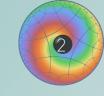
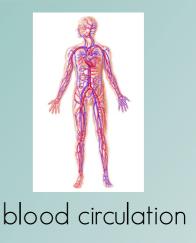
ABC's of Holographic Hydrodynamics

Why should we care about Hydro?



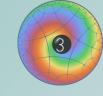






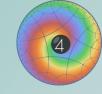
- describes a vast amount of different systems
- equations (Navier-Stokes) poorly understood (e.g. turbulence, CLAY Millennium Prize)

Outline



- Hydrodynamics
- Gravity and Hydrodynamics
 - Membrane Paradigm
 - ▶ Fluid/Gravity Correspondence

What is Hydrodynamics?



Def: effective long-distance description of classical or quantum many-body system at finite temperature

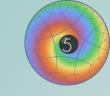
Conservation laws:

$$\partial_{\alpha}T^{\alpha\beta} = 0$$
 $\partial_{\alpha}J^{\alpha} = 0$

Constitutive relations:

$$T_{\alpha\beta} \equiv T_{\alpha\beta}(T, u_{\sigma}, \mu) \quad J_{\alpha} \equiv J_{\alpha}(T, u_{\sigma}, \mu)$$

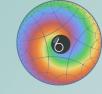
What do we want?



Goal: - systematic construction of $T_{lphaeta}$ and J_{lpha}

- in terms of $T(x), \ u^{\alpha}(x), \ \mu(x)$
- in agreement with the symmetries

Conserved Quantities

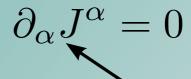


$$\partial_{\alpha} T^{\alpha\beta} = 0$$

spacetime symmetries:

translations, rotations and boosts

$$d+1$$
 equations



internal symmetries:

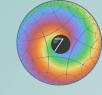
e.g. U(1) baryon number symmetry

1 equation

d+2 equations

Noether theorem:
conservation laws related to
continuous symmetries of the
fundamental microscopic theory

Why can we use Hydro Fields?



- Assumption: many body systems always equilibrate locally over a finite time scale t_m
- ightharpoonup for $t > t_m$ local equilibrium at every point x
- Parameters characterising local equilibrium: $T(x),\ u^{\alpha}(x),\ \mu(x)$
- riangle these parameters vary on length scales $L{>}l_m$
- effective dynamical fields at late times
- evolution governed by hydro equations

Relation btw. Equations and Fields

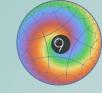


$$>$$
 $T_{lphaeta}$ $(d+1)(d+2)/2$ components $>$ J_{lpha} $(d+1)$

- Hydrodynamical limit (long-wavelength limit):

 To and I expressed in terms of delay fields
 - $T_{lphaeta}$ and J_{lpha} expressed in terms of d+2 fields:
 - local temperature T(x)
 - local fluid velocity $u_{lpha}(x)$ (Note: $u_{lpha}u^{lpha}=-1$)
 - local chemical potential $\mu(x)$

Constitutive Relations

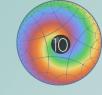


Now we can systematically construct $T_{\alpha\beta}$ and J_{α} out of the hydrodynamic variables:

$$T^{\alpha\beta} = \epsilon u^{\alpha} u^{\beta} + P(\eta^{\alpha\beta} + u^{\alpha} u^{\beta}) + \Pi^{\alpha\beta}$$
$$J^{\alpha} = q u^{\alpha} + \Upsilon^{\alpha}$$

- Equations hold only locally in spacetime
- $\triangleright \epsilon, P, q$ depend on T, μ, u^{α}
- $\triangleright \Pi^{\alpha\beta}$, Υ^{α} depend on gradients of T, μ , u^{α}

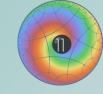
Higher-Order Hydrodynamics



- ho $\Pi_{(n)}^{lphaeta}$ n-th order in derivative of hydro fields
- $ilde{ ilde{}}$ every derivative suppresses fields by 1/L
- \triangleright magnitude of n-th term relative to 0th order:

$$\left(rac{l_m}{L}
ight)^n$$
 with $L\gg l_m$

First-Order Hydrodynamics



$$\Pi_{(1)}^{ab} = -\eta \left(\partial^a u^b + \partial^b u^a - \frac{2}{d} \delta^{ab} \partial \cdot u \right) - \zeta \delta^{ab} \partial \cdot u$$

$$\Upsilon^a = -\sigma T \partial^a \left(\frac{\mu}{T} \right)$$

Constraint by: Dorentz invariance

Local form of 2nd law

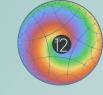
of thermodynamics: $\partial_{\alpha}J_{s}^{\alpha}\geq0$ entropy current

$$\eta \ge 0, \ \zeta \ge 0, \ \sigma \ge 0$$

transport

coefficients

Transport Coefficients



- Determined by the underlying microscopic theory
- weakly coupled theory:

Computed using linear response and fluctuation-dissipation theorem

quantifies the relation between thermal fluc. of a system in equilibrium and response of the system to ext. pert.

strongly coupled theory: AdS/CFT

Ringberg2013

describes how an observable

changes due to small ext. pert.

Hydrodynamics and Gravity



Dynamics of membranes in curved spacetime are governed by fluid dynamic equations in long-wavelength limit

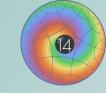
Einstein's equations in d+1 dim contain fluid equations in d dim

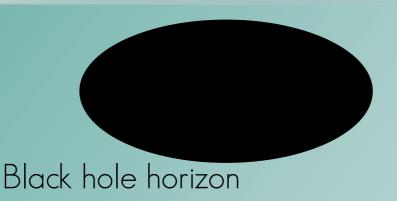
In the following:

1st: Membrane Paradigm (membrane close to b.h. hor.)

2nd: Fluid/Gravity Duality (AdS Boundary = membrane)

Membrane Paradigm





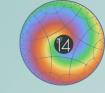
At the b.h. horizon:

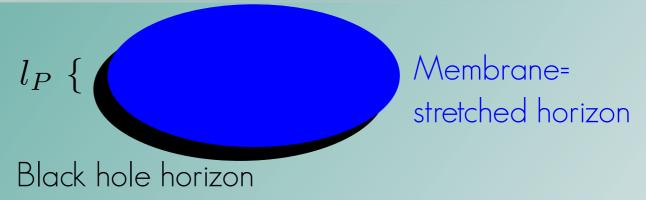
entropy density computed using QFT diverges

$$\neq \frac{A_H}{4G\hbar} = S_{BH}$$

- OFT not valid at this point, need Quant. Grav.
- Stretched horizon = cutoff for QFT

Membrane Paradigm





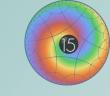
At the b.h. horizon:

entropy density computed using QFT diverges

$$\neq \frac{A_H}{4G\hbar} = S_{BH}$$

- DQFT not valid at this point, need Quant. Grav.
- Stretched horizon = cutoff for QFT

Stretched Horizon



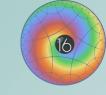
- Stretched horizon has dynamics of its own
- Behaves like a hot conducting viscous fluid
- Governed by Navier-Stokes Equation (non-relativistic hydrodynamic equation)

Aside:
$$\frac{\eta}{s} = \frac{1}{4\pi}$$



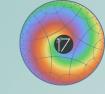
Black hole horizon

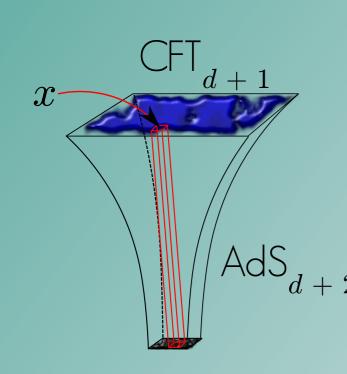
Fluid/Gravity Correspondence



- Low energy dynamics of planar black holes in asymptotically AdS given by relativistic conformal fluids
- Unperturbed black hole solution to Einstein equations describes thermal equilibrium of dual theory
- Long-wavelength perturbation of black hole corresponds to hydrodynamic expansion

Fluid dynamics in N=4 SYM





$$\partial_{\mu}T^{\mu\nu} = 0$$

$$T^{\alpha\beta} = \epsilon u^{\alpha} u^{\beta} + P(\eta^{\alpha\beta} + u^{\alpha} u^{\beta}) + \Pi^{\alpha\beta}$$

$$P = \frac{\pi^{d-1}}{16\pi G_N^{(d+2)}} T^{d-1}$$

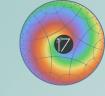
$$\eta = \frac{1}{16\pi G_N^{(d+2)}} \left(\frac{4\pi}{d}T\right)^d$$

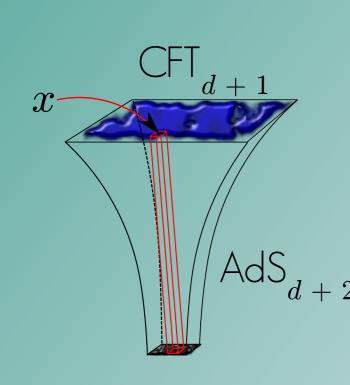
$$\epsilon = \frac{\pi^{d-1}}{16\pi G_N^{(d+2)}} (d-2) T^{d-1}$$



$$R_{MN} - \frac{1}{2}Rg_{MN} + \Lambda g_{MN} = 0$$

Fluid dynamics in N=4 SYM





$$\partial_{\mu}T^{\mu\nu} = 0$$

$$T^{\alpha\beta} = \epsilon u^{\alpha} u^{\beta} + P(\eta^{\alpha\beta} + u^{\alpha} u^{\beta}) + \Pi^{\alpha\beta}$$

$$P = \frac{\pi^{d-1}}{16\pi G_N^{(d+2)}} T^{d-1}$$

$$\epsilon = \frac{\pi^{d-1}}{16\pi G_N^{(d+2)}} (d-2) T^{d-1}$$

$$\eta = \frac{1}{16\pi G_N^{(d+2)}} \left(\frac{4\pi}{d}T\right)^d$$

$$\frac{\eta}{s} = \frac{1}{4\pi}$$

$$\uparrow$$

claimed to be lower bound for fluids found in nature (Koytun et.al.)

$$R_{MN} - \frac{1}{2}Rg_{MN} + \Lambda g_{MN} = 0$$

Further Results from Fluid/Gravity



- Systematic treatment of superfluids (i.e. Goldstone modes)
- Anomaly effects found in hydrodynamics:



- stir really fast

Further Results from Fluid/Gravity

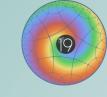


Systematic treatment of superfluids (i.e. Goldstone modes)

Anomaly effects found in hydrodynamics:

- stir really fast
- sugar jumps out of the mug

Conclusion



- Hydrodynamics is found in a vast number of different systems.
- There is a systematic way to construct hydrodynamic constitutive relations.
- There is a close connection to gravity.
- New results were found in the context of AdS/CFT correspondence
- Chance to win 1mio Dollars!



Hot Chocolate Effect

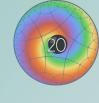
pour hot milk into a mug



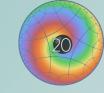




- pour hot milk into a mug
- stir in chocolate powder



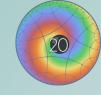




Hot Chocolate Effect

- pour hot milk into a mug
- stir in chocolate powder
- tap the bottom of the mug with a spoon while the milk is still in motion





Hot Chocolate Effect

- pour hot milk into a mug
- stir in chocolate powder
- tap the bottom of the mug with a spoon while the milk is still in motion
- pitch of taps will increase progressively with no relation to speed or force of tapping

MHX55

