

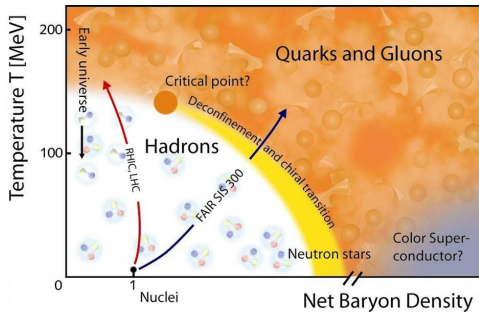
Holographic Analysis of Energy Loss Processes in Strongly Coupled Plasmas

Ling Lin

Nov 26, 2012

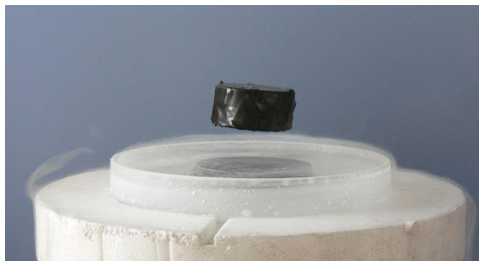
Motivation

- Quark-gluon plasma is strongly coupled.



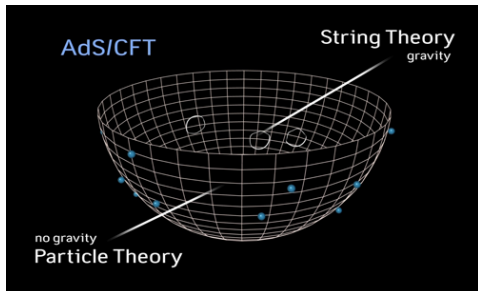
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- Quark-gluon plasma is strongly coupled.
- Condensed matter systems may have strong coupling.
- Theoretical interest in holographic principle.



The Holographic Principle

- Realization: AdS/CFT correspondence (Maldacena, 1997):

$\mathcal{N} = 4$ supersymmetric
 $SU(N_c)$ gauge theory in
4d Minkowski space

is dual to

type IIB string theory in
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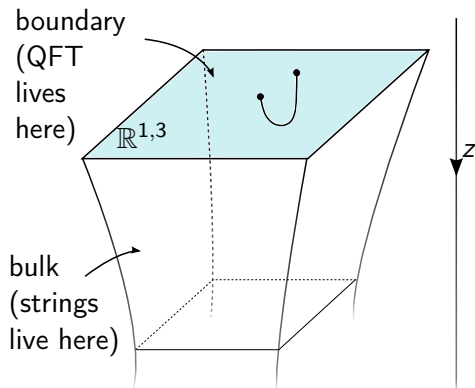
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- AdS/CFT 'dictionary' translates gauge theory quantities to gravity calculations.

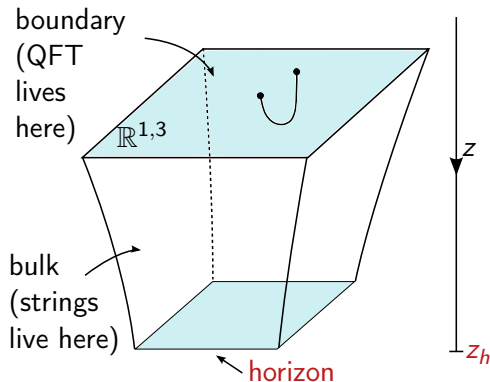
The Holographic Principle



$$\text{Bulk metric: } G_{\alpha\beta} dx^\alpha dx^\beta = \frac{R^2}{z^2} \left(-h(z) dt^2 + d\vec{x}^2 + \frac{dz^2}{h(z)} \right)$$

$$h(z) = 1$$

The Holographic Principle

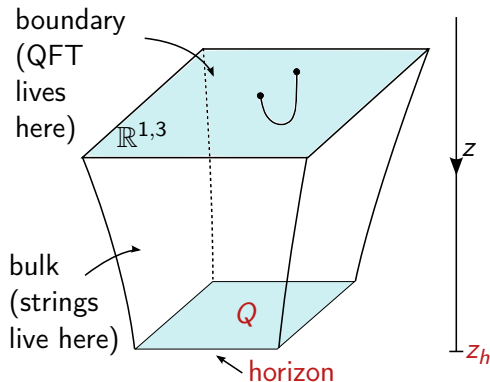


$$T = \frac{1}{\pi z_h}$$

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horizon function: $h(z) = 1 - \left(\frac{z}{z_h} \right)^4$

The Holographic Principle



$$T = \frac{1}{\pi z_h} \left(1 - \frac{Q^2}{2} \right)$$

$$\mu = \sqrt{3} \frac{Q}{z_h}$$

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The Holographic Principle

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- AdS-Reissner-Nordström (AdS-RN) solution is dual to $\mathcal{N} = 4$ theory at finite T and μ .
- Bottom-up approach: Propose a gravity theory and use dictionary to **define** the boundary theory.
- **Guideline:** Investigate dual theory by analysing observables.

Extended Models

- CGN-model [Colangelo, Giannuzzi, Nicotri] with *ad hoc* deformation c :

$$G_{\alpha\beta}^{\text{CGN}} dx^\alpha dx^\beta = \frac{R^2 e^{c^2 z^2}}{z^2} \left(-h(z) dt^2 + d\vec{x}^2 + \frac{dz^2}{h(z)} \right),$$

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- Exhibits 'confining phase' at low μ/T .
- Does not solve any known gravitational equations of motion.

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- '1-parameter model' with metric ansatz

$$ds^2 = e^{2A(z)} (-h(z)dt^2 + d\vec{x}^2) + \frac{e^{2B(z)}}{h(z)}dz^2$$

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$$S = \frac{1}{16\pi G_N} \int d^5x \sqrt{-g} \left(\mathcal{R} - \frac{1}{2}(\partial\phi)^2 - V(\phi) - \frac{f(\phi)}{4} F_{\mu\nu} F^{\mu\nu} \right)$$

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- Scalar ϕ may or may not be dilaton.
- Stability argument *forbids* low μ/T area. No confinement observed.

Jet Quenching Parameter \hat{q}

- Hard parton in plasma gets transverse kicks: $\hat{q} \equiv \frac{\langle p_{\perp}^2 \rangle}{L}$

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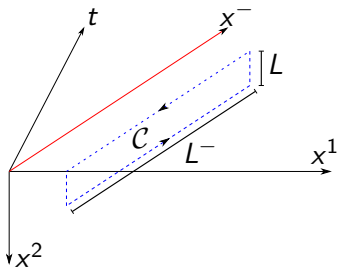
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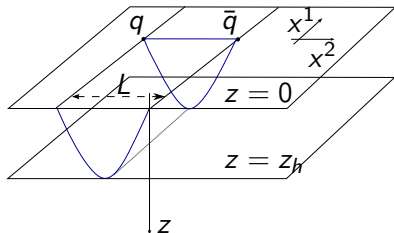
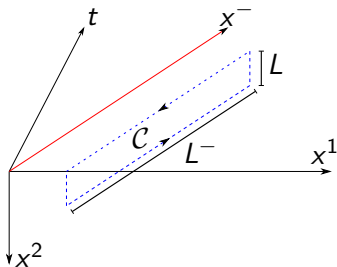
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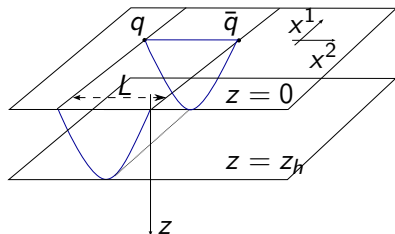
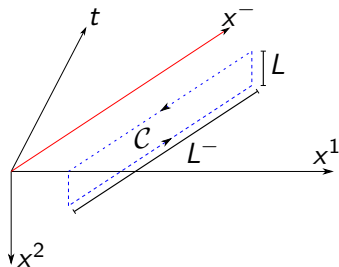
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- Applying dictionary:

$$\langle W(C) \rangle = \exp \left(-2S_{\text{reg}}^{\text{on-shell}} \right)$$

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- In general metric ansatz

$$ds^2 = e^{2A(z)} (-h(z)dt^2 + d\vec{x}^2) + \frac{e^{2B(z)}}{h(z)}dz^2$$

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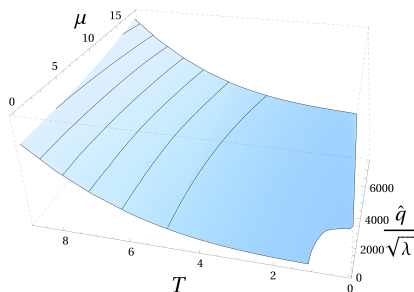
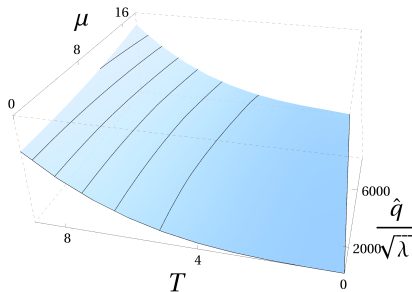
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- At $\mu = 0$ and no deformation [Liu, Rajagopal, Wiedemann]:

$$\hat{q} = \sqrt{\lambda} \frac{\pi^{3/2} \Gamma(3/4)}{\Gamma(5/4)} T^3 \approx \sqrt{\lambda} \cdot 7.528 T^3$$

\hat{q} in Deformed Models

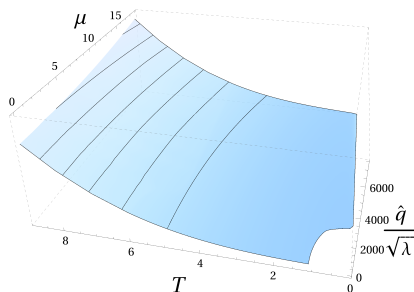
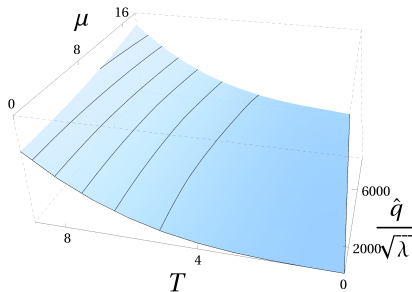
- In all our models \hat{q} increases with growing T and μ



Examples: Behaviour of \hat{q} in $\mathcal{N} = 4$ (left) and in 1-parameter model at $\kappa = 10$ (right)

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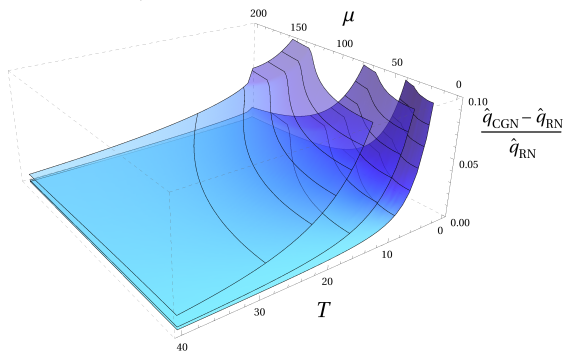
Examples: Behaviour of \hat{q} in $\mathcal{N} = 4$ (left) and in 1-parameter model at $\kappa = 10$ (right)

- Agrees with intuition of how plasma affects jets.

\hat{q} in Deformed Models

- Deformation unimportant at higher μ/T .

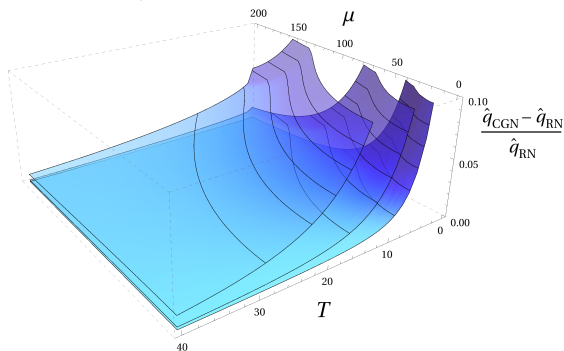
Relative deviation from $\mathcal{N} = 4$ in CGN model at different deformations ($c = 5, 10, 20$):



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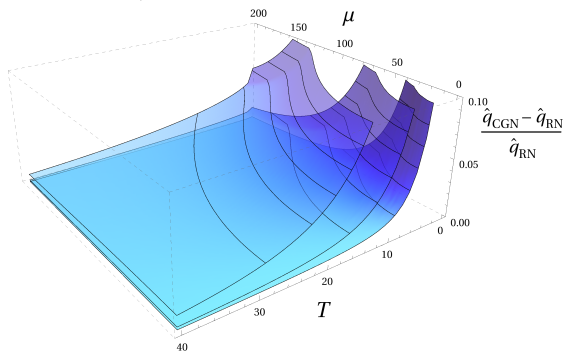


- \hat{q} is a very 'robust' observable.

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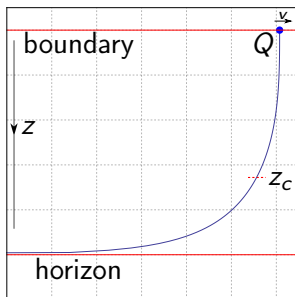
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- \hat{q} is a very 'robust' observable.
- Apply calculation to other (measurable) quantities, e.g. R_{AA} (?).

Moving Quark

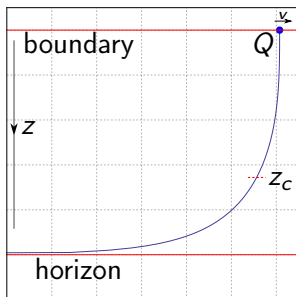
Linear motion [Gubser]:



$$\frac{dE}{dt} = -\frac{1}{2\pi\alpha'} e^{2A(z_c)} v^2$$

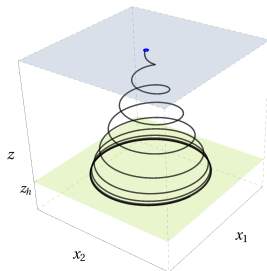
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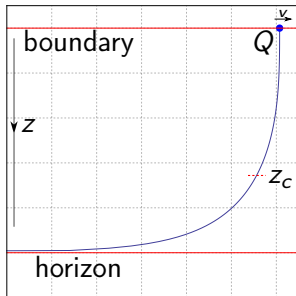
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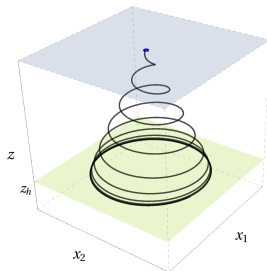
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Drag force shows robust behaviour.

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- Applied to QGP: Analyse jet quenching.
- In models we discussed: Observables are robust.
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Thank you for your attention!