

LHC phenomenology – made in Munich

“Munich MPI Project Review 2014”.

Jan Winter

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Max Planck Institute for Physics

Munich

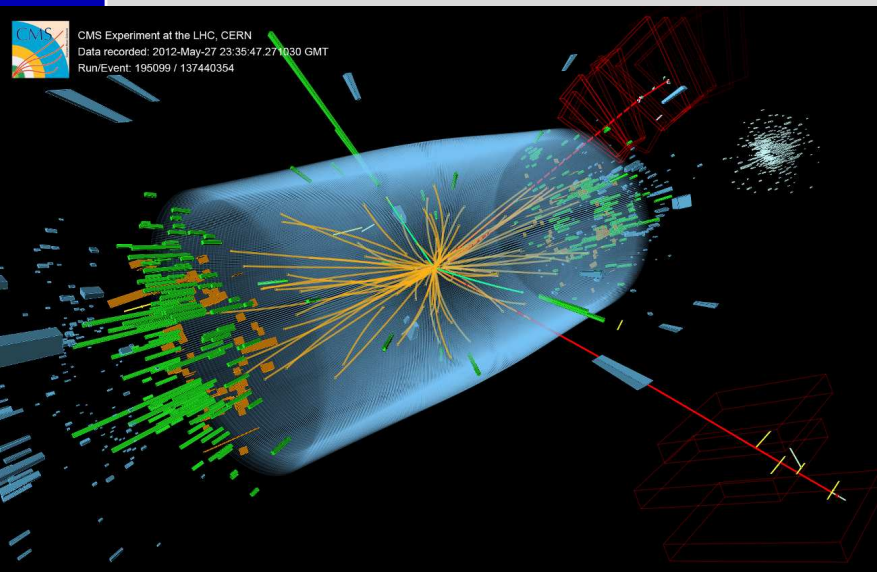


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[Munich MPI Project Review 2014]

Jan Winter

– MPP Munich, Germany –

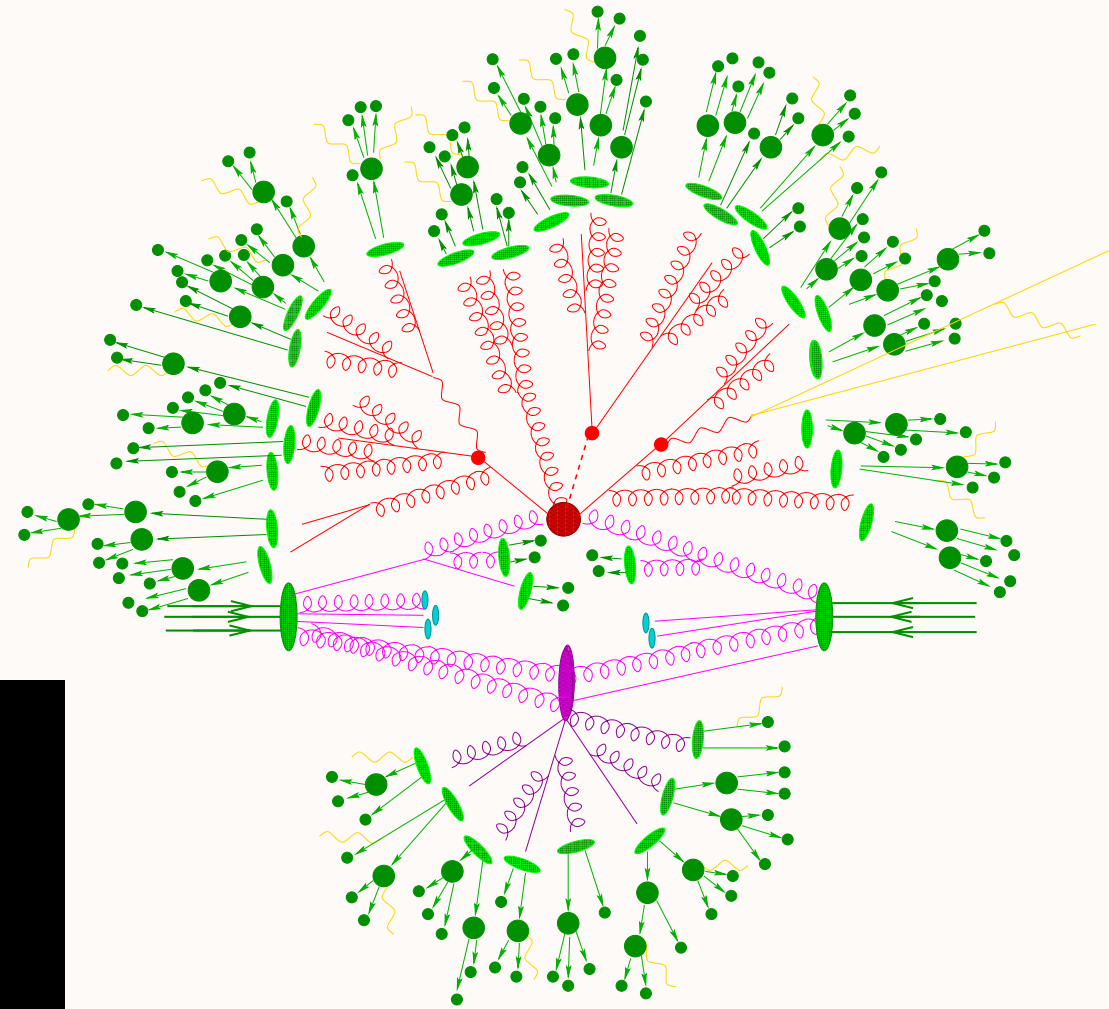
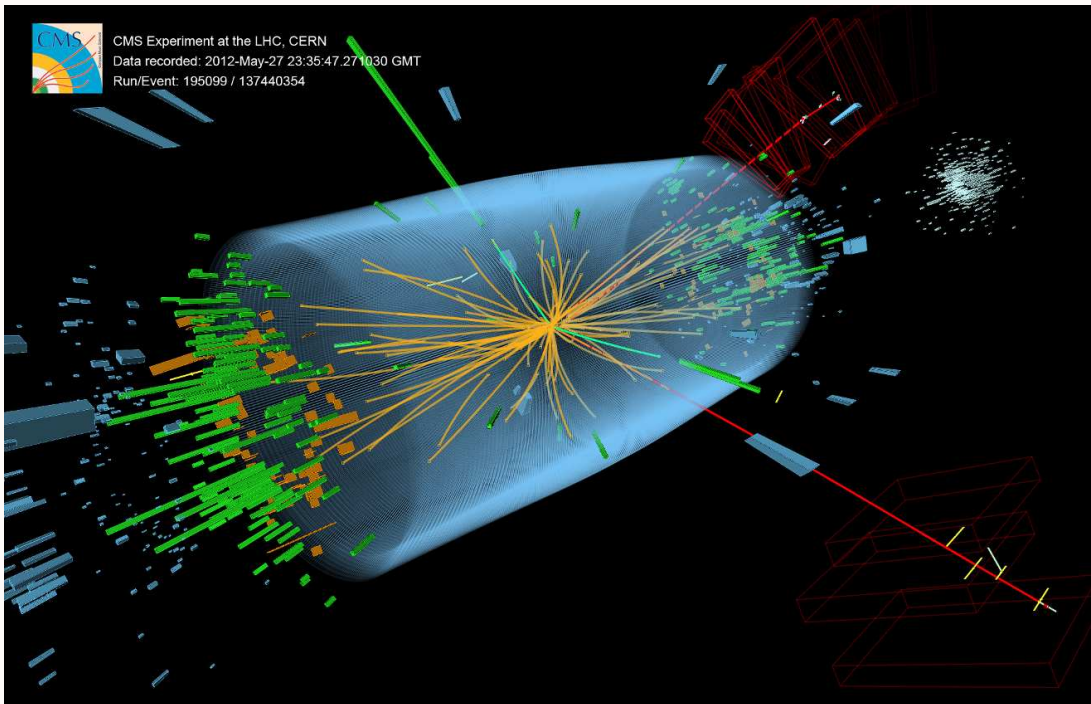


➔ *Introduction:
Monte Carlo event generation.*

- *The TH Pheno Group at the MPI.*
- *Examples of SM and BSM applications.*
- *Summary.*

➔ *Visualizing the events.*

- As we have seen already, this is how a four-lepton event looks like in the CMS detector.



- Illustration of the various phases in the evolution of an event. To simulate an event, we use **Monte Carlo techniques**.

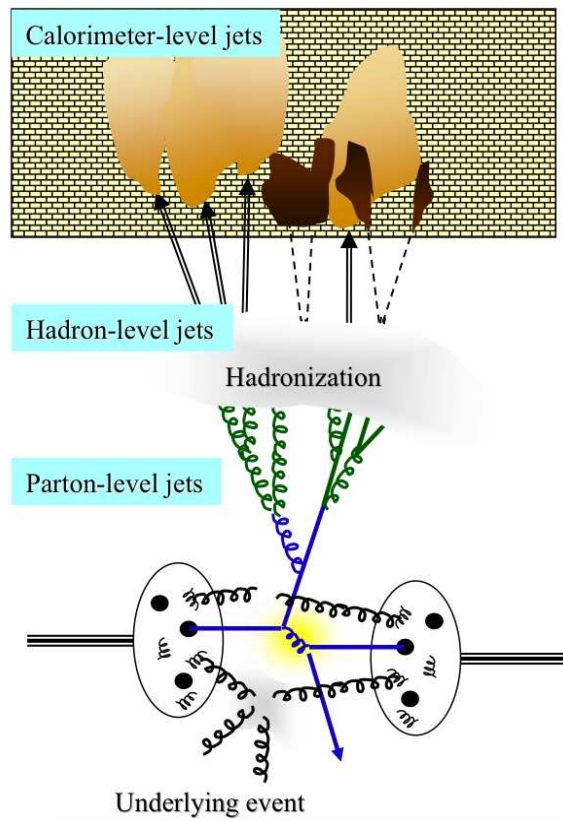


Event evolution

➔ *Incoming protons ready for head-on collision.*

Protons are made of (valence/sea) quarks, (sea) antiquarks and gluons (= partons). It is the individual partons that collide, with only a fraction of the proton's energy being available to a single parton collision.

Quark, gluons, partons? We only ever see JETS.

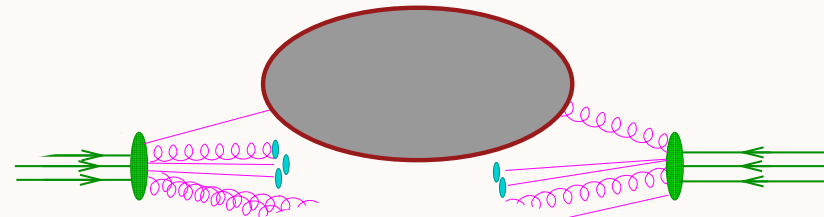
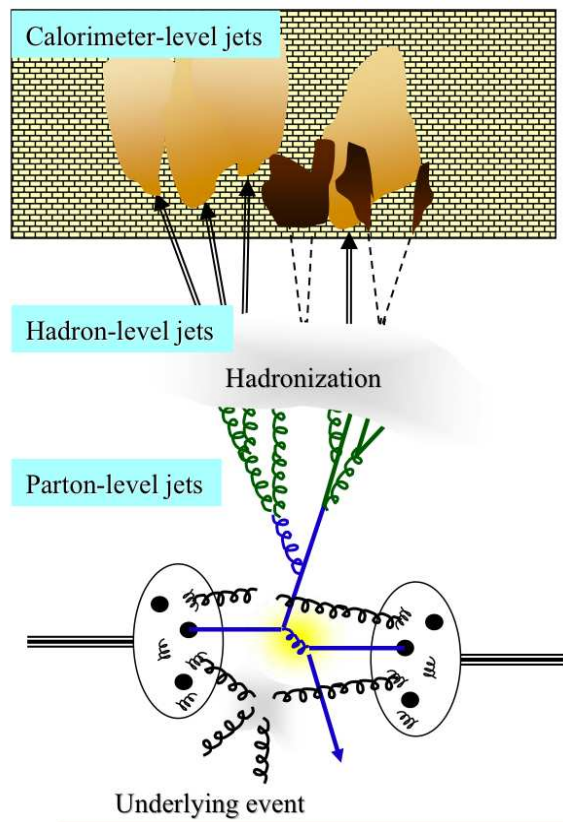


Event evolution

→ *The hard interaction.*

Partonic interaction calculated from “phase space times matrix element squared”. Convolution with PartonDensityFunctions. Numerical integration method using multidimensional Monte Carlo.

Quark, gluons, partons? We only ever see JETS.

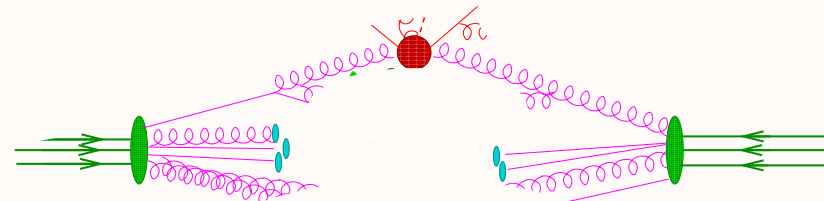
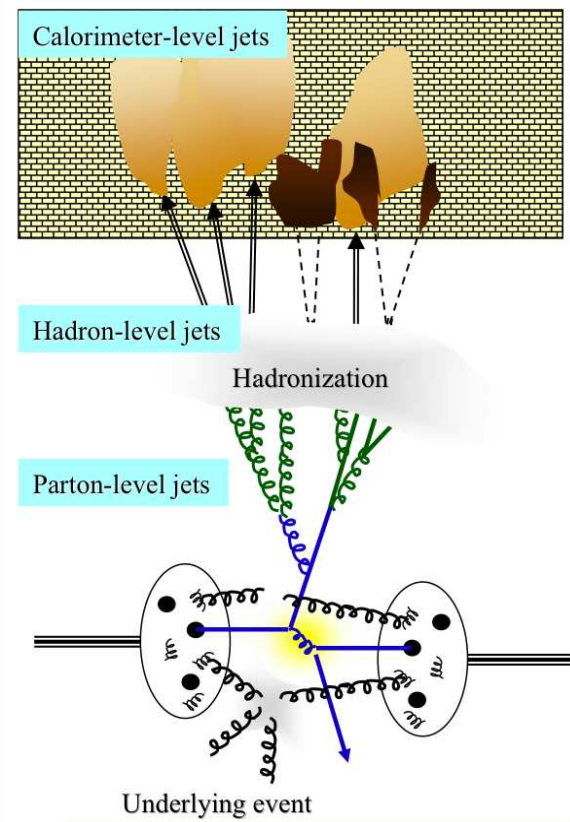


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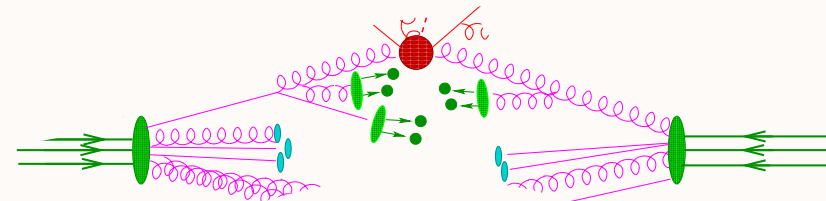
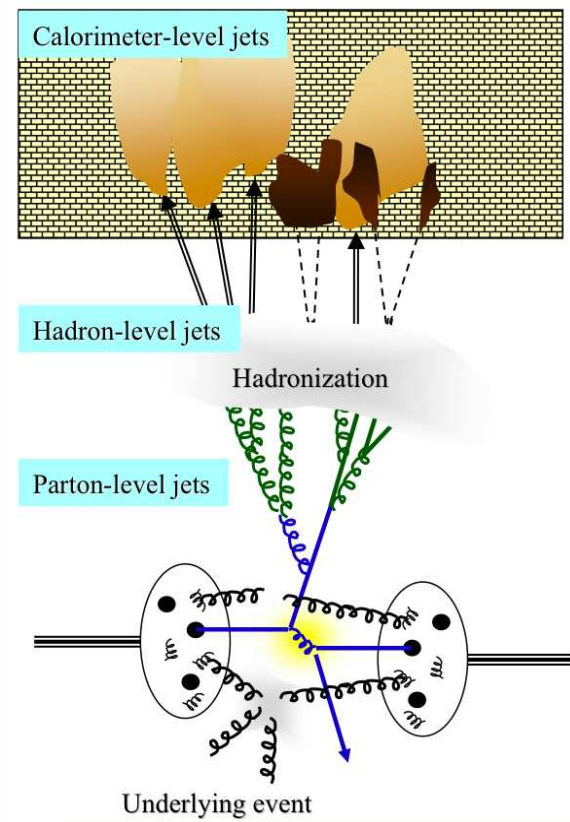


Event evolution

→ *Initial state radiation.*

Multiple soft/collinear emission from incoming partons. Backward evolution since both ends of evolution are fixed.

Quark, gluons, partons? We only ever see JETS.

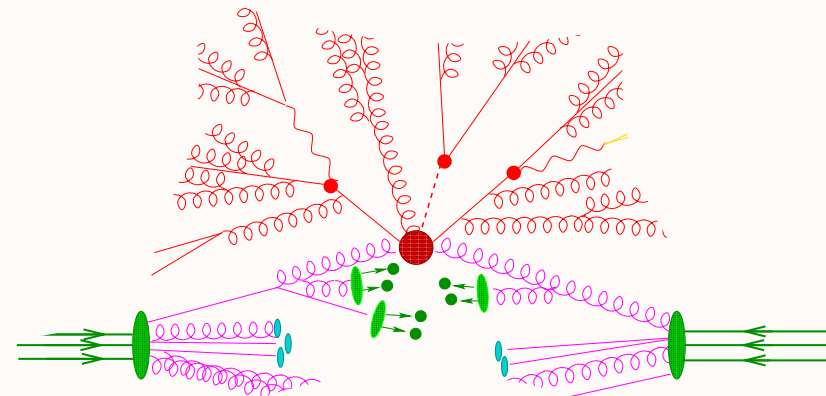
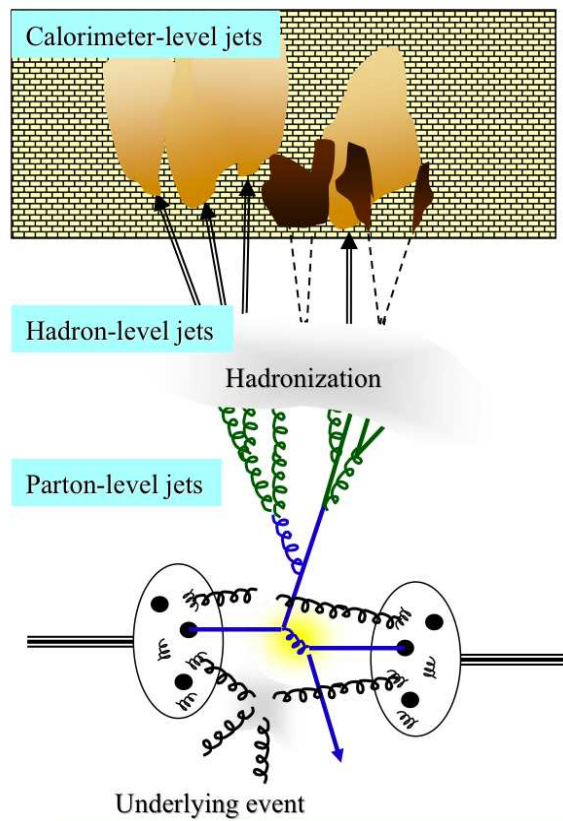


Event evolution

→ *Final state radiation.*

In pQCD, probability that parton emits a gluon is $\sim \alpha_s \frac{dE}{E} \frac{d\theta}{\theta}$, which diverges for small gluon energies E and/or small angles θ . A quark never survives unchanged, it always emits a (low-energy) gluon (at small angles). Each gluon radiates a further gluon, and so forth. Numerically treated through Metropolis-type algorithms.

Quark, gluons, partons? We only ever see JETS.

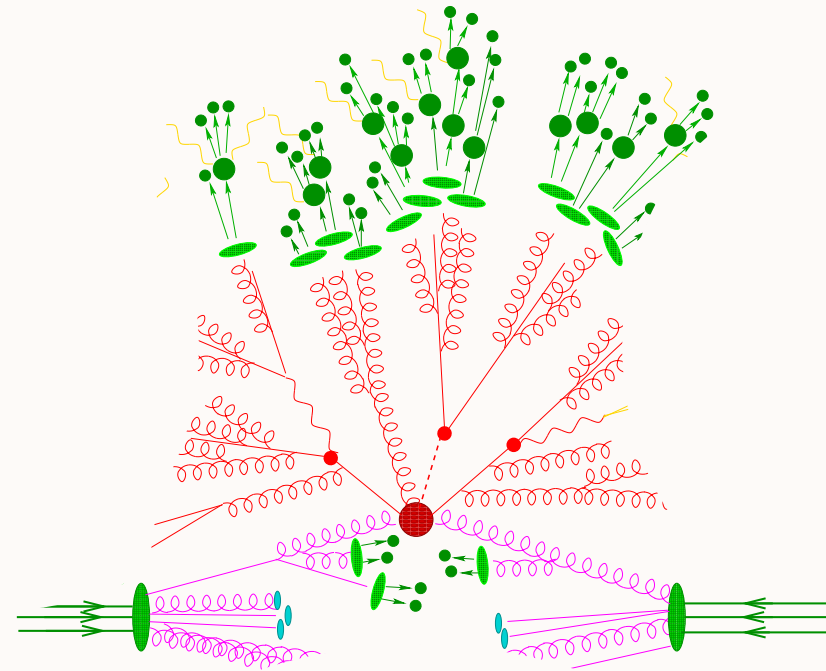
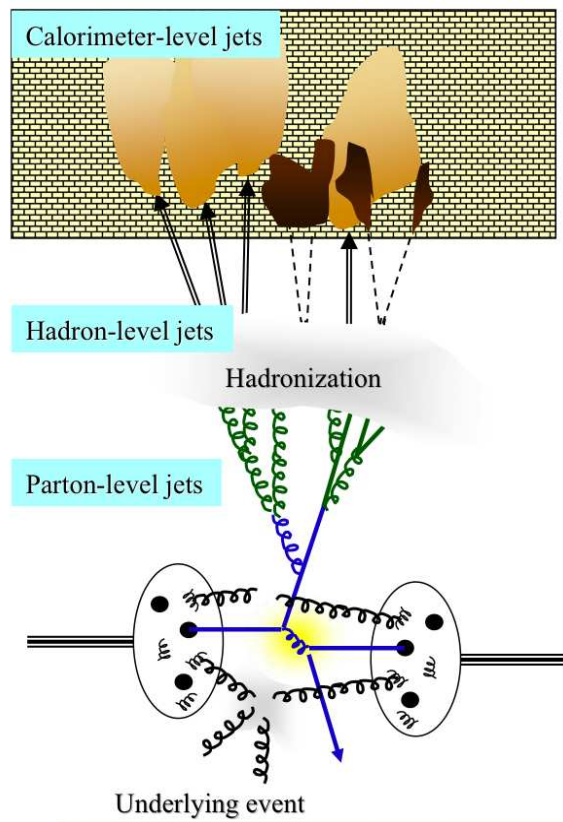


Event evolution

→ *Parton-to-hadron transition of jets.*

Phenomenological models tuned to data based on universality of the parton-to-hadron transition setting in at scales of the order of 1 GeV.

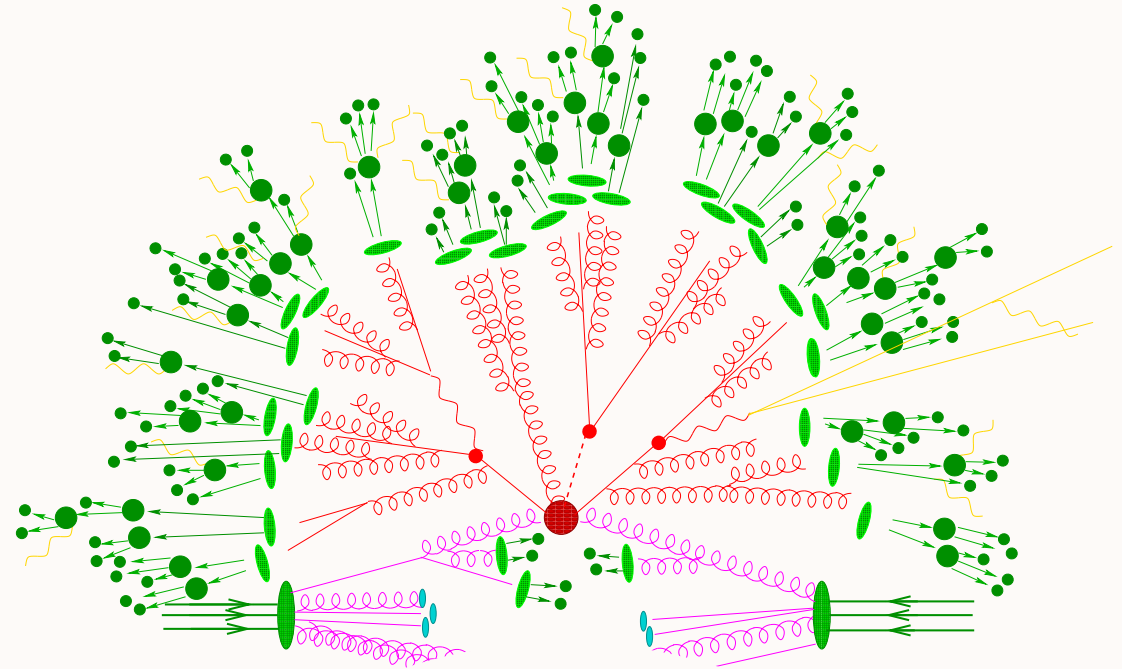
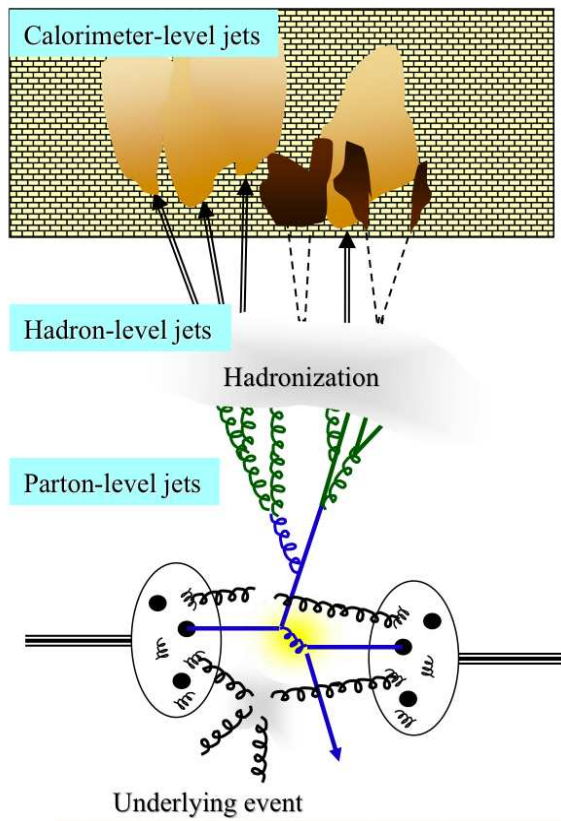
Quark, gluons, partons? We only ever see JETS.



Event evolution

- ➔ *Hadronization and decay of primary hadrons.*
- Photon emissions.*

Quark, gluons, partons? We only ever see JETS.

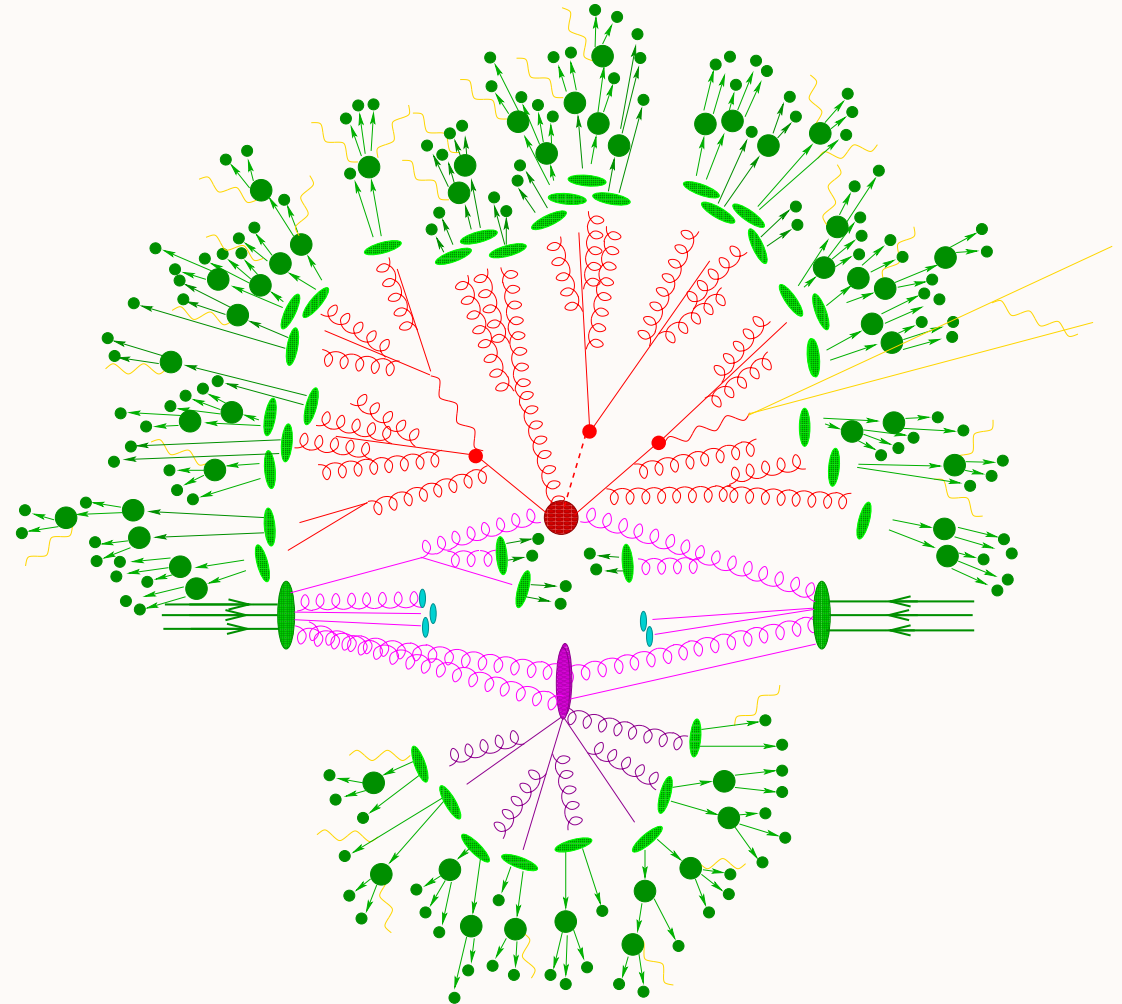
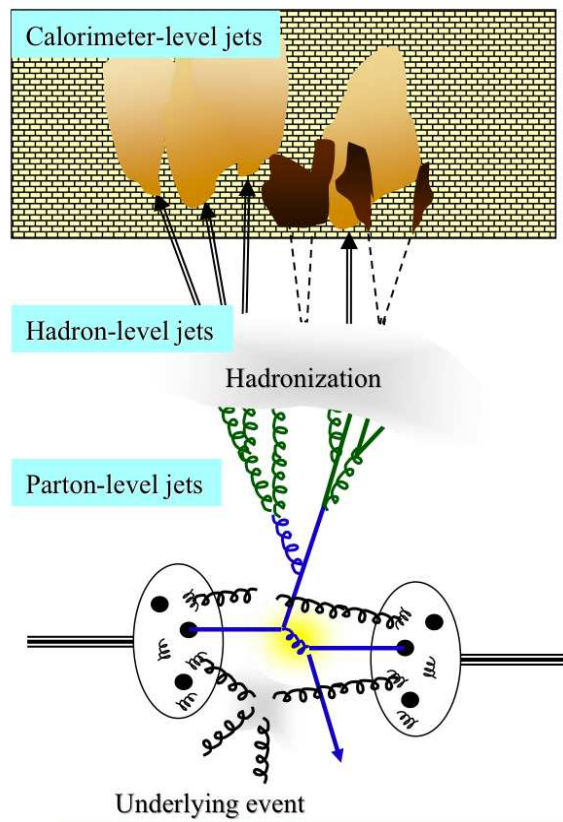


Event evolution

➔ *Multiple parton interactions/underlying event.*

Secondary interactions at lower energies compared to hard event.

Quark, gluons, partons? We only ever see JETS.



THX TO FRANK KRAUSS



Monte Carlo event generation

*Event generators are used to model **multi-hadron** final states of high-energy particle collisions. Factorization approach: divide jet simulation into different phases – use Monte Carlo methods.*

➔ *Perturbative Phases: [parton jets]*

● **Hard process/interaction (hard jet production)**

exact matrix elements $|\mathcal{M}|^2$

● **QCD bremsstrahlung (soft/coll multiple emissions)**

initial- and final-state parton showering

● **Multiple/Secondary interactions**

modelling the underlying event

➔ *Non-perturbative Phases: [jet confinement – particle jets]*

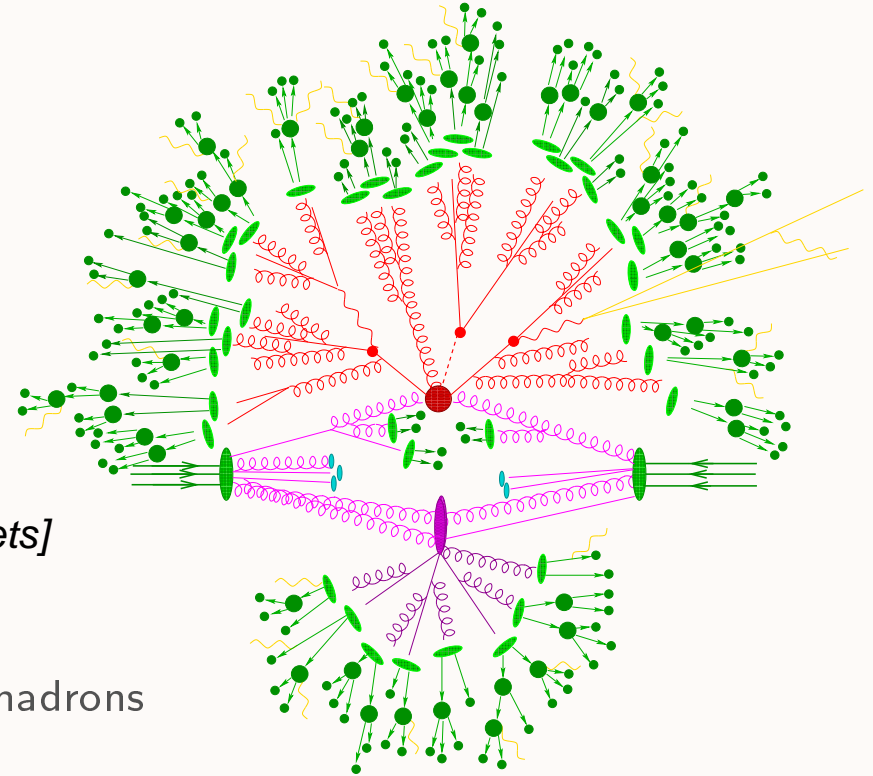
● **Hadronization**

phenomenological models to convert partons into primary hadrons

● **Hadron decays**

phase-space or effective models to decay unstable into stable hadrons as observed in detectors

➔ predictions at hadron level – comparable to experimental data if corrected for detector effects



Munich pheno group developments aim at improving the description of hard interactions.

Often these have the largest phenomenological impact, and we can study signatures at the parton level to already get most of the idea.

Luckily the other perturbative and the non-perturbative phases can be handled in a more universal manner. They are essential ingredients of Monte Carlo event generators, and we have learned (are learning) how to match parton showers with the higher-order calculations in a consistent way.



MPI Pheno Group Members.

➔ **Director:** *Wolfgang Hollik*

➔ **Staff members:** *Thomas Hahn, Gudrun Heinrich, Pierpaolo Mastrolia*

- **Postdoctoral Researchers:**

Stefano Di Vita, Stephen Jones, Matthias Kerner, Gionata Luisoni, Edoardo Mirabella, Jan Winter, Valery Yundin, Tom Zirke

- **PhD Students:**

Tao-Li Cheng, Hans van Deurzen, Stephan Hessenberger, Viktor Papara, Cyril Pietsch, Johannes Schlenk, Ulrich Schubert, Johann Felix von Soden-Fraunhofen

- **Congratulations! Successfully graduated this year (now off to their 1st PostDoc):**

Sophia Borowka (U Zurich), Sebastian Passehr (DESY Hamburg), Tiziano Peraro (Edinburgh)

- **Best of luck! Left the PostDoc team:** Nicolas Greiner

- **Welcome! New PostDocs to the team:** Stephen Jones, Matthias Kerner, Tom Zirke

➔ **Sofja Kovaleskaja Award Team, PI:** *Pierpaolo Mastrolia*

Gionata Luisoni, Valery Yundin, Hans van Deurzen, Tiziano Peraro, Ulrich Schubert



MPI Pheno Group Tools.

➔ *LHC phenomenology needs tools for precision calculations.*

- **GoSam** [= Golem + Samurai]

automated generation of one-loop amplitudes within & beyond the Standard Model ⇒ GoSam 2.0
[V.DEURZEN, GREINER, HEINRICH, LUISONI, MASTROLIA, MIRABELLA, PERARO, SCHLENK, V.SODEN-FRAUNHOFEN, ...]

- **Ninja** [C++ library] [MASTROLIA, MIRABELLA, PERARO]

interfaced with GoSam to provide faster & more stable integrand reduction for multi-scale processes

- **NJet** [C++ library] [YUNDIN, ...]

multi-parton one-loop MEs in massless QCD (many legs!), NLO predictions through interface to Sherpa

- **Sherpa** [Monte Carlo event generator (written in C++)] [WINTER, ...]

using OLPs, it provides NLO results, as well as matched/merged, fully showered/hadronized events

- **SecDec** [BOROWKA, HEINRICH, SCHLENK, ...]

numerical integration of multi-loop integrals (and other integrals)

- **FeynHiggs** [HAHN, HOLLIK, ...]

calculation of Higgs masses and cross sections in the MSSM

- **FeynArts & FormCalc** [HAHN, ...]

Feynman diagram and amplitude generation & simplification of Feynman diagrams up to one loop

- **Cuba & LoopTools** [HAHN, ...]

multi-dimensional numerical integration & library of one-loop integrals



MPI Pheno Group New Methods.

➔ *LHC phenomenology improves with our ability to handle more complex computations. This is a call for simplifying HEP calculations and trying out new methods. The new frontier is NNLO calculations (for everything), i.e. two-loop integral calculations attract much attention.*

- **Multi-loop recursive integrand reduction**

Reduction of n -denominator to $(n-1)$ -denominator integrand plus remainder (=residue).

Motivated by quest for unique mathematical framework for loop amplitudes and search for hidden properties in QFT. [MASTROLIA, MIRABELLA, PERARO, ...]

- **Method of Differential Equations for the evaluation of Feynman integrals**

Recipes to find the canonical form of systems of equations for master integrals.

(canonical form → separation of dependence on dimensional parameter from kinematics)

“Magnus and Dyson Series for Master Integrals”

[JHEP 1403 (2014) 082; DI VITA, MASTROLIA, MIRABELLA, SCHLENK, SCHUBERT, ...]

“Three-loop master integrals for ladder-box diagrams with one massive leg”

[JHEP 1409 (2014) 148; DI VITA, MASTROLIA, SCHUBERT, YUNDIN]



Examples of phenomenological applications and tool developments.

Thanks to all group members for providing me with plenty of material !!



GoSam

Slides from N. Greiner, G. Heinrich and G. Luisoni.



GoSAm-2.0: a tool for automated one-loop calculations within the Standard Model and Beyond

Gavin Cullen¹, Hans van Deurzen^{a,2},
Nicolas Greiner^{b,2}, Gudrun Heinrich^{c,2},
Gionata Luisoni^{d,2}, Pierpaolo Mastrolia^{e,2,3},
Edoardo Mirabella^{f,2}, Giovanni Ossola^{g,4,5},
Tiziano Peraro^{h,2}, Johannes Schlenk^{i,2},
Johann Felix von Soden-Fraunhofen^{j,2},
Francesco Tramontano^{k,6}

¹Deutsches Elektronen-Synchrotron DESY, Zeuthen, Germany

²Max-Planck-Institut für Physik, München, Germany

³Dipartimento di Fisica, Università di Padova, Italy

⁴New York City College of Technology, City University of New York, USA

⁵The Graduate School and University Center, City University of New York, USA

⁶Dipartimento di Scienze Fisiche, Università di Napoli and INFN, Sezione di Napoli, Italy

Abstract We present the version 2.0 of the program package GoSAm for the automated calculation of one-loop amplitudes. GoSAm is devised to compute one-loop QCD and/or electroweak corrections to multi-particle processes within and beyond the Standard Model. The new code contains improvements in the generation and in the reduction of the amplitudes, performs better in computing time and numerical accuracy, and has an extended range of applicability. The extended version of the “BinOth-Les-Houches-Accord” interface to Monte Carlo programs is also implemented. We give a detailed description of installation and usage of the code, and illustrate the new features in dedicated examples.

Keywords NLO calculations · automation · hadron colliders

PACS 12.38.-t · 12.38.Bx · 12.60.-i

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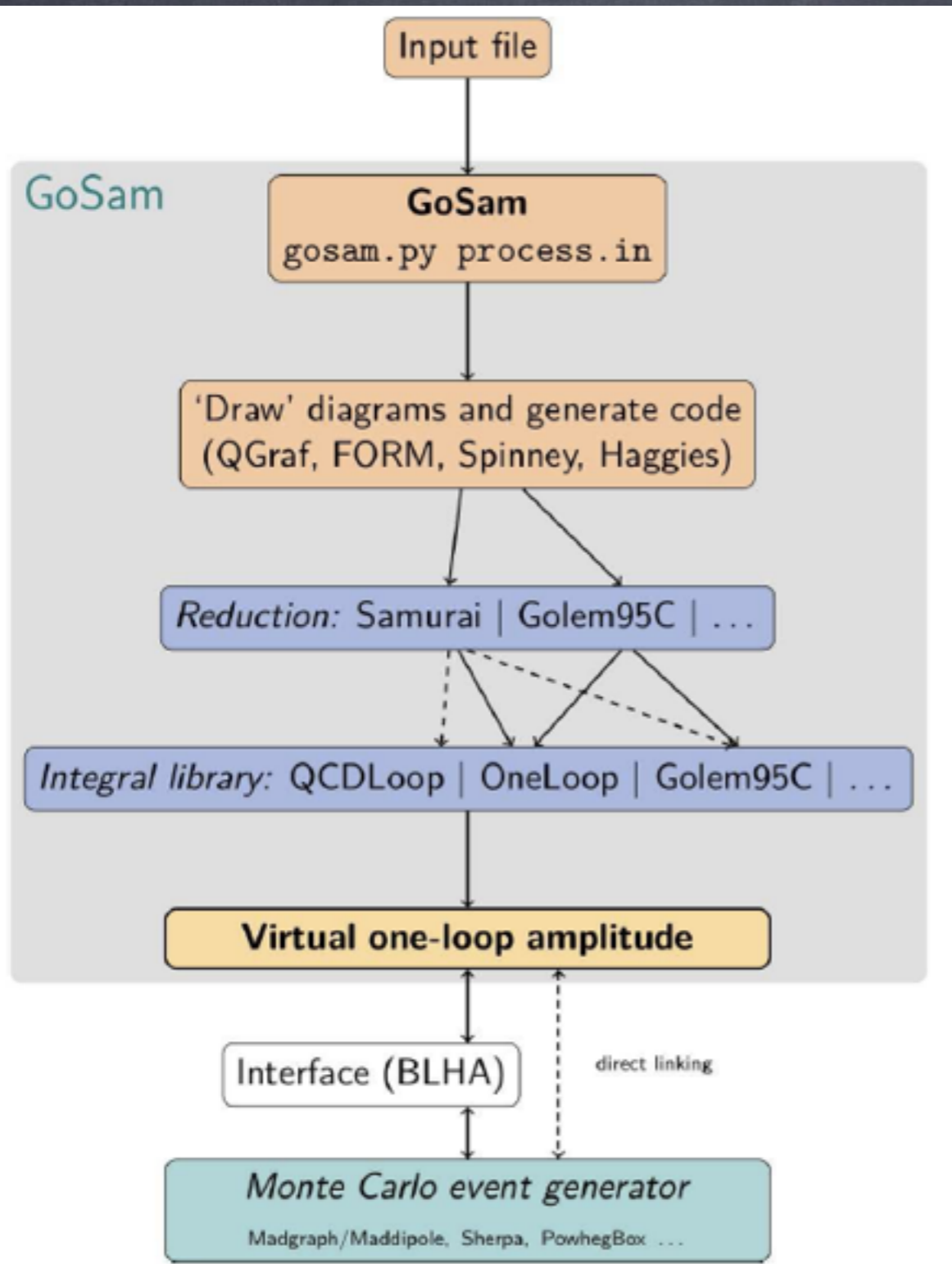
^hE-mail: peraro@mpp.mpg.de

ⁱE-mail: jschlenk@mpp.mpg.de

^jE-mail: jfsoden@mpp.mpg.de

^kE-mail: Francesco.Tramontano@cern.ch

GoSam for pedestrians



1. **Input file process.in**
Specify process ($ud^{\sim} \rightarrow W+gggg$) and options/parameters.
2. **Generation of Feynman diagrams**
QGraf, Form, Spinney
Translation into Fortran code
Haggies

3. **At runtime: Reduction of Tensor Integrals**
Samurai, Golem95
4. **Evaluation of scalar integrals**
QCDLoop, OneLoop, Golem95

5. **Interface with numerical Phase space integration for xsec and distributions**

GoSam-2.0

program available at

<http://gosam.hepforge.org>

very simple usage

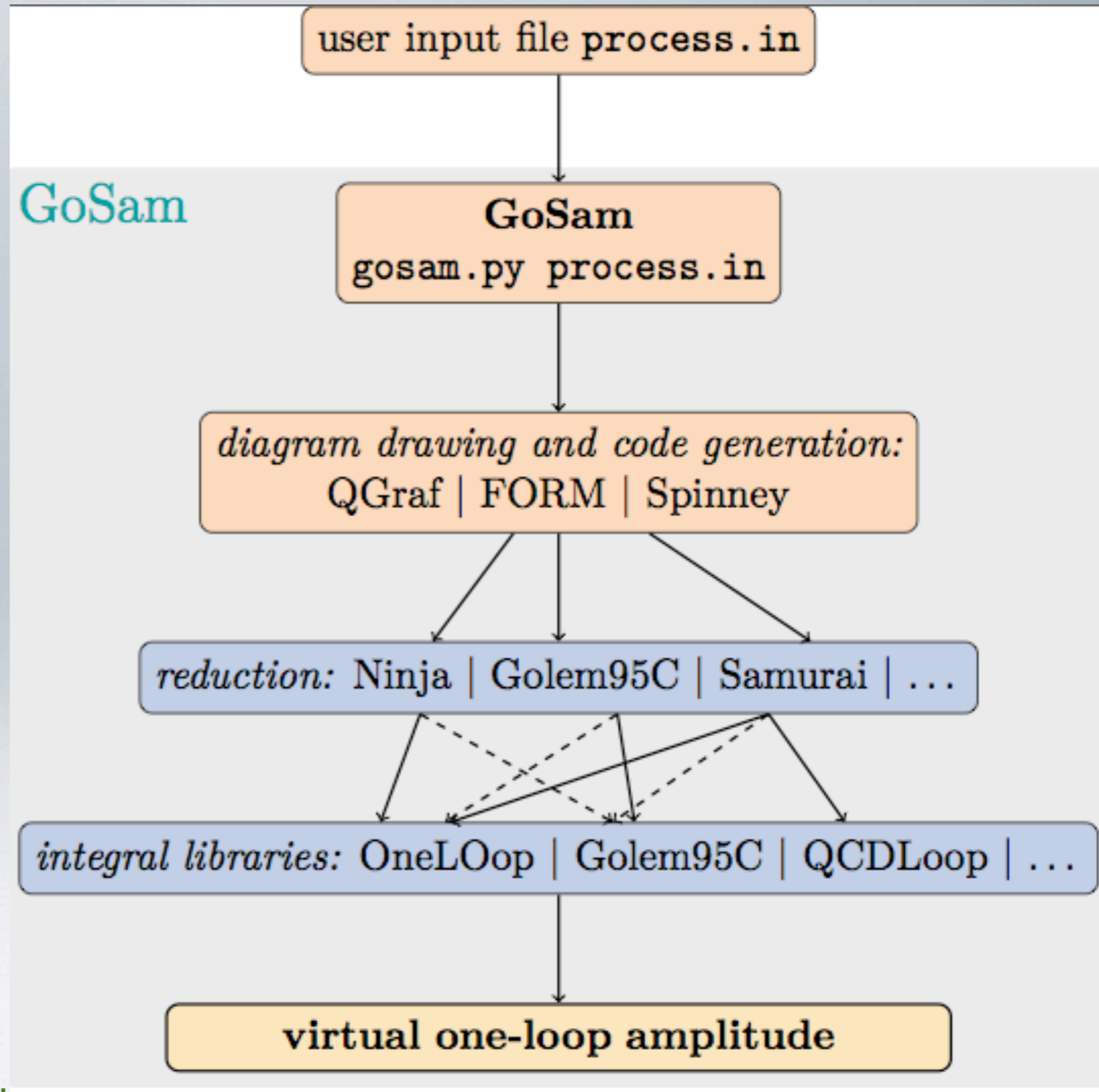
example input file for

$$e^+ e^- \rightarrow t \bar{t}$$

```
process_path=eett
in=      e+, e-
out=     t, t~
order=   gs, 0, 2
```

arXiv:1404.7096

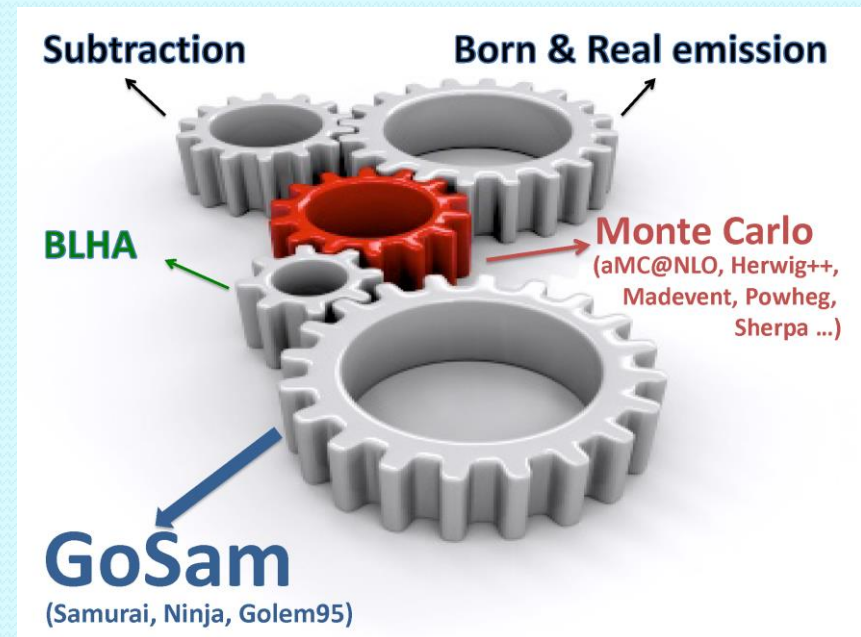
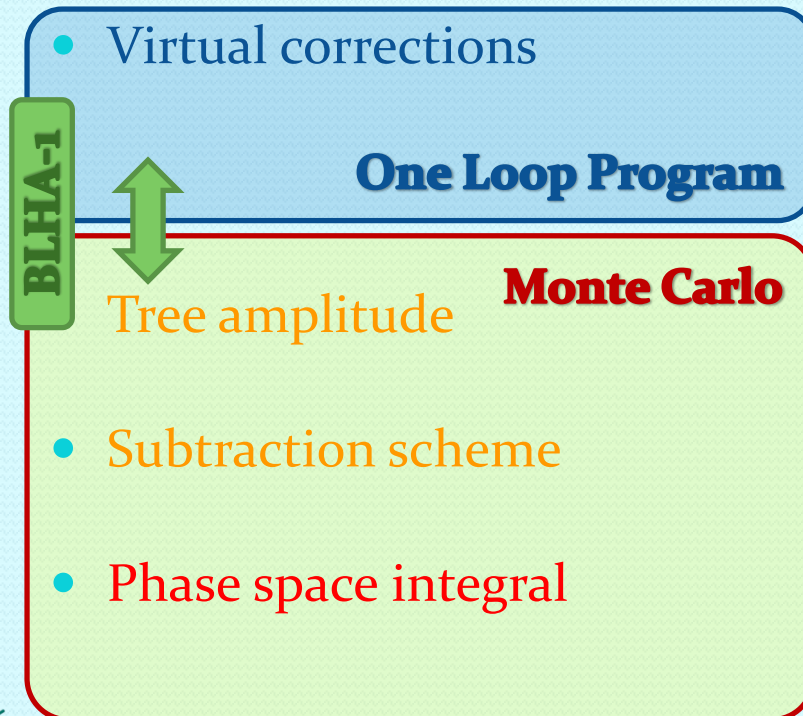
Cullen, Greiner, GH, Luisoni, Mastrolia, Mirabella,
Ossola, Peraro, Schlenk, van Deurzen,
von Soden-Fraunhofen, Tramontano



NLO Results

- Ingredients for a full NLO calculation:

$$\sigma_{\text{NLO}} = \int d\Phi_m d\sigma_{\text{Born}} + \int d\Phi_{m+1} (d\sigma_{\text{NLO}}^{\text{R}} - d\sigma_{\text{NLO}}^{\text{S}}) + \int d\Phi_m \left[\int d\Phi_1 d\sigma_{\text{NLO}}^{\text{S}} + d\sigma_{\text{NLO}}^{\text{V}} \right]$$



Recent NLO results using GoSam

- **GoSam + MadGraph/MadDipole/MadEvent**

- $pp \rightarrow b\bar{b}b\bar{b}$ [Greiner, Guffanti, Reiter, Reuter]
- $pp \rightarrow W^+ W^- jj$ [Greiner, Heinrich, Mastrolia, Ossola, Reiter, Tramontano]
- $pp \rightarrow \tilde{\chi}^0 \tilde{\chi}^0 j$ [Cullen, Greiner, Heinrich]
- $pp \rightarrow \gamma\gamma j / \gamma\gamma jj$ [Gehrmann, Greiner, Heinrich] <http://gosam.hepforge.org/diphoton>
- $pp \rightarrow G (-> \gamma\gamma) j$ [Greiner, Heinrich, Reichel, v. Soden-Fraunhofer]

- **GoSam + Powheg**

- $pp \rightarrow HW j / HZ j$ [G.L., Nason, Oleari, Tramontano]

- **GoSam + Sherpa**

- $pp \rightarrow H jj$ [in ggf] [v. Deurzen, Greiner, G.L., Mastrolia, Mirabella, Ossola, Peraro, v. Soden-Fraunhofer, Tramontano]
- $pp \rightarrow t\bar{t} (j)$ [Höche, Huang, G.L., Schönherr, Winter]
- $pp \rightarrow H t\bar{t} (j)$ [v. Deurzen, G.L., Mastrolia, Mirabella, Ossola, Peraro]
- $pp \rightarrow W^+ W^- b\bar{b}$ [Heinrich, Schlenk, Winter]

Other process packages available at:
<http://gosam.hepforge.org/proc/>

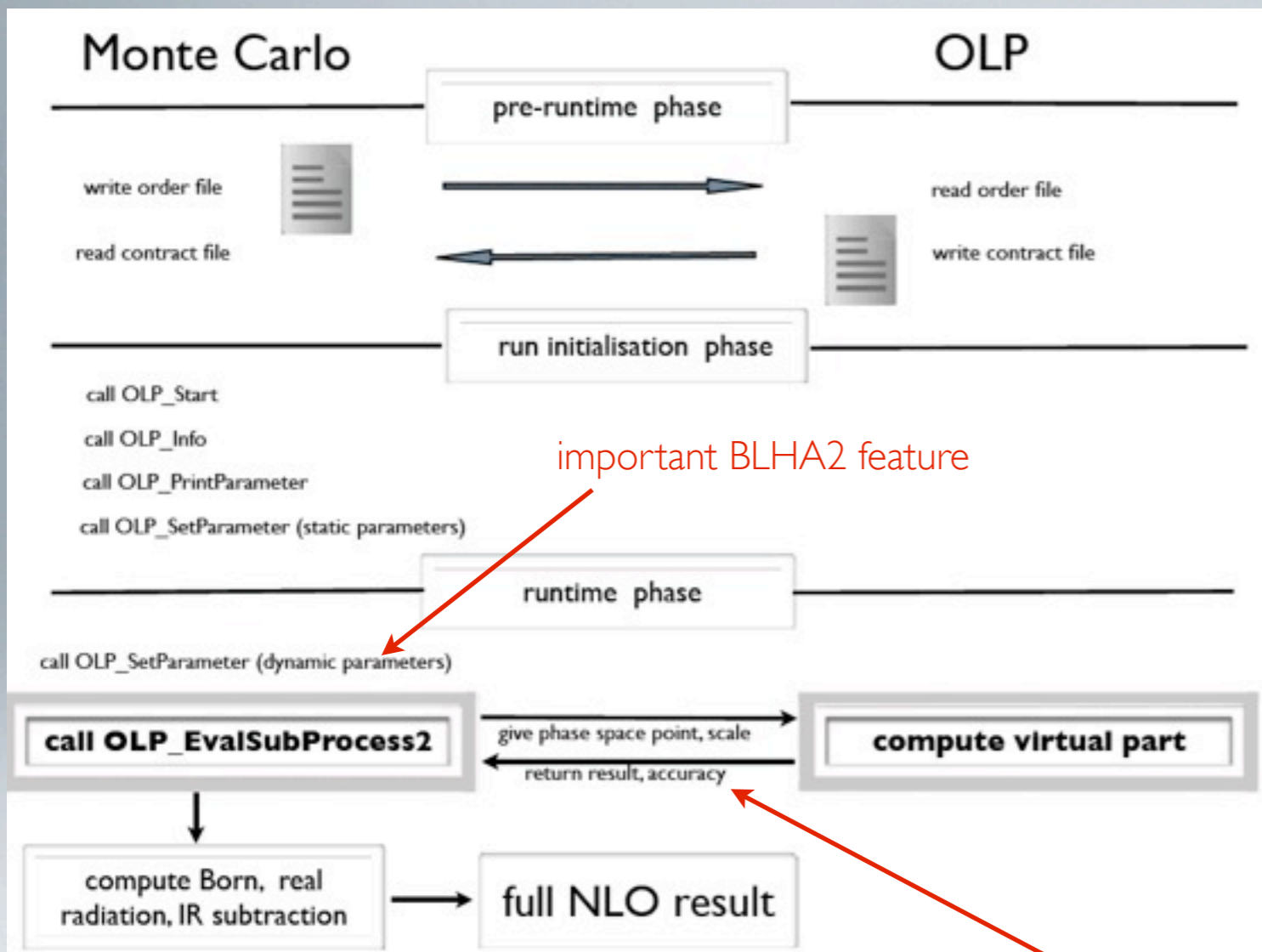
- **GoSam + MadGraph/MadDipole/MadEvent + Sherpa**

- $pp \rightarrow H jjj$ [in ggf] [Cullen, v. Deurzen, Greiner, G.L., Mastrolia, Mirabella, Ossola, Peraro, Tramontano]



Interface to Monte Carlo programs

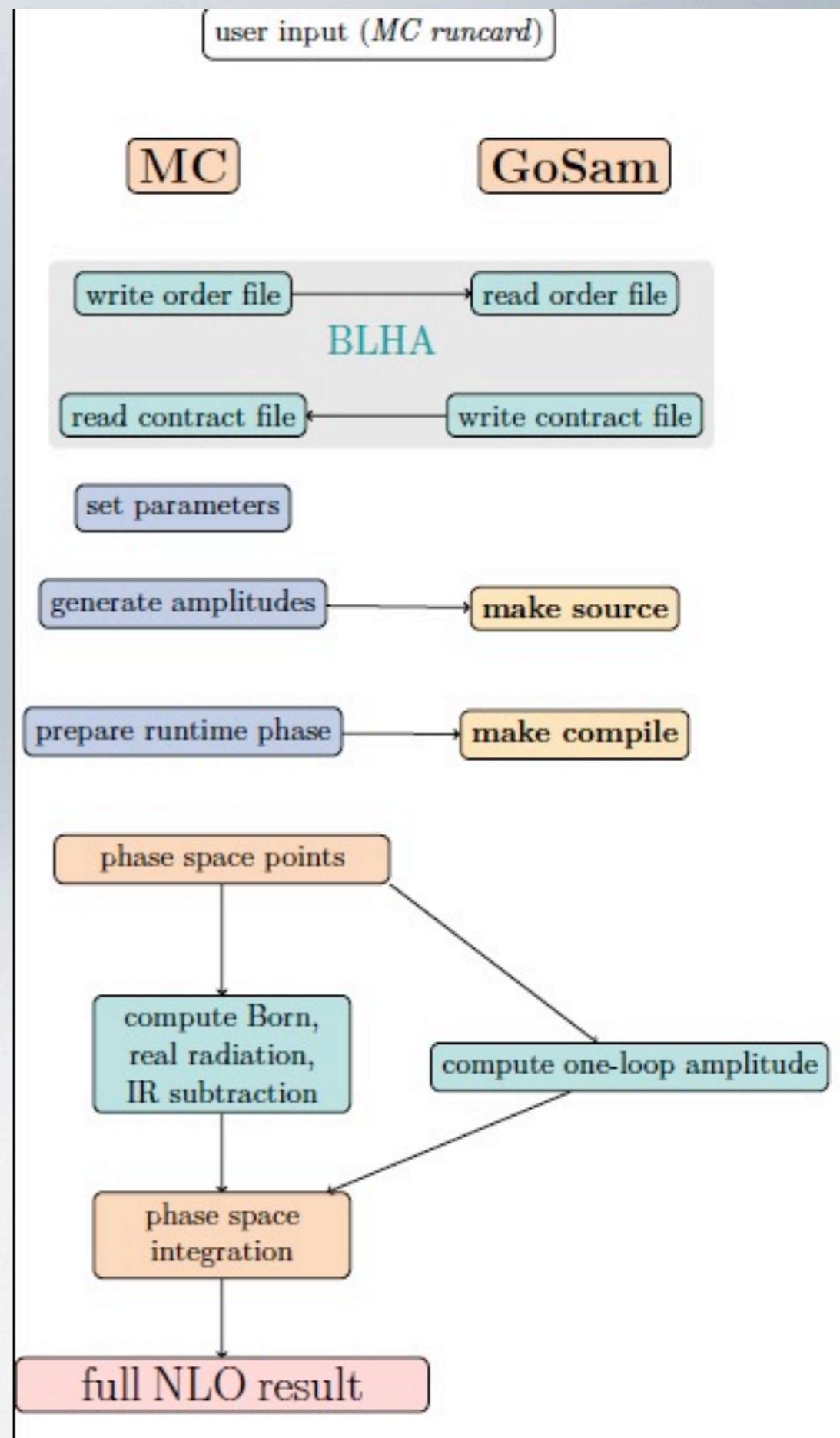
both original Binoth-Les-Houches-Accord and extended standards [CPC 185 (2014)] are supported



important BLHA2 feature

important BLHA2 feature

allows combination with different MC programs



Interface with Monte Carlo Programs

Universal interface that enables us to combine one-loop programs (OLP) with Monte Carlo programs (MC).

→ **Full NLO automation.**

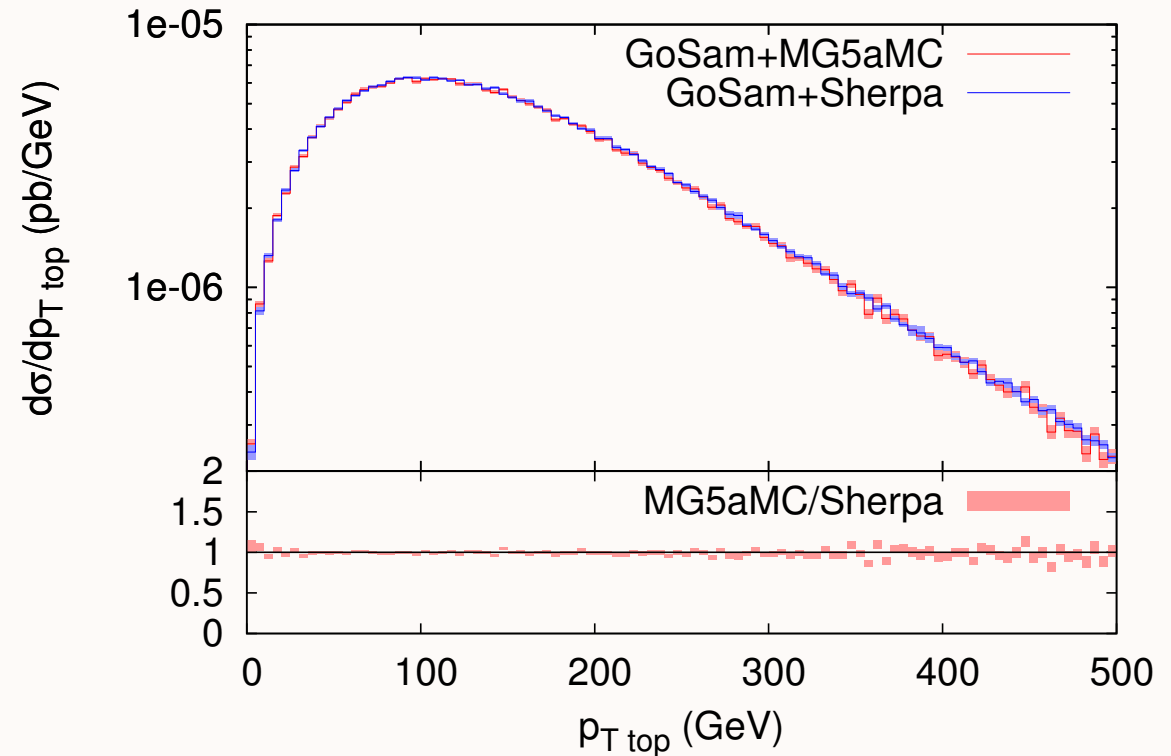
→ **GoSam interfaces exist with:**

- MadGraph/MadDipole/Madevent.
- Sherpa.
- Powheg.

Successfully tested for several benchmark processes also for:

- Herwig.
- aMC@NLO.

e.g. in top quark pair production associated with two jets,
[V.DEURZEN, FREDERIX, HIRSCHI, LUISONI, MASTROLIA, OSSOLA]



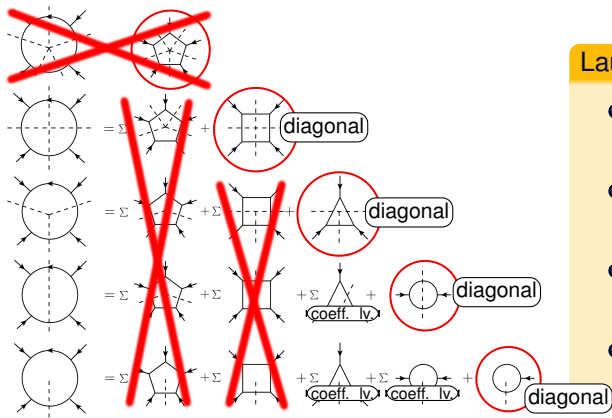
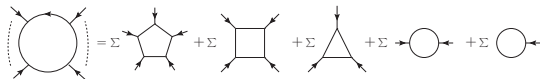
Ninja

Slide from T. Peraro.



Integrand reduction via Laurent expansion (NINJA)

Integrand decomposition:



Laurent-expansion method

- pentagons not needed
- boxes never subtracted
- diagonal systems of equations
- subtractions at coefficient level

Higgs boson plus jets in gluon-gluon fusion

Slides from G. Luisoni.

(in prep; Greiner, Höche, Luisoni, Schönherr, Winter, Yundin)



H+jets in gluon-gluon fusion

H+3 jets

- Calculation setup so far:

- B amplitudes: **Sherpa (Amebic)**
- V amplitudes: **GoSam**

- IRS amplitudes: **MG4/MadDipole**

} PS integration: **Sherpa** (BLHA)

} PS integration: **MadEvent**

↳ **Full NLO**

- New ongoing calculation:

- B amplitudes: **Sherpa (Comix)**
- V amplitudes: **GoSam**

- IRS amplitudes: **Sherpa (Comix)**

} PS integration: **Sherpa** (BLHA)

↳ **Full NLO + merging + shower**

↳ **NLO Events as NTuples**



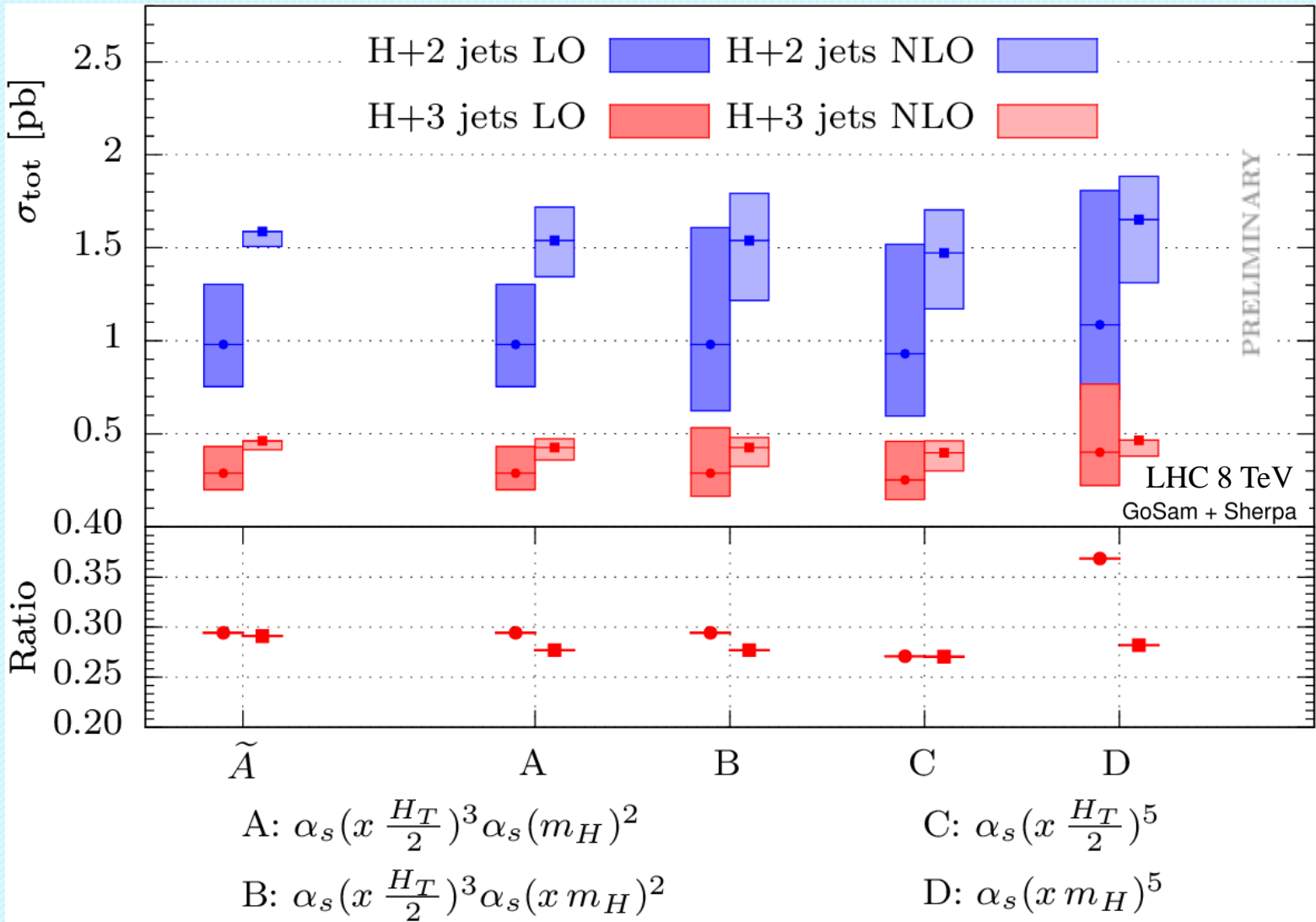
H+jets: virtual corrections

	Processes	# Diagrams	# Helicities	# Groups	Timing (col.+hel. summed)
H+0 jets	$g + g \longrightarrow H$	1	1	1	< 1 ms
H+1 jets	$q + \bar{q} \longrightarrow H + g$	14	4	3	~ 3 ms
	$g + g \longrightarrow H + g$	48	8	3	~ 7 ms
		62			
H+2 jets	$q + \bar{q} \longrightarrow H + q' + \bar{q}'$	32	4	6	~ 9 ms
	$q + \bar{q} \longrightarrow H + q + \bar{q}$	64	6	8	~ 15 ms
	$q + \bar{q} \longrightarrow H + g + g$	179	8	12	~ 56 ms
	$g + g \longrightarrow H + g + g$	651	16	12	~ 309 ms
		926			
H+3 jets	$q + \bar{q} \longrightarrow H + q' + \bar{q}' + g$	467	8	32	~ 68 ms
	$q + \bar{q} \longrightarrow H + q + \bar{q} + g$	868	12	44	~ 157 ms
	$q + \bar{q} \longrightarrow H + g + g + g$	2519	16	60	~ 999 ms
	$g + g \longrightarrow H + g + g + g$	9325	32	60	~ 8'960 ms
		13179			



Dependence on scale choice

PRELIMINARY

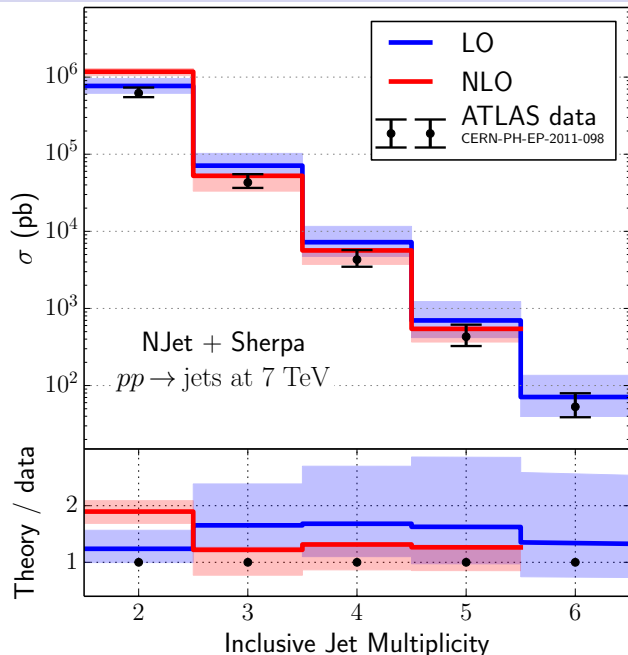


NJet

Slide from V. Yundin.



NJet+Sherpa: total XS for 2, 3, 4, 5 jets at 7 TeV vs ATLAS measurements



Cuts

anti-kt $R = 0.4$

$p_T^{\text{1st}} > 80$ GeV

$p_T^{\text{other}} > 60$ GeV

$|\eta| < 2.8$

NLO

$\mu_R = \mu_F = \hat{H}_T/2$

vars. $\hat{H}_T/4$ and \hat{H}_T

$\alpha_s(M_Z) = 0.118$

NNPDF23 PDF set

FormCalc & Friends / SecDec / BSM applications

Slides from T. Hahn, G. Heinrich and W. Hollik.



Summary

FeynArts 3.9:

- MSSMCT model file incl. complete 1-loop renormalization.

FormCalc 8.4:

- Support for run-time renormalization-scheme selection.
- Automated vectorization of helicity loop.
- Suppression of negligible helicity combinations.
- Ninja interface.

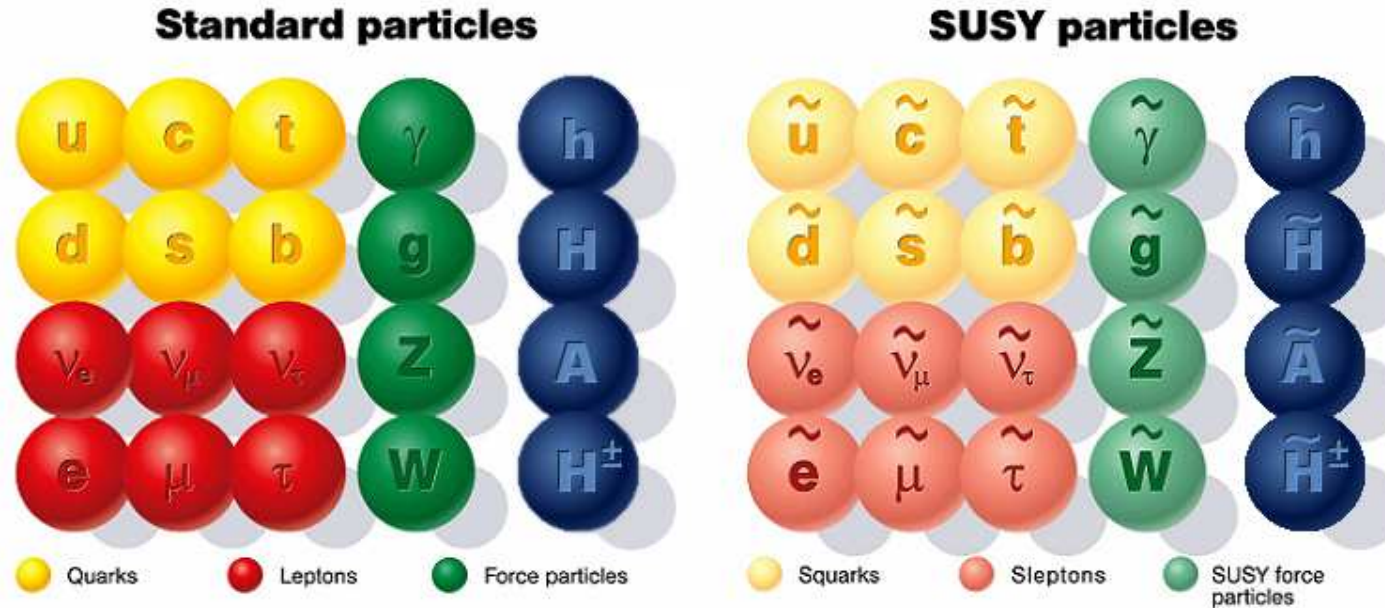
LoopTools 2.10:

- Cache concurrency issues solved.

Cuba 3.3:

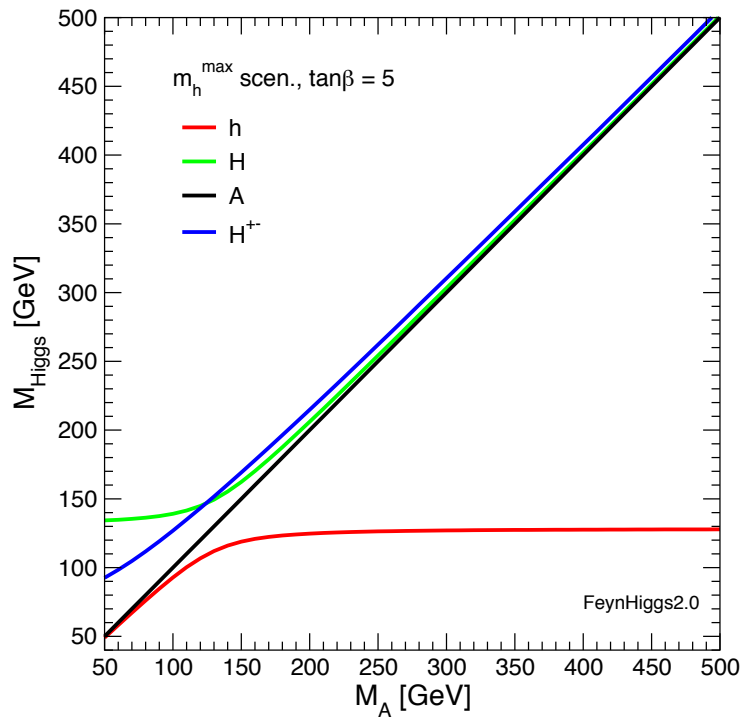
- Ways to fine-tune distribution of points to cores suited for vectorization/GPU 'on top' of Cuba's fork/wait.

Supersymmetry



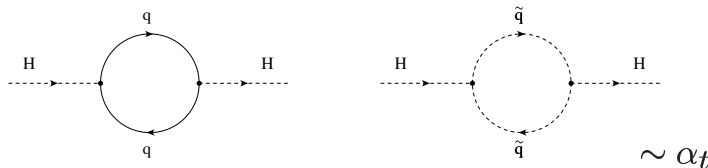
- gauge coupling unification
- dark matter candidate (lightest SUSY particle, LSP)
- physical Higgs bosons: h^0, H^0, A^0, H^\pm
lightest Higgs boson $h^0 < 130 \text{ GeV}$

Higgs bosons in the MSSM: h^0, H^0, A^0, H^\pm



- *light Higgs boson h^0 predicted*
- *for heavy A^0, H^0, H^\pm :
 h^0 like Standard Model Higgs boson*

m_h^0 strongly influenced by quantum effects, e.g. t, \tilde{t}



public tool:

FEYNHIGGS

Hahn, Hollik [MPP] + Heinemeyer, Rzehak, Weiglein

recent progress in precision calculations

- full momentum dependence of self-energies at 2-loop order $O(\alpha_t\alpha_s)$

$$\Sigma(0) \text{ [effective potential]} \rightarrow \Sigma(q^2)$$

Borowka, Hahn, Heinrich, Heinemeyer, Hollik EPJC 74(2014)2994

- MSSM with complex parameters ($\rightarrow CP$ -violation): contributions of 2-loop order $O(\alpha_t^2)$ with complex phases

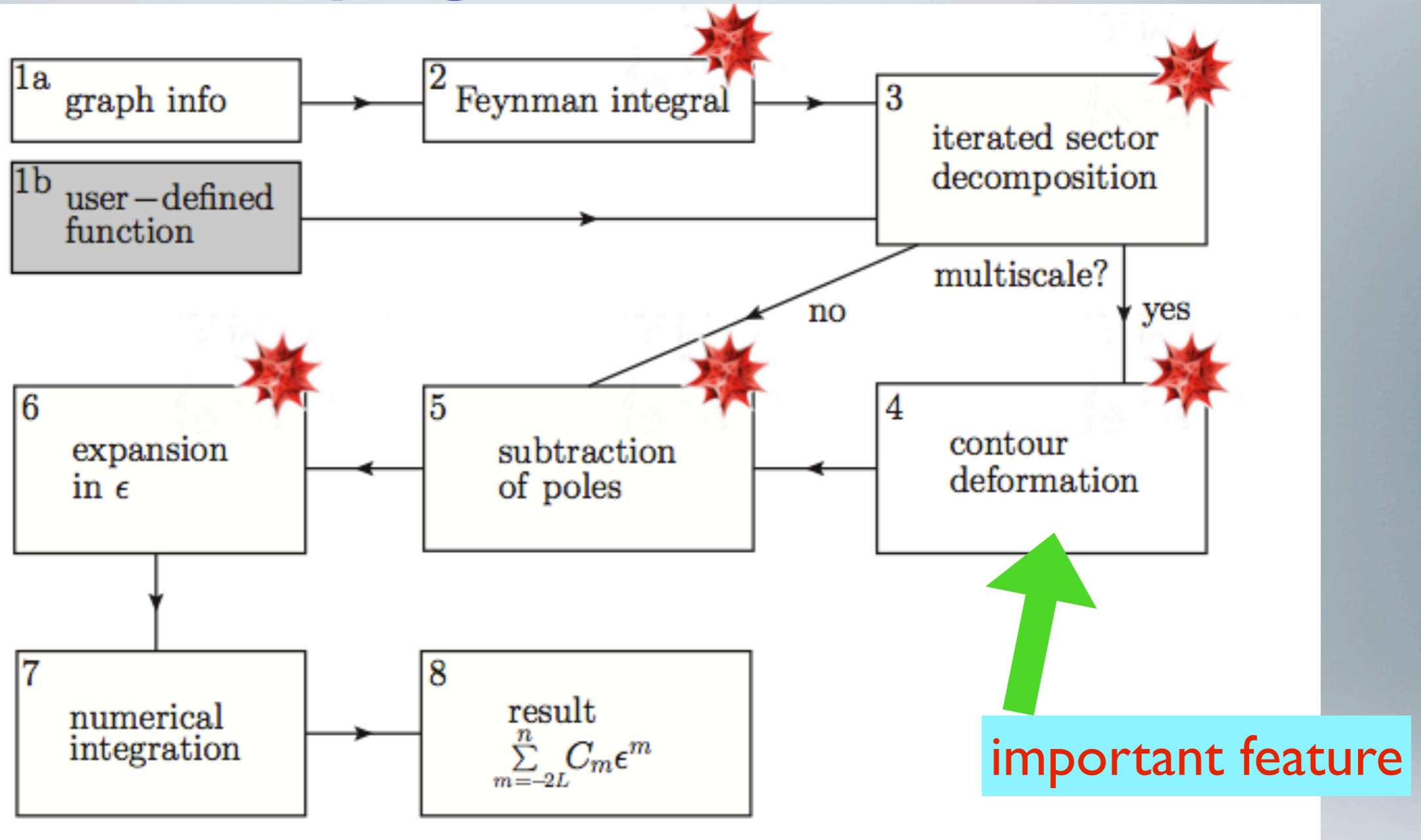
Hollik, Paßehr PLB 733(2014)144; JHEP 10(2014)171

- resummation of large logarithms $\log(M_S/m_t)$ at LL and NLL for large SUSY mass scales M_S

\Rightarrow allows predictions also for high $M_S \rightarrow$ TeV scale

Hahn, Heinemeyer, Hollik, Rzehak, Weiglein PRL 112(2014)14180

The program SecDec



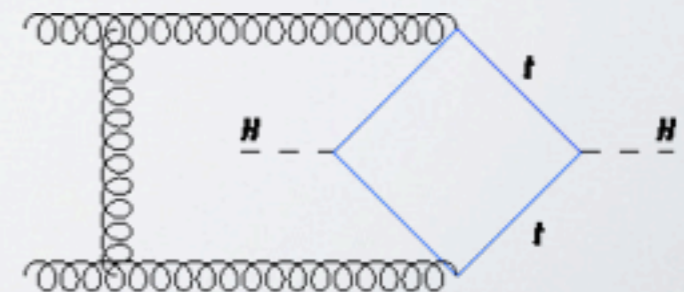
[S. Borowka, J. Carter, G.Heinrich '12,'13]

- allows calculation of multi-loop integrals with several mass scales

- **current project:**

two-loop integrals entering **double Higgs production** with full top mass dependence

(analytical calculation not in reach)



Recent application of SecDec

momentum dependent $\mathcal{O}(\alpha_s \alpha_t)$

corrections to neutral Higgs boson masses in the MSSM

[S. Borowka, T. Hahn, S. Heinemeyer, GH, W. Hollik 2014]

- requires calculation of Higgs boson self-energies up to two-loop level

$$\Gamma \equiv \Delta_{\text{Higgs}}^{-1} = \begin{pmatrix} p^2 - m_{H,\text{tree}}^2 + \hat{\Sigma}_H(p^2) & \hat{\Sigma}_{hH}(p^2) \\ \hat{\Sigma}_{hH}(p^2) & p^2 - m_{h,\text{tree}}^2 + \hat{\Sigma}_h(p^2) \end{pmatrix}$$

- find the complex solutions of $\det \Gamma = 0$
- the masses are identified with the real parts of the solutions

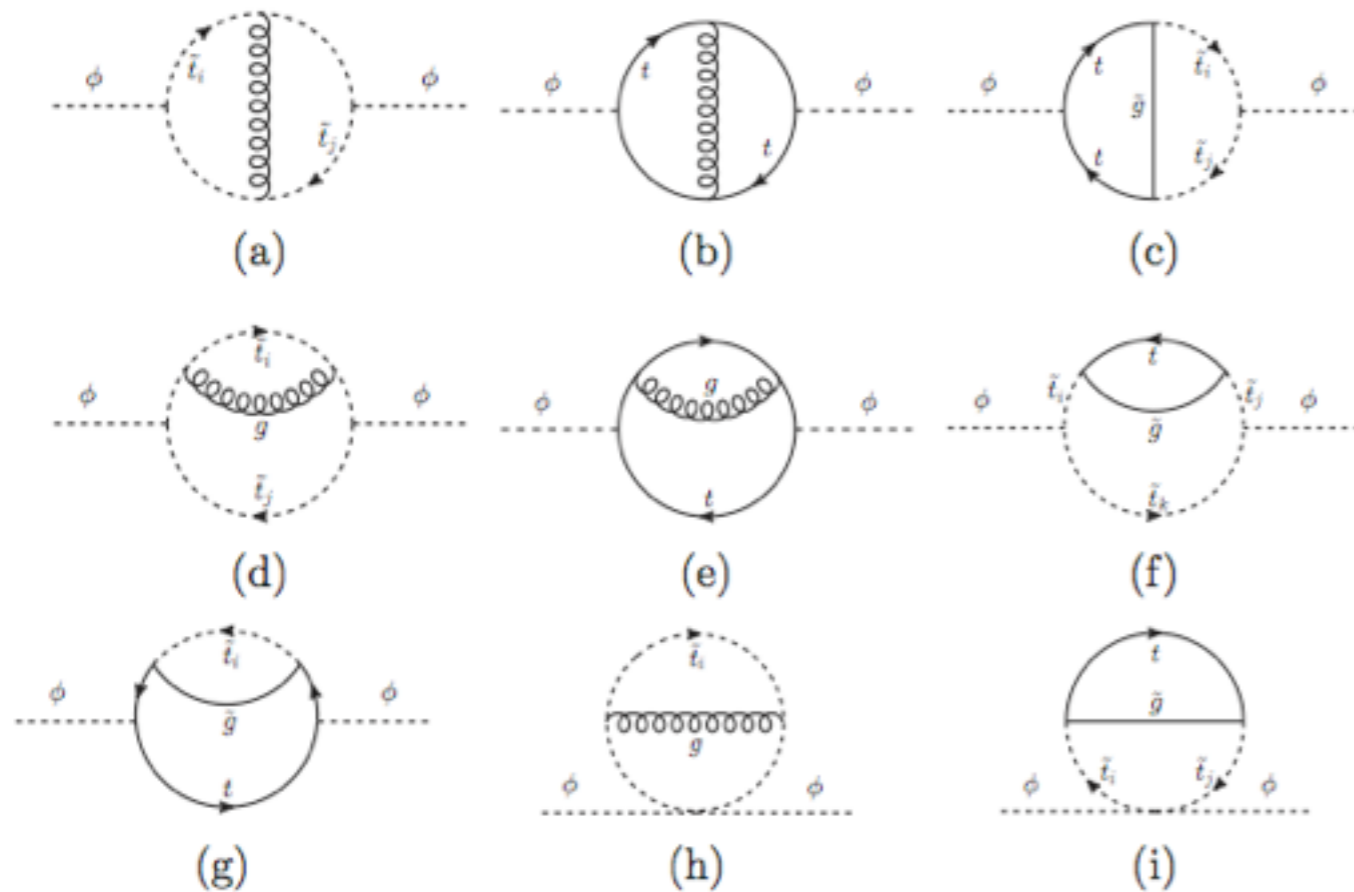
note: experimental precision on Higgs boson mass:

ATLAS: $M_H^{\text{exp}} = 125.5 \pm 0.4 \pm 0.2 \text{ GeV}$

CMS: $M_H^{\text{exp}} = 125.7 \pm 0.3 \pm 0.3 \text{ GeV}$

theoretical precision (MSSM): $\Delta M_h \simeq 3 \text{ GeV}$

two-loop Higgs boson selfenergy diagrams

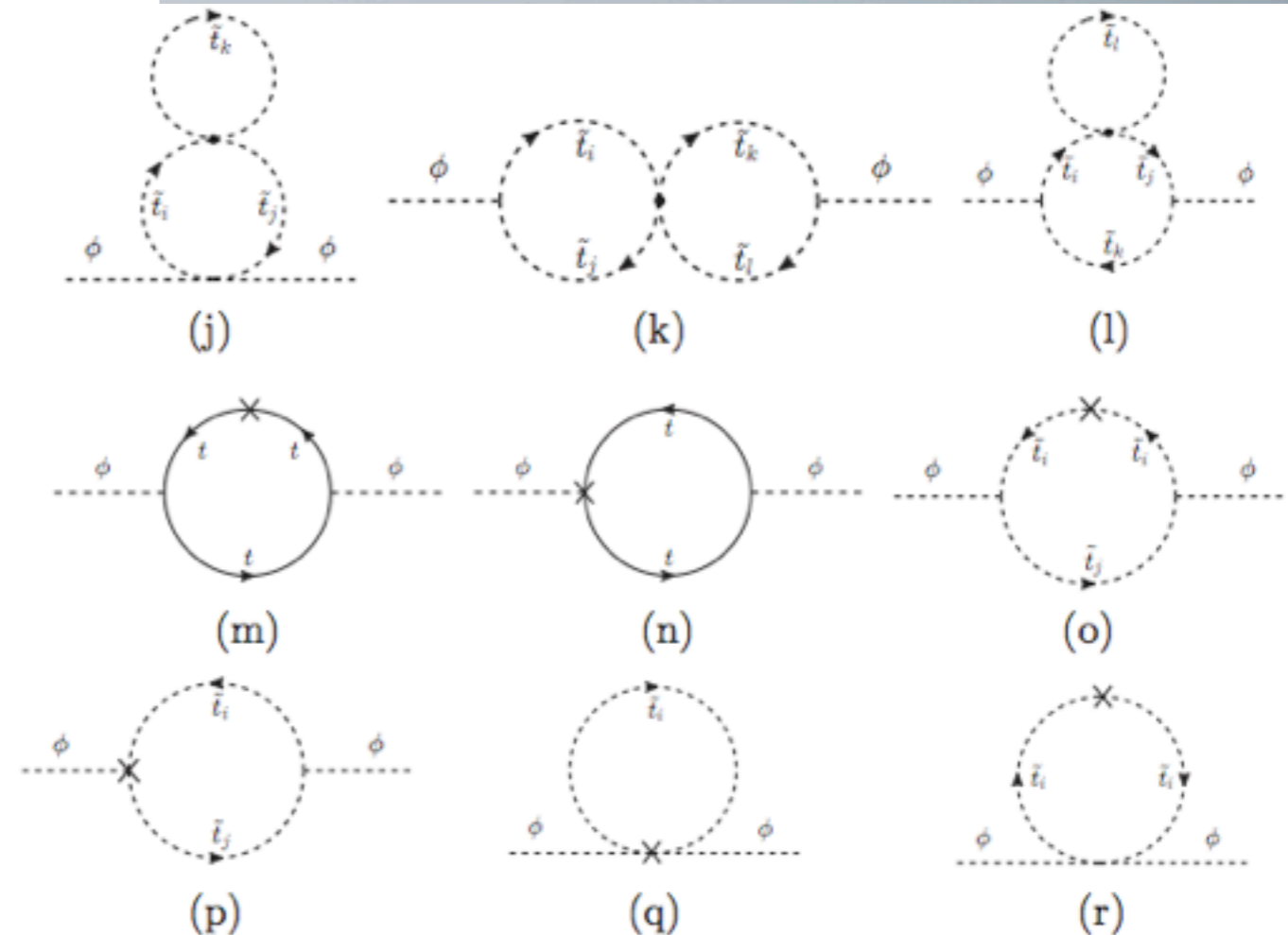


- only few of these integrals are known analytically
- up to 4 different masses
- was technically not feasible so far:
momentum dependence

numerical calculation possible with

SecDec-2.1

[S. Borowka, GH]



Status of corrections in real MSSM

public programs:

FeynHiggs Frank, Hahn, Heinemeyer, Hollik, Rzehak, Weiglein '00 '03 '07 '14

SoftSusy Allanach '02 **SPheno** Porod '03 Staub '11

CPsuperH Carena, Choi, Drees, Ellis, Lee, Pilaftsis, Wagner '04 '09 '12

Suspect Djouadi, Kneur, Moutaka '07

H3m Kant, Harlander, Mihaila, Steinhauser '10

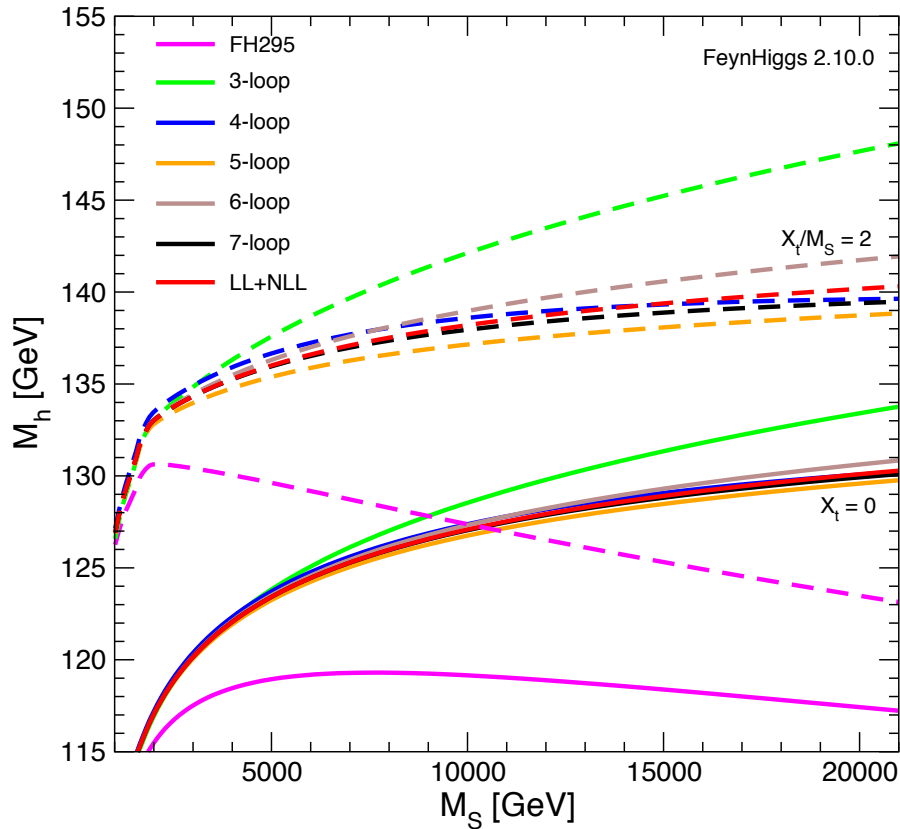
implemented corrections:

- **1-loop** : complete
- **2-loop** : $\mathcal{O}(\alpha_s \alpha_t)$, $\mathcal{O}(\alpha_t^2)$, $\mathcal{O}(\alpha_s \alpha_b)$, $\mathcal{O}(\alpha_b^2)$, $\mathcal{O}(\alpha_t \alpha_b)$ at $p^2 = 0$
- **3-loop** : $\mathcal{O}(\alpha_s^2 \alpha_t)$ at $p^2 = 0$

now: **2-loop** $\mathcal{O}(\alpha_s \alpha_t)$ at $p^2 \neq 0$

[Borowka, Hahn, Heinemeyer, GH, Hollik 04/2014]

[Degrassi, Di Vita, Slavich 10/2014]



— fixed order

— full resummation of LL and NLL

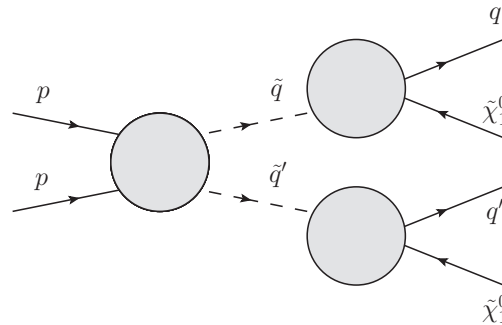
\tilde{t} -squark mass matrix $\mathcal{M}_{\tilde{t}}^2 =$

$$\begin{pmatrix} M_S^2 + m_t^2 & m_t X_t \\ m_t X_t & M_S^2 + m_t^2 \end{pmatrix}$$

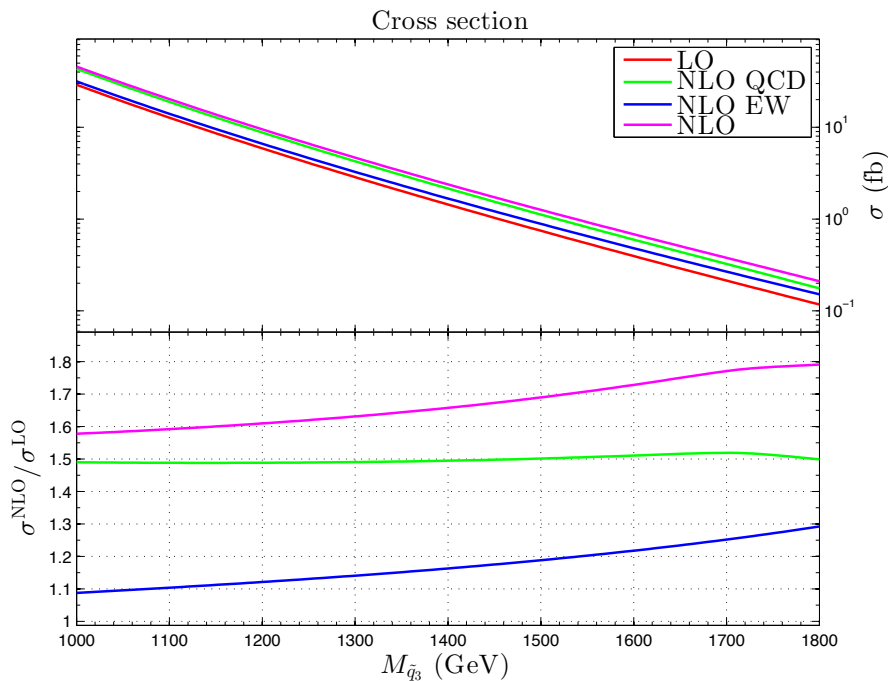
- sizeable upward shift for heavy top-squarks
- large impact for confronting constrained models with the Higgs signal at 125 GeV

SUSY searches at the LHC

- searches need predictions for production and decays of SUSY particles



- NLO calculations required for stable predictions



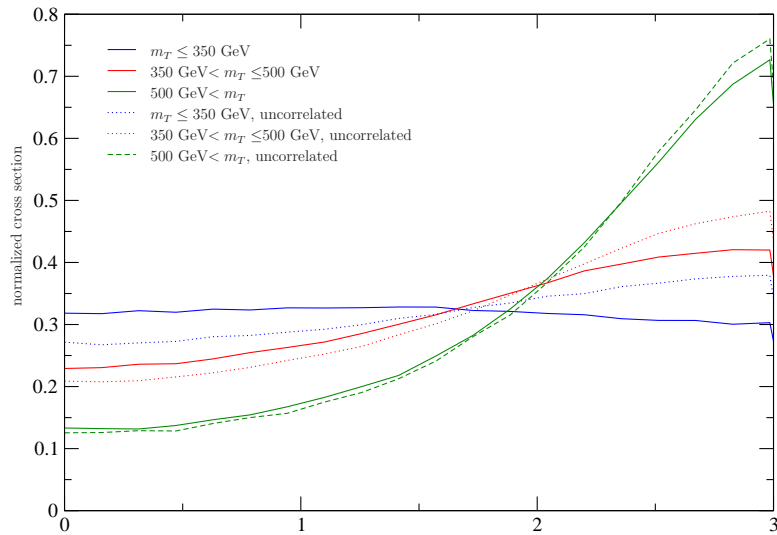
top-squark pair production at the LHC
 $pp \rightarrow \tilde{t}_1 \tilde{t}_1^*$

QCD and electroweak NLO

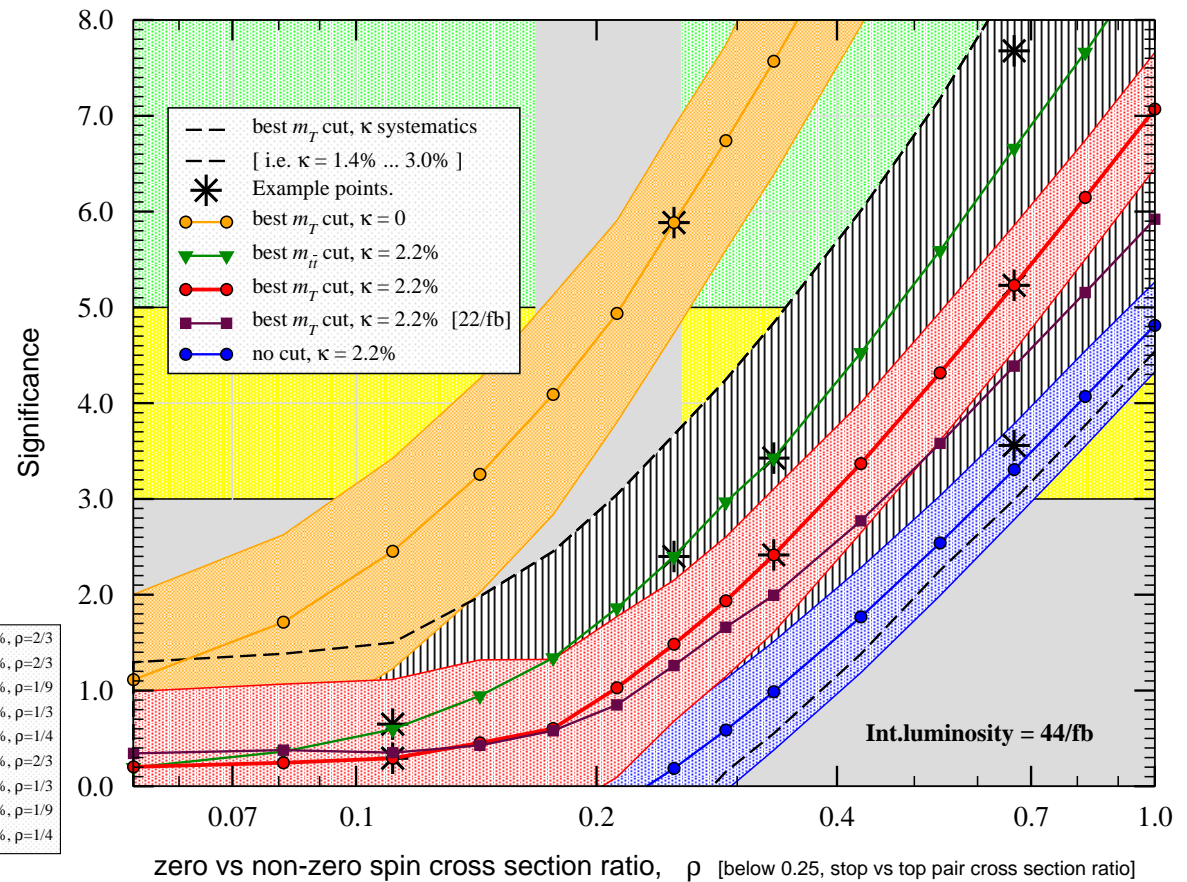
*[Hollik, Lindert, Mirabella
 JHEP 09(2014)022]*

Spin correlations in top quark pair production

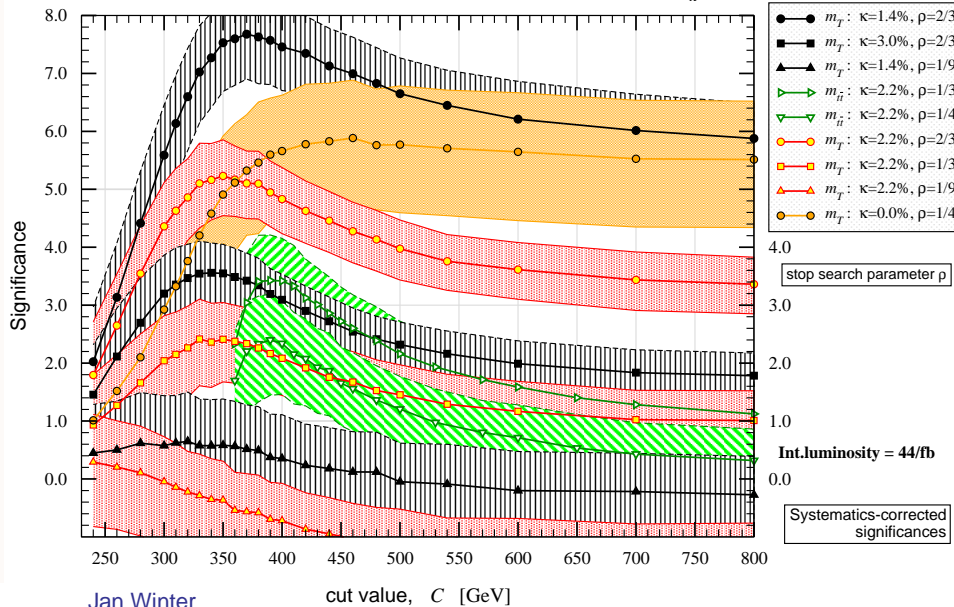
Use the shape of the lepton azimuthal angle separation in the dilepton channel to search for contamination from light stop quarks.



Shape analysis reach regarding stop scenario exclusion @ LHC 8 TeV, $m_{T,\bar{t}} < C$



Optimization of transverse and invariant mass constraints @ LHC 8 TeV, $m_x < C$



[IN PREP; MAHBUBANI, PARKE, RZEHAK, WINTER]



Summary.

The activities of the group cover a large spectrum of LHC phenomenology in and beyond the Standard Model.

Precision calculations: automated NLO corrections, also for multi-scale and multi-leg processes, Higgs masses and cross sections in the MSSM, top-squark pair production.

Development of computational tools that are useful for and being used by the entire community.

Development of new analytical/numerical methods at the calculational frontier (NNLO/multi-loops).

Great collaboration with Monte Carlo generator experts.

There is nice cross-talk between MPI experimentalists and pheno-group members already. An increase in this exchange between EXP & TH would be something we are definitely looking forward to.



The end. (:o) Thank You.



Magnus/Dyson series for Master Integrals

[Argeri, Di Vita, Mastrolia, Mirabella, Schlenk, Schubert, Tancredi JHEP 1403 (2014) 082]

A modern multi-loop and -leg calculation is done in three steps

- 1 Find integral basis for the process **IBP, Multiloop integrand reduction (MIR)**
- 2 Determine “coordinates” in such basis **Feynman diags., Generalized unitarity, OPP, MIR**
- 3 Calculate basis elements **Feynman/Schwinger, Mellin-Barnes, Differential Equations**

Differential Equations for d -dimensional Feynman Integrals of a given topology

IBP relations \Rightarrow **most integrals linearly dep.** Choose indep “Master Integrals”, functions of:

- 1 the kinematic invariants x_i built with the external momenta
- 2 the internal masses
- 3 the number of spacetime dimensions d

IBP relations \Rightarrow MIs obey **system of first order linear differential equations** in ∂_{x_i}

Conjecture: basis choice can simplify the solution of such systems! [Henn 13]

bad choice

$$\partial_{x_i} \vec{g}(\vec{x}, \epsilon) = A_{x_i}(\vec{x}, \epsilon) \vec{g}(\vec{x}, \epsilon)$$

good choice (if \exists)

$$\partial_{x_i} \vec{g}(\vec{x}, \epsilon) = \epsilon A_{x_i}(\vec{x}) \vec{g}(\vec{x}, \epsilon)$$

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with a good (**canonical**) choice:

- ▶ integration (order by order in ϵ) is just algebra, iterated \int over kernels
- ▶ system dictates form of the integration kernels and analytic properties
- ▶ most of the boundary conditions can be obtained *algebraically* by exploiting analytical properties of the solution (regularity at pseudothresholds etc)

Magnus/Dyson series for Master Integrals

[Argeri, Di Vita, Mastrolia, Mirabella, Schlenk, Schubert, Tancredi JHEP 1403 (2014) 082]

A modern multi-loop and -leg calculation is done in three steps

- 1 Find integral basis for the process [IBP](#), [Multiloop integrand reduction \(MIR\)](#)
- 2 Determine “coordinates” in such basis [Feynman diags.](#), [Generalized unitarity](#), [OPP](#), [MIR](#)
- 3 Calculate basis elements [Feynman/Schwinger](#), [Mellin-Barnes](#), [Differential Equations](#)

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How to find good basis? Is it possible to proceed in *algorithmic* way? Still open question.

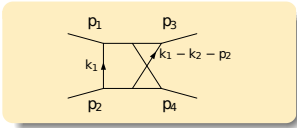
A first proposal: use **Magnus Exponential** to turn ϵ -linear systems into canonical systems

[Argeri, Di Vita, Mastrolia, Mirabella, Schlenk, Schubert, Tancredi JHEP 1403 (2014) 082]

Two-loop non-planar box [Tausk 99; Anastasiou, Gehrmann, Oleari, Remiddi, Tausk 00]

[Argeri, Di Vita, Mastroli, Mirabella, Schlenk, Schubert, Tancredi JHEP 1403 (2014) 082]

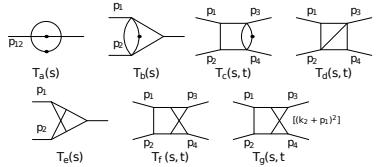
12 MI's for the crossed topology



$$x = -\frac{t}{s}, \quad s > 0, t < 0, |s| > |t|$$

This choice of the MI's f 's obeys an ϵ -linear DE

$$\begin{aligned}
 f_1 &= \epsilon^2 s T_a(s), & f_2 &= \epsilon^2 t T_a(t), & f_3 &= \epsilon^2 u T_a(u), \\
 f_4 &= \epsilon^3 s T_b(s), & f_5 &= \epsilon^3 st T_c(s, t), & f_6 &= \epsilon^3 su T_c(s, u), \\
 f_7 &= \epsilon^4 u T_d(s, t), & f_8 &= \epsilon^4 s T_d(t, u), & f_9 &= \epsilon^4 t T_d(u, s), \\
 f_{10} &= \epsilon^4 s^2 T_\theta(s), \\
 f_{11} &= \epsilon^4 st u T_f(s, t) - \frac{3}{4s(4\epsilon + 1)} \left[\epsilon^2 (s^2 T_a(s) + t^2 T_a(t) + u^2 T_a(u)) \right. \\
 &\quad \left. - 4\epsilon^4 (u^2 T_d(s, t) + s^2 T_d(t, u) + t^2 T_d(u, s)) \right], \\
 f_{12} &= \epsilon^4 st T_g(s, t) - \frac{3}{8u(4\epsilon + 1)} \left[\epsilon^2 (s^2 T_a(s) + t^2 T_a(t) + u^2 T_a(u)) \right. \\
 &\quad \left. - 4\epsilon^4 (u^2 T_d(s, t) + s^2 T_d(t, u) + t^2 T_d(u, s)) \right],
 \end{aligned}$$



$$\partial_x f(\epsilon, x) = (A_0(x) + \epsilon A_1(x)) f(\epsilon, x) \quad A_i(x) = \frac{M_{f,1}}{x} + \frac{M_{f,2}}{1-x}$$

Magnus Exponential \Rightarrow get rid of A_0 , g 's obey a canonical DE

$$\partial_x g(\epsilon, x) = \epsilon \hat{A}_1(x) g(\epsilon, x) \quad \hat{A}_1(x) = \frac{M_1}{x} + \frac{M_2}{1-x}$$

$$M_1 = \begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & -2 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & -\frac{1}{2} & 0 & 0 & -2 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -3 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -2 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 2 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -4 & -2 & -18 & -12 & -12 & 1 & 1 & -2 & -2 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 3 & 2 & -3 & 12 & -6 & -18 & 0 & 0 & -2 \end{pmatrix}$$

$$M_2 = \begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 2 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & -\frac{1}{2} & 0 & 0 & 3 & -1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 2 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -2 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 2 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 2 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -4 & -2 & -18 & -12 & -12 & 1 & 1 & -2 & -2 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -6 & 2 & -4 & 12 & -6 & -24 & 1 & -1 & 0 \end{pmatrix}$$

Three-loop Higgs + 1Jet ladder-box [Di Vita, Mastrolia, Schubert, Yundin JHEP 1409 (2014) 148]

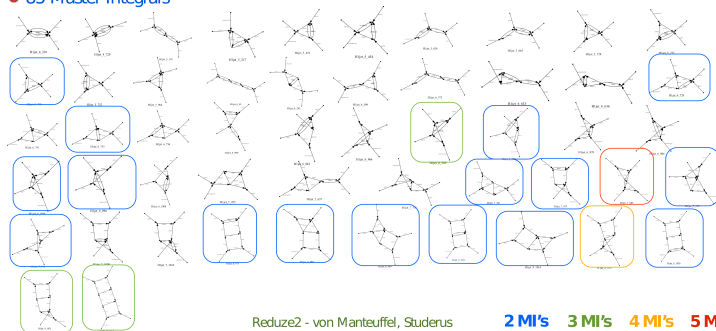


4 MI's [Kindergarten]

18 MI's [Gehrmann, Remiddi 00]

85 MI's!!

● 85 Master Integrals



Reduze2 - von Manteuffel, Studerus

2 MI's 3 MI's 4 MI's 5 MI's

now 2 integration variables, not just two "letters" $\{x, 1 - x\}$ but a bigger *alphabet* $\{x, 1 - x, y, 1 - y, 1 - x - y, x + y\}$

own GHPL code + REDUZE [Studerus, Studerus and von Manteuffel] + GiNaC [Bauer, Frink and Kreckel; Vollinga and Weinzierl]

Three-loop Higgs + 1Jet ladder-box [Di Vita, Mastrolia, Schubert, Yundin JHEP 1409 (2014) 148]

