

# Implications of the Higgs discovery

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- **The Higgs in the Standard Model and beyond**
  - **The Higgs at the LHC**
  - **First implications of the discovery**
    - **What next?**

# 1. The Higgs in the Standard Model and beyond

We have a theory, the Standard Model, which describes microscopic world.

the interaction of  $s = \frac{1}{2}$  matter particles via exchange of  $s = 1$  force particles.

It is based on a gauge symmetry:

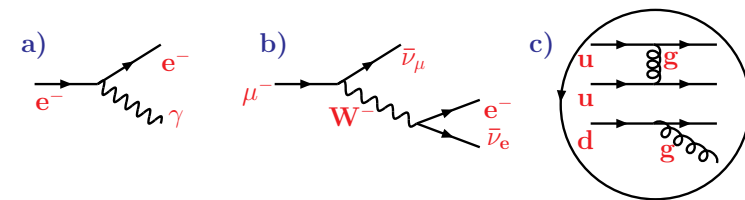
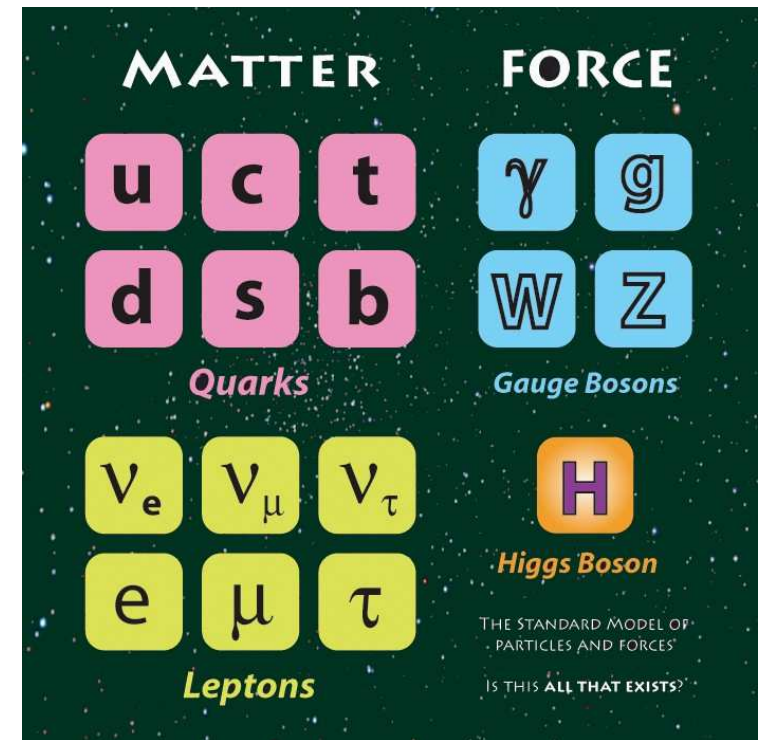
$$SU(3)_C \times SU(2)_L \times U(1)_Y$$

- relativistic quantum field theory,
- perturbative, renormalisable,
- and most of all, very successful:
  - ⇒ infinitely precise predictions,
  - ⇒ high precision experimental tests.

But true only if particles are massless<sup>a</sup>: putting naively masses for W/Z/fermions spoils gauge invariance and therefore the nice properties of the theory above.

Problem: how to generate particle masses in a gauge invariant way?

⇒ the Brout–Englert–Higgs mechanism for EW symmetry breaking!



<sup>a</sup>This has nothing to do with mass of macroscopic objects due to binding energy..

# 1. The Higgs in the Standard Model and beyond

Introduce a doublet of complex scalar fields  $\Phi = \begin{pmatrix} \Phi^+ \\ \Phi^0 \end{pmatrix}$ : 4 degrees of freedom.

Scalar potential:  $V_S = \mu^2 \Phi^\dagger \Phi + \lambda (\Phi^\dagger \Phi)^2$ ,  $SU(2)_L \times U(1)_Y$  invariant.

$\mu^2 > 0$ : minimum of  $V_S$  at  $\langle 0 | \Phi^0 | 0 \rangle = 0$ :

4 new scalar particles with mass  $m_S = \mu$ .

$\mu^2 < 0$ : (via quantum fluctuations?):

the field  $\Phi$  develops a non-zero vev

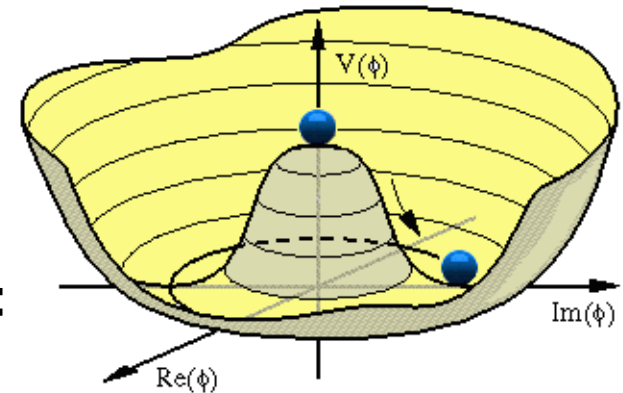
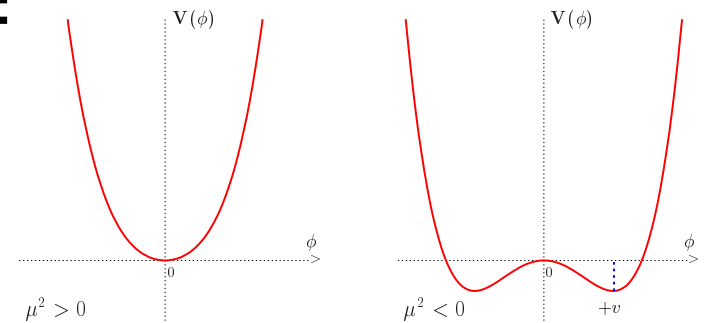
$$\langle 0 | \Phi^0 | 0 \rangle = v = \sqrt{\frac{-\mu^2}{\lambda^2}} (= 246 \text{ GeV})$$

fields/interactions still  $SU(2) \times U(1)$  symmetric  
but vacuum not  $\Rightarrow$  spontaneous EW breaking.

$\Rightarrow$  three d.o.f. for  $M_{W^\pm}$  and  $M_Z$ .

Introduce interaction of fermions with same  $\Phi$ :

fermions masses  $m_f$  also generated!



Residual d.o.f corresponds to spin-0 Higgs particle.

- Unique particle: spin zero, not matter particle and not force particle,
- couples to all particles  $\propto$  their masses:  $g_{Hff} \propto m_f$ ,  $g_{HVV} \propto M_V$ ,
- couples to itself,  $g_{HHH} \propto M_H^2$  with the relation  $M_H^2 = 2\lambda v^2$ .

# 1. The Higgs in the Standard Model and beyond

Since  $v$  is known, the only free parameter in the SM is  $M_H$  (or  $\lambda$ ).

Pre-LHC constraints on  $M_H$ :

- **Experimental constraints:**

- direct searches at LEP/Tevatron:

$M_H > 114 \text{ GeV}$ ,  $M_H \neq 160 \text{ GeV}$

- quantum effects in EW data:

$M_H < 160 \text{ GeV}$  @95% confidence.

- **The Higgs unitarizes the theory:**

without H:  $|A_0(VV \rightarrow VV)| \propto E^2$

including H:  $|A_0| \propto M_H^2/v^2$

theory unitary but  $M_H \lesssim 700 \text{ GeV}$ ...

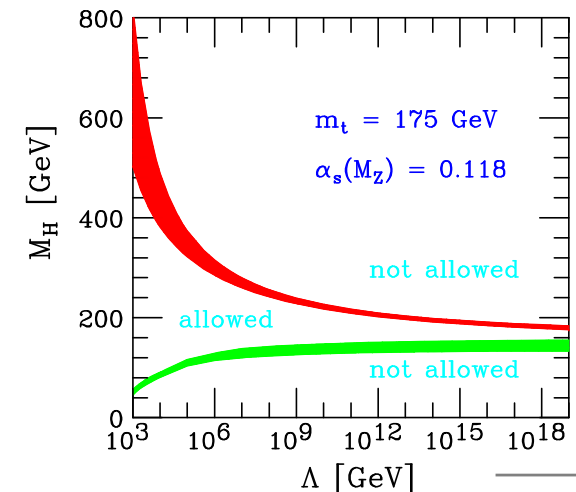
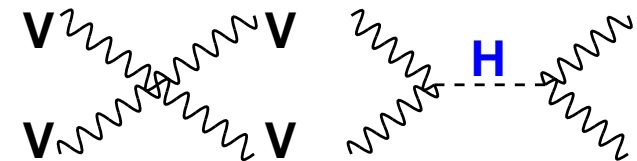
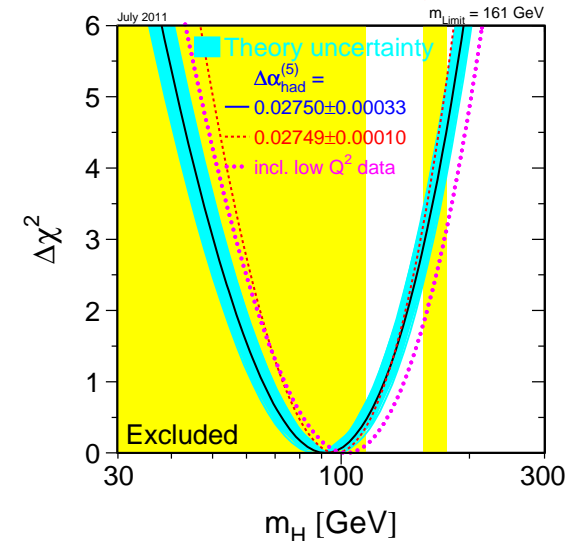
- **Triviality and stability bounds:**

coupling evolves with energy  $\lambda \equiv \lambda(Q^2)$

$\lambda \gg 1$ : becomes infinite (no perturbation)

$\lambda \ll 1$ : potential unstable (no EWSB)

$\Lambda \sim M_{\text{Pl}} : 120 \lesssim M_H \lesssim 180 \text{ GeV}!$

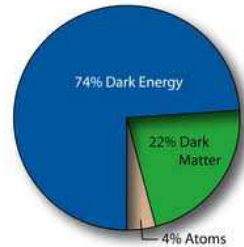
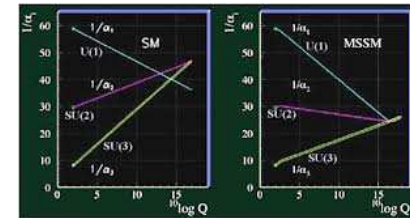




# 1. The Higgs in the Standard Model and beyond

## 3) Supersymmetry: doubling the world.

- SUSY = most attractive SM extension:
  - links  $s = \frac{1}{2}$  fermions to  $s = 1$  bosons
  - links internal and space-time symmetry
  - if made local, it provides link to gravity
  - naturally present in string theory (TOE?)
  - natural  $\mu^2 < 0$ : radiative EWSB
  - fixes gauge coupling unification
  - ideal candidate for Dark Matter...

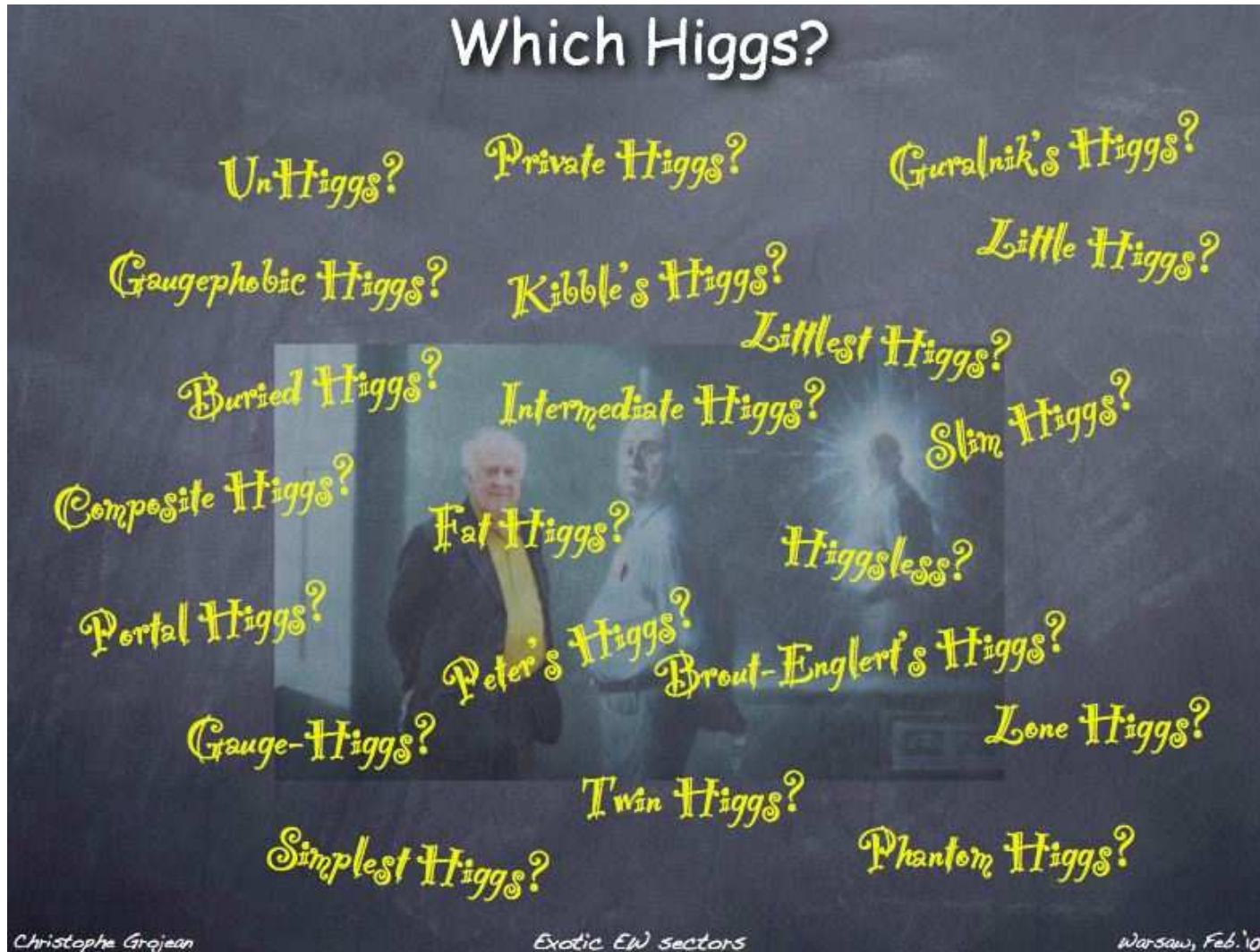


- Needs two doublets  $\Phi_1, \Phi_2$  for EWSB:
  - ⇒ extended Higgs sector:  $h, H, A, H^\pm$  with  $h \oplus H \approx H_{SM}$
  - SUSY ⇒ only two inputs at tree level:  $\tan\beta = v_2/v_1, M_A$
  - SUSY ⇒ hierarchy spectrum:  $M_h \approx 100 \text{ GeV}, M_H \approx M_A \approx M_{H^\pm}$  (SUSY scale  $M_{SUSY}$  pushes via radiative corrections  $M_h$  to 130 GeV).
  - Most often decoupling regime:  $h \equiv H_{SM}$ , others decouple from W/Z.



# 1. The Higgs in the Standard Model and beyond

and along the avenues, many possible streets, paths, corners...



**Which scenario chosen by Nature? The LHC was devised to tell!**

## 2. The Higgs at the LHC

**Physics at hadron machines is a nightmare...**

- Protons non–elementary: difficult environment
- Huge cross sections for QCD processes
- Small cross sections for EW Higgs signal

**$S/B \gtrsim 10^{10} \Rightarrow$  a needle in a haystack!**

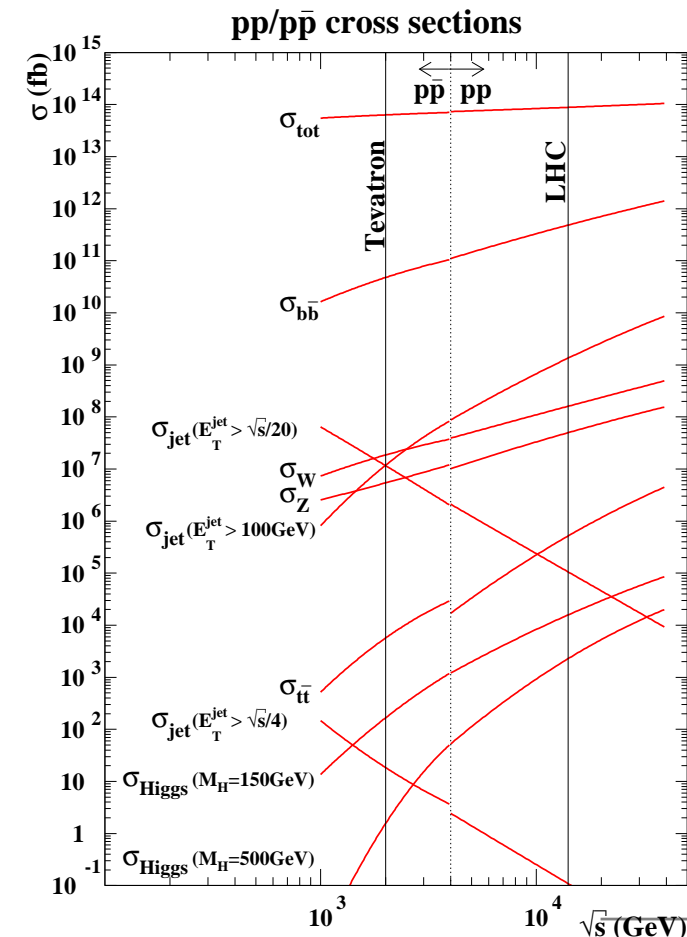
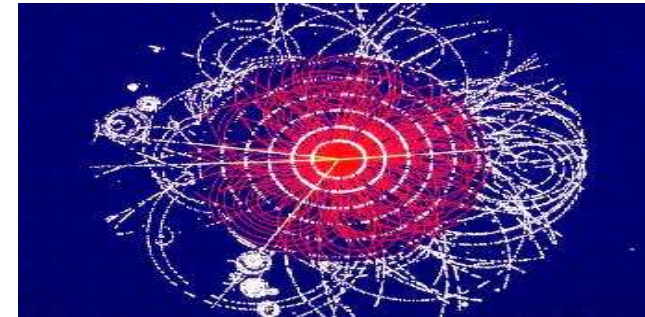
- Need some strong selection criteria:
  - trigger: get rid of uninteresting events...
  - select clean channels: leptons and photons
  - use specific kinematic features of Higgs

● Combine # decay/production channels  
(and eventually several experiments...)

● Have a precise knowledge of S and B rates  
(higher/quantum effects can be factor of 2!)

● Gigantic experimental + theoretical efforts  
(more than 30 years of extremely hard work!)

**to make sure that the Higgs will not escape!**





## 2. The Higgs at the LHC

Since  $v$  is known, the only free parameter in the SM is  $M_H$  (or  $\lambda$ ).  
Once  $M_H$  known, all the properties of the Higgs boson are fixed.

### Example: Higgs decays in the SM

• As  $g_{HPP} \propto m_P$ , H will decay into heaviest particle phase-space allowed:

•  $M_H \lesssim 130 \text{ GeV}$  :

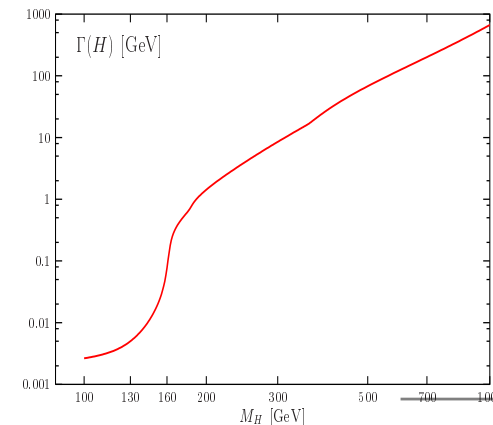
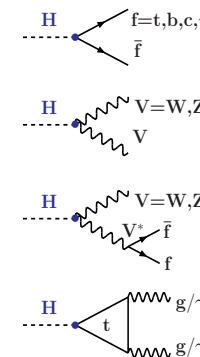
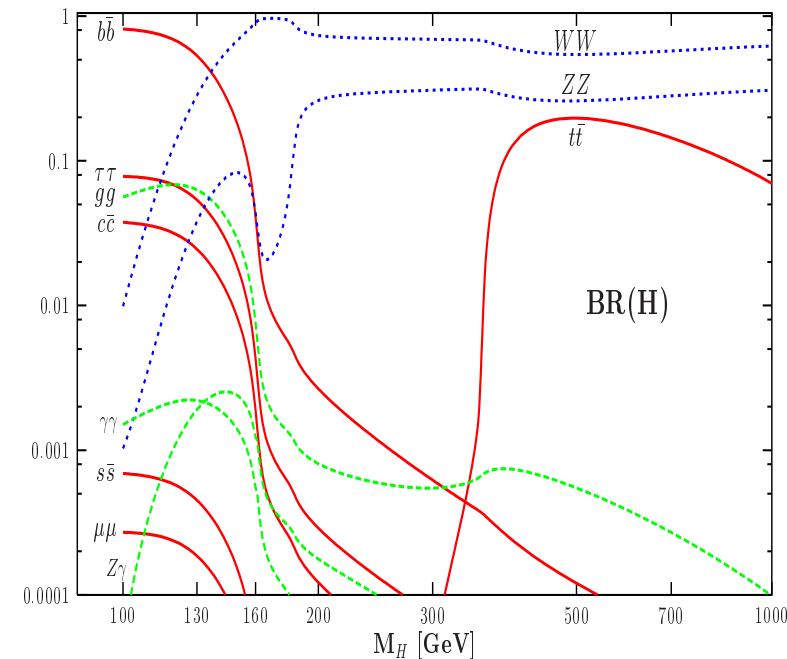
- $H \rightarrow b\bar{b}$ : dominant decay
- $H \rightarrow cc, \tau^+\tau^-, gg = \mathcal{O}(\text{few } \%)$
- $H \rightarrow \gamma\gamma, Z\gamma = \mathcal{O}(0.1\%)$

•  $M_H \gtrsim 130 \text{ GeV}$ :

- $H \rightarrow WW, ZZ$  dominant
- decays into  $t\bar{t}$  for heavy Higgs

• Total Higgs decay width:

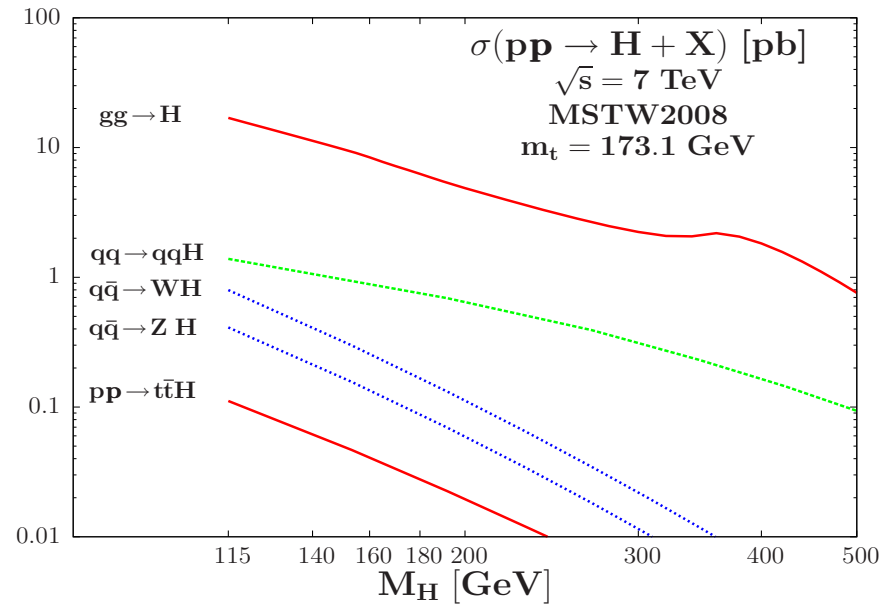
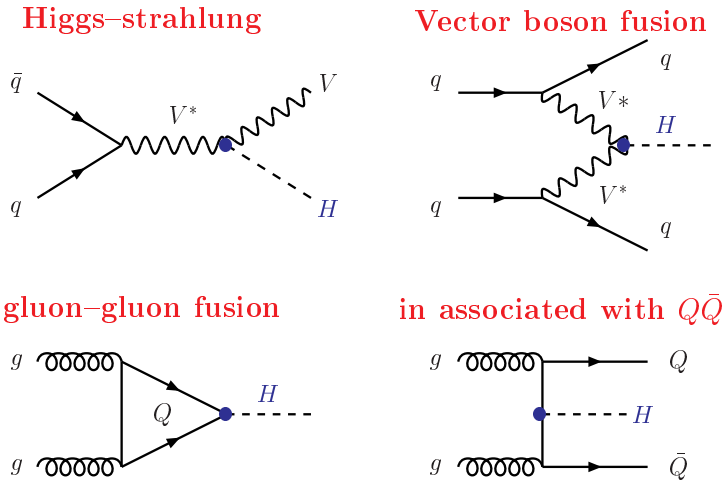
- very small for a light Higgs
- comparable to mass if heavy



HDECAY  $\Rightarrow$

# 2. The Higgs at the LHC

## Main Higgs production channels



## Large production cross sections

with  $gg \rightarrow \text{H}$  by far dominant process

$1 \text{ fb}^{-1} \Rightarrow \mathcal{O}(10^4)$  events@LHC

$\Rightarrow \mathcal{O}(10^3)$  events@Tevatron

but eg  $\text{BR}(\text{H} \rightarrow \gamma\gamma, \text{ZZ} \rightarrow 4\ell) \approx 10^{-3}$

... a small # of events at the end...

with a huge QCD-jet background.

$\Rightarrow$  an extremely challenging task!

## Main sensitive channels:

$gg \rightarrow \text{H} \rightarrow \gamma\gamma$

$gg \rightarrow \text{H} \rightarrow \text{ZZ} \rightarrow 4\ell, 2\ell 2\nu, 2\ell 2\gamma$

$gg \rightarrow \text{H} \rightarrow \text{WW} \rightarrow \nu\nu + 0, 1j$

also help from other channels:

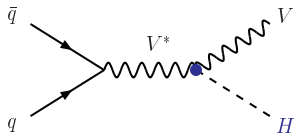
-  $\text{VBF} + gg \rightarrow \text{H} \rightarrow \tau\tau$

-  $q\bar{q} \rightarrow \text{HV} \rightarrow b\bar{b}\ell X$

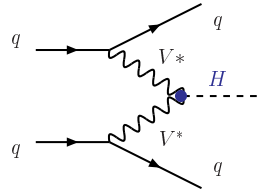
# 2. The Higgs at the LHC

Things are even more complicated/challenging in BSM: MSSM case

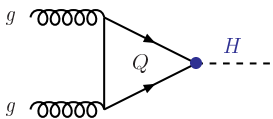
Higgs-strahlung



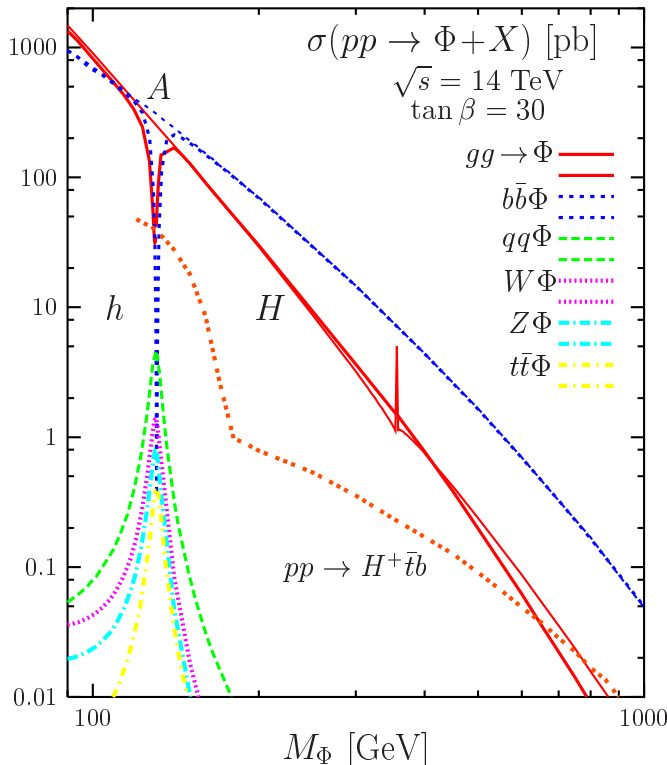
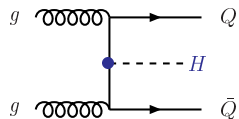
Vector boson fusion



gluon-gluon fusion



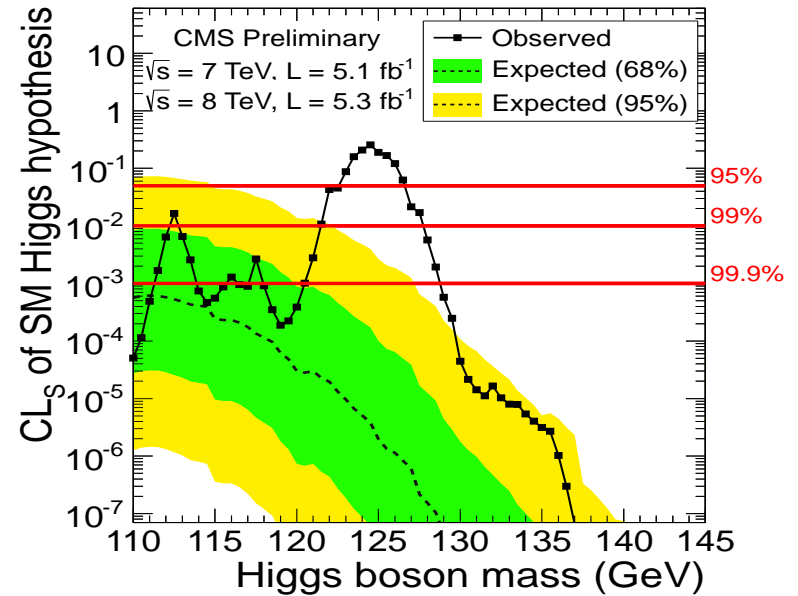
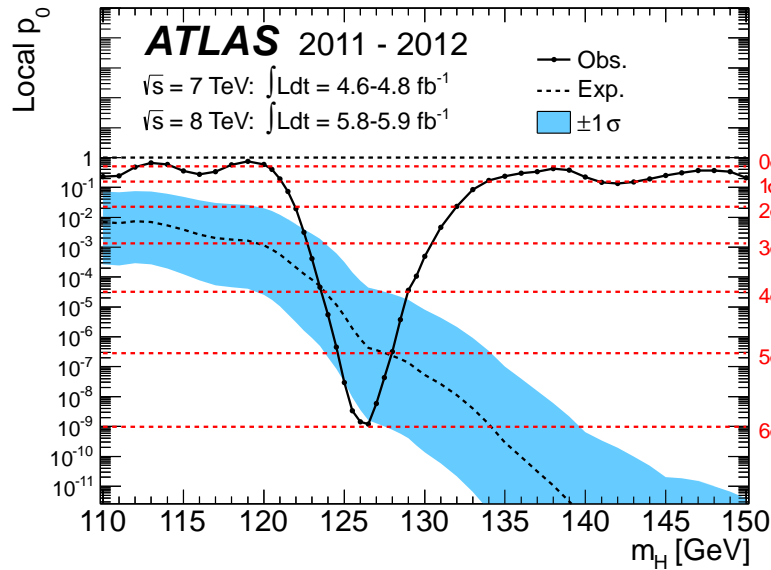
in associated with  $Q\bar{Q}$



- More Higgs particles:  $\Phi = h, H, A, H^\pm$ 
  - some couple almost like the SM Higgs,
  - but some are more weakly coupled.
- In general same production as in SM but also new/more complicated processes (rates can be smaller or larger than in SM)
- Possibility of different decay modes (and clean decays eg into  $\gamma\gamma$  suppressed)
- Impact of light SUSY particles?
  - $\Rightarrow$  In general very complicated situation!
- But simpler in the decoupling regime:
  - h as in SM with  $M_h = 115 - 130$  GeV
  - dominant mode:  $gg, b\bar{b} \rightarrow H/A \rightarrow \tau\tau$
- It is even more tricky in beyond MSSM! and also in some non-SUSY extensions..

# 2. The Higgs at the LHC

... a challenge met the 4th of July, when the Higgs was discovered at LHC.



# 3. Implications of the discovery: is it a Higgs?

**Spin:** the state decays into  $\gamma\gamma$

- not spin-1: Landau-Yang
- could be spin-2 like graviton? **Ellis et al.**
- miracle that couplings fit that of H,
- “prima facie” evidence against it:

e.g.:  $c_g \neq c_\gamma, c_V \gg 35c_\gamma$

many th. analyses (no suspense).

**CP:** is it CP-even or CP-odd?

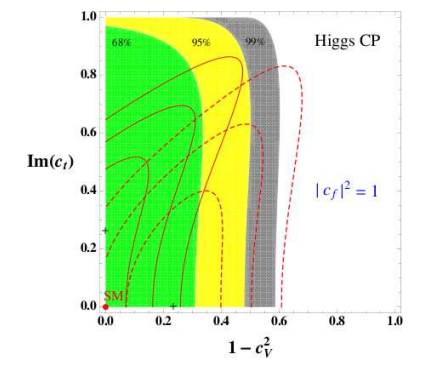
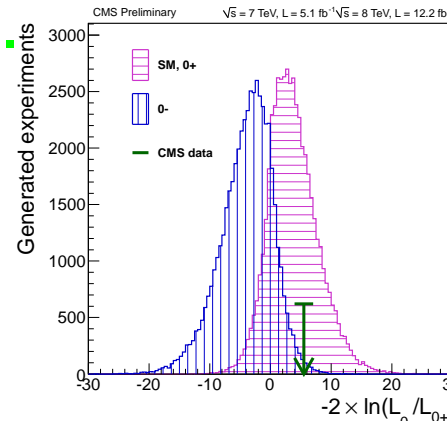
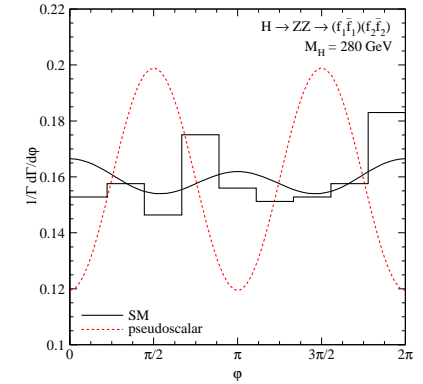
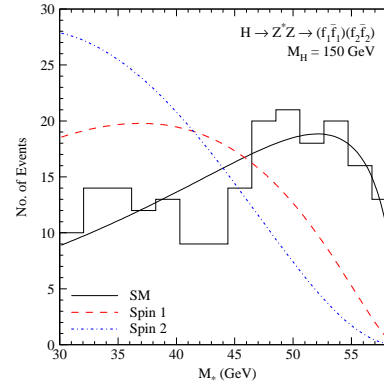
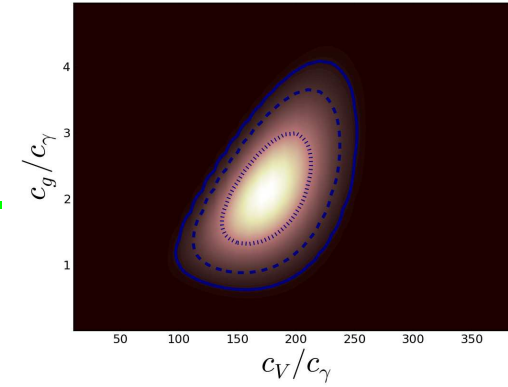
$$H V_\mu V^\mu \text{ vs } H \epsilon^{\mu\nu\rho\sigma} Z_{\mu\nu} Z_{\rho\sigma}$$

$$\Rightarrow \frac{d\Gamma(H \rightarrow ZZ^*)}{dM_*} \text{ and } \frac{d\Gamma(H \rightarrow ZZ)}{d\phi}$$

**ATLAS/CMS:**  $\approx 3\sigma$  for CP-even.

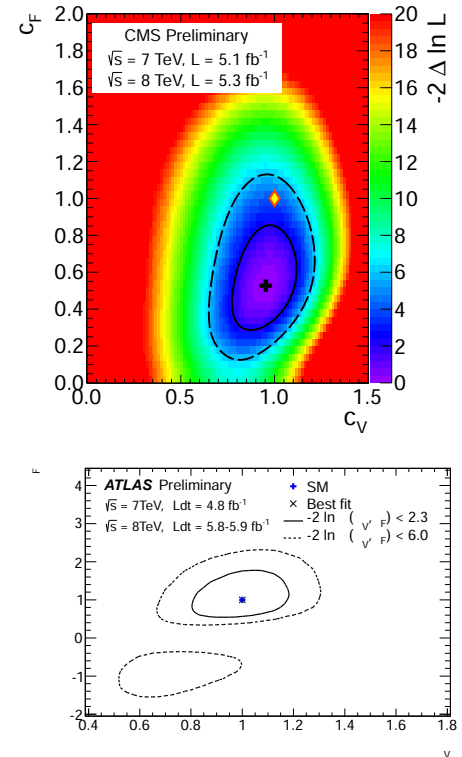
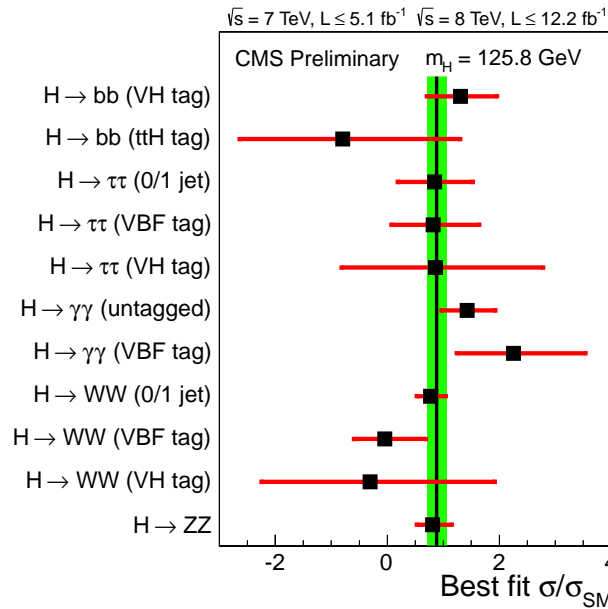
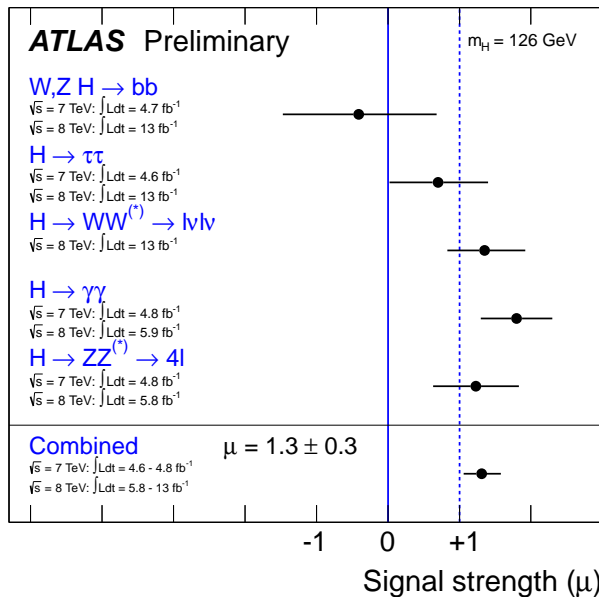
**Problem:** if H is CP mixture, only  $0^+$  component is projected out!  
(or very large  $0^-$  VV loop cplg).

$\Rightarrow$  better probe:  $\hat{\mu}_{ZZ} = 1.1 \pm 0.4!$





# 3. Implications of the discovery: is it a Higgs?



From ATLAS/CMS results:

Higgs couplings to elementary particles as predicted by Higgs mechanism

- couplings to  $WW, ZZ, \gamma\gamma$  roughly as expected for a CP-even Higgs
- couplings proportional to masses as expected for the Higgs boson

So, it is not only a “new particle”, the “126 GeV boson”, a “new state”...

**IT IS A HIGGS BOSON!**

But is it **THE** SM Higgs boson or **A** Higgs boson from some extension?

To check this you need very precise measurements to see small deviations...

# 3. Implications of the discovery: the SM

The Higgs looks like expected in SM  $\Rightarrow$   
 a triumph for high-energy physics!

Indirect constraints from EW data <sup>a</sup>

H contributes to RC to W/Z masses:

$$\begin{array}{c}
 \text{wavy line} \quad \text{H} \quad \text{wavy line} \\
 \text{W/Z} \quad \quad \quad \text{W/Z}
 \end{array}
 \propto \frac{\alpha}{\pi} \log \frac{M_H}{M_W} + \dots$$

Fit the EW precision measurements,  
 one obtains  $M_H = 92^{+34}_{-26}$  GeV, or

$$M_H \lesssim 160 \text{ GeV at 95\% CL}$$

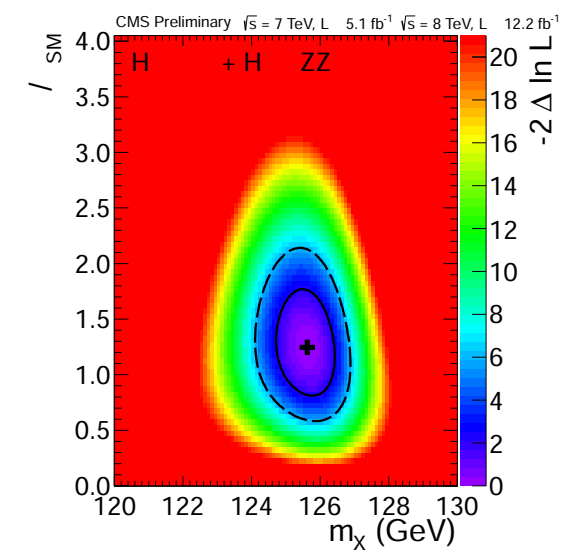
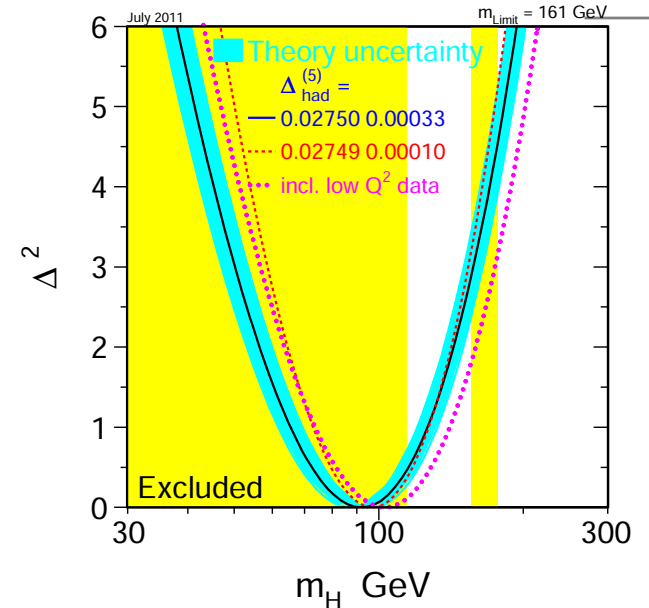
compared with the measured mass

$$M_H \approx 126 \text{ GeV.}$$

A very non-trivial consistency check!

(remember the story of the top quark!).

**The SM is a very successful theory!**



<sup>a</sup> Still some problems with  $A_{\text{FB}}^b$  (LEP),  $A_{\text{FB}}^t$  (TeV) and  $g-2$  but not severe...

# 3. Implications of the discovery: the SM

- **The theory preserves unitarity** as we have  $M_H \ll 700 \text{ GeV}$ ...
- **Particle spectrum complete:** Fourth generation excluded by  $H \rightarrow ZZ, WW, \gamma\gamma, bb$  rates... (as well as by direct searches@LHC...)

- **Extrapolable up to highest scales.**

$$\frac{\lambda(Q^2)}{\lambda(v^2)} \approx 1 + 3 \frac{2M_W^4 + M_Z^4 - 4m_t^4}{16\pi^2 v^4} \log \frac{Q^2}{v^2}$$

tops make  $\lambda < 0$ : unstable vacuum

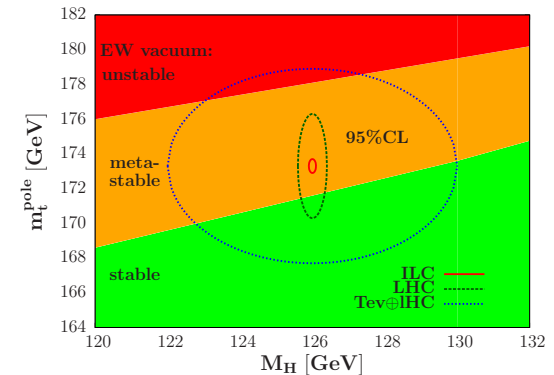
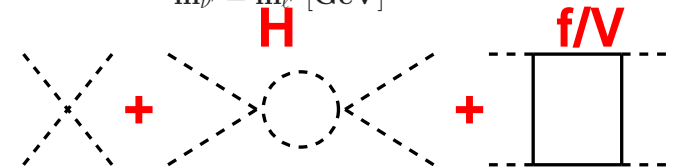
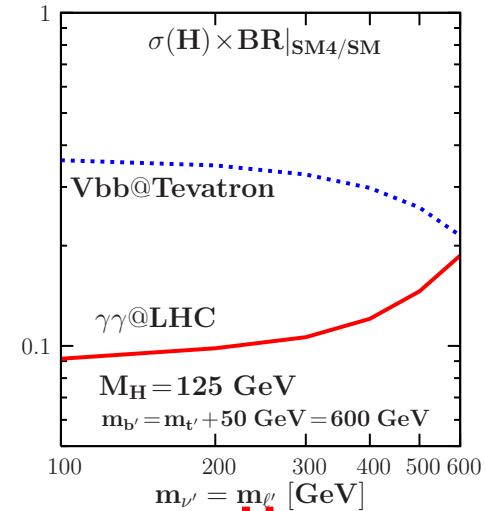
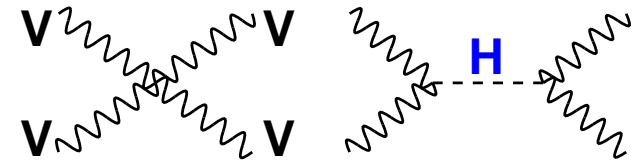
$$\Lambda_C \sim M_{Pl} \Rightarrow M_H \gtrsim 129 \text{ GeV!}$$

at 2 loops for  $m_t^{\text{pole}} = 173 \text{ GeV}$ .....

(Degrassi et al., Bezrukov et al.)

but what is measured  $m_t$  value?

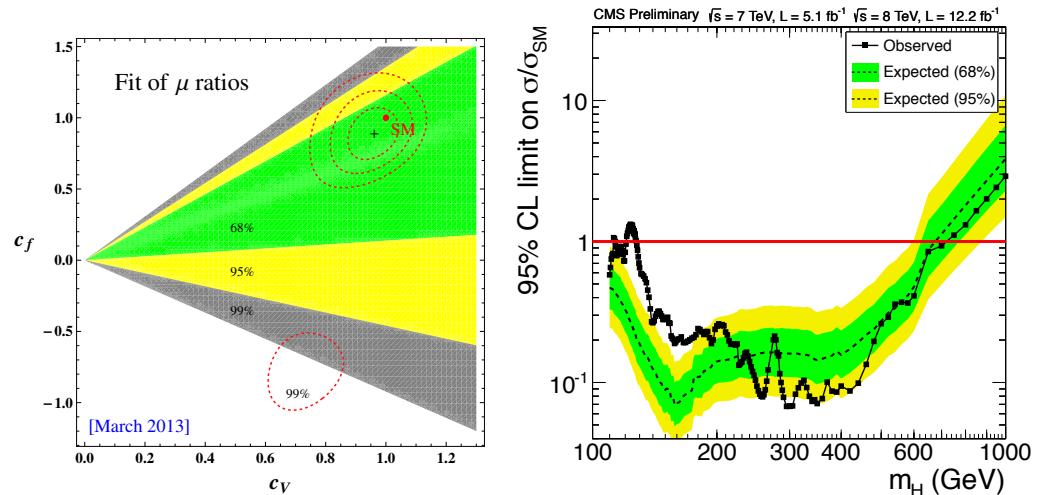
- **SM = TOE? Maybe not (?):**
- $m_\nu$ , DM, GUT OK with extensions
- but about the hierarchy problem?



# 3. Implications of the discovery: beyond the SM

- Rates compatible with SM fit of all data  $\Rightarrow$  OK at  $\approx 20\%$
- No other resonance found in many search channels....

**Huge implications for BSM!**



Some beyond the SM scenarios are in ‘mortuary’:

- Higgsless models, extreme Technicolor and composite scenarios, ..
- fermiophobic Higgs, gauge-phobic Higgs, 4th generation, ...

Some beyond the SM scenarios are in ‘hospital’:

- ‘light’ versions of Technicolor and composite models...
- many other extended Higgs scenarios: private, portal, ....

Other BSM scenarios are strongly constrained...

**and the best example is Supersymmetry and the MSSM.**

### 3. Implications of the discovery: the MSSM

In MSSM, two doublets  $H_1, H_2 \Rightarrow$  5 physical states:  $h, H, A, H^\pm$

only two parameters at tree-level:  $\tan\beta, M_A$  but rad. cor. important:

$$M_h \lesssim M_Z |\cos 2\beta| + RC \lesssim 130 \text{ GeV}, \quad M_H \approx M_A \approx M_{H^\pm} \lesssim M_{\text{EWSB}}$$

**126 GeV is large for MSSM:  $\Rightarrow M_h$  needs to be maximal from start...**

$$M_h \xrightarrow{M_A \gg M_Z} M_Z |\cos 2\beta| + \frac{3\bar{m}_t^4}{2\pi^2 v^2 \sin^2 \beta} \left[ \log \frac{M_S^2}{\bar{m}_t^2} + \frac{X_t^2}{M_S^2} \left( 1 - \frac{X_t^2}{12M_S^2} \right) \right]$$

- decoupling regime with  $M_A \sim \mathcal{O}(\text{TeV})$ ;  $h$  is SM-like
- large values of  $\tan\beta \gtrsim 10$  to maximize tree-level value;
- maximal mixing scenario:  $X_t = \sqrt{6}M_S$ ;
- heavy stops, i.e. large  $M_S = \sqrt{\bar{m}_{\tilde{t}_1} \bar{m}_{\tilde{t}_2}}$ ; but  $M_S \lesssim 3 \text{ TeV} \dots$

**Scan parameter space with all corrections and full SUSY spectrum**

**Constrained MSSMs are interesting from model building point of view:**

- concrete schemes: SSB occurs in hidden sector  $\xrightarrow{\text{gravity, ...}}$  MSSM fields
- provide solutions to some MSSM problems: CP, flavor, etc..
- parameters obey boundary conditions  $\Rightarrow$  small number of inputs...

**the prototype model is mSUGRA:  $\tan\beta, m_{1/2}, m_0, A_0, \text{sign}(\mu)$**

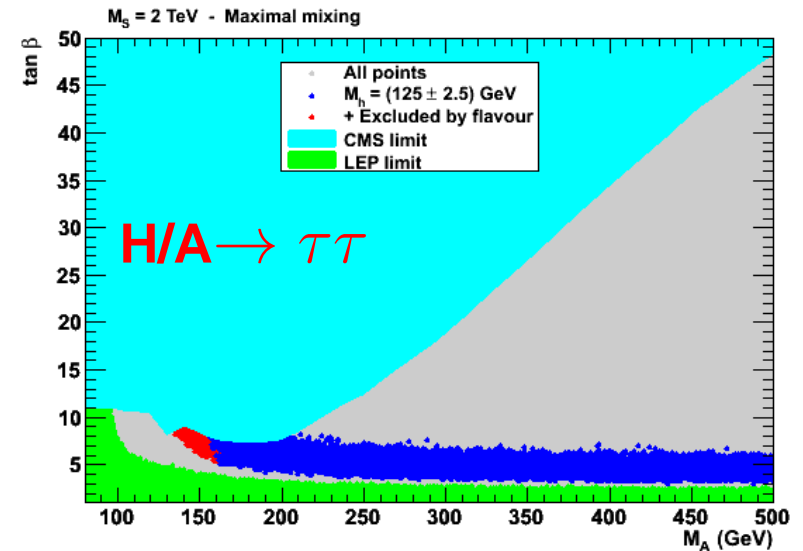
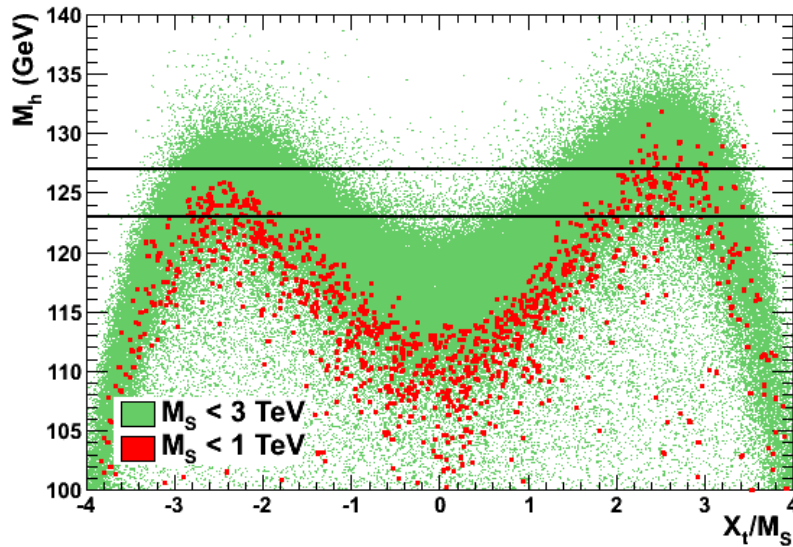
**full scan of the model parameters with  $123 \text{ GeV} \leq M_h \leq 129 \text{ GeV}$**



# 3. Implications of the discovery: the MSSM

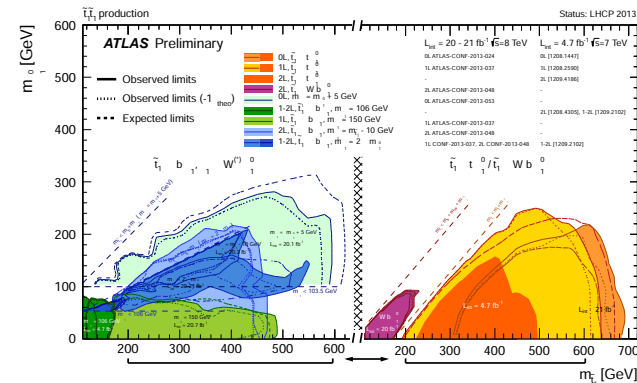
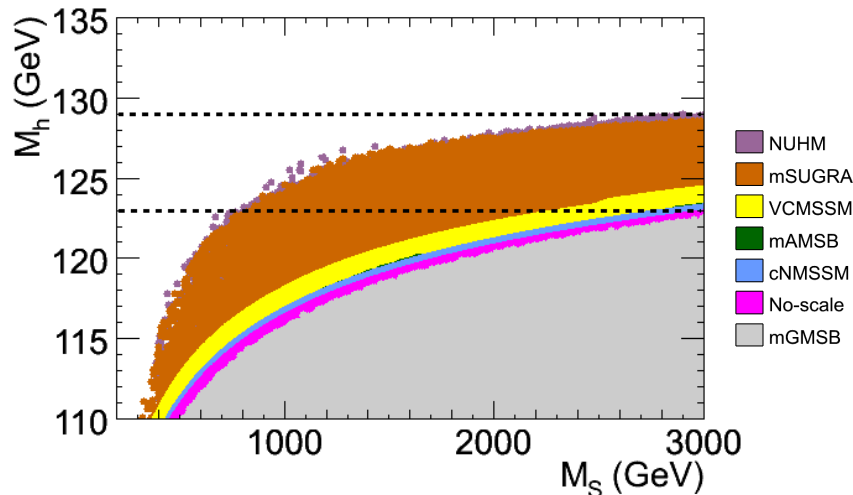
⇒ SUSY scale rather large...

... backed up by direct searches



especially in constrained MSSMs ...

especially squarks/gluinos...



# 3. Implications of the discovery: the MSSM

**A 126 GeV Higgs provides information on BSM and SUSY in particular:**

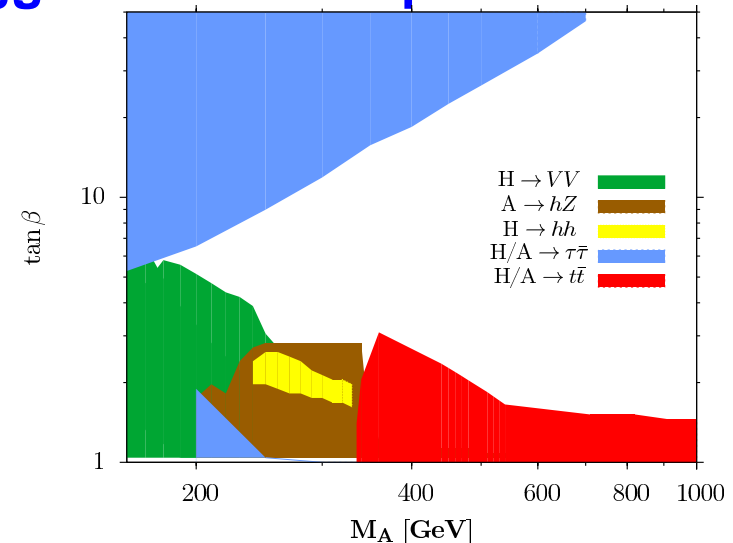
- $M_H = 119$  GeV would have been a boring value: everybody OK..
- $M_H = 145$  GeV would be a devastating value: mass extinction..
- $M_H \approx 126$  GeV is Darwinian: (natural) selection among models..

**SUSY spectrum heavy; except maybe for weakly interacting sparticles and also stops  $\Rightarrow$  more focus on them in SUSY searches!**

**One has to refine all other MSSM Higgs searches in particular:**

- $gg, bb \rightarrow H/A \rightarrow \tau\tau, \mu\mu$
- $t \rightarrow H^+b, gg \rightarrow tH^-$
- $H \rightarrow WW, ZZ$  as in SM
- $gg, H/A \rightarrow tt$
- $H \rightarrow hh, A \rightarrow Zh....$

**and of course sparticle searches...**



**7–8 TeV LHC for the lightest h and 13–14 TeV LHC for H/A/H<sup>+</sup> ?  
and maybe some supersymmetric particles will show up?**

## 4. What next?

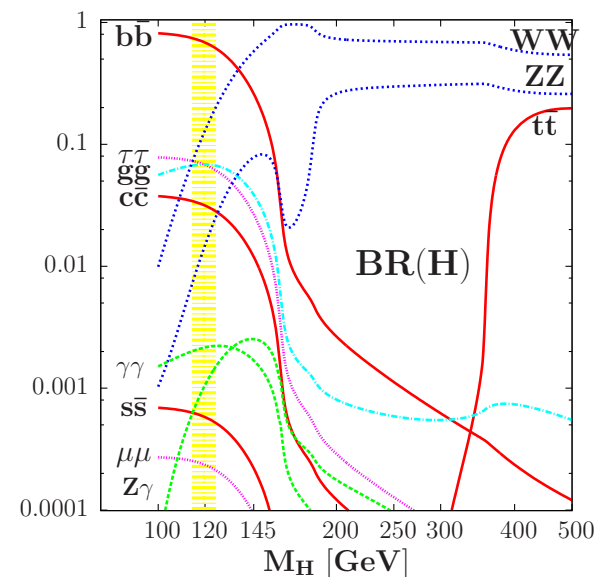
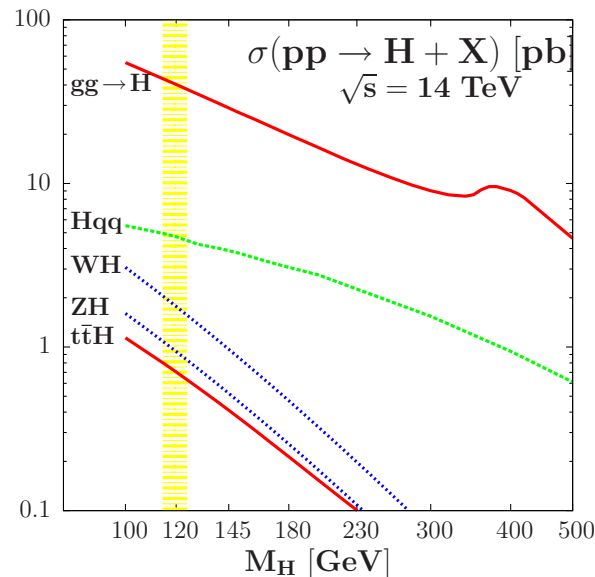
Even if no sign of BSM physics is seen: is Particle Physics “closed”?

**No! Need to check that H is indeed responsible of sEWSB (and SM-like?)**

**Measure its fundamental properties in the most precise way:**

- its mass and total decay width (invisible width due to dark matter?),
- its spin–parity quantum numbers and check SM prediction for them,
- its couplings to fermions and gauge bosons and check that they are indeed proportional to the particle masses (fundamental prediction!),
- its self–couplings to reconstruct the potential  $V_H$  that makes EWSB.

**Possible for  $M_H \approx 126$  GeV as all production/decay channels useful!**



# 4. What next? Couplings

- Look at various H production/decay channels and measure  $N_{ev} = \sigma \times BR$
- But large errors mainly due to:
  - experimental: stats, system., lumi...
  - theory: PDFs, HO/scale, jetology...
- total error about 20–30% in  $gg \rightarrow H$
- Hjj contaminates VBF (now 30%)..

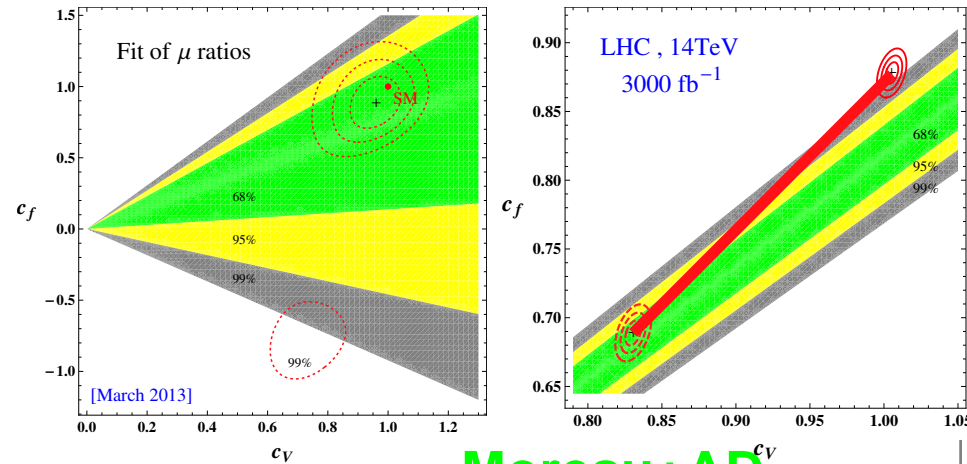
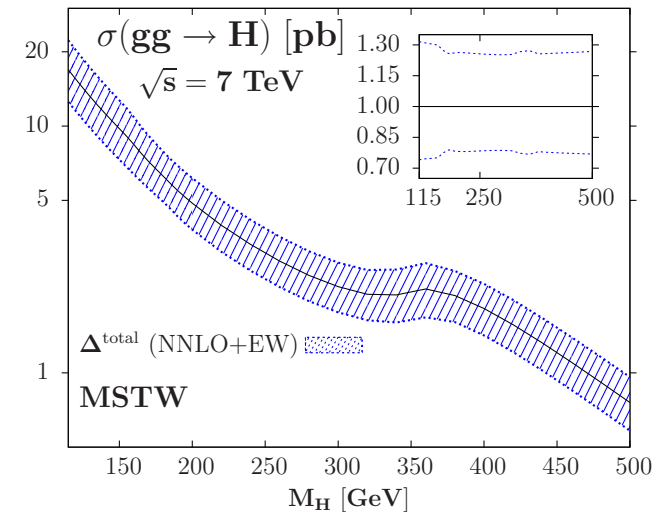
⇒ ratios of  $\sigma \times BR$ : many errors out!

Deal with width ratios  $\Gamma_X/\Gamma_Y$

- TH on  $\sigma$  and some EX errors
- parametric errors in BRs
- TH ambiguities from  $\Gamma_H^{tot}$
- Achievable accuracy:
  - now: 20–30% on  $\gamma\gamma/VV, \tau\tau/VV$
  - future: few % at HL–LHC!

Sufficient to probe BSM physics?

Baglio+AD

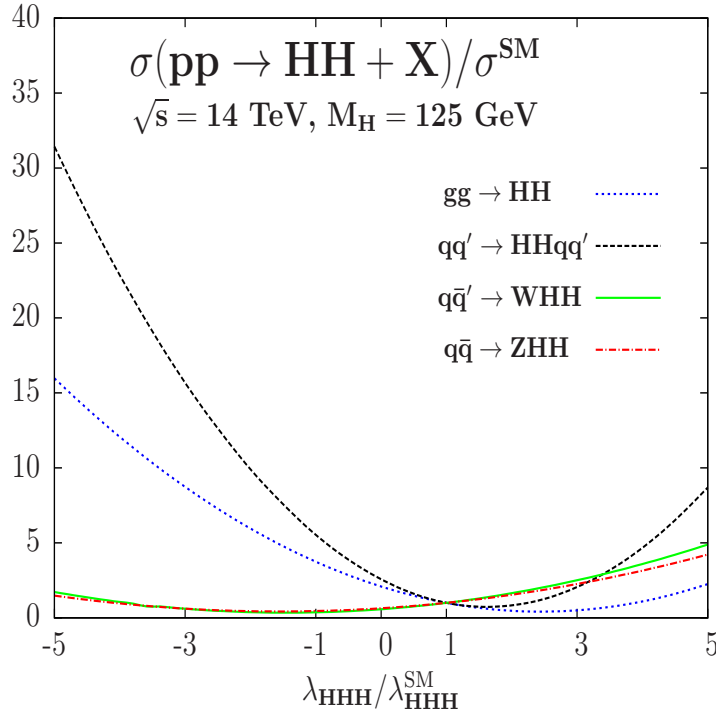
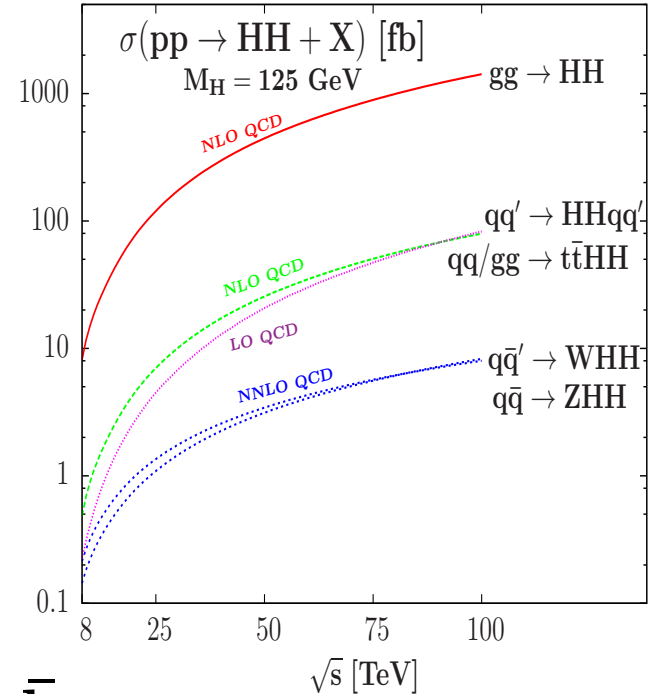
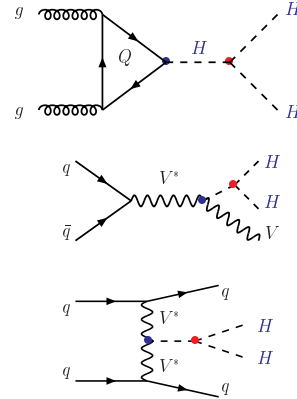


Moreau+AD

# 5. What next? Self-coupling

Challenge: measurement of Higgs self-couplings and access to  $V_H$ .

- $g_{H^3}$  from  $pp \rightarrow HH + X \Rightarrow$
  - $g_{H^4}$  from  $pp \rightarrow 3H + X$ , hopeless.
- Various processes for HH prod:  
only  $gg \rightarrow HHX$  relevant...



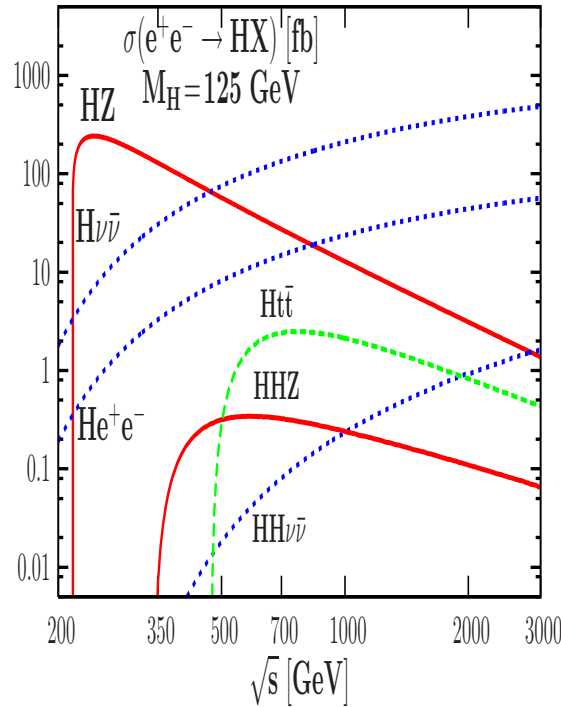
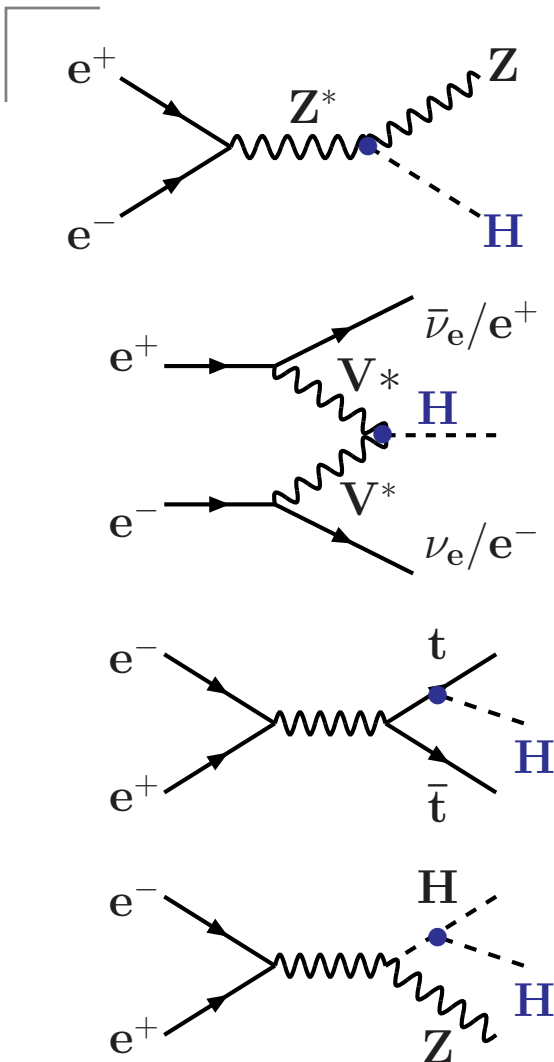
- $H \rightarrow b\bar{b}$  decay alone not clean
- $H \rightarrow \gamma\gamma$  decay very rare,
- $H \rightarrow \tau\tau$  would be possible?
- $H \rightarrow WW$  not useful?
- $bb\tau\tau, bb\gamma\gamma$  viable?
- but needs <sup>1</sup>very large luminosity.

Maybe even needs an ILC.....

Baglio et al., arXiv:1212.5581



# 4. What next? ILC



Very precise measurements mostly at  $\sqrt{s} \lesssim 500$  GeV and mainly in  $e^+e^- \rightarrow ZH$  (with  $\sigma \propto 1/s$ ) and  $ZHH$ ,  $ttH$

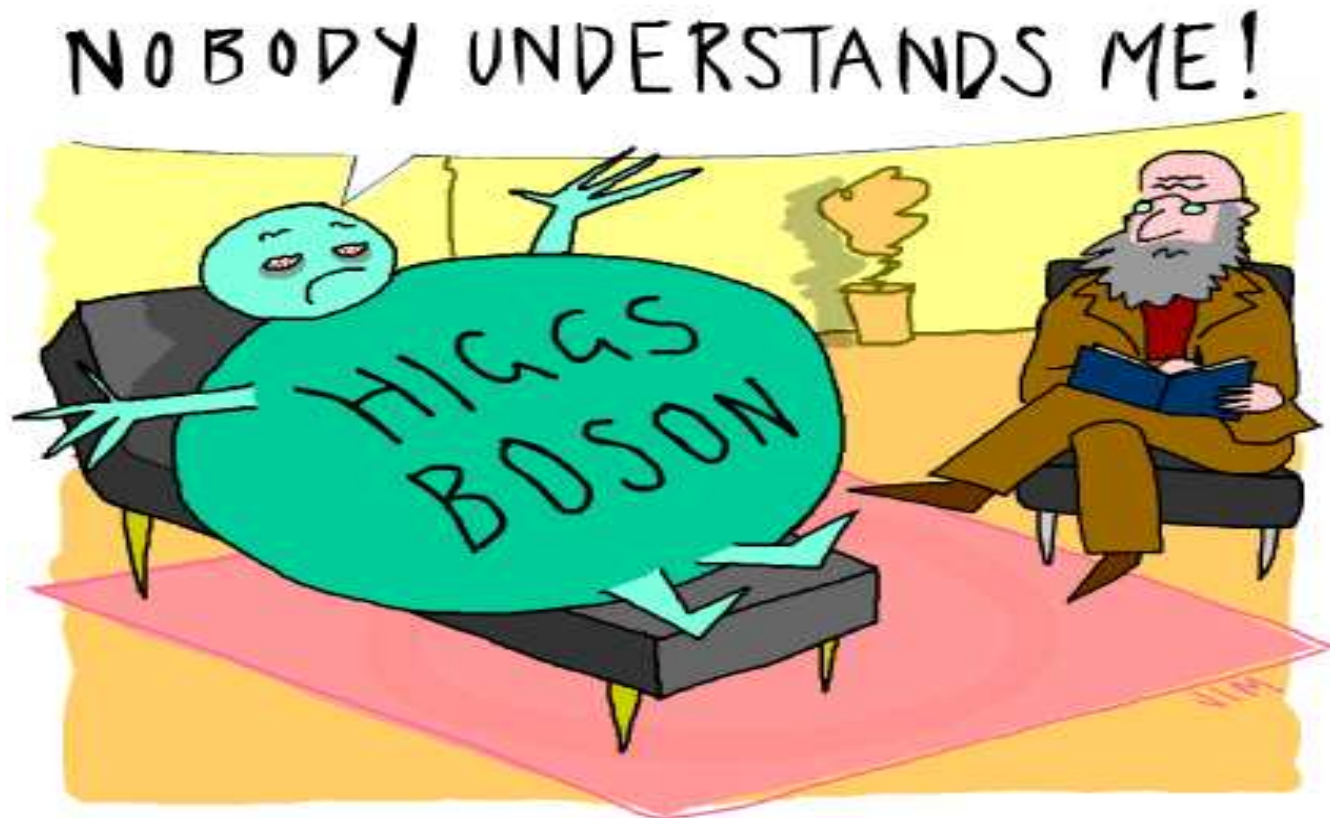
$g_{HWW}$	$\pm 0.012$
$g_{HZZ}$	$\pm 0.012$
$g_{Hbb}$	$\pm 0.022$
$g_{Hcc}$	$\pm 0.037$
$g_{H\tau\tau}$	$\pm 0.033$
$g_{Htt}$	$\pm 0.030$
$\lambda_{HHH}$	$\pm 0.22$
$M_H$	$\pm 0.0004$
$\Gamma_H$	$\pm 0.061$
CP	$\pm 0.038$

⇒ difficult to be beaten by anything else for  $\approx 125$  GeV Higgs

⇒ welcome to the ILC!

## 4. What next?

We hope that we will finally understand the Higgs mechanism...



... but there is a long way until we get there....

... and there might be many surprises waiting for us...