Supersymmetric Holography on AdS₃

Michael Kech

IMPRS workshop

March 18, 2013

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Pigher Spin Duality on AdS₃

- The Original Proposal
- Supersymmetric Holography on AdS₃

Partition Functions

- CFT Partition Function
- Gravity Partition Function

Section 1

Introduction

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Duality between higher spin theories and CFTs Simplified version of AdS/CFT duality

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Duality between higher spin theories and CFTs

- Simplified version of AdS/CFT duality
- Conceptual insight in AdS/CFT

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Duality between higher spin theories and CFTs

- Simplified version of AdS/CFT duality
- Conceptual insight in AdS/CFT
- CFT methods become available for quantum gravity

Section 2

Higher Spin Duality on AdS_3

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Consider higher spin duality in three dimensions2d CFTs are well understood

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Consider higher spin duality in three dimensions

- 2d CFTs are well understood
- Higher spin theories simplify in three dimensions

Subsection 1

The Original Proposal

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Duality between one of Vasiliev's higher spin theories on AdS_3 containing a complex scalar and the \mathcal{W}_N^k minimal models in the t'Hooft limit

 \bullet CFT coupling given by t'Hooft parameter λ

Duality between one of Vasiliev's higher spin theories on AdS_3 containing a complex scalar and the \mathcal{W}_N^k minimal models in the t'Hooft limit

- \bullet CFT coupling given by t'Hooft parameter λ
- Mass of the scalar given by $M^2 = -(1-\lambda^2)$

Minimal models are the coset theories

$$\mathcal{W}_N^k = rac{\mathfrak{su}(N)_k \oplus \mathfrak{su}(N)_1}{\mathfrak{su}(N)_{k+1}}$$

where $\lambda = rac{N}{N+k}$.

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Image: A match a ma

There is a description of higher spin theories analogous to the Chern-Simons formulation of gravity

$$S_{CS} = S[A] - S[\bar{A}], \ S[A] = \frac{\hat{k}}{4\pi} \int Tr[A \wedge dA + \frac{2}{3}A \wedge A \wedge A]$$

with gauge group $hs(\lambda)$ instead of $\mathfrak{sl}(2)$.

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Much non-trivial evidence in favour of proposal based on Symmetry analysis

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Much non-trivial evidence in favour of proposal based on

- Symmetry analysis
- Correlation functions

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Much non-trivial evidence in favour of proposal based on

- Symmetry analysis
- Correlation functions
- Spectrum

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Subsection 2

Supersymmetric Holography on AdS₃

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CFT is based on the supersymmetric coset

$$\frac{\mathfrak{su}(N+1)_{k+N+1}^1}{\mathfrak{su}(N)_{k+N+1}^1 \oplus \mathfrak{u}(1)_{\kappa}^1}$$

This coset is one of the $\mathcal{N} = 2$ Kazama-Suzuki cosets.

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Image: A match a ma

Higher spin theory is based on the superalgebra $shs(\lambda)$. $shs(\lambda)$ can be obtained from the associative algebra $A(\nu)$ generated by

$$[y_{\alpha}, y_{\beta}] = 2i\epsilon_{\alpha\beta}(1+\nu k), \ \{k, y_{\alpha}\} = 0, \ \alpha, \beta = 1, 2$$

by introducing a grading according to the powers of $y_{\alpha,\beta}$. ν is related to λ by $\lambda = \frac{1+\nu}{2}$.

Section 3

Partition Functions

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Subsection 1

CFT Partition Function

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Orthogonal version of supersymmetric coset is

$$rac{\mathfrak{so}(2\mathsf{N}+2)^1_{k+2\mathsf{N}+2}}{\mathfrak{so}(2\mathsf{N})^1_{k+2\mathsf{N}+2}\oplus\mathfrak{u}(1)^1_\kappa}.$$

This is isomorphic to the bosonic subalgebras and free Majorana fermions

$$\frac{\mathfrak{so}(2\mathsf{N}+2)_k\oplus\mathfrak{so}(4\mathsf{N})_1}{\mathfrak{so}(2\mathsf{N})_{k+2}\oplus\mathfrak{u}(1)_\kappa}$$

The coset representations are given by the multiplicity spaces (Λ, Φ, I) in

$$\Lambda\otimes NS = \bigoplus_{I,\Phi} (\Lambda, \Phi, I) \otimes \Phi \otimes I.$$

The Hilbert space is defined by charge conjugation

$$\mathcal{H}_{CFT} = \bigoplus_{\Lambda, \Phi, I} (\Lambda, \Phi, I) \otimes \overline{(\Lambda, \Phi, I)}.$$

To calculate partition function one has to find the characters of the multiplicity spaces (Λ, Φ, I) . Main ideas

• In t'Hooft limit affine characters simplify

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- Problem reduces to counting $\mathfrak{so}(2N)$ representations in a free theory

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- In t'Hooft limit affine characters simplify
- Problem reduces to counting $\mathfrak{so}(2N)$ representations in a free theory
- \bullet Result can be expressed in terms of $\mathfrak{gl}(\infty|\infty)_+$ characters

Subsection 2

Gravity Partition Function

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• CFT partition function determines field content of higher spin theory

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Spectrum

- CFT partition function determines field content of higher spin theory
- In the gauge sector the dual higher spin theory has to contain two N = 2 multiplets with lowest spin s for each spin s = 1, 2, 3, ... and s = 1/2, 5/2, 9/2, ...

Spectrum

- CFT partition function determines field content of higher spin theory
- In the gauge sector the dual higher spin theory has to contain two N = 2 multiplets with lowest spin s for each spin s = 1, 2, 3, ... and s = 1/2, 5/2, 9/2, ...
- Sums over gl(∞|∞)₊ characters can be interpreted in terms of massive fields. There are two complex scalars, two real scalars, two complex fermions and two real fermions

- Goal is to construct a higher spin algebra that posses this spectrum.
- Such an algebra can be obtained by restricting the Lie superalgebra obtained from

 $\mathfrak{gl}(2)\otimes A(\nu)$

by an automorphism.

• Higher spin dualities on *AdS*₃ allow for strong consistency tests

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- Higher spin dualities on *AdS*₃ allow for strong consistency tests
- Prove duality in t'Hooft limit (better understand scalar on gravity side)

- Higher spin dualities on *AdS*₃ allow for strong consistency tests
- Prove duality in t'Hooft limit (better understand scalar on gravity side)
- Understand relation between higher spin theories and string theories