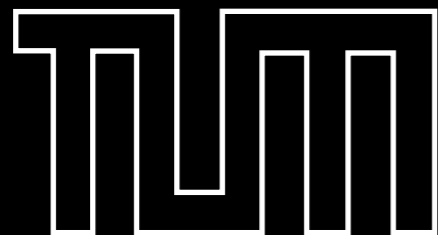


24th IMPRS Workshop
Munich

Dark Radiation from a hidden $U(1)$



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Outline

1. Implications of a hidden $U(1)$
2. Dark Radiation
3. Summary

Kinetic Mixing

- Lagrangian density for $U(1)_Y$

$$\mathcal{L} = -\frac{1}{4} F^{\mu\nu} F_{\mu\nu}$$



γ

Kinetic Mixing

- Lagrangian density for $U(1)_Y \times U(1)_h$

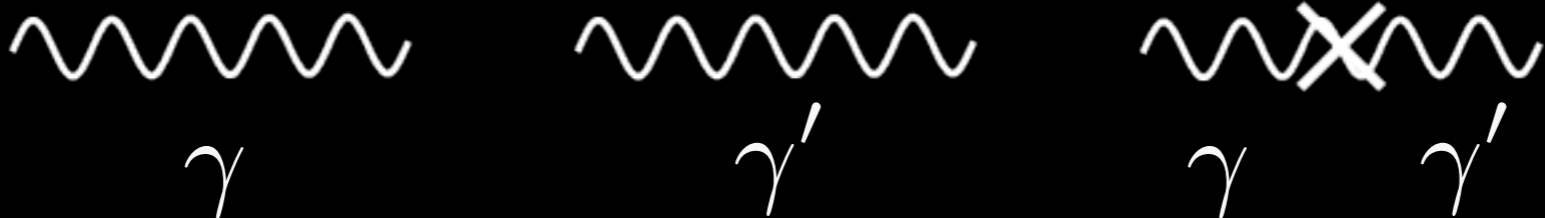
$$\mathcal{L} = -\frac{1}{4} F^{\mu\nu} F_{\mu\nu}$$



γ

Kinetic Mixing

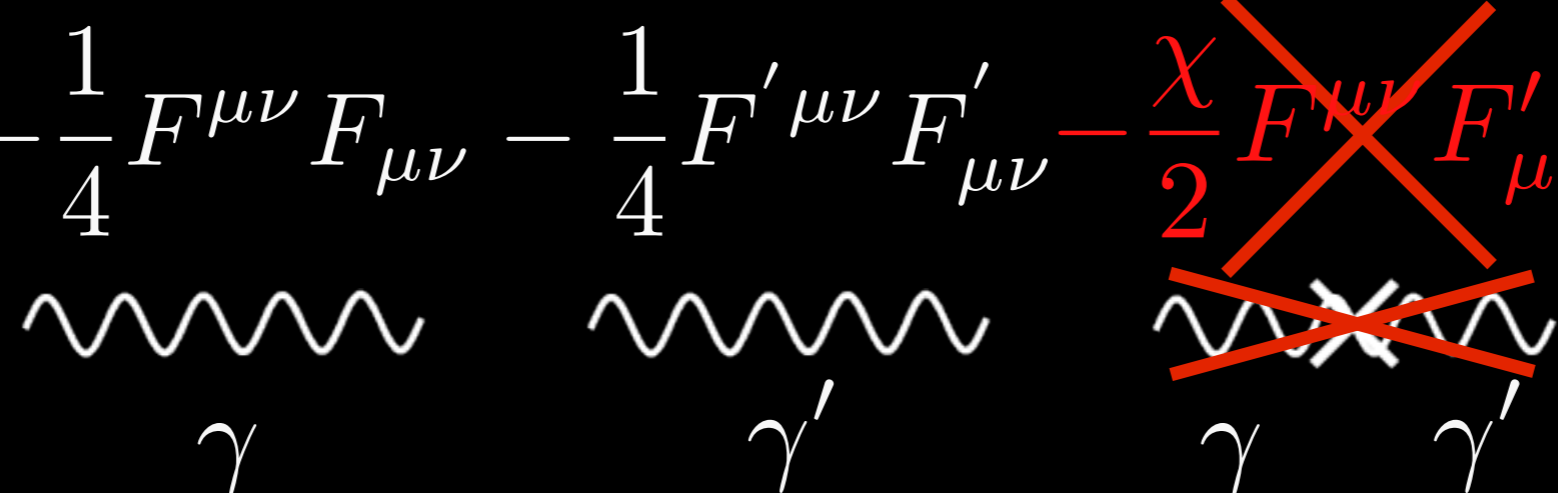
- Lagrangian density for $U(1)_Y \times U(1)_h$

$$\mathcal{L} = -\frac{1}{4}F^{\mu\nu}F_{\mu\nu} - \frac{1}{4}F'^{\mu\nu}F'_{\mu\nu} - \frac{\chi}{2}F^{\mu\nu}F'_{\mu\nu}$$


The diagram below the equation illustrates the terms. The first term, $-\frac{1}{4}F^{\mu\nu}F_{\mu\nu}$, is represented by a single wavy line labeled γ . The second term, $-\frac{1}{4}F'^{\mu\nu}F'_{\mu\nu}$, is represented by a single wavy line labeled γ' . The third term, $-\frac{\chi}{2}F^{\mu\nu}F'_{\mu\nu}$, is represented by two wavy lines, one labeled γ and one labeled γ' , with a large 'X' over them, indicating mixing between the two fields.

Kinetic Mixing

- Lagrangian density for $U(1)_Y \times U(1)_h$

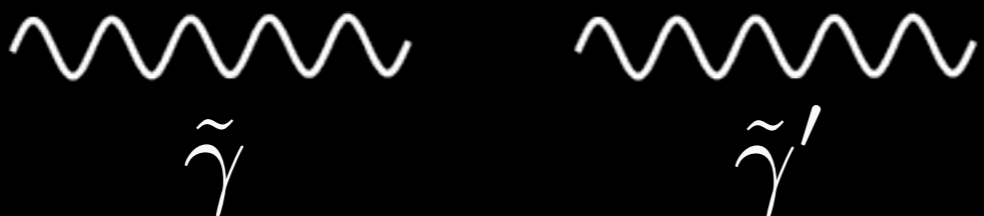
$$\mathcal{L} = -\frac{1}{4}F^{\mu\nu}F_{\mu\nu} - \frac{1}{4}F'^{\mu\nu}F'_{\mu\nu} - \frac{\chi}{2}F^{\mu\nu}F'_{\mu\nu}$$


The diagram illustrates the Lagrangian density with Feynman diagrams for each term. The first term, $-\frac{1}{4}F^{\mu\nu}F_{\mu\nu}$, is associated with a wavy line labeled γ . The second term, $-\frac{1}{4}F'^{\mu\nu}F'_{\mu\nu}$, is associated with a wavy line labeled γ' . The third term, $-\frac{\chi}{2}F^{\mu\nu}F'_{\mu\nu}$, is crossed out with a large red X and is associated with two wavy lines, one labeled γ and one labeled γ' .

- **Redefinition:** $A'_\mu \rightarrow A'_\mu - \chi A_\mu$
 $A_\mu \rightarrow \frac{1}{\sqrt{1-\chi^2}}A_\mu$

Kinetic Mixing

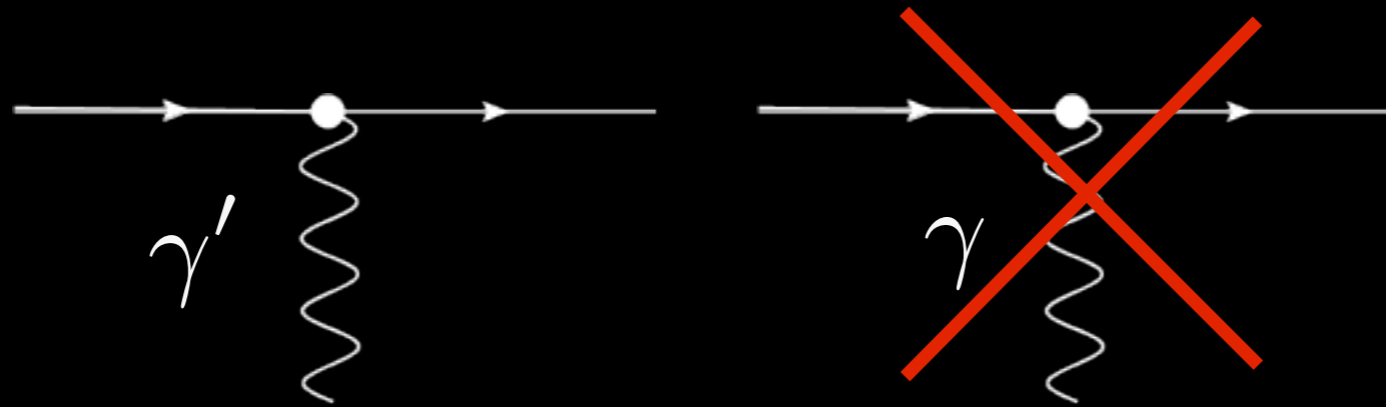
- Lagrangian density for $U(1)_Y \times U(1)_h$

$$\mathcal{L} = -\frac{1}{4}\tilde{F}^{\mu\nu}\tilde{F}_{\mu\nu} - \frac{1}{4}\tilde{F}'^{\mu\nu}\tilde{F}'_{\mu\nu}$$


- **Redefinition:** $A'_\mu \rightarrow A'_\mu - \chi A_\mu$
 $A_\mu \rightarrow \frac{1}{\sqrt{1-\chi^2}}A_\mu$

Kinetic Mixing (+ fermions)

- Dirac fermion Ψ

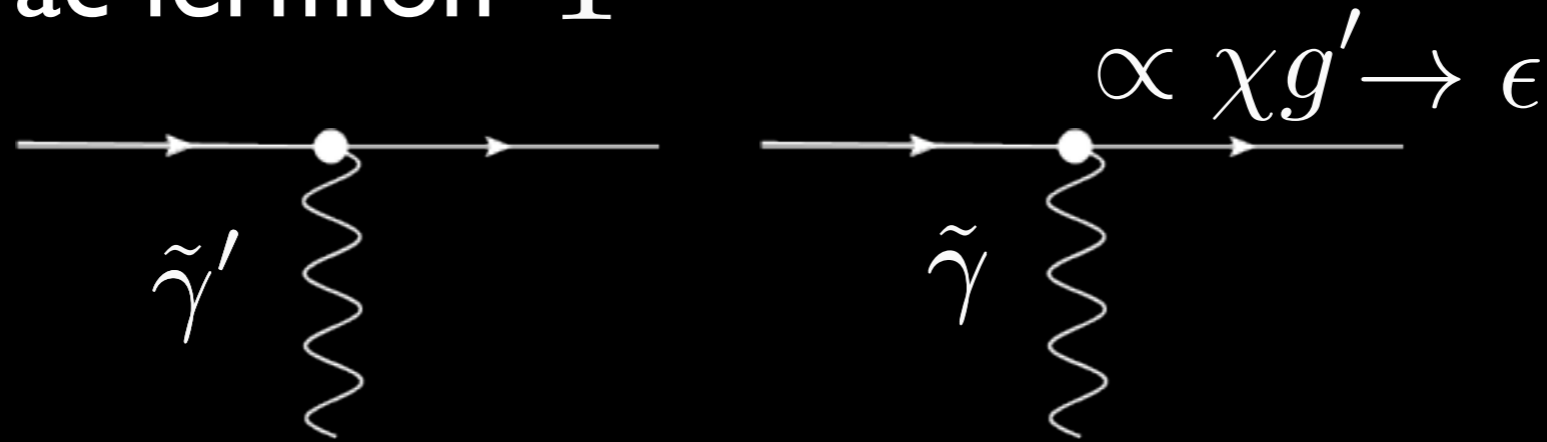


$$D_{\mu}\Psi = (\partial_{\mu} - ig' A'_{\mu})\Psi$$

Kinetic Mixing

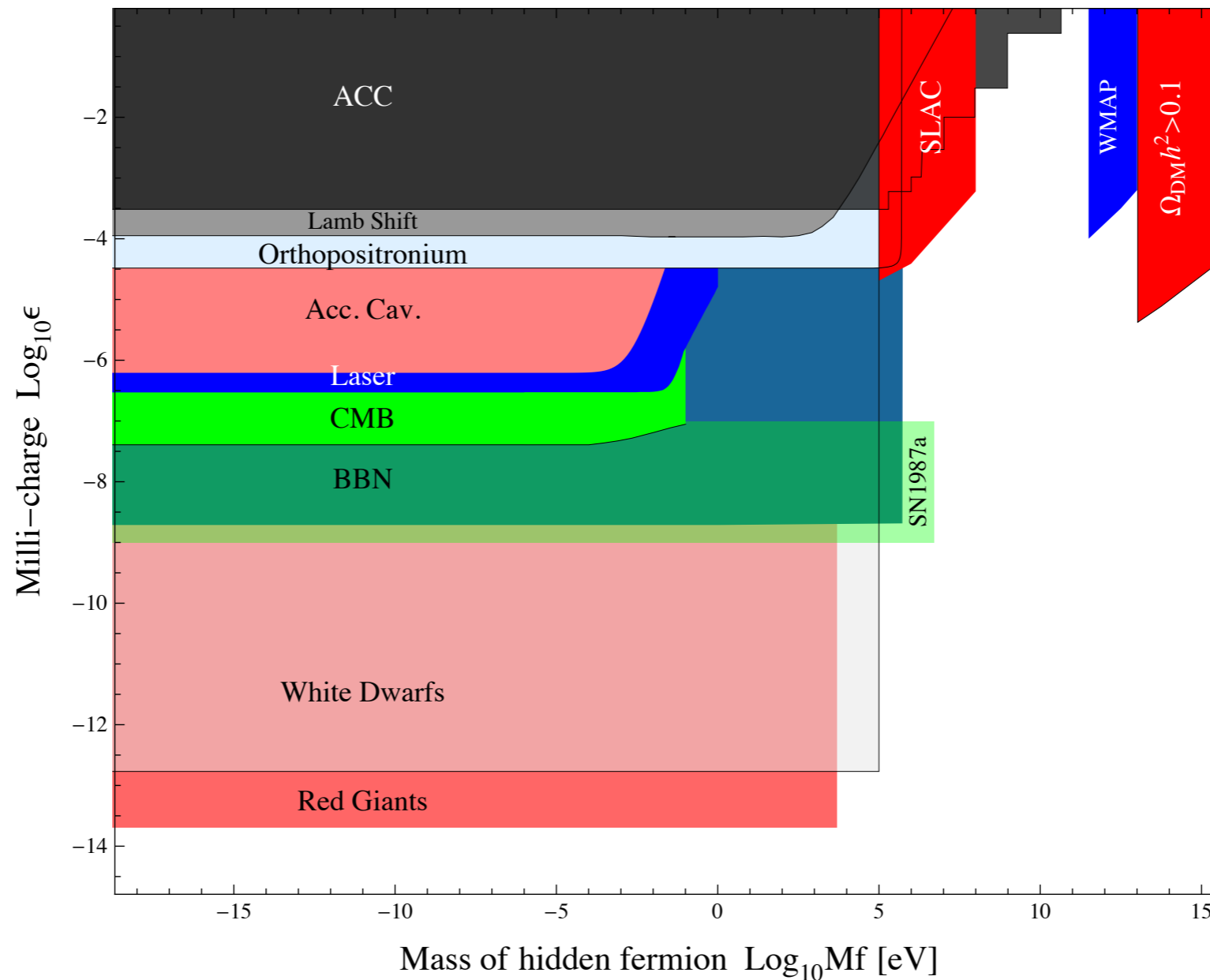
(+ fermions)

- Dirac fermion Ψ



$$D_{\mu}\Psi = \left(\partial_{\mu} - ig' \tilde{A}'_{\mu} - ig' \frac{\chi}{\sqrt{1 - \chi^2}} \tilde{A}_{\mu}\right)\Psi$$

Milli-Charged Particles

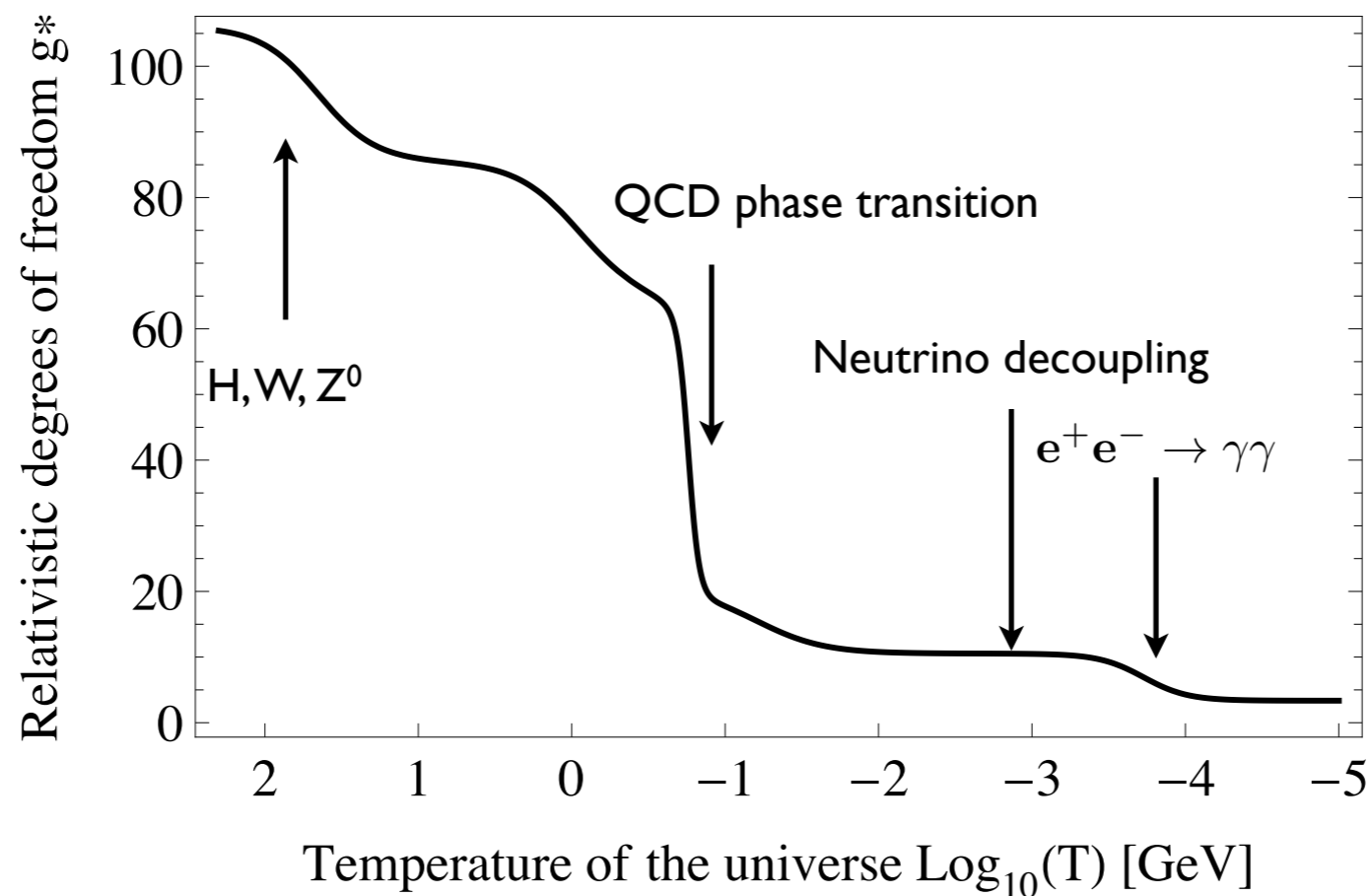


Davidson, Hannestad, Raffelt 2000, modified by Javier Redondo

2. Dark Radiation

Expansion of the Universe

Hubble parameter: $H \propto \sqrt{\rho} \propto g_{\star}^{1/2} T^2$



Fit by Wantz & Shellard (2009): 0910.1066

N_{eff}

Effective neutrino degrees of freedom N_{eff}

$$N_{\text{eff}} = 3$$

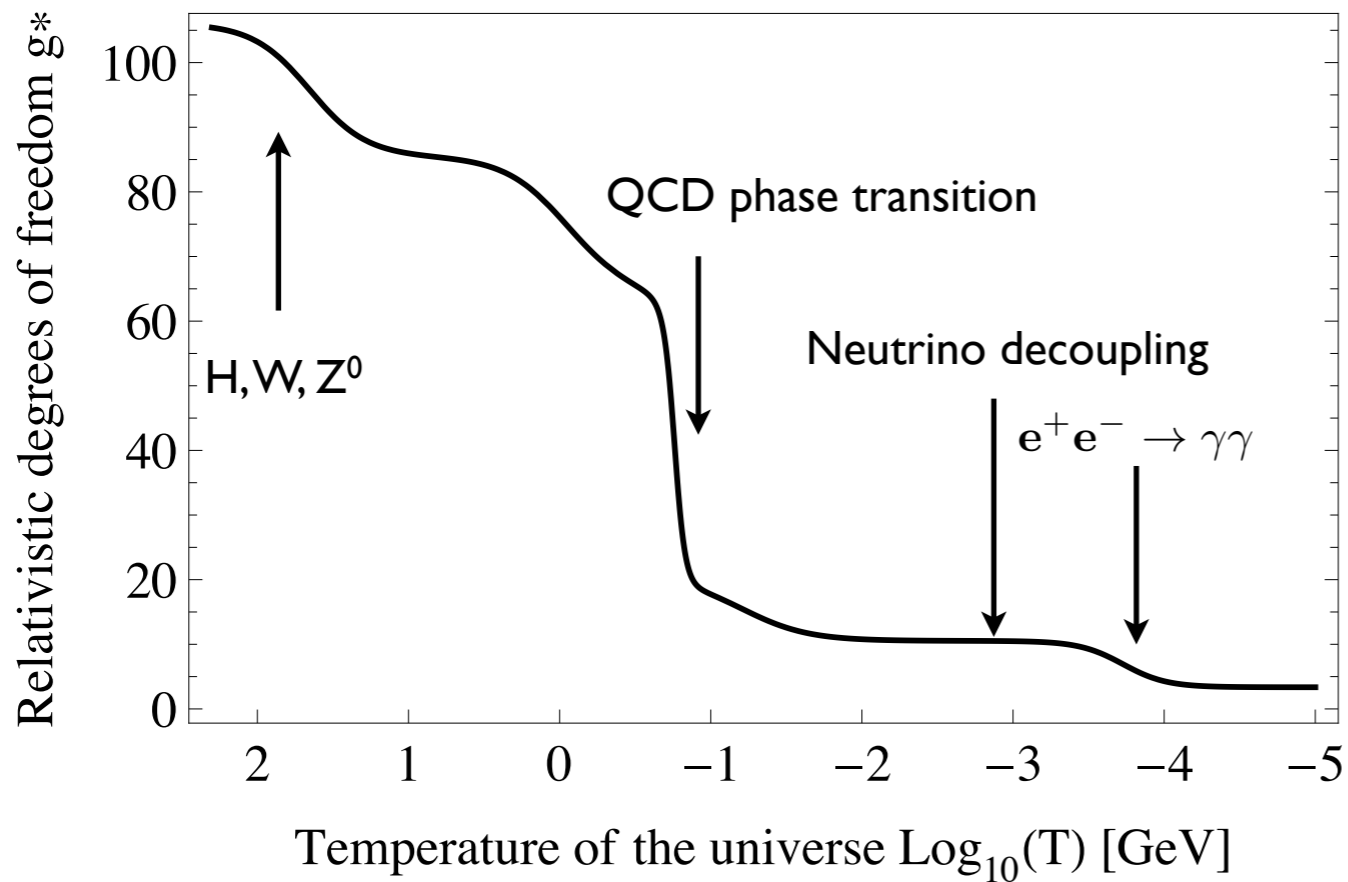
Standard Model

$$N_{\text{eff}} = 3.85 \pm 0.84 \text{ 95\% c.l.}$$

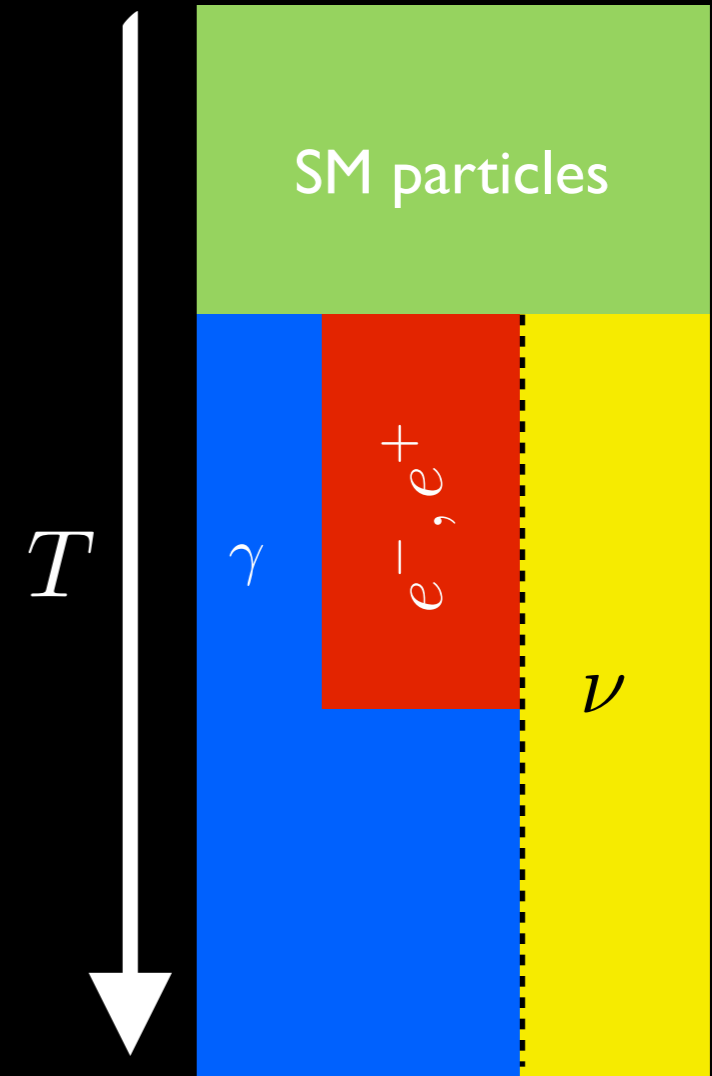
for CMB (Keisler et al. (2011): 1108.4136)

$$\Delta N_{\text{eff}} = \frac{\rho_R}{\rho_\nu} \propto \left(\frac{T_R}{T_\nu} \right)^4 = \left(\frac{T_R}{\left(\frac{4}{11} \right)^{\frac{1}{3}} T_\gamma} \right)^4$$

Neutrino Temperature

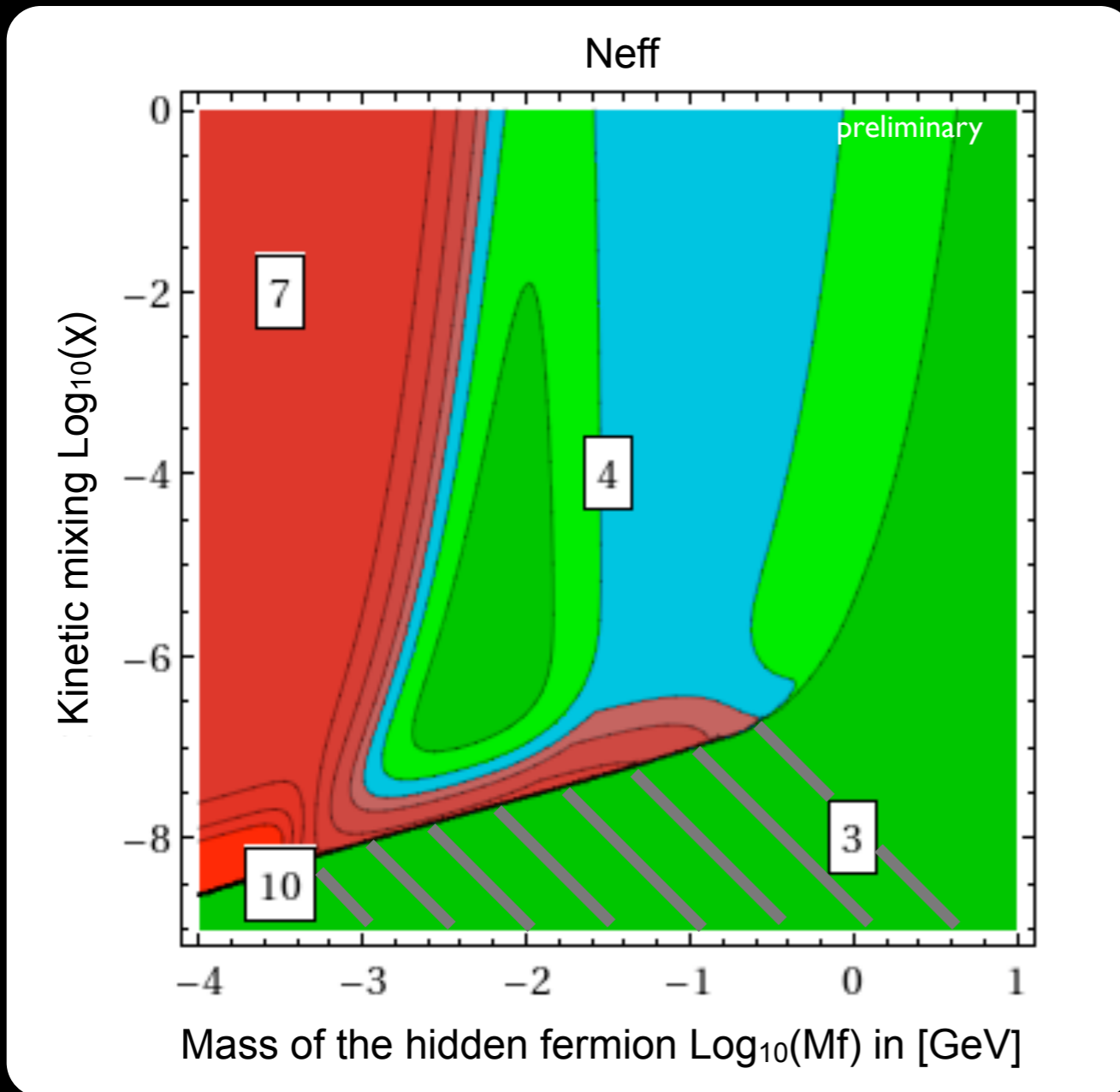


Fit by Wantz & Shellard (2009): 0910.1066

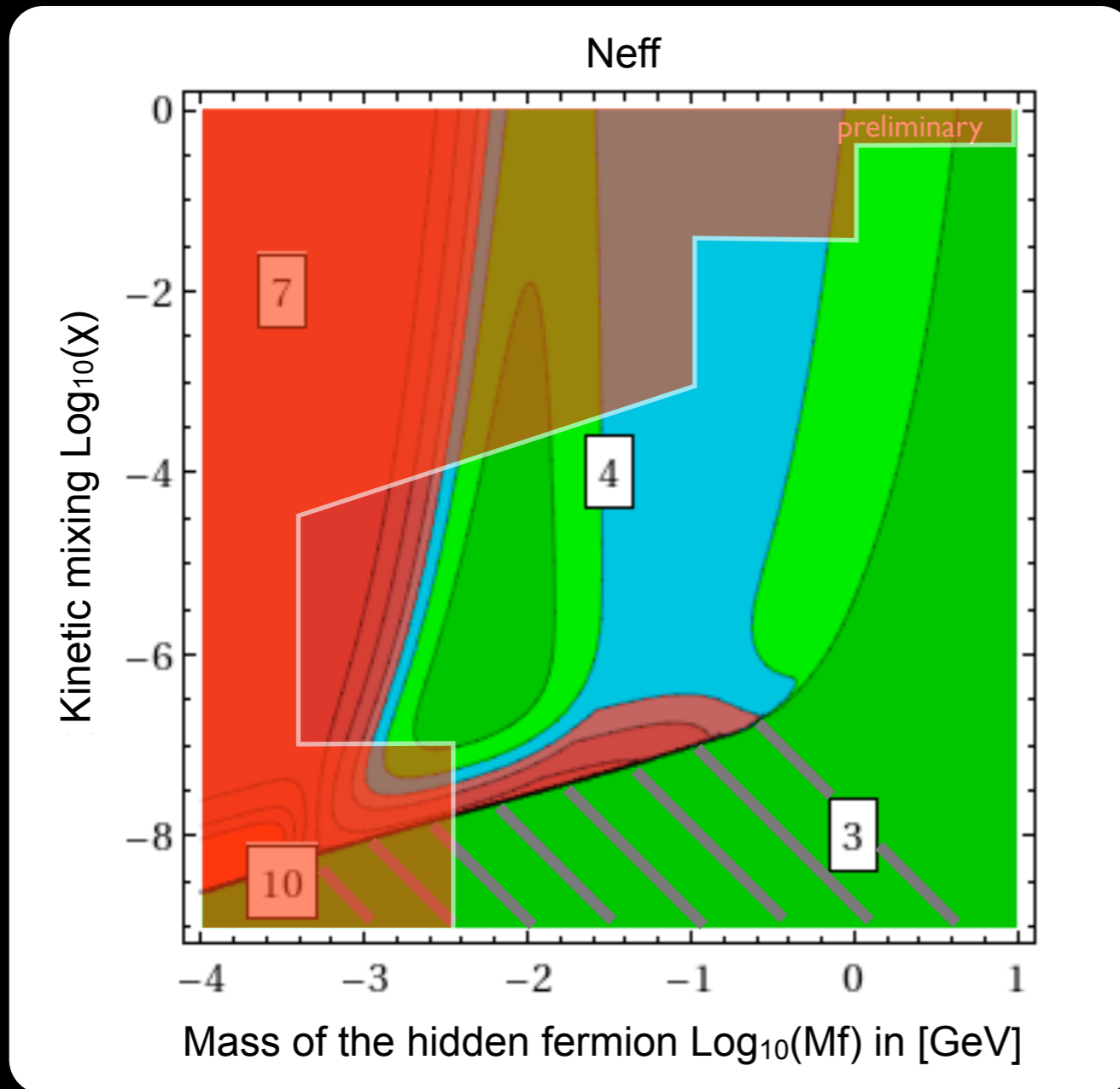


$$\frac{T_\nu}{T_\gamma} = \left(\frac{4}{11}\right)^{\frac{1}{3}}$$

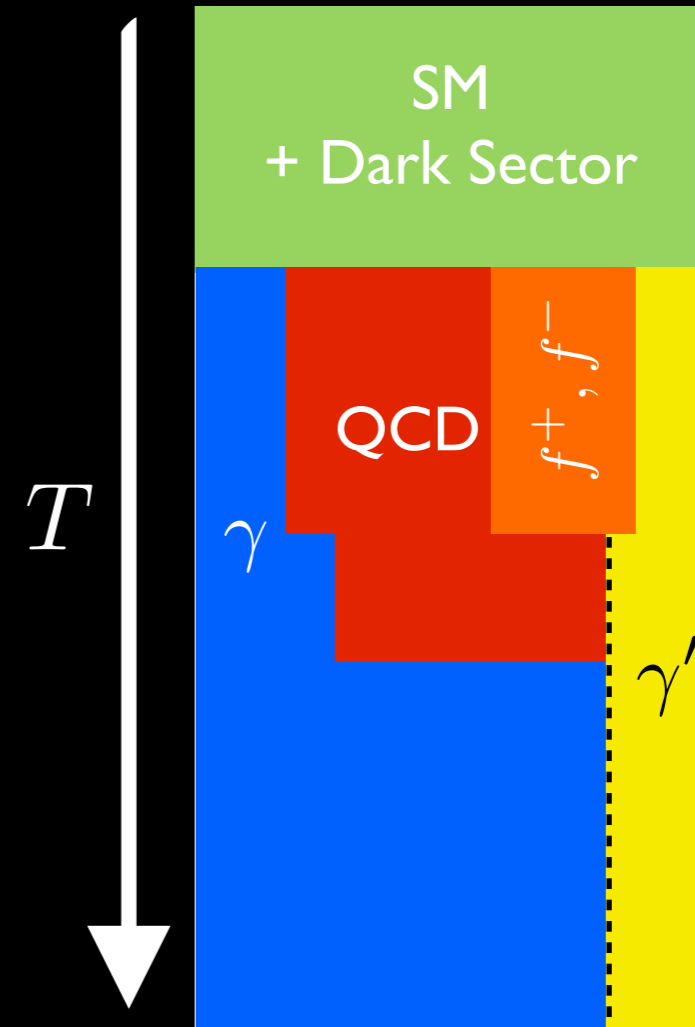
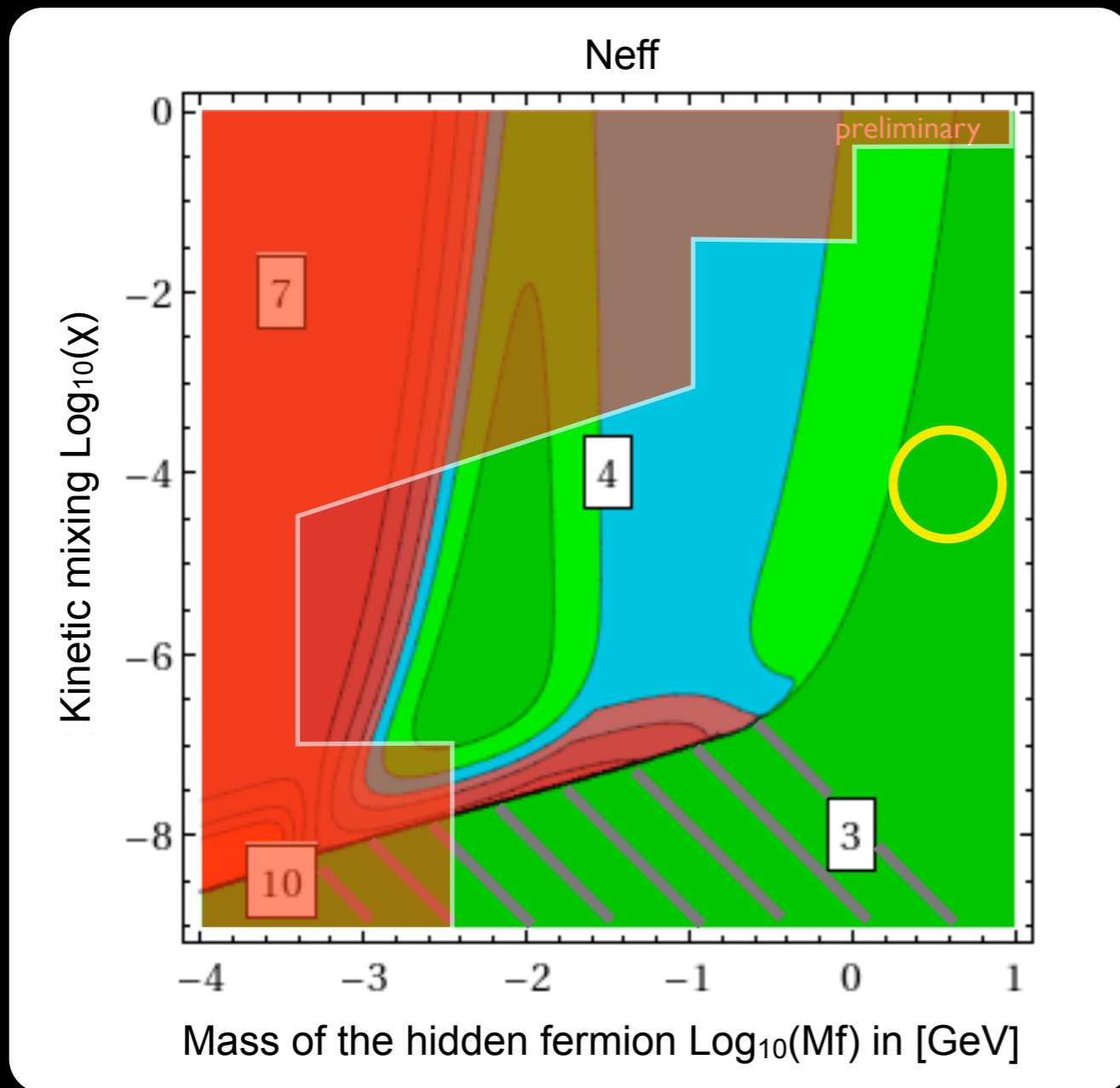
Results



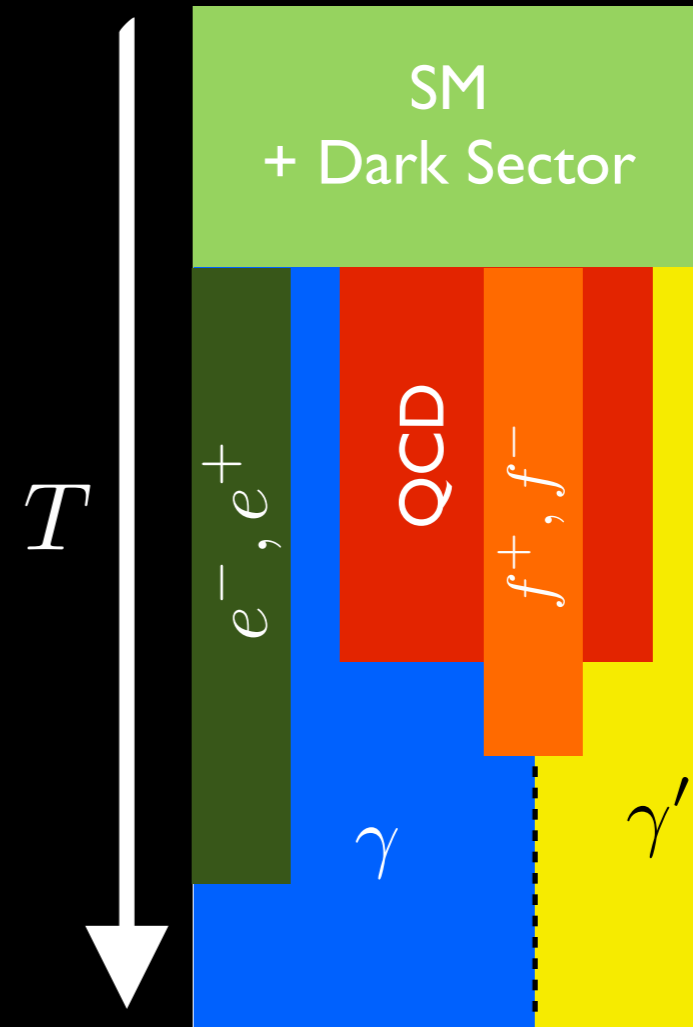
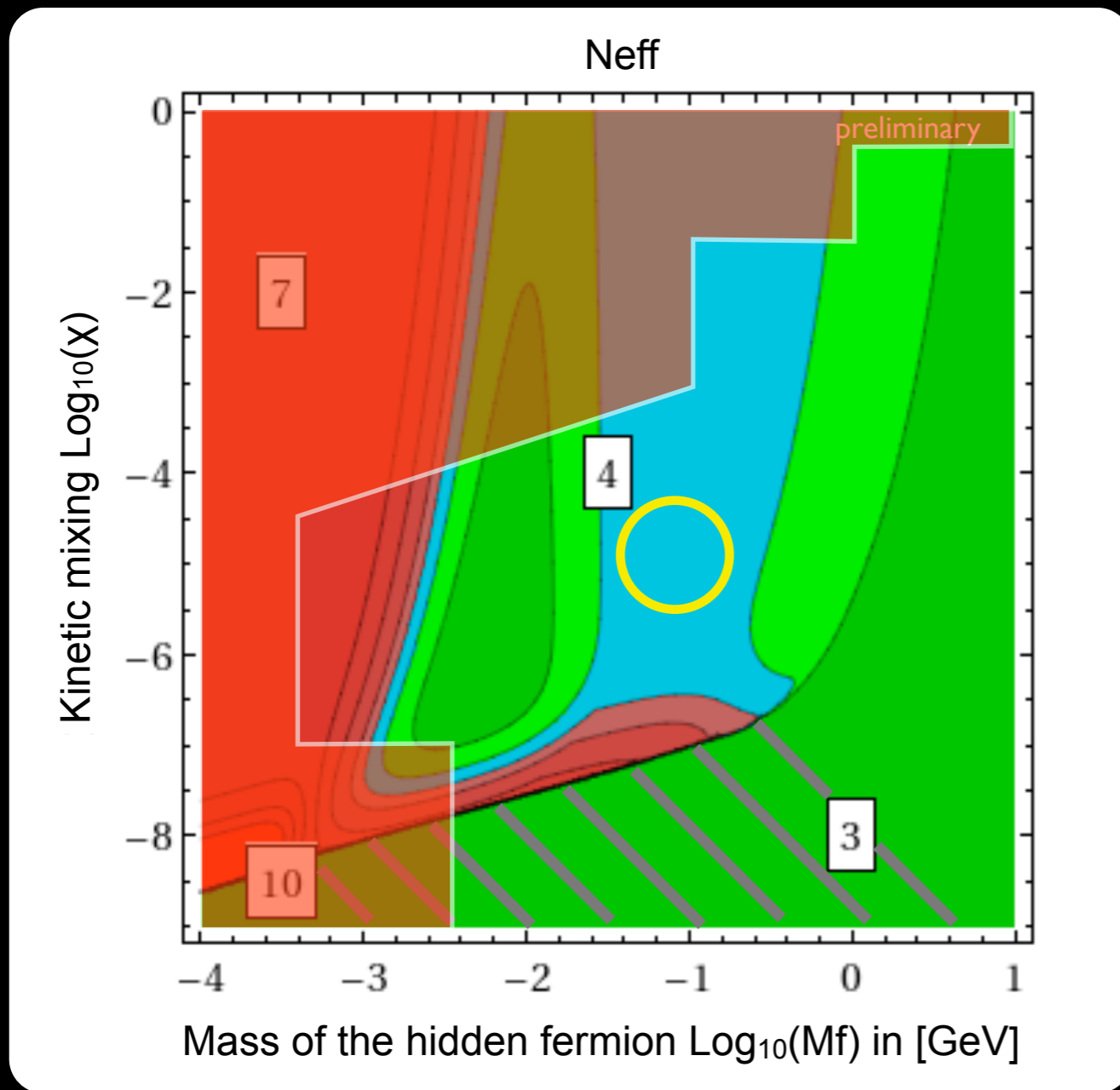
Results



Results



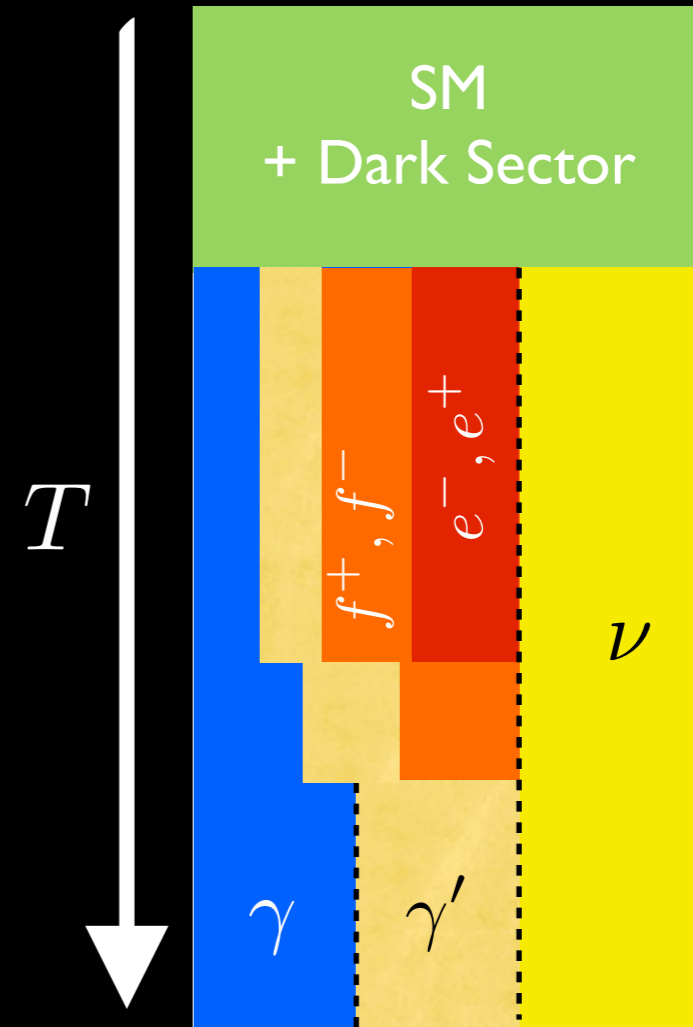
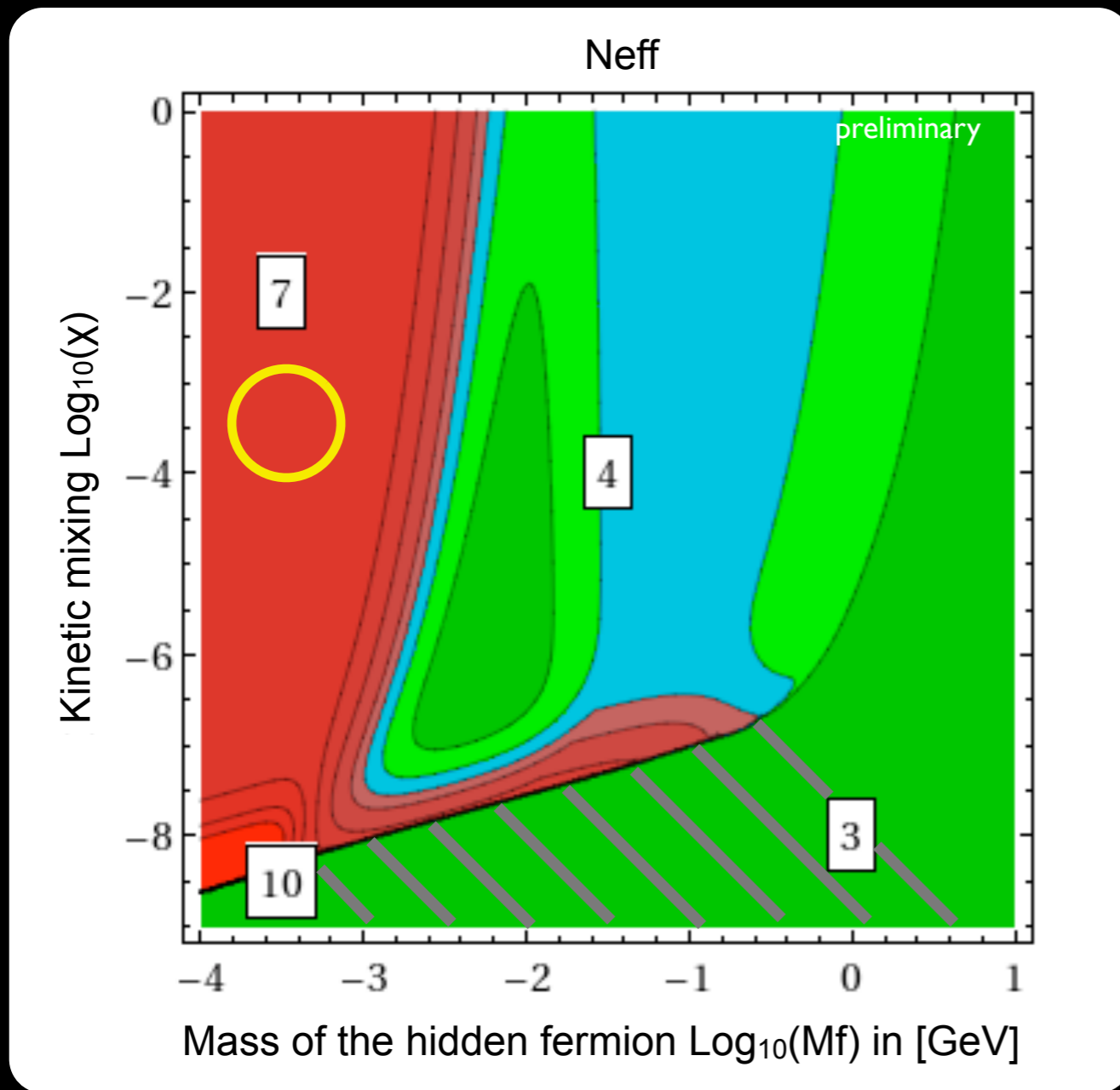
Results



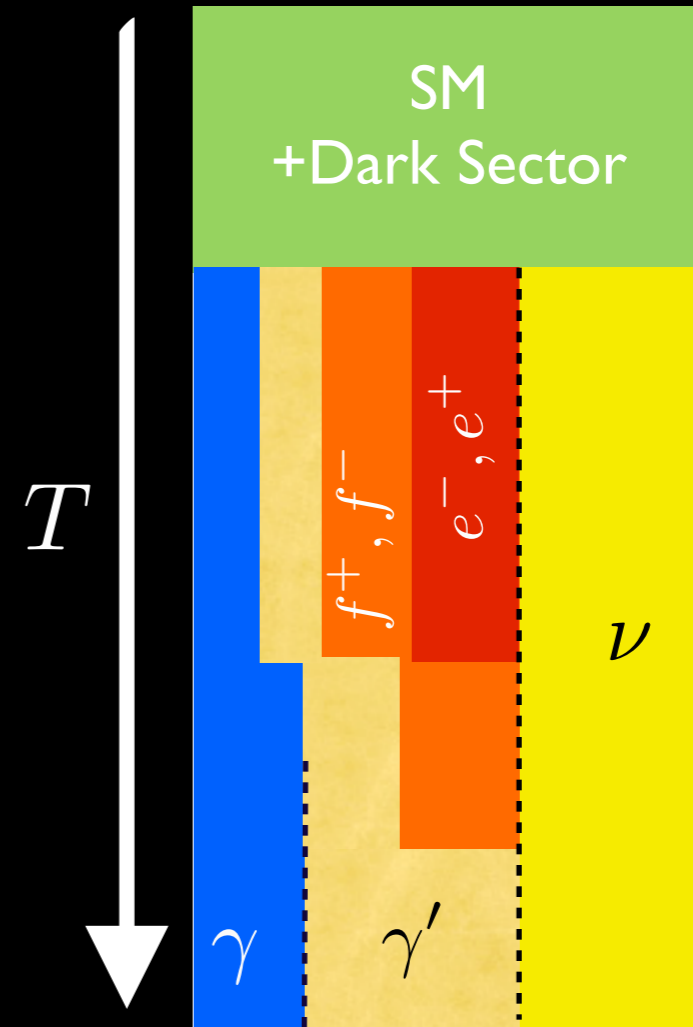
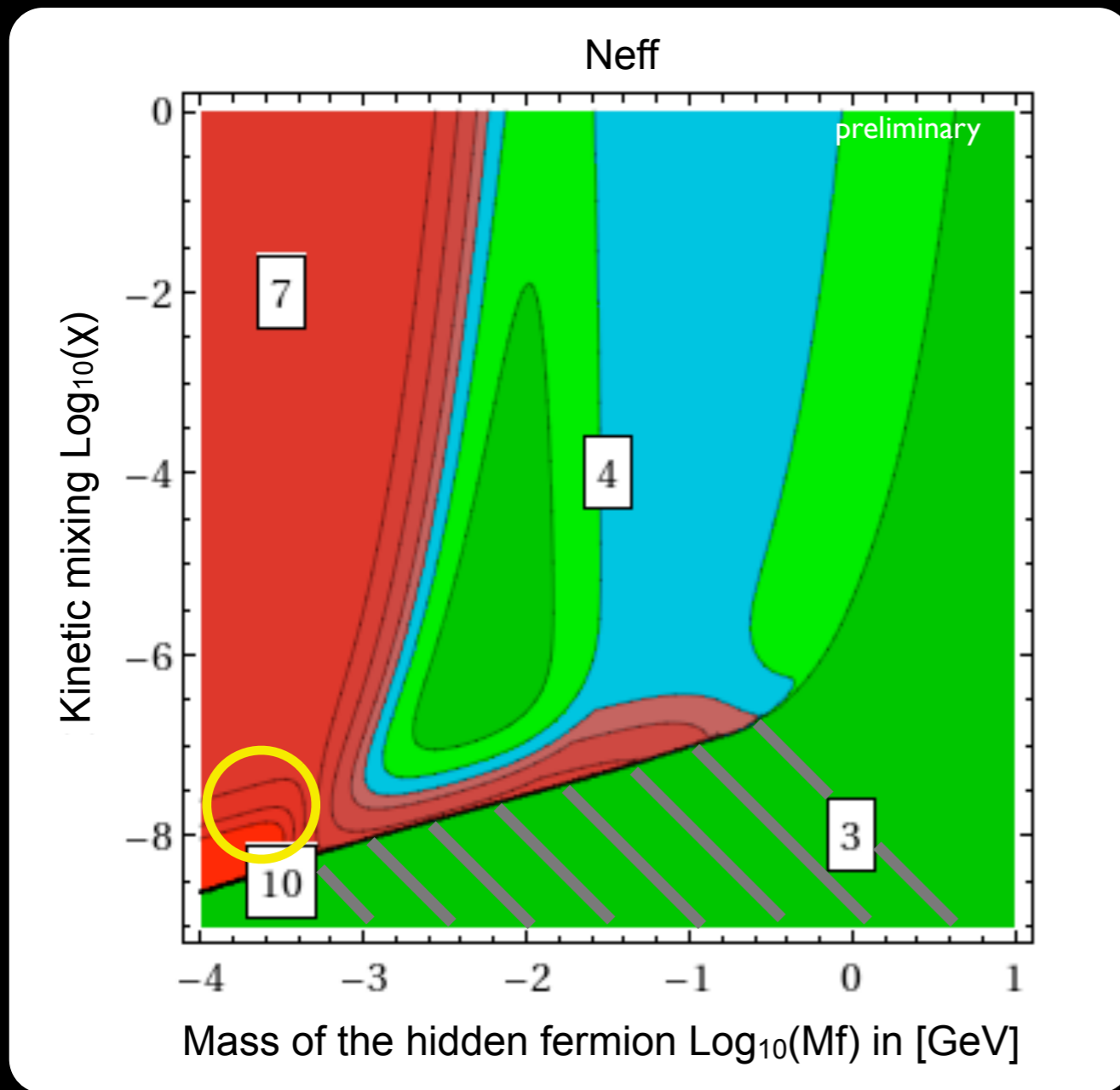
Summary

- Photons of an additional **unbroken U(1)** alone are **unobservable**
- Adding fermions leads to **milli-charged particles**
- Coupling these fermions to the Standard Model accelerates expansion of the universe (**Dark Radiation**)
- **Observations** favor an additional radiative component ($N_{\text{eff}} > 3$)
- This **can be explained** with an **additional photon and an additional fermion**

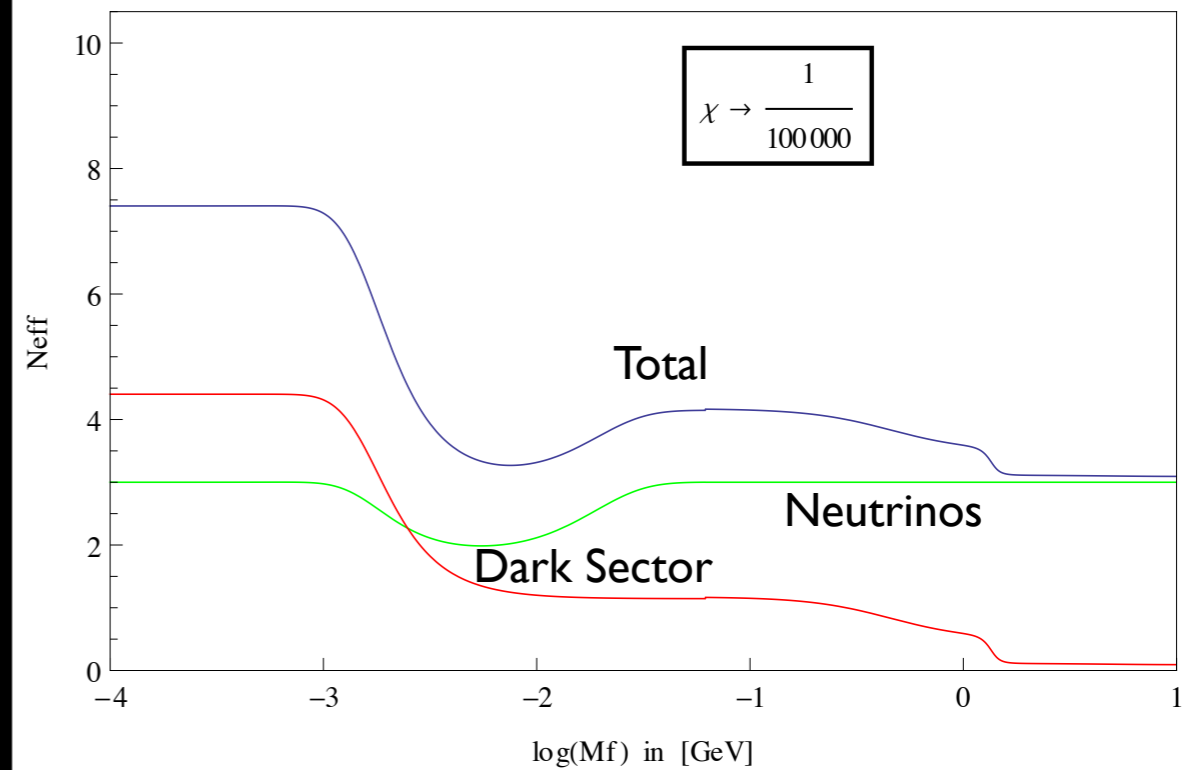
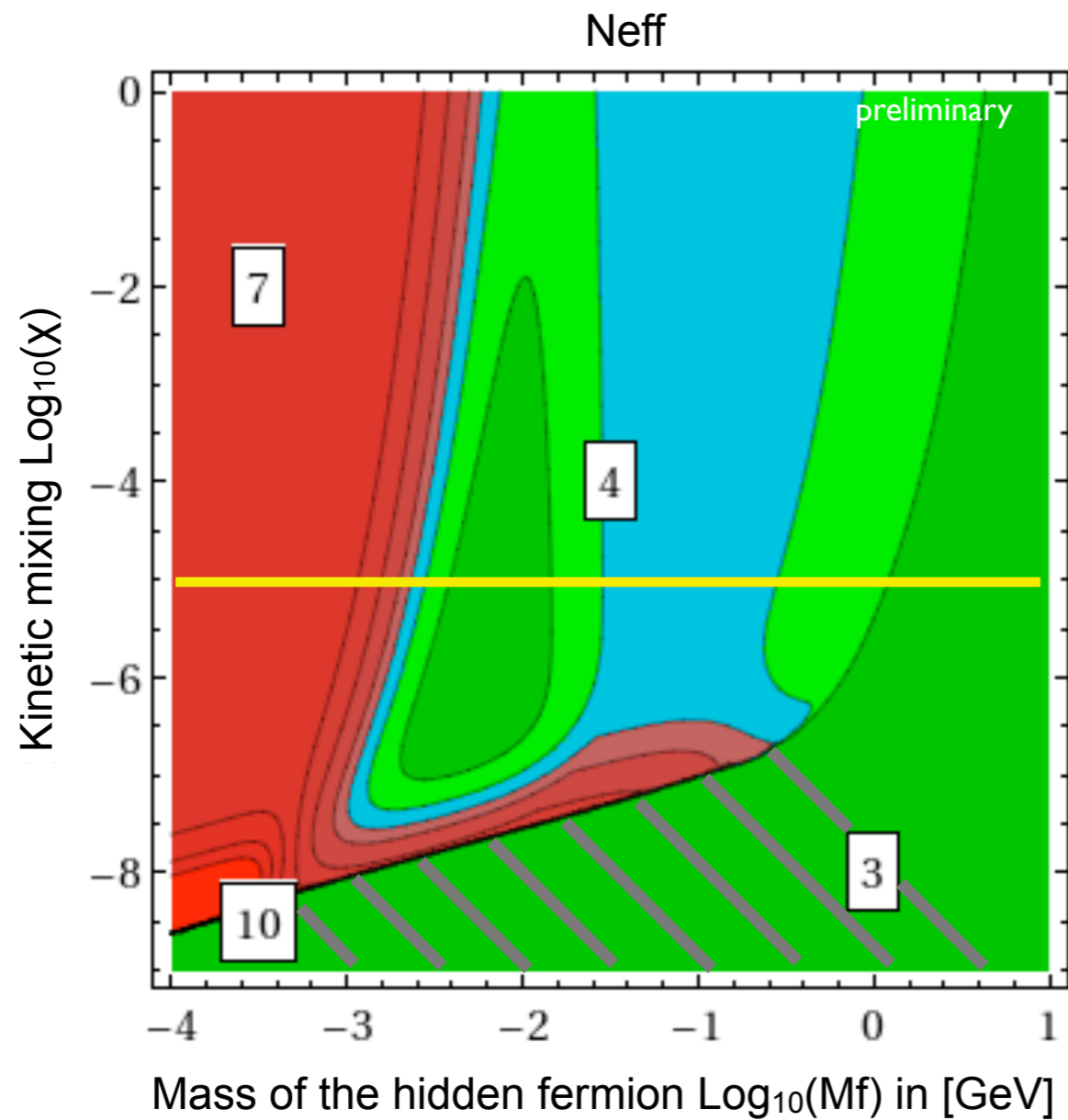
Results



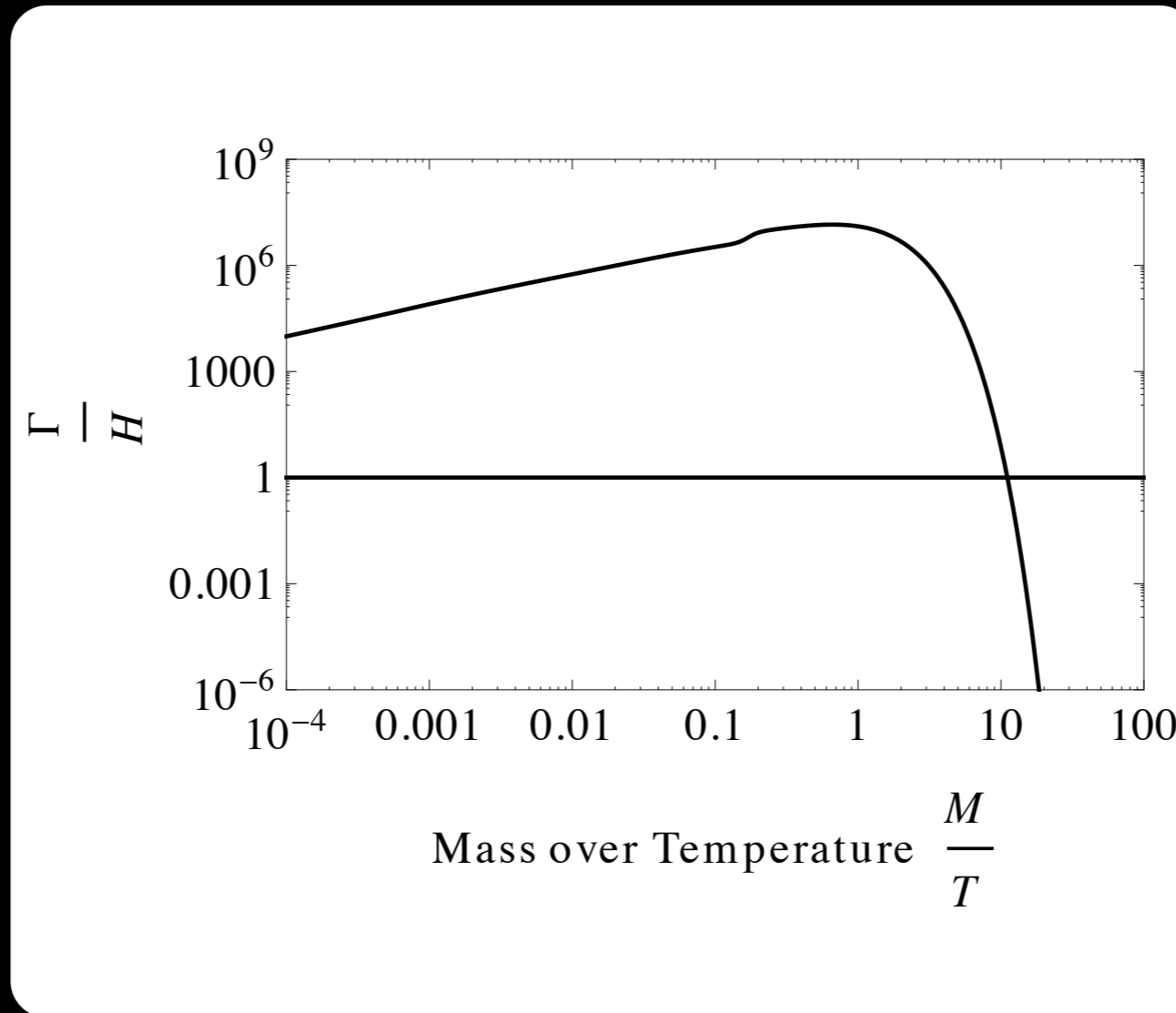
Results



Results



Decoupling



$$\Gamma = \langle \sigma v \rangle n_f$$