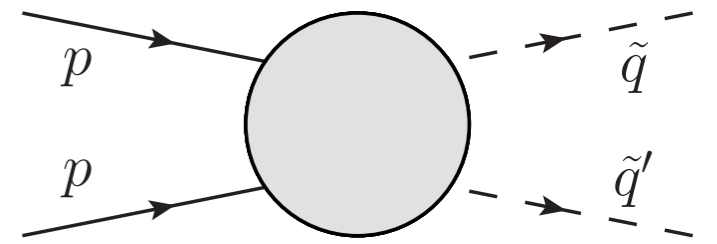


Higher Order Corrections to Production



Higher Order Corrections to Decay





All counterterms, but the one for the QCD coupling $\delta g_s = g_s \delta Z_{g_s}$ are renormalized according to the **on-shell** scheme.

Choice of scheme for the renormalization of the QCD coupling is fixed by definition of α_s in the PDF distributions: $\overline{MS} + 5$ flavour scheme.

$$\delta Z_{g_s} = -\frac{\alpha_s}{4\pi} \left[\Delta \frac{\beta_0}{2} + \frac{1}{3} \log \frac{m_t^2}{\mu_F^2} + \log \frac{m_{\tilde{g}}^2}{\mu_F^2} + \frac{1}{12} \sum_{\tilde{q}} \log \frac{m_{\tilde{q}}^2}{\mu_F^2} \right]$$

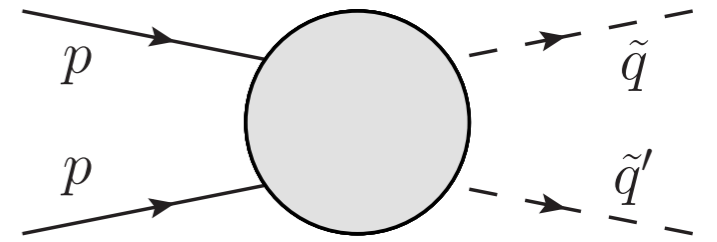
Using \overline{MS} and Dim. Reg. breaks supersymmetric Slavnov-Taylor identity, that relates the QCD coupling in the qqg QCD vertex and the \hat{g}_s coupling in the $q\tilde{q}\tilde{g}$ SQCD vertex.

Can be restored:

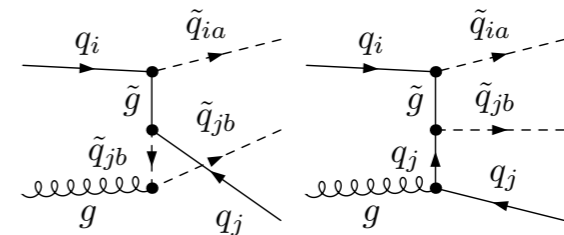
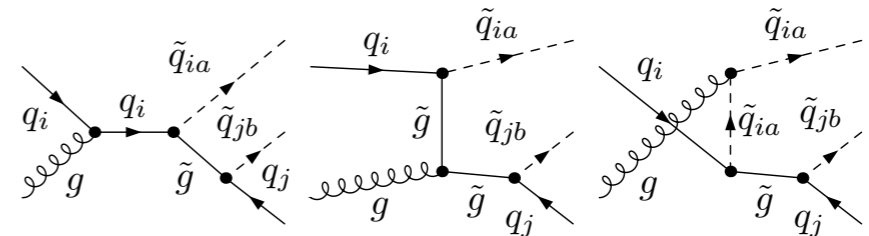
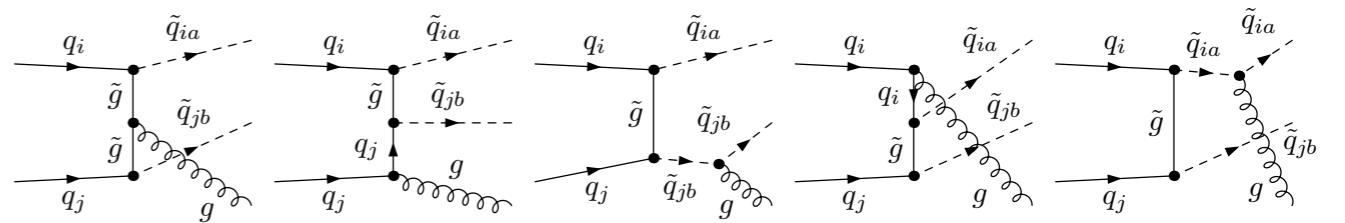
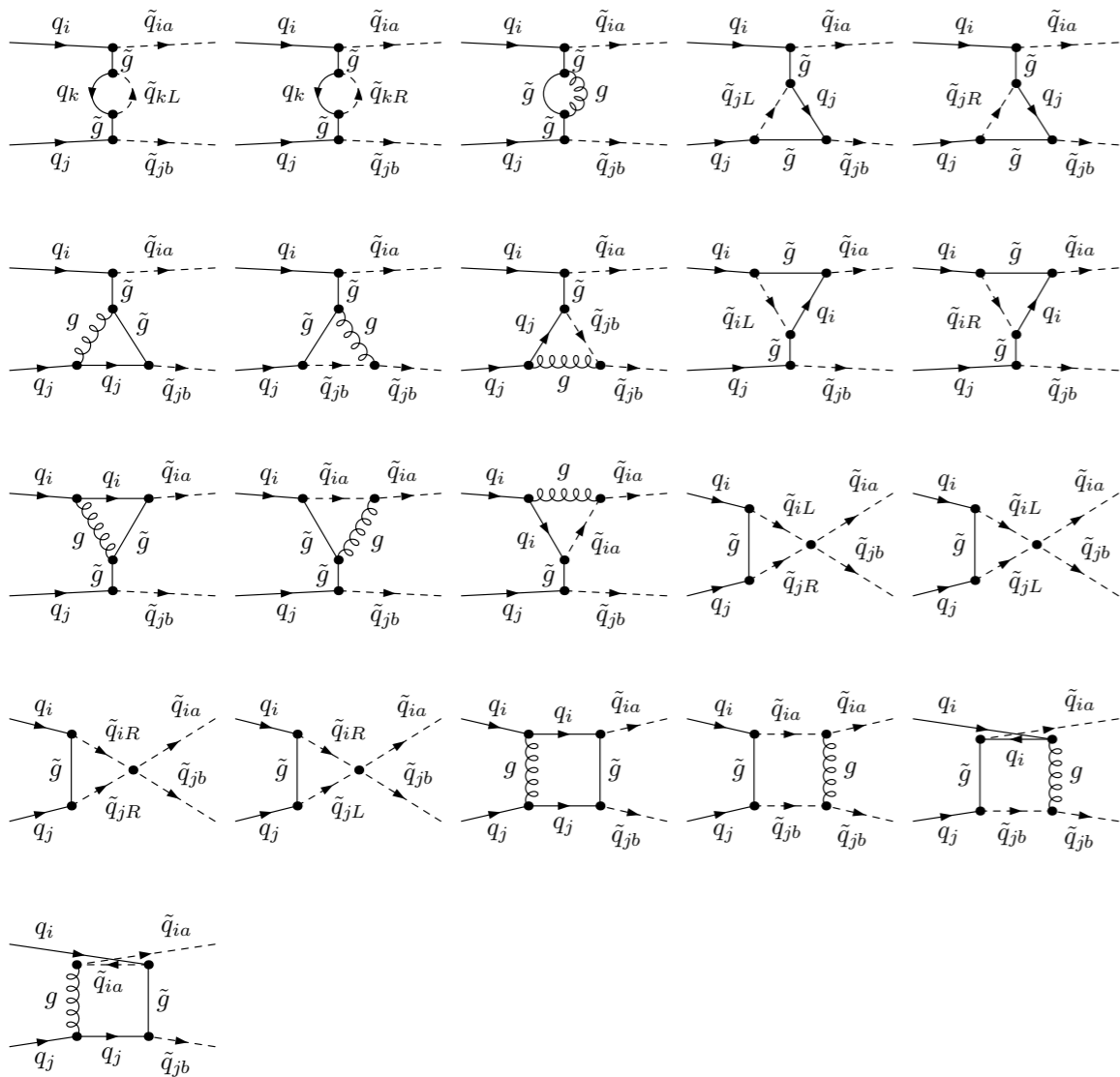
$$\delta Z_{\hat{g}_s} = \delta Z_{g_s} + \frac{\alpha_s}{3\pi}$$

[Beenakker et al. '96;
Hollik, Stöckinger '01]

NLO production

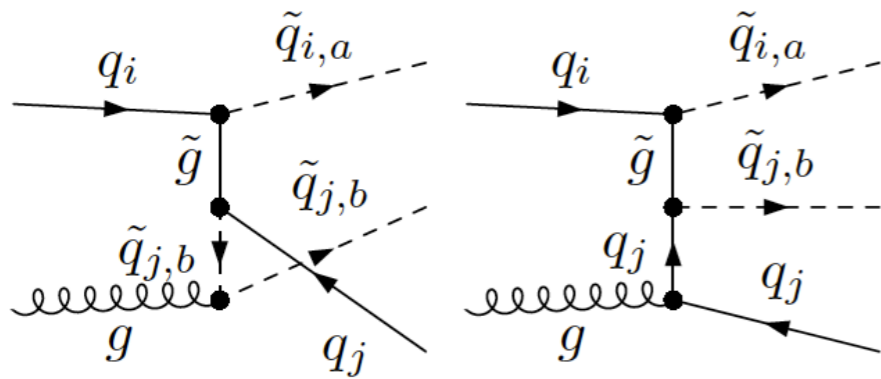
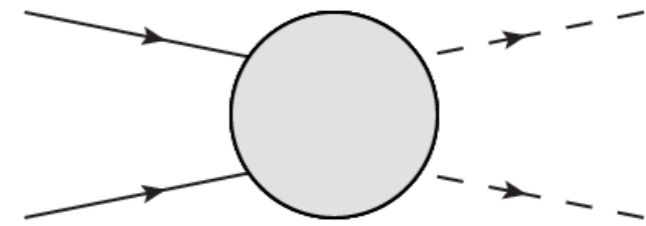


$$d\sigma_{pp \rightarrow \tilde{q}\tilde{q}'}^{(1)}(+X) = d\sigma_{pp \rightarrow \tilde{q}\tilde{q}'}^{\text{virtual+soft}}(g) + d\sigma_{pp \rightarrow \tilde{q}\tilde{q}'}^{\text{coll}}(g) + d\sigma_{pp \rightarrow \tilde{q}\tilde{q}'g}^{\text{hard}} + d\sigma_{pp \rightarrow \tilde{q}\tilde{q}'\bar{q}(\prime)}^{\text{real-quark}}$$

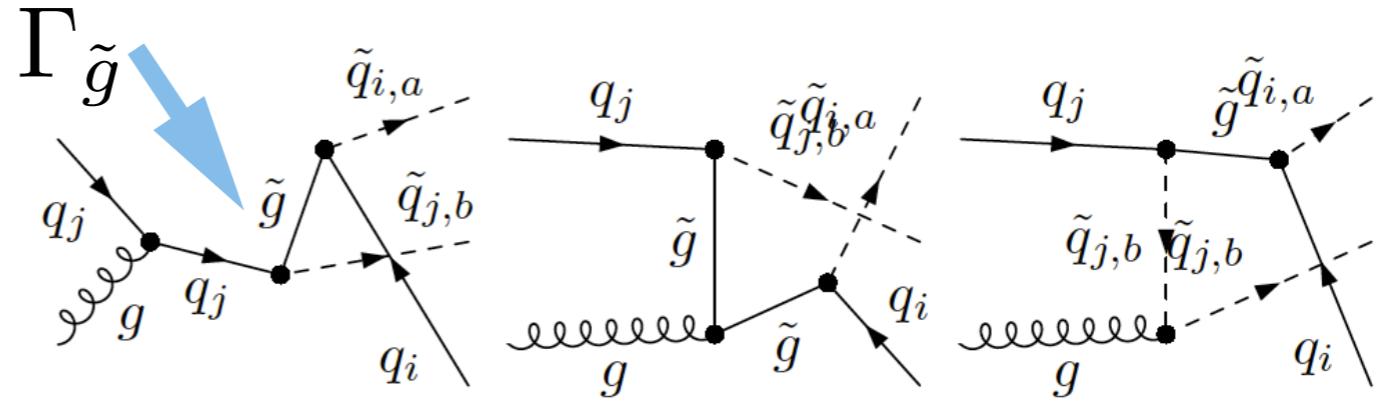


Real quark radiation

DR scheme



non-resonant



resonant

$$d\hat{\sigma}(q_i g \rightarrow \tilde{q}_{i,a} \tilde{q}_{i,b} q_i) = \frac{1}{\Phi} \left[|\mathcal{M}_{\text{nonres}}|^2 + 2\text{Re}(\mathcal{M}_{\text{nonres}} \mathcal{M}_{\text{res}}^*) \right]$$

DS scheme

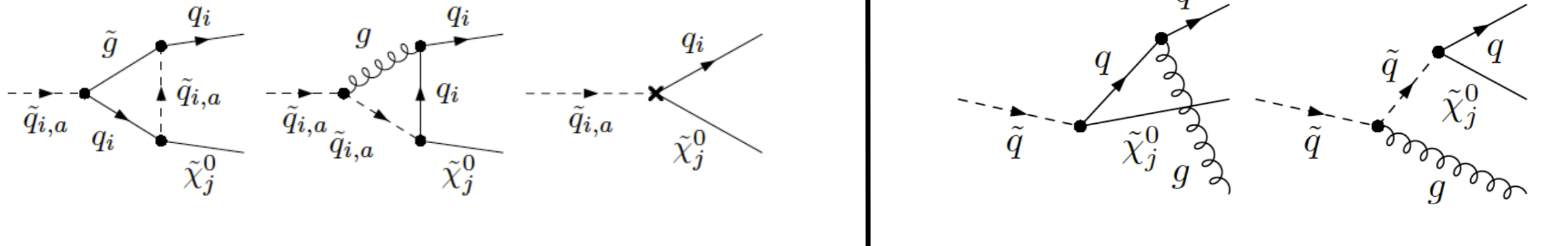
[Binoth et. al.; '11]

$$\frac{|\mathcal{M}|^2(s_{q\tilde{q}})}{(s_{q\tilde{q}} - m_{\tilde{g}}^2)^2 + m_{\tilde{g}}^2 \Gamma_{\tilde{g}}^2} \rightarrow \frac{|\mathcal{M}|^2(s_{q\tilde{q}})}{(s_{q\tilde{q}} - m_{\tilde{g}}^2)^2 + m_{\tilde{g}}^2 \Gamma_{\tilde{g}}^2} - \frac{|\mathcal{M}|^2(m_{\tilde{g}}^2)}{(s_{q\tilde{q}} - m_{\tilde{g}}^2)^2 + m_{\tilde{g}}^2 \Gamma_{\tilde{g}}^2}$$

and usually: $\Gamma \rightarrow 0$ numerically.

NLO decay

$$d\Gamma_{\tilde{q} \rightarrow q \tilde{\chi}_j^0}^{(1)} = d\Gamma_{\tilde{q} \rightarrow q \tilde{\chi}_j^0}^{\text{virtual}} + d\Gamma_{\tilde{q} \rightarrow q \tilde{\chi}_j^0}^{\text{soft}}(g) + d\Gamma_{\tilde{q} \rightarrow q \tilde{\chi}_j^0}^{\text{coll}}(g) + d\Gamma_{\tilde{q} \rightarrow q \tilde{\chi}_j^0}^{\text{hard}}g$$



NLO total decay

$$\Gamma_{\tilde{q} \rightarrow q \tilde{\chi}_j^0}^{(0+1)} = \Gamma^{(0)} \left[1 + \frac{4}{3} \frac{\alpha_s}{\pi} F^{QCD} \left(\frac{m_{\tilde{\chi}_j^0}}{m_{\tilde{q}}}, \frac{m_{\tilde{q}}}{m_{\tilde{g}}} \right) \right]$$

[Djouadi, Hollik, Jünger; '97]

analytical universal form factor,
recalculated with independent regulators