## Non-geometry in string theory

based on 1211.6437 in collaboration with<br>D. Andriot, M. Larfors and D. Lüst

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Particle Physics School Colloquium
March 15 ${ }^{\text {th }}, 2013$


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Innocent idea:

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2. Canonical quantisation: $\left[X^{\mu}, \Pi^{\nu}\right]=\mathrm{i}$

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2. Canonical quantisation: $\left[X^{\mu}, \Pi^{\nu}\right]=\mathrm{i}$
3. T-duality to a non-geometric situation: $X^{\mu} \rightarrow Z^{\mu}$
4. What are the commutators $\left[Z^{\mu}, Z^{\nu}\right]$ ?

## String theory as a field theory

- Consider a two-dimensional field theory with fields

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S=-\frac{1}{4 \pi \alpha^{\prime}} \int_{\Sigma} \mathrm{d}^{2} \sigma\left(G_{\mu \nu}(X) \eta^{\alpha \beta}+B_{\mu \nu}(X) \varepsilon^{\alpha \beta}\right) \partial_{\alpha} X^{\mu} \partial_{\beta} X^{\nu}
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- Define three-dimensional torus as "target space"

$$
X^{\mu}(\tau, \sigma+2 \pi)=X^{\mu}(\tau, \sigma)+2 \pi N^{\mu}
$$

## String theory as a field theory 2

- Choose simple metric and "H-flux"

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G=\left(\begin{array}{lll}
1 & 0 & 0 \\
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- Highly non-trivial equations of motion

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$$

- Our ansatz: dilute flux approximation

$$
X^{\mu}(\tau, \sigma)=X_{0}^{\mu}(\tau, \sigma)+H X_{H}^{\mu}(\tau, \sigma)+\mathcal{O}\left(H^{2}\right)
$$

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Solution to EOM and target space BC , order by order

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- $\mathcal{O}\left(H^{1}\right)$ : much more involved

$$
\begin{aligned}
X_{H}^{\mu}(\tau, \sigma)= & x_{H}^{\mu}+p_{H}^{\mu} \tau+\frac{i}{2} \sum_{n \neq 0} \frac{1}{n}\left(\widetilde{\gamma}_{n}^{\mu} e^{-\mathrm{i} n \sigma_{+}}+\gamma_{n}^{\mu} e^{-\mathrm{i} n \sigma_{-}}\right) \\
& -\epsilon^{\mu}{ }_{\nu \rho} p^{\rho} N^{\nu} \frac{\tau^{2}}{2} \\
& -\epsilon^{\mu}{ }_{\nu \rho} \frac{1}{2} \tau\left(\left.N^{\nu} X_{0}^{\rho}\right|_{\Sigma}-p^{\nu} \tilde{X}_{0}^{\rho} \mid \Sigma\right) \\
& -\left.\left.\epsilon^{\mu}{ }_{\nu \rho} \frac{1}{4} \tilde{X}_{0}^{\nu}\right|_{\Sigma} X_{0}^{\rho}\right|_{\Sigma}
\end{aligned}
$$

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- Impose canonical equal- "time" commutation relations

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with momentum

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\Pi_{\mu} \equiv \frac{\delta \mathcal{L}}{\delta \partial_{\tau} X^{\mu}}=\frac{1}{\pi}\left(G_{\mu \nu}(X) \partial_{\tau} X^{\nu}+B_{\mu \nu}(X) \partial_{\sigma} X^{\nu}\right)
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- Again, proceed order by order


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- Plug in mode expansion and read off

$$
\begin{aligned}
& {\left[\widetilde{\alpha}_{m}^{\mu}, \widetilde{\alpha}_{n}^{\nu}\right]=\left[\alpha_{m}^{\mu}, \alpha_{n}^{\nu}\right]=m \delta_{m,-n} \eta^{\mu \nu}} \\
& {\left[x^{\mu}, p^{\nu}\right]=\frac{i}{2} \eta^{\mu \nu}}
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- Instead, two terms

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\begin{aligned}
0 & =\left.\left[X^{\mu}(\tau, \sigma), X^{\nu}\left(\tau, \sigma^{\prime}\right)\right]\right|_{H} \\
& =H\left[X_{0}^{\mu}(\tau, \sigma), X_{H}^{\nu}\left(\tau, \sigma^{\prime}\right)\right]+H\left[X_{H}^{\mu}(\tau, \sigma), X_{0}^{\nu}\left(\tau, \sigma^{\prime}\right)\right]
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- Finally, read off some - but by far not all - commutators

$$
\begin{array}{r}
{\left[\gamma^{1}, \alpha^{2}\right]-\left[\gamma^{2}, \alpha^{1}\right] \neq 0} \\
{\left[\gamma^{1}, p^{2}\right]-\left[p_{H}^{2}, \alpha^{1}\right] \neq 0} \\
{\left[\gamma^{1}, N^{2}\right]=0}
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- No exhaustive list of such commutators $\rightarrow$ eventually, not possible to judge whether quantisation is fully consistent
- Difficulties come from solving order by order and cannot be overcome in that procedure, even in principle


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- Always necessary: isometry in the target space
- Here: metric and $B$-field only depend on $X^{3} \rightarrow \mathbf{2}$ isometries


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- $f$ makes this "non-geometric": going around a circle $Z^{3} \rightarrow Z^{3}+2 \pi$ cannot be absorbed by gauge transformation or diffeomorphism (for $X$-frame it could!)


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- Integrate! $\rightarrow$ super-complicated expressions $Z^{\mu}(\tau, \sigma)=\ldots$
- Unknown integration constants (operators)


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\left[Z^{1}(\tau, \sigma), Z^{2}\left(\tau, \sigma^{\prime}\right)\right]=\text { long and nasty computation }
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- Strategy: use obtained expressions for $Z^{\mu}$ and known commutators for expansion coefficients
- Plug in!

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- Had to choose particular commutators for the unknown integration constants

Non-commutativity 2

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Results

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