

# Stringy Geometry

## Geometrization of $O(d,d)$ -transformations

**Felix Rennecke**

Based on

arXiv: 1202.4934, 1205.1522, 1210.1591, 1211.0030, 1304.2784



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Max-Planck-Institut für Physik  
(Werner-Heisenberg-Institut)

# MOTIVATION

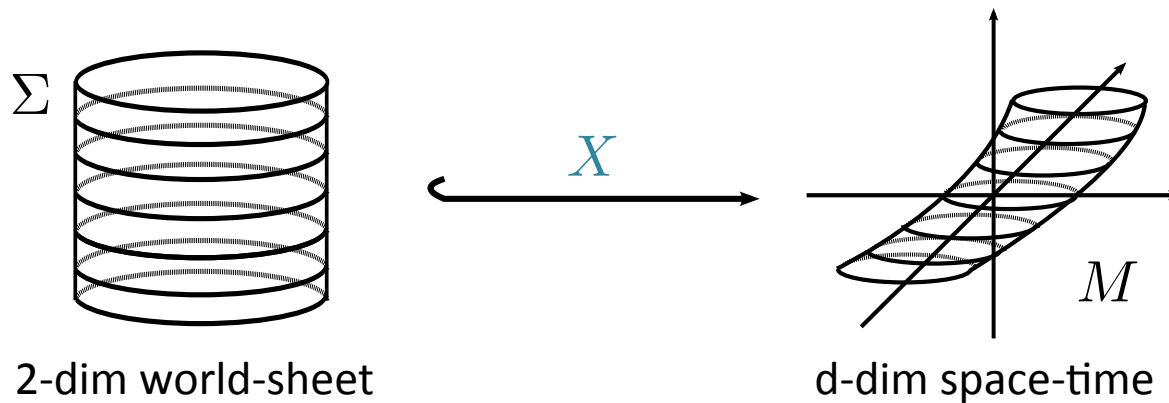
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$$S(G, B) = \frac{1}{2\pi\alpha'} \int_{\Sigma} d^2z (G_{ab} + B_{ab}) \partial X^a \bar{\partial} X^b$$

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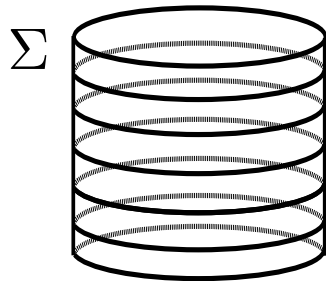


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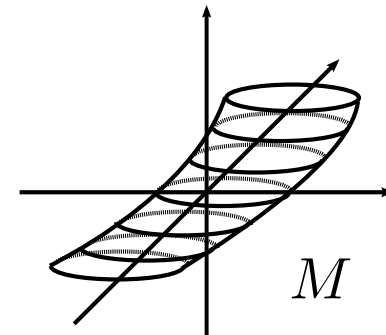
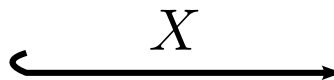
Metric determines  
*shape* of spacetime

$B$ -field gives  
*magnetic field*

$$S(G, B) = \frac{1}{2\pi\alpha'} \int_{\Sigma} d^2z (G_{ab} + B_{ab}) \partial X^a \bar{\partial} X^b$$



2-dim world-sheet

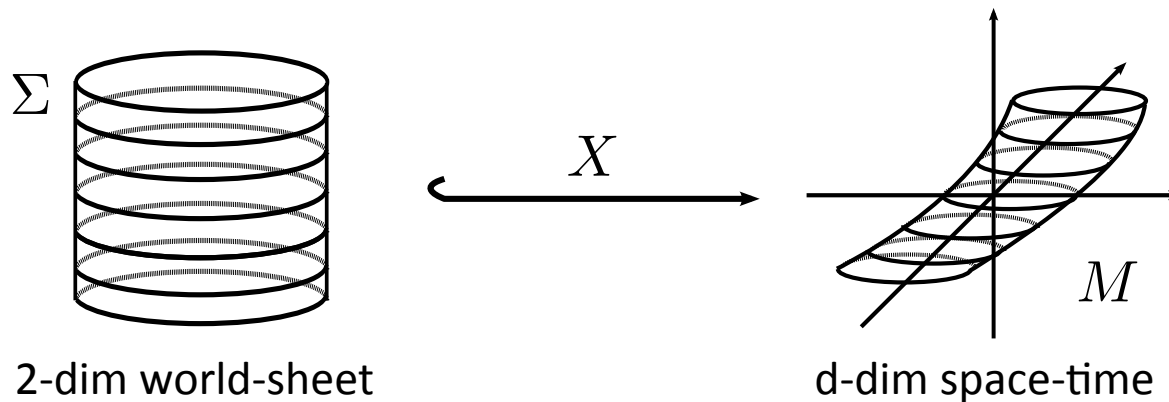


d-dim spacetime

# MOTIVATION

The **2-dim. sigma model**  $S$  describes motion of closed string in  $d$ -dim. space-time: geodesic motion of a string on Riemann-Cartan space  $M$

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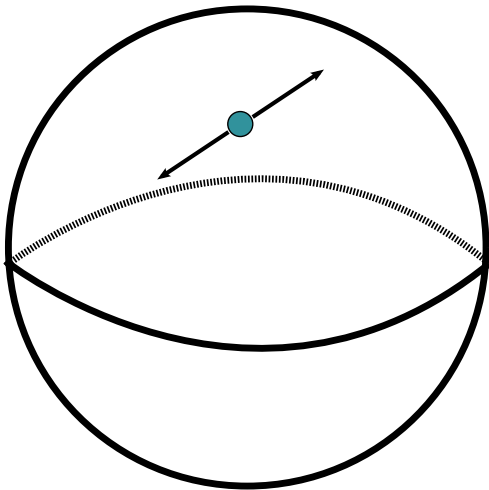


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point particle



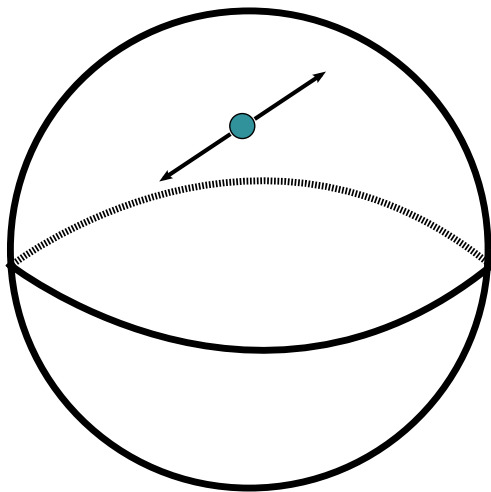
closed string

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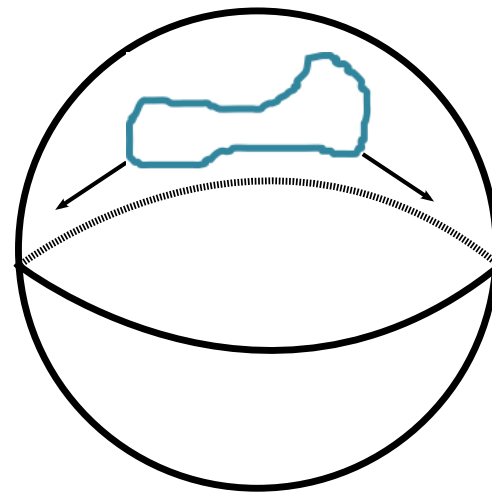
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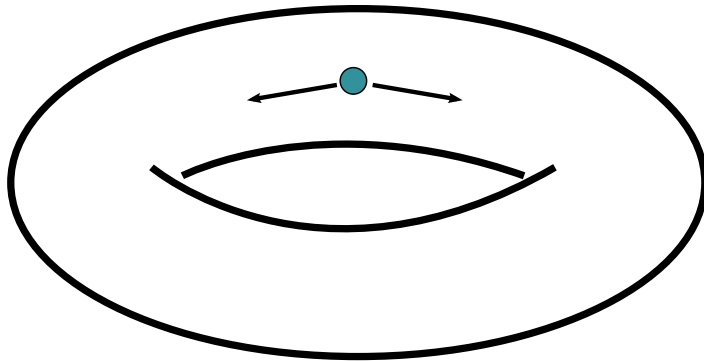
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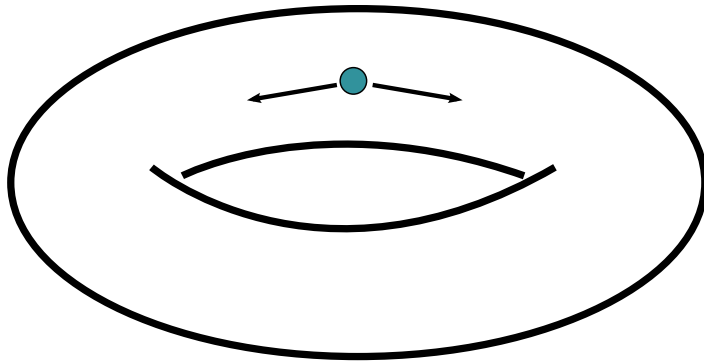


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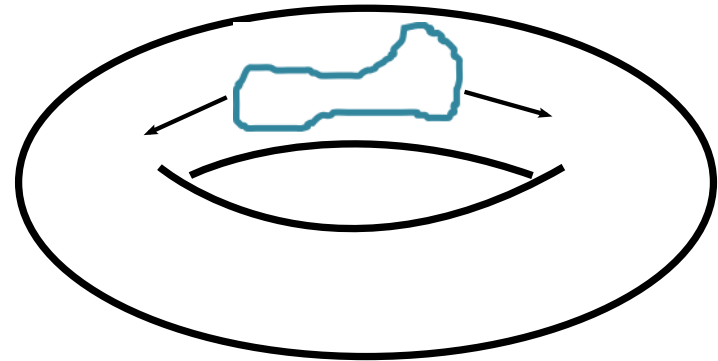
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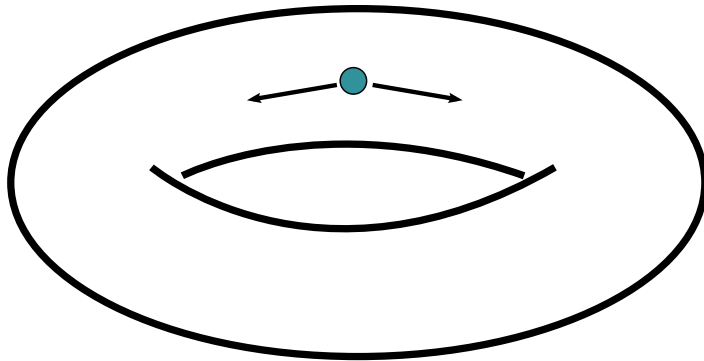


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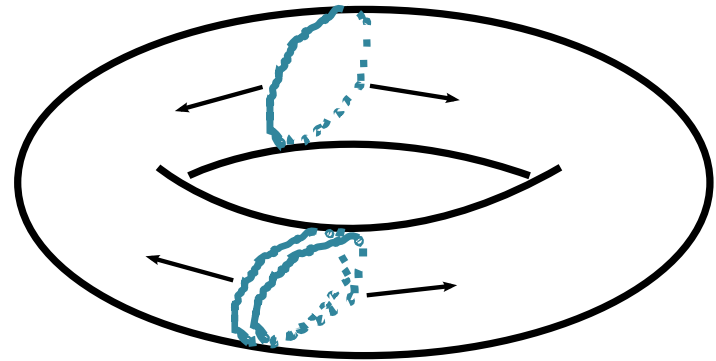
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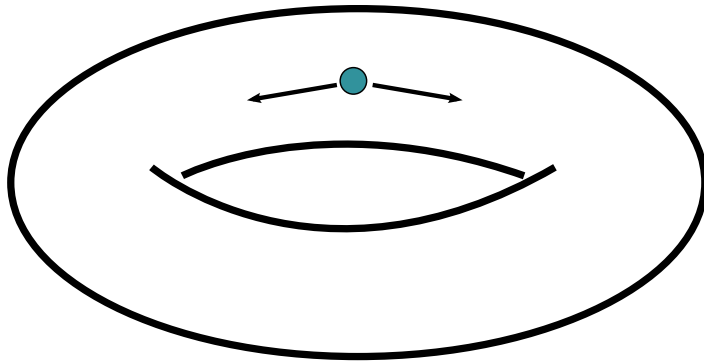


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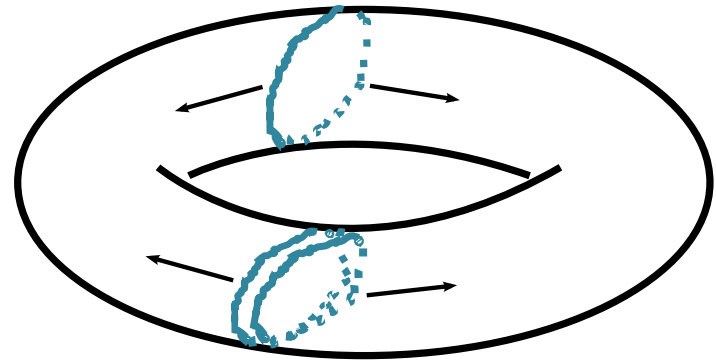
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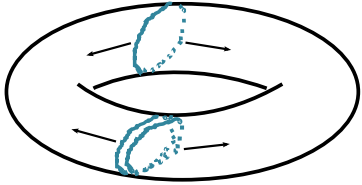
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String probes spacetime very differently – it can wind around certain directions

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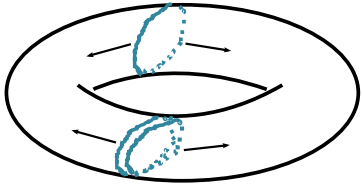
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String **probes** spacetime very differently –  
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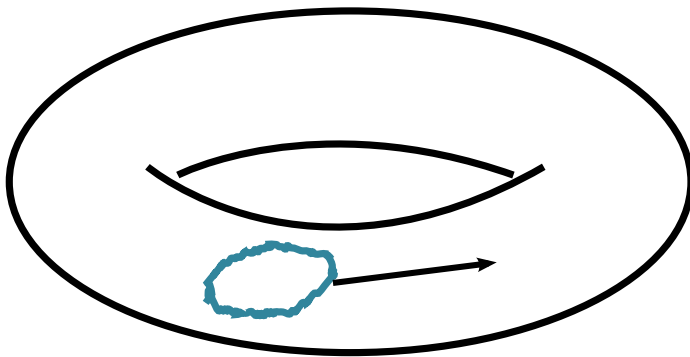
There is a specialty:

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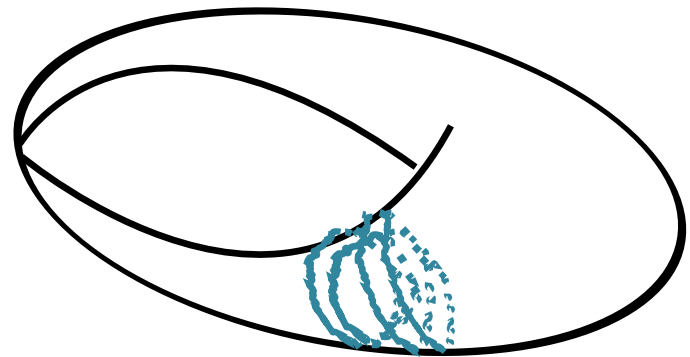
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momentum  $m$  and winding  $n$

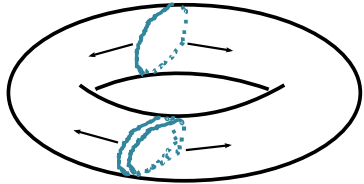
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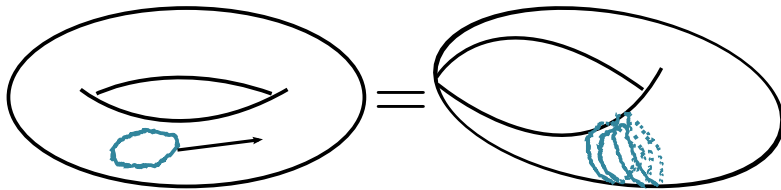
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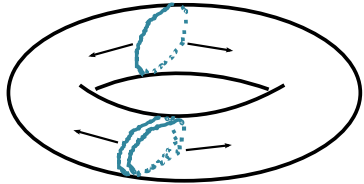
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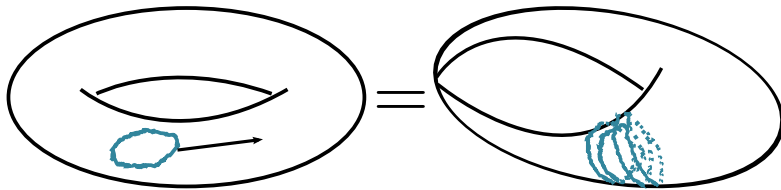
Very different spacetimes can give same  
theory: **T-duality**

# MOTIVATION

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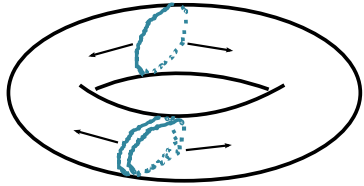
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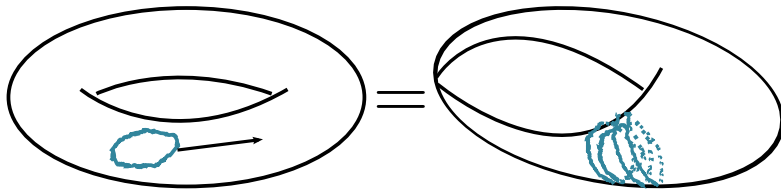
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$$(G, B) \iff (\tilde{G}, \tilde{B})$$

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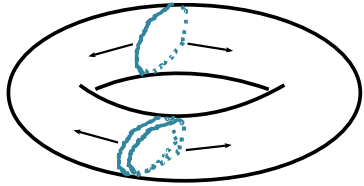
$$(G, B) \iff (\tilde{G}, \tilde{B})$$

related by the **Buscher rules**

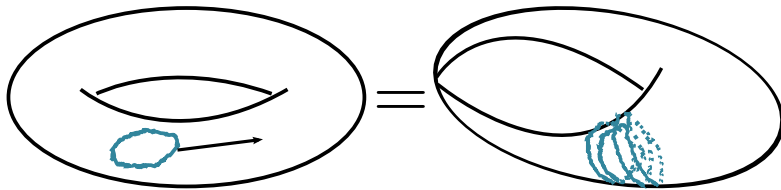
$$\begin{aligned} \tilde{G}_{00} &= \frac{1}{G_{00}}, & \tilde{G}_{0m} &= \frac{B_{0m}}{G_{00}}, & \tilde{G}_{mn} &= G_{mn} - \frac{G_{m0}G_{n0} + B_{m0}B_{0n}}{G_{00}} \\ \tilde{B}_{0m} &= \frac{G_{0m}}{G_{00}}, & \tilde{B}_{mn} &= B_{mn} - \frac{G_{m0}B_{0n} + B_{m0}G_{0n}}{G_{00}} \end{aligned}$$



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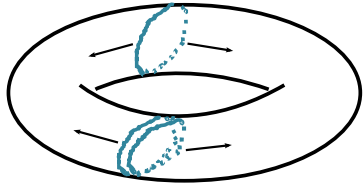
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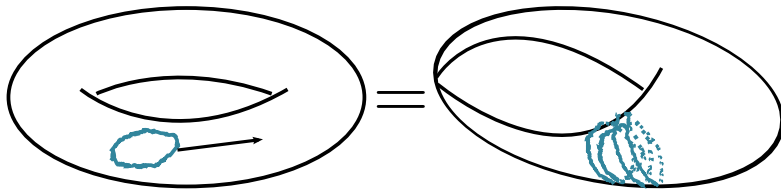
related by the **Buscher rules**  
in terms of **generalized metric**

$$\mathcal{H}(G, B) = \begin{pmatrix} G - BG^{-1}G & BG^{-1} \\ -G^{-1}B & G^{-1} \end{pmatrix}$$

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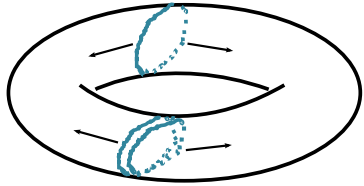
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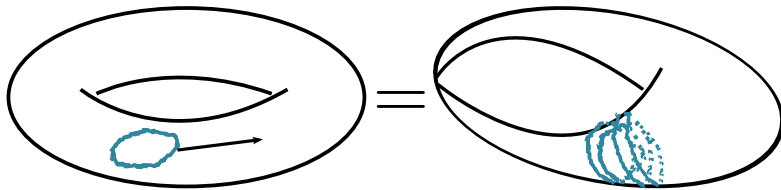
$$\mathcal{H}(\tilde{G}, \tilde{B}) = \mathcal{H}'(G, B) = \mathcal{T}^t \mathcal{H}(G, B) \mathcal{T} \quad \text{with} \quad \mathcal{T} = \left( \begin{array}{ccc|ccc} 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 \\ \hline 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 \end{array} \right)$$

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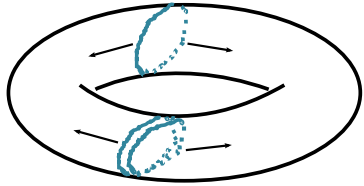
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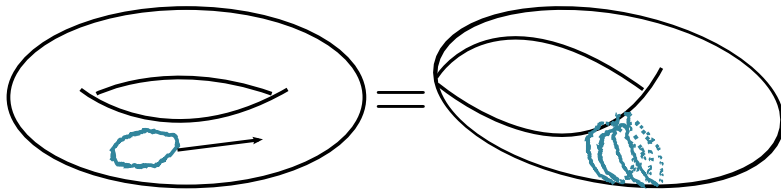
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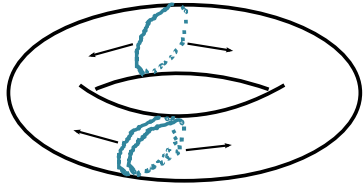
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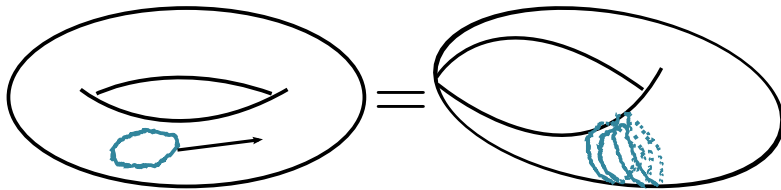
related by the **Buscher rules**

Simple Example:  $\mathbb{R}^2 \times \mathbb{R}^{1,d-3}$

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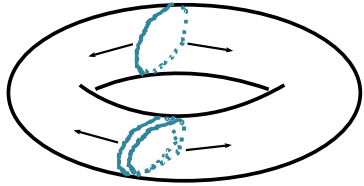
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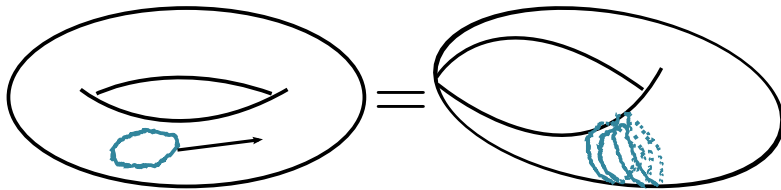
$$ds^2 = dr^2 + r^2 d\phi^2 + d\mathbf{x}^2$$

$$R = 0$$

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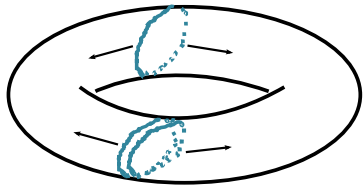
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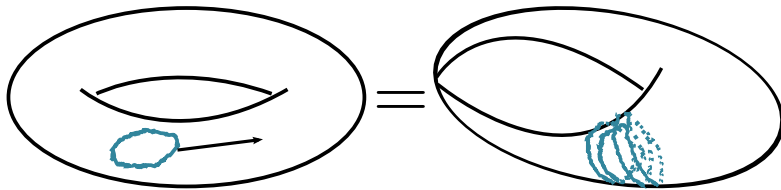
$$ds^2 = dr^2 + r^2 d\phi^2 + d\mathbf{x}^2 \quad \xleftrightarrow{T_\phi} \quad d\tilde{s}^2 = dr^2 + \frac{1}{r^2} d\tilde{\phi}^2 + d\mathbf{x}^2$$

$$R = 0 \quad \quad \quad \tilde{R} = -\frac{2}{r^2}$$

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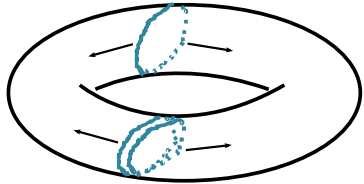
Simple Example:  $\mathbb{R}^2 \times \mathbb{R}^{1,d-3}$

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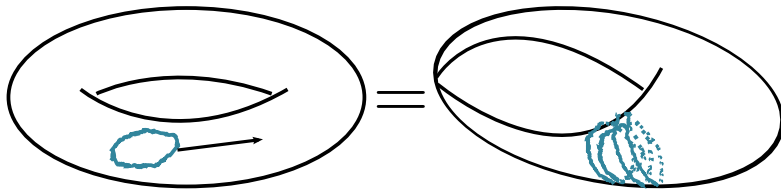
$$R = 0 \quad \quad \quad \tilde{R} = -\frac{2}{r^2} \quad \text{singular at origin}$$

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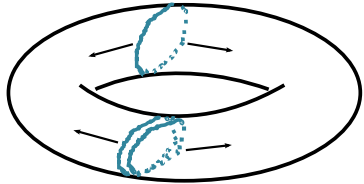
Very different spacetimes can give same  
theory: **T-duality**

T-duality can **change topology** and cause troubles

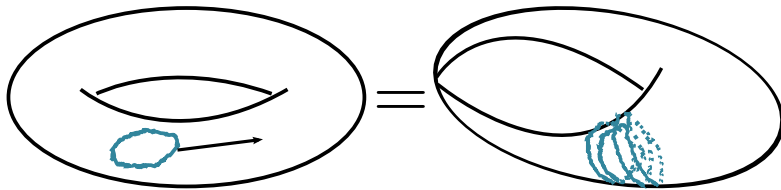


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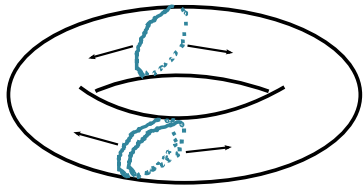
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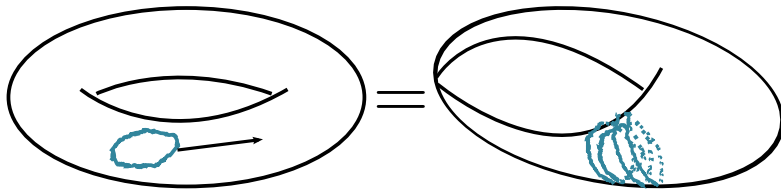
$B$  transforms under **gauge transformations**  $B \rightarrow B + d\xi$

Better Object: **flux**  $H = dB$

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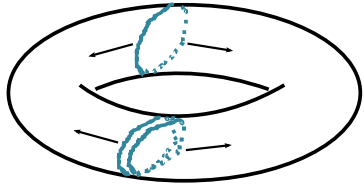
Better Object: **flux**  $H = dB$

T-duality relates backgrounds with different metric and flux; e.g. for 3-torus

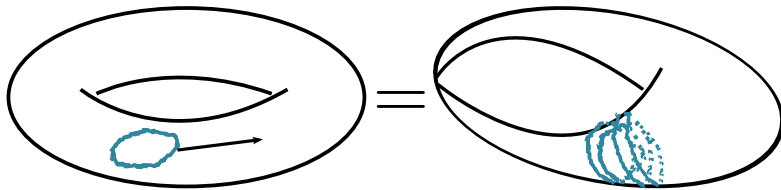
$$H_{abc} \xleftrightarrow{T_1} f^a{}_{bc} \xleftrightarrow{T_2} Q^{ab}{}_c \xleftrightarrow{?} R^{abc}$$

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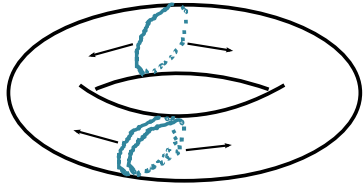


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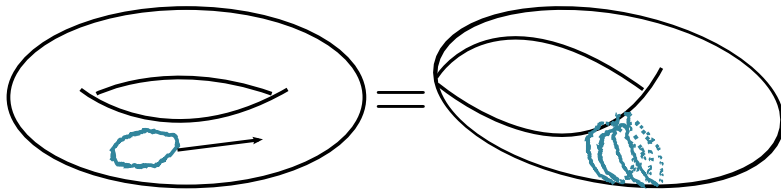
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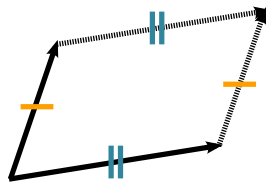


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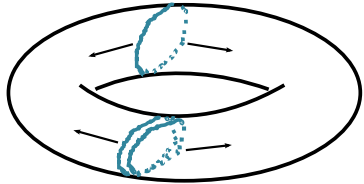
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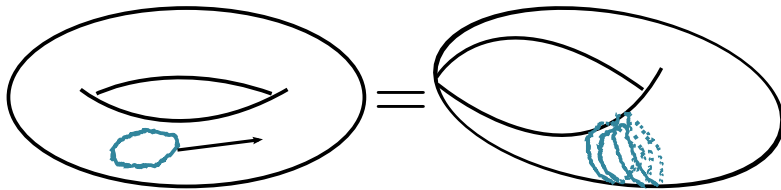
torus



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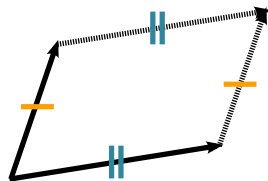


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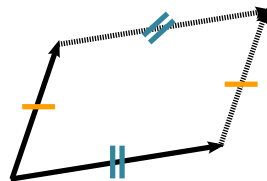
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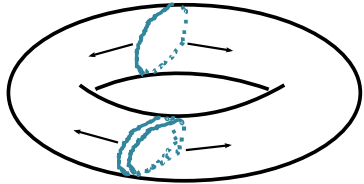
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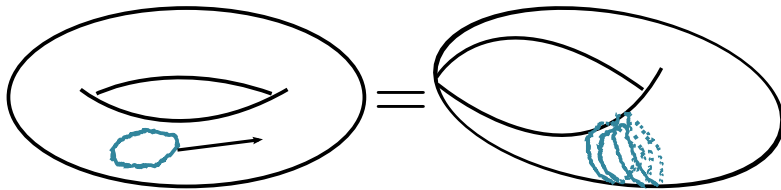
twisted torus



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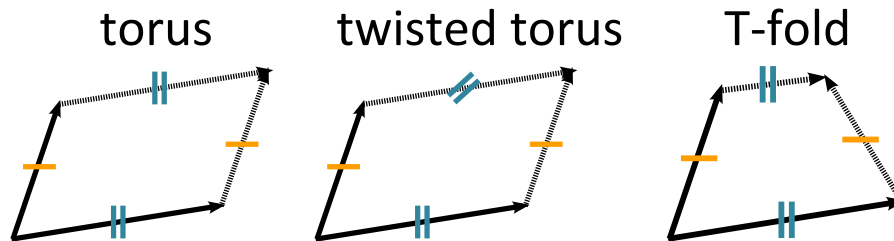
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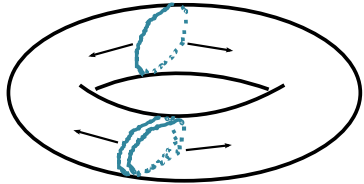
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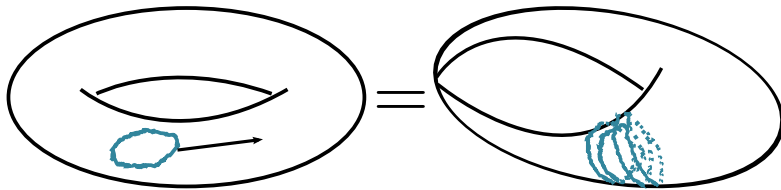
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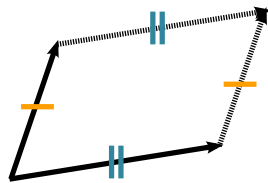


Very different spacetimes can give same  
theory: **T-duality**

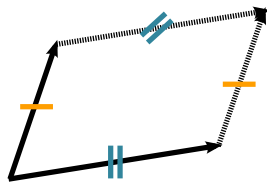
T-duality can **change topology** and cause troubles

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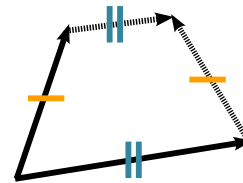
torus



twisted torus

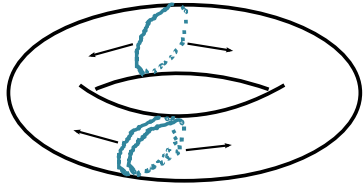


T-fold

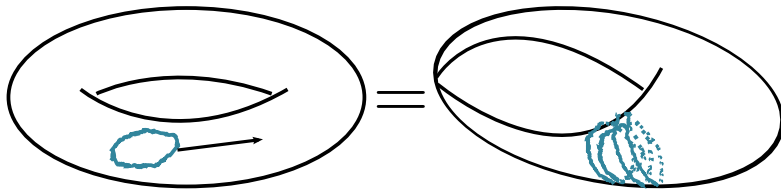


unknown

# MOTIVATION



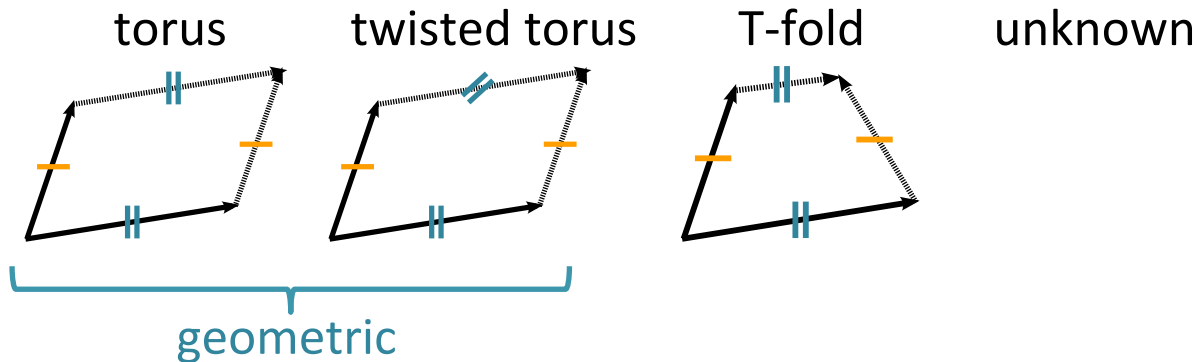
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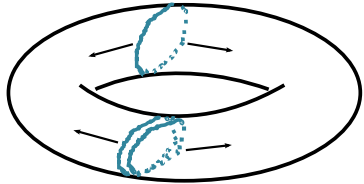
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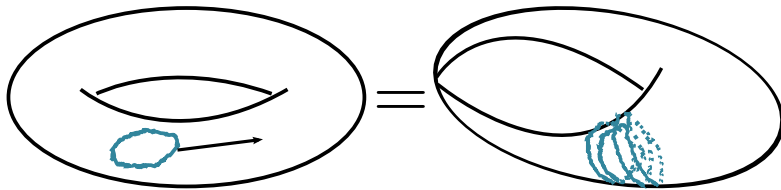




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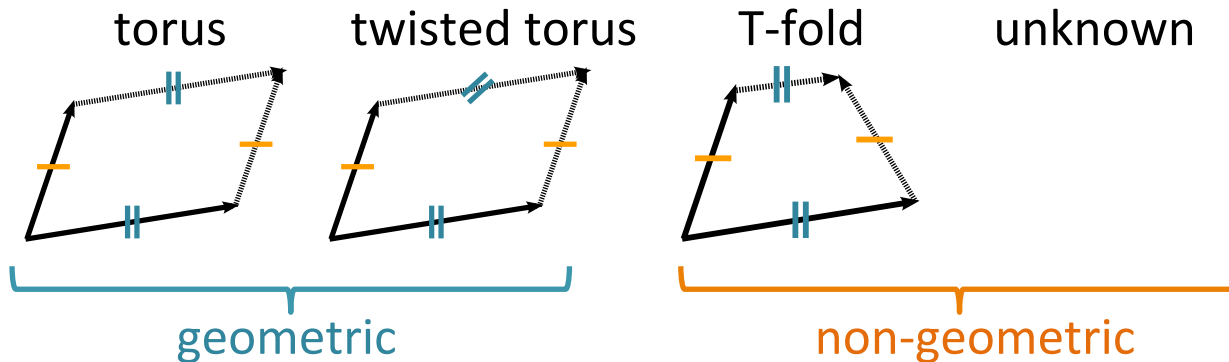
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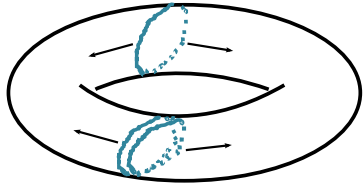
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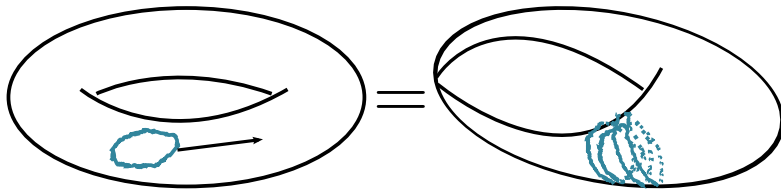
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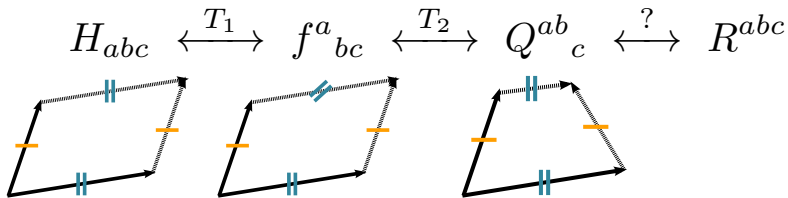
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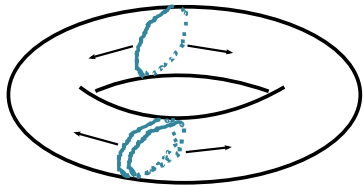


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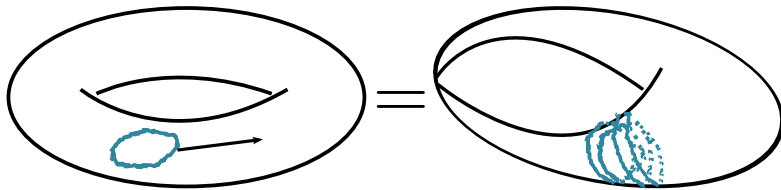


**Non-geometric** backgrounds

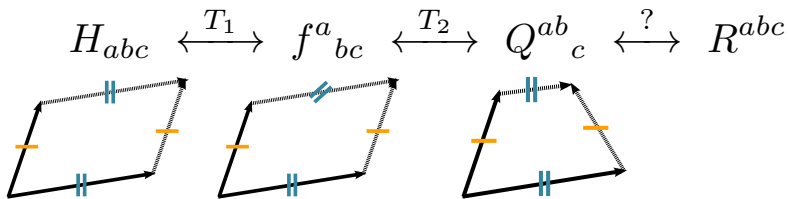
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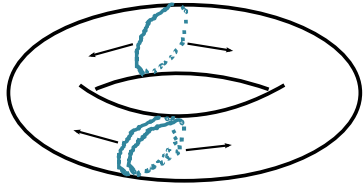


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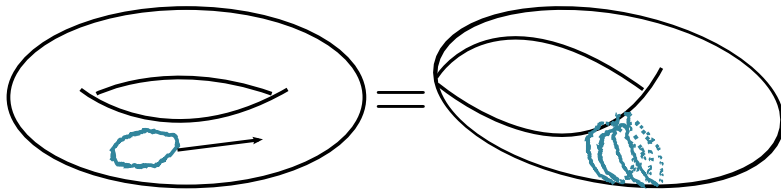
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$$S_{\text{eff}}^d(G, B) = \frac{1}{2\kappa^2} \int d^d x \sqrt{|G|} e^{-2\phi} \left( R - \frac{1}{12} H_{abc} H^{abc} \right) + \dots$$

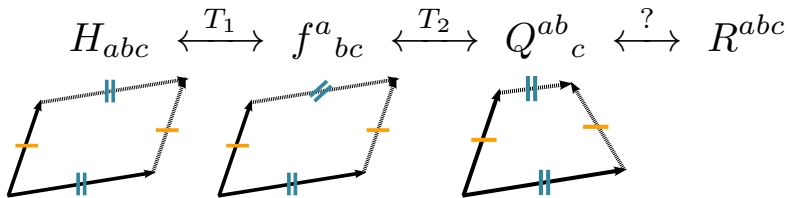
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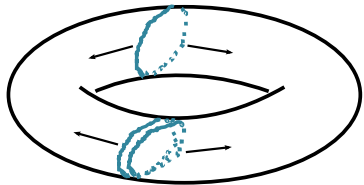


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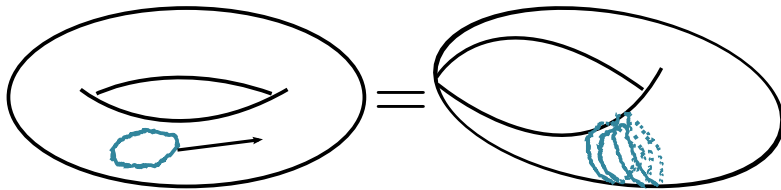
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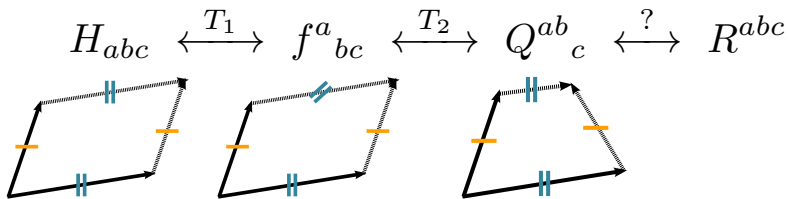
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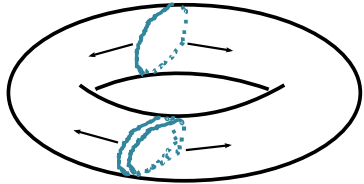
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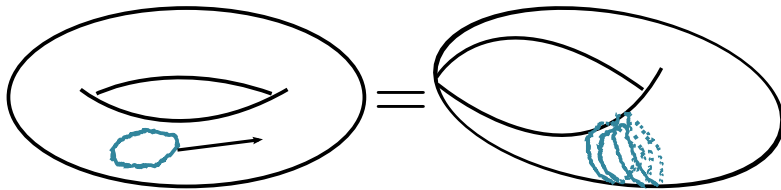
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$H = dB$

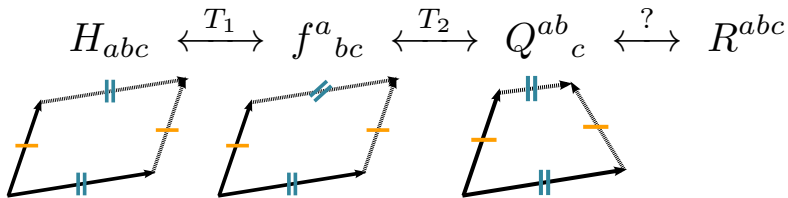
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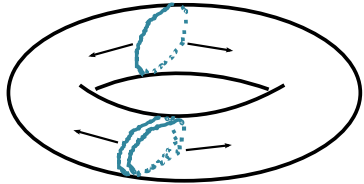
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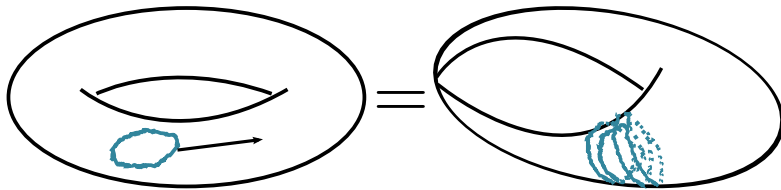
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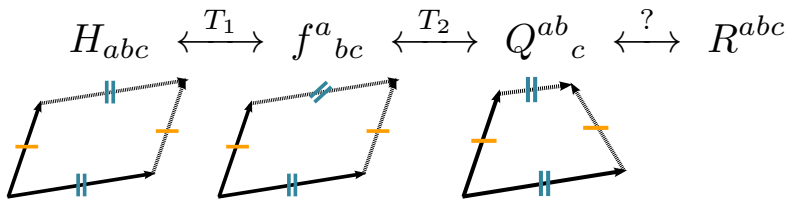
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T-duality and non-geometry can be understood via ***O(d,d)-transformations***

# all at once: GENERALIZED GEOMETRY

---

$$H_{abc} \xleftrightarrow{T_1} f^a{}_{bc} \xleftrightarrow{T_2} Q^{ab}{}_c \xleftrightarrow{?} R^{abc}$$

All possible positions of indices



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quasi-Lie algebroids describing  $f$ - and  $Q$ -flux with deviation from Jacobi- identity given by  $H$ - and  $R$ -flux respectively

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$$[a, [b, c]] + [c, [a, b]] + [b, [c, a]] = 0$$

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The combination of both describes all fluxes: Dorfmann bracket

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This defines a Courant algebroid

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Yields algebra  $\llbracket e_a, e_b \rrbracket = f^c{}_{ab} e_c + H_{abc} e^c$

$$\llbracket e_a, e^b \rrbracket = Q^{bc}{}_a e_c - f^b{}_{ac} e^c$$

$$\llbracket e^a, e^b \rrbracket = R^{abc} e_c + Q^{ab}{}_c e^c$$

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?

There is a sub-structure allowing for gravity theory: [Lie algebroids](#)

# O(d,d) and LIE ALGEBROIDS

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$$\mathcal{H}(G, B) = \begin{pmatrix} G - BG^{-1}G & BG^{-1} \\ -G^{-1}B & G^{-1} \end{pmatrix}$$



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Non-geometric backgrounds can be patched-up using this construction

# OUTLOOK

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There is a nice geometric structure behind stringy fluxes

Various interesting routes:

- better understanding of **non-geometry**: classical aspects  
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- conformal field theory approaches: T-duality simple
- **non-associative geometry**

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Thank You