

MIAPP Workshop 2014

Top Quark Physics Day

**Probing top quark electroweak couplings
at the LHC**

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- Our understanding of the top quark as an elementary particle is very solid
- Basic properties were explored at the Tevatron
- Electroweak couplings are relatively unexplored in collider experiments
- P.Uwer (ICHEP'14): „*The time for predicted discoveries might be over for a while. We need to turn every stone and look below.*“

Electroweak top quark couplings

We want to study $t\bar{t}$ +electroweak gauge bosons

$$\begin{array}{ll} t\bar{t} + \gamma & t\bar{t} + H \\ t\bar{t} + Z & t\bar{t} + W^\pm \end{array}$$

additional sensitivity: single top, top decay dynamics

Our framework:

NLO QCD corrections in top production and decay.

Top quarks treated in narrow width approximation.

Features:

- NLO spin correlations
- Any top quark decay channel at NLO
- Photon radiation off top decay products

Electroweak top quark couplings

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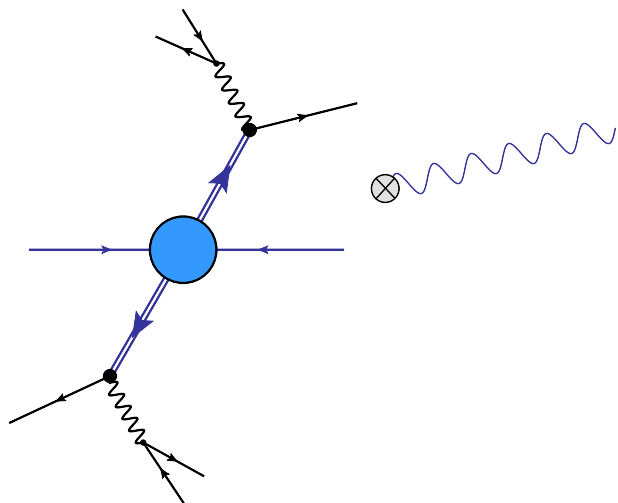
Top quarks treated in narrow width approximation.

Neglect:

- Parton Showering (taken care of by NLO decay)
- Threshold corrections (expected to be small beyond NLO; Phasespace $\sim \beta^4$)
- Top off-shell effects (parametrically small $\sim \Gamma/m$, explicitly verified in $t\bar{t}$)

$$t\bar{t} + \gamma$$

ttbar + photon



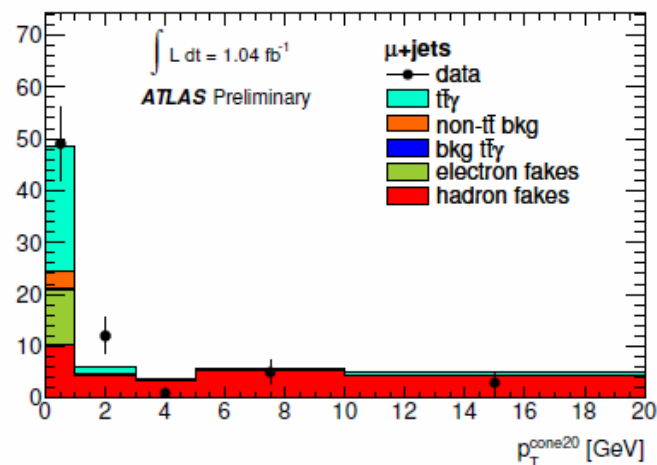
- Directly sensitive to top quark electric charge
- At LHC gg dominated (small ISR contamination)
- Has FB asymmetry already at LO
- Serves as control sample for $t\bar{t}b+H$

[Aguilar-Saavedra]

$$\mathcal{L}_{\gamma tt} = -eQ_t \bar{t} \gamma^\mu t A_\mu - e \bar{t} \frac{i\sigma^{\mu\nu} q_\nu}{m_t} (d_V^\gamma + i d_A^\gamma \gamma_5) t A_\mu.$$

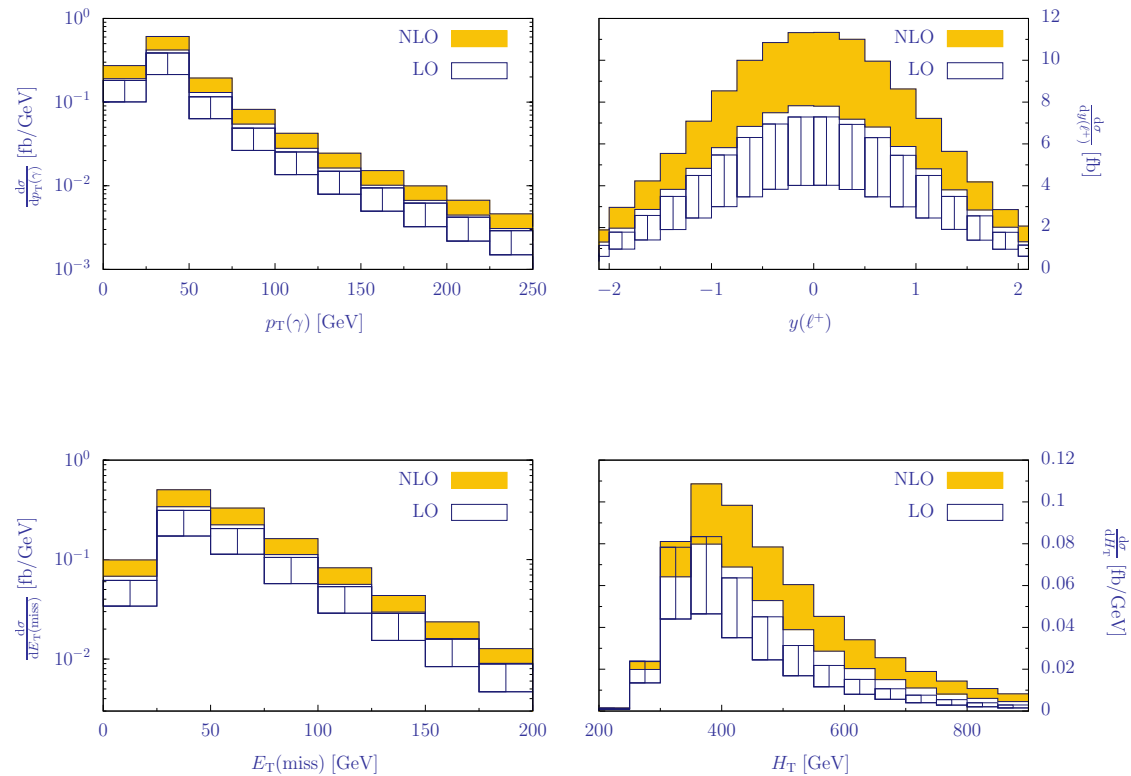
$$\delta d_V^\gamma = \frac{\sqrt{2}}{e} \text{Re} [c_W C_{uB\phi}^{33} + s_W C_{uW}^{33}] \frac{vm_t}{\Lambda^2},$$

$$\delta d_A^\gamma = \frac{\sqrt{2}}{e} \text{Im} [c_W C_{uB\phi}^{33} + s_W C_{uW}^{33}] \frac{vm_t}{\Lambda^2}.$$



ttbar + photon

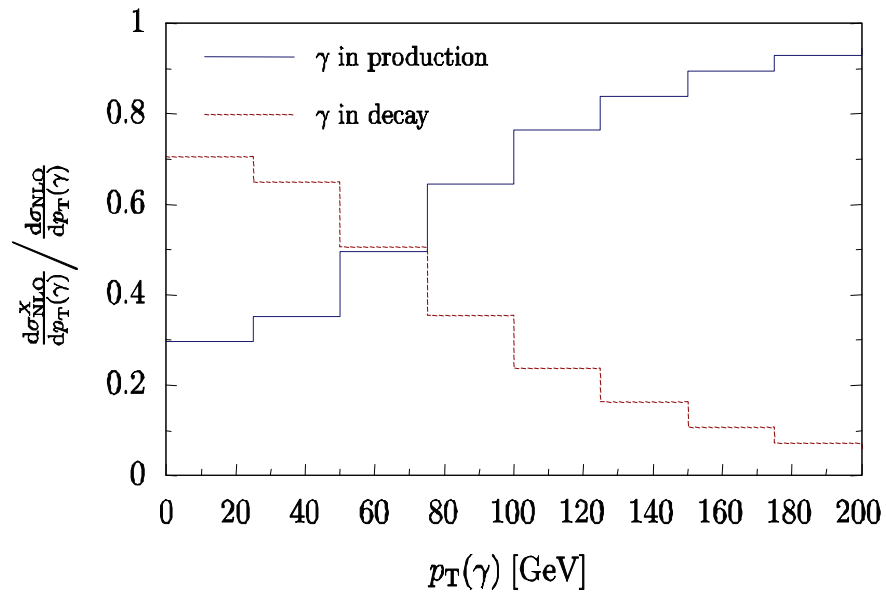
$$pp(7\text{TeV}) \rightarrow t\bar{t} + \gamma \rightarrow b\bar{b} \ell\nu jj + \gamma$$



$$\sigma_{t\bar{t}\gamma}^{\text{LO}} = 15.6 \pm 4.6 \text{ fb} \quad \sigma_{t\bar{t}\gamma}^{\text{NLO}} = 26.8 \pm 5.0 \text{ fb}$$

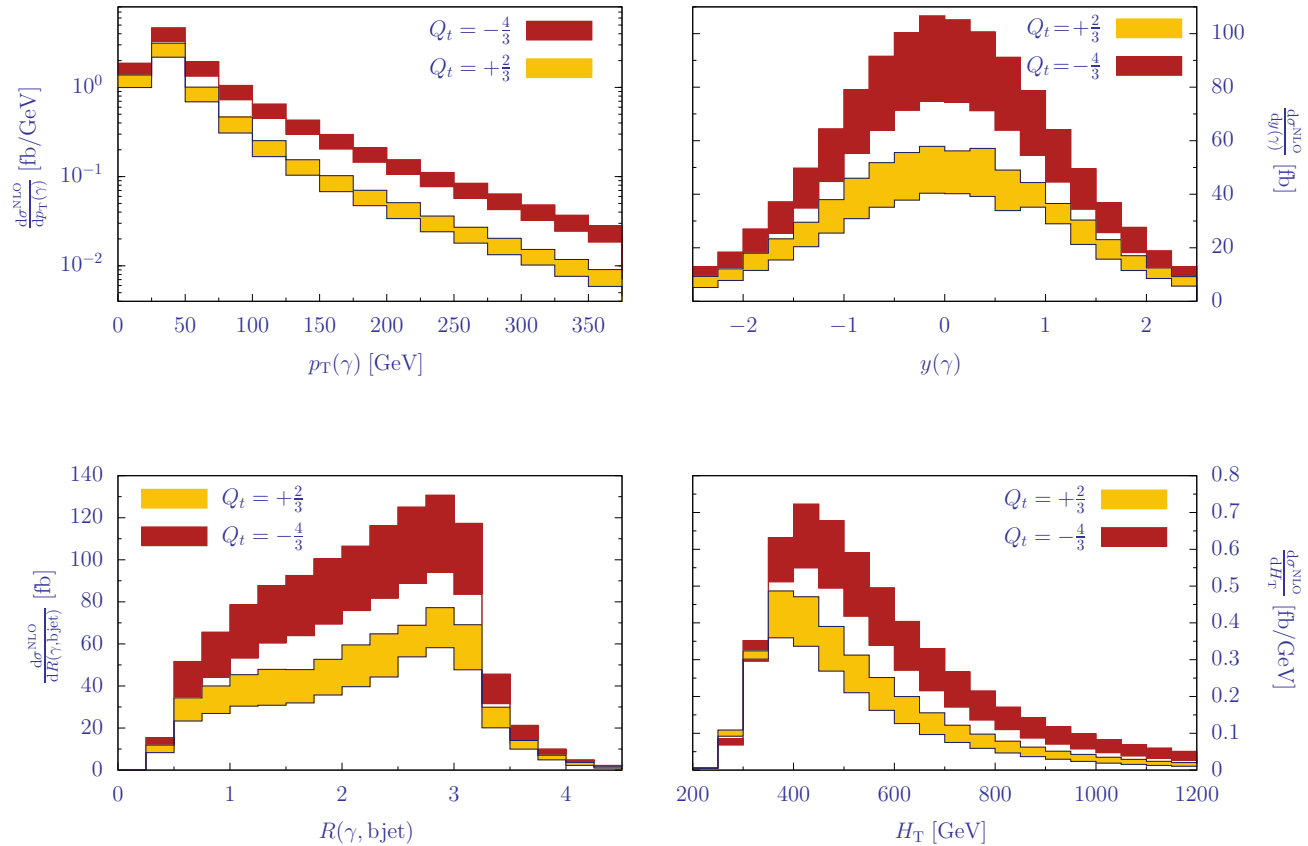
[Melnikov,Scharf,M.S]

ttbar + photon



Most of the photons with $p_T^\gamma < 50$ GeV are radiated in the top quark decay.

ttbar + photon



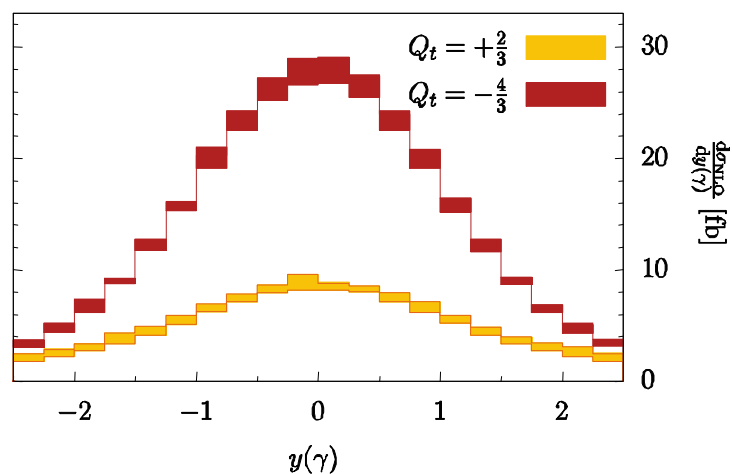
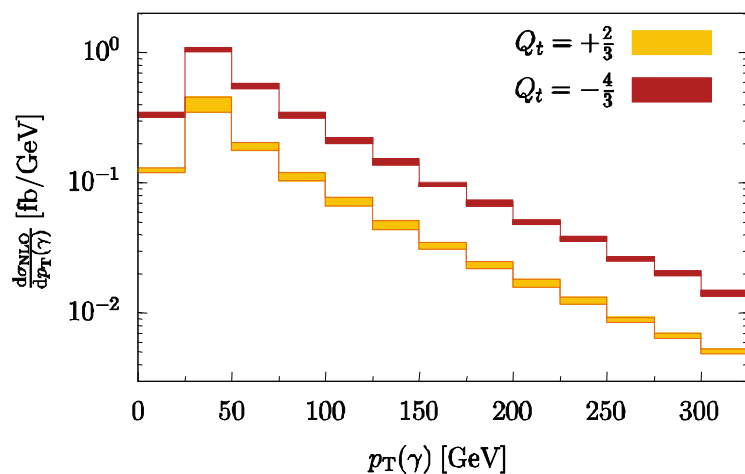
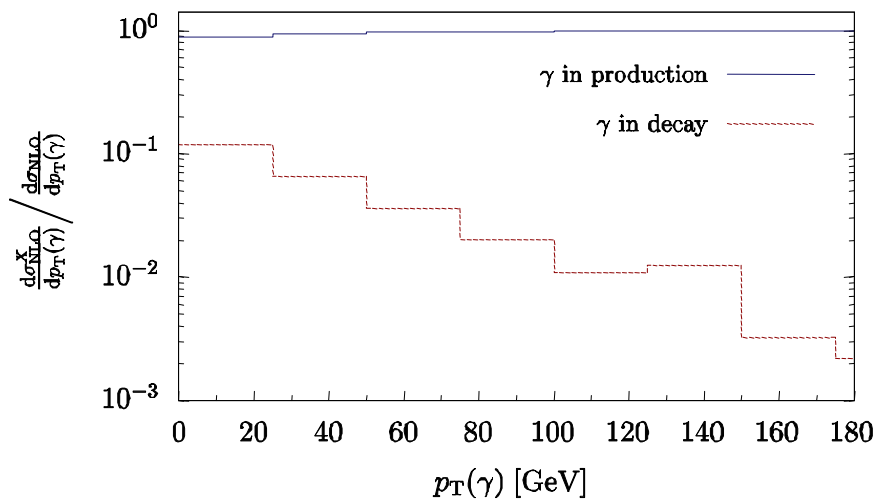
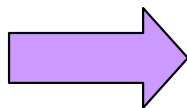
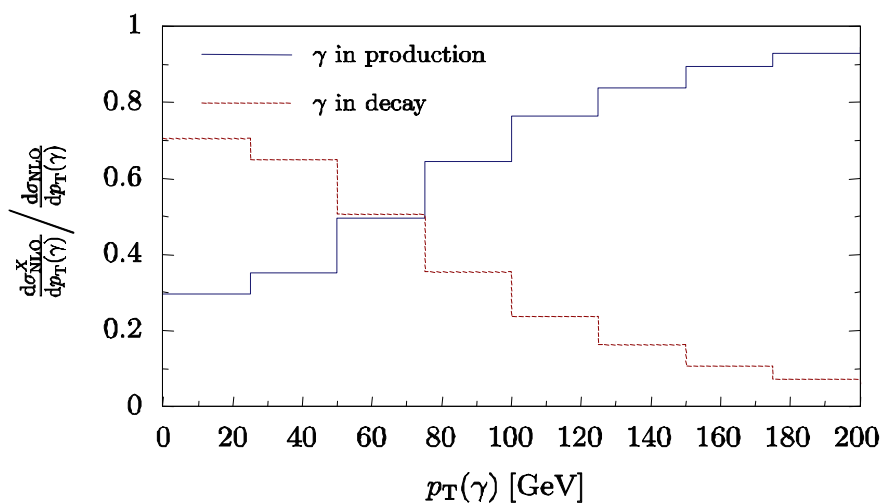
$$\sigma_{t\bar{t}\gamma}^{\text{NLO}} = 138 \text{ fb} \xrightarrow{Q_t = \frac{2}{3} \rightarrow -\frac{4}{3}} \sigma_{t\bar{t}\gamma}^{\text{NLO}} = 243 \text{ fb}$$

Naive expectation of Q_t^2 scaling fails.

ttbar + photon

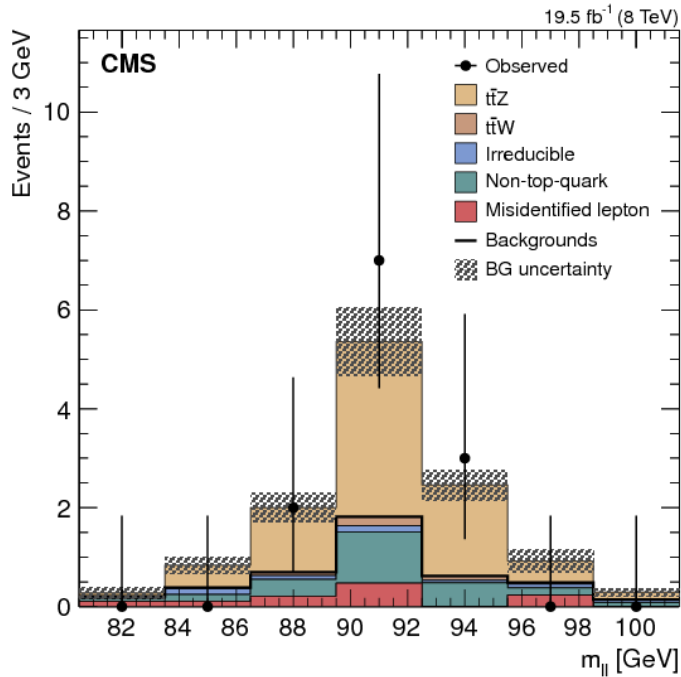
$$m_T(bl\gamma; E_T^{\text{miss}}) > 180 \text{ GeV}, \quad m_T(l\gamma; E_T^{\text{miss}}) > 90 \text{ GeV},$$

$$160 \text{ GeV} < m(bjj) < 180 \text{ GeV}, \quad 70 \text{ GeV} < m(j, j) < 90 \text{ GeV}$$



$$t\bar{t} + Z$$

$t\bar{t} + Z$



- Direct probe of $t\bar{t}b$ - Z interactions
- At LHC, gg dominated, $Z \rightarrow ll$ has small background
- Never observed at the Tevatron

$$\mathcal{L}_{t\bar{t}Z} = ie\bar{u}(p_t) \left[\gamma^\mu (C_{1,V} + \gamma_5 C_{1,A}) + \frac{i\sigma_{\mu\nu} q_\nu}{M_Z} (C_{2,V} + i\gamma_5 C_{2,A}) \right] v(p_{\bar{t}}) Z_\mu,$$

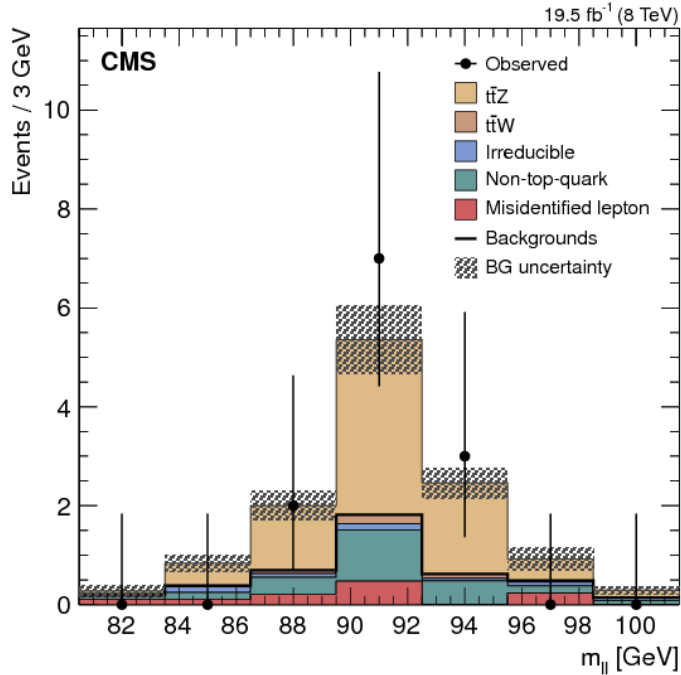
$$C_V^{\text{SM}} = \frac{T_t^3 - 2Q_t \sin^2 \theta_w}{2 \sin \theta_w \cos \theta_w},$$

$$C_{1,V} = C_{1,V}^{\text{SM}}$$

$$C_A^{\text{SM}} = \frac{-T_t^3}{2 \sin \theta_w \cos \theta_w},$$

$$C_{1,A} = C_{1,A}^{\text{SM}}$$

$t\bar{t} + Z$



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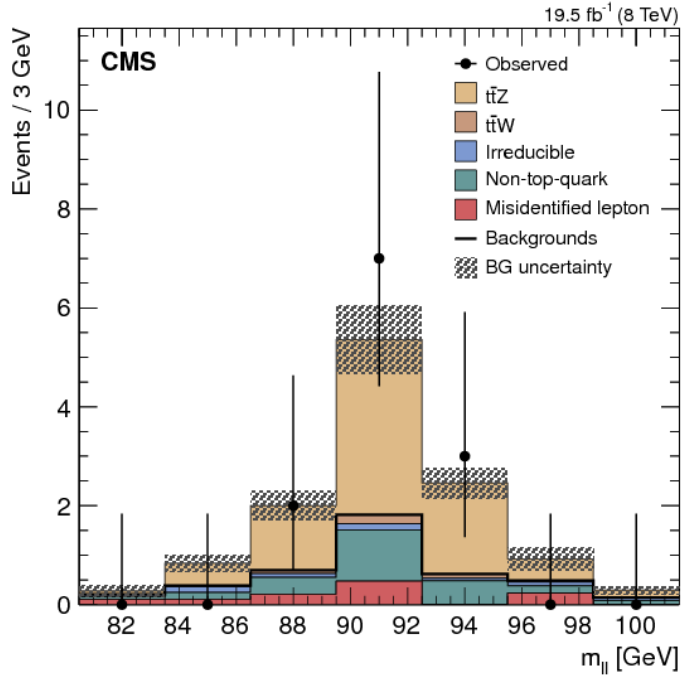
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$$C_{1,V} = C_{1,V}^{\text{SM}} + \Delta C_{1,V}$$

$$C_{1,A} = C_{1,A}^{\text{SM}} + \Delta C_{1,A}$$

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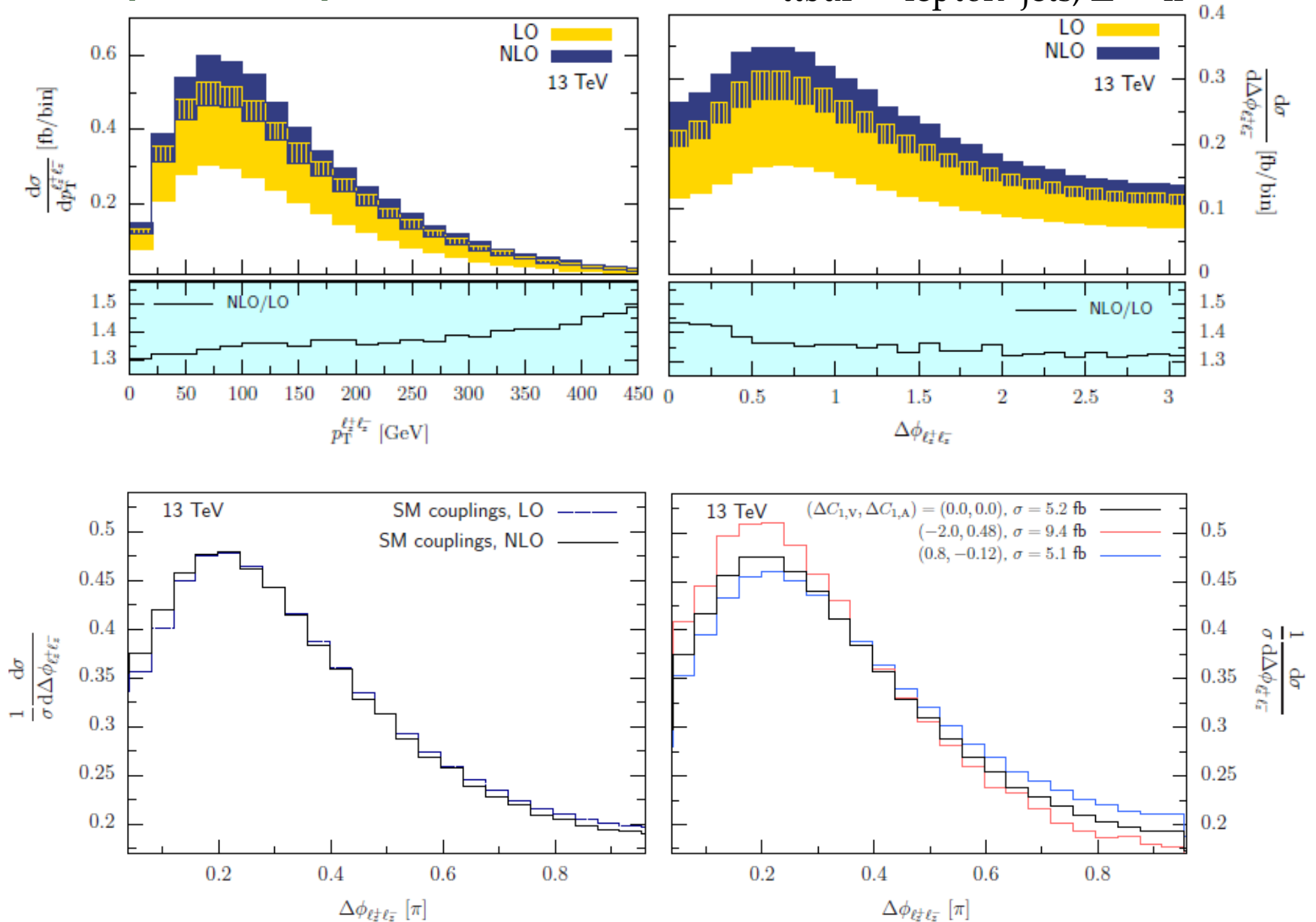
$$C_{1,V} = C_{1,V}^{\text{SM}} + \left(\frac{v^2}{\Lambda^2} \right) \text{Re} \left[C_{\phi q}^{(3,33)} - C_{\phi q}^{(1,33)} - C_{\phi u}^{33} \right],$$

$$C_{1,A} = C_{1,A}^{\text{SM}} + \left(\frac{v^2}{\Lambda^2} \right) \text{Re} \left[C_{\phi q}^{(3,33)} - C_{\phi q}^{(1,33)} + C_{\phi u}^{33} \right],$$

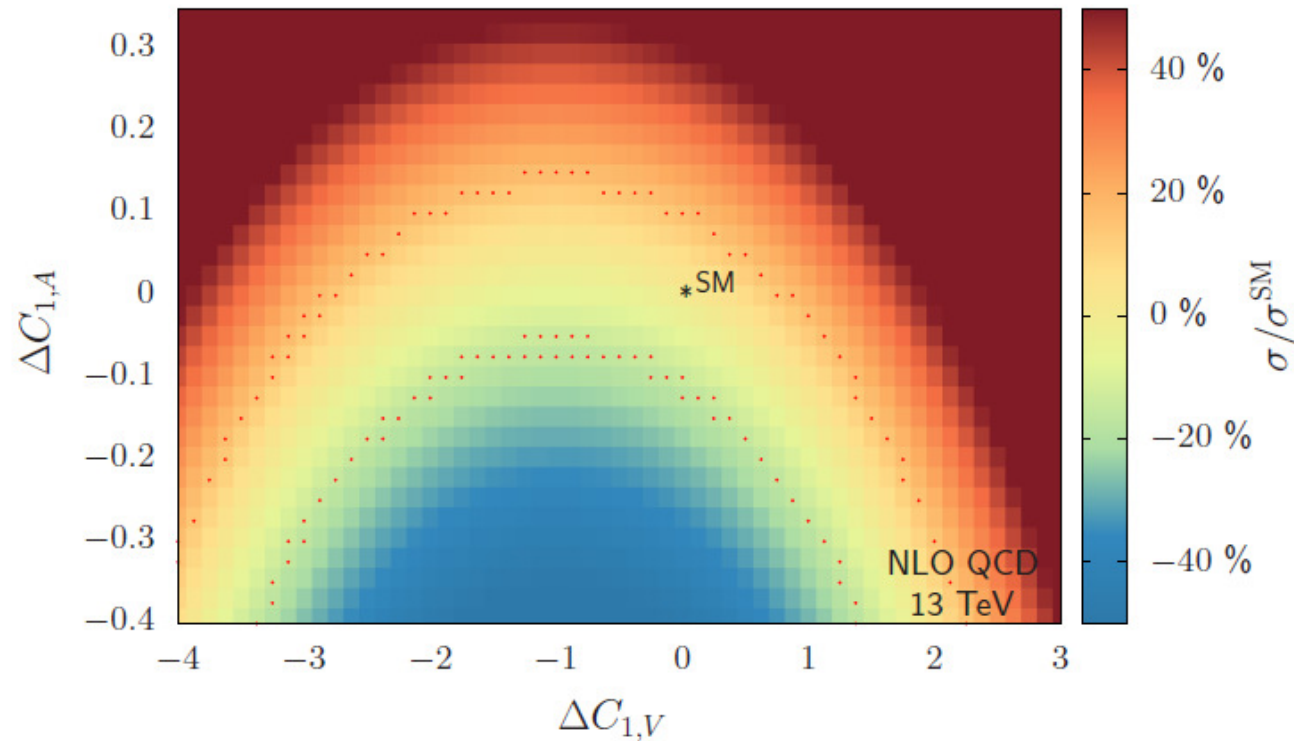
ttbar + Z

[Röntsch, M.S.]

ttbar → lepton+jets, Z → ll



$t\bar{t} + Z$



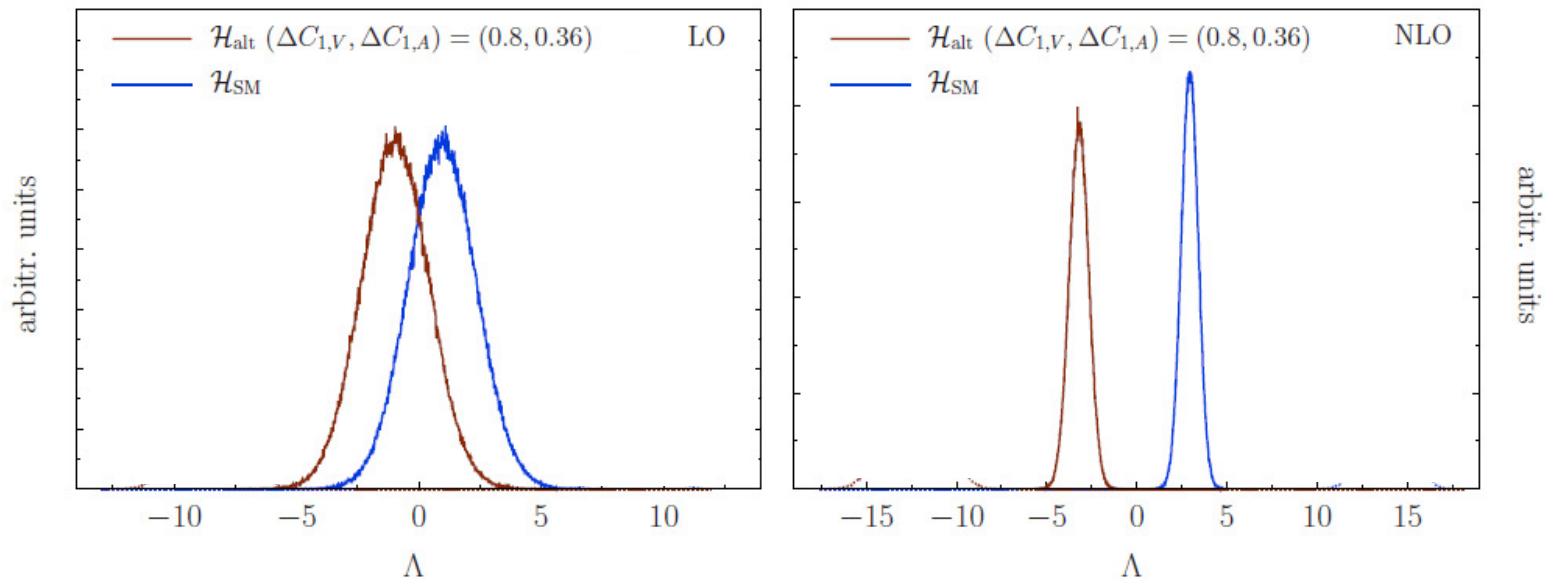
- Non-SM coupling combinations within the 20% band cannot be distinguished from the SM.
- Additional shape information is needed

ttbar + Z

- We employ a log-likelihood ratio test which accounts for statistical uncertainties, experimental systematics and theoretical uncertainties from scales+pdfs

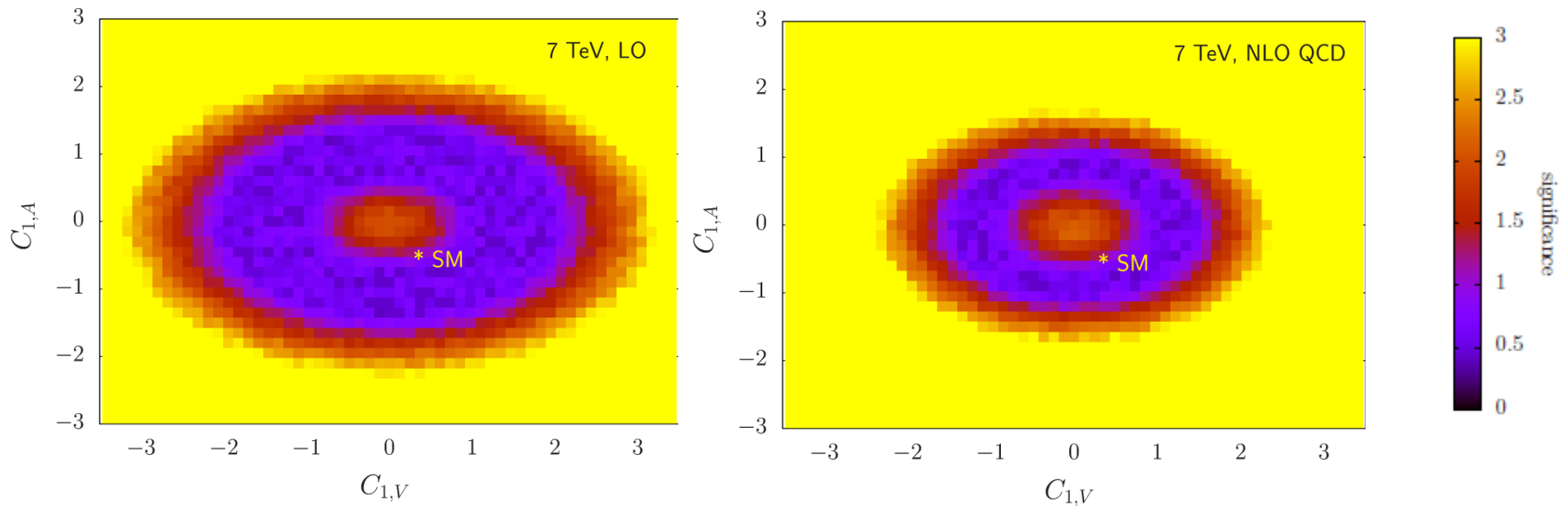
$$\mathcal{L}(\mathcal{H}|\vec{n}) = \prod_{i=1}^{N_{\text{bins}}} P_i(n_i|\nu_i^{\mathcal{H}})$$

$$\Lambda(\vec{n}_{\text{obs}}) = \log\left(\frac{\mathcal{L}(\mathcal{H}_{\text{SM}}|\vec{n}_{\text{obs}})}{\mathcal{L}(\mathcal{H}_{\text{alt}}|\vec{n}_{\text{obs}})}\right)$$



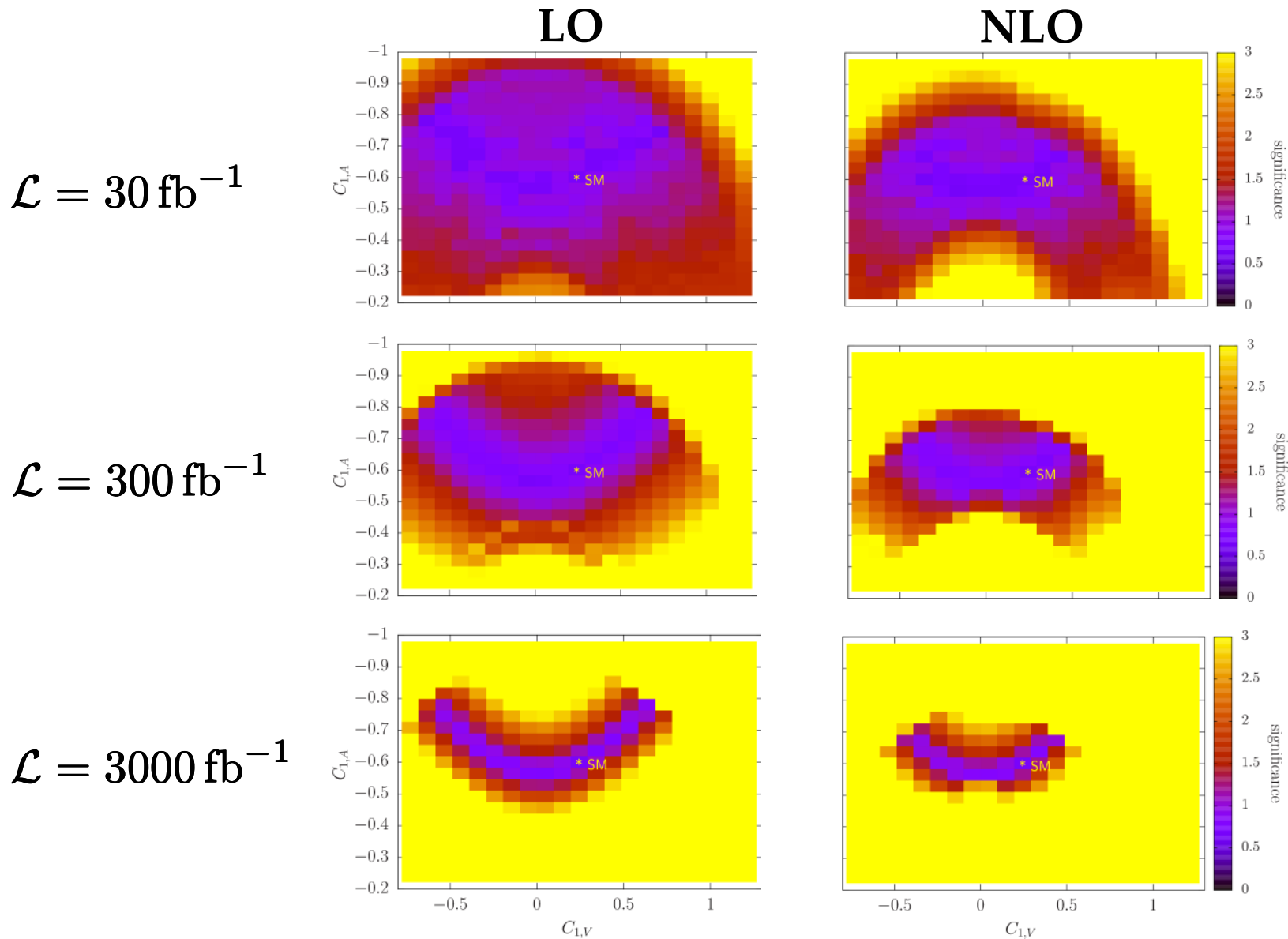
$t\bar{t} + Z$

- First measurement of total cross section by CMS
- $H_{\text{null}} = \text{CMS data}$ vs. $H_{\text{alt}} = \text{some BSM coupling choice}$
- Uncertainty treatment:
 - Theoretical*: scale + pdfs 40% at LO and 15% at NLO (uniform/flat distr.)
 - Experimental*: statistical (Poisson) + 20% systematics (Gaussian)

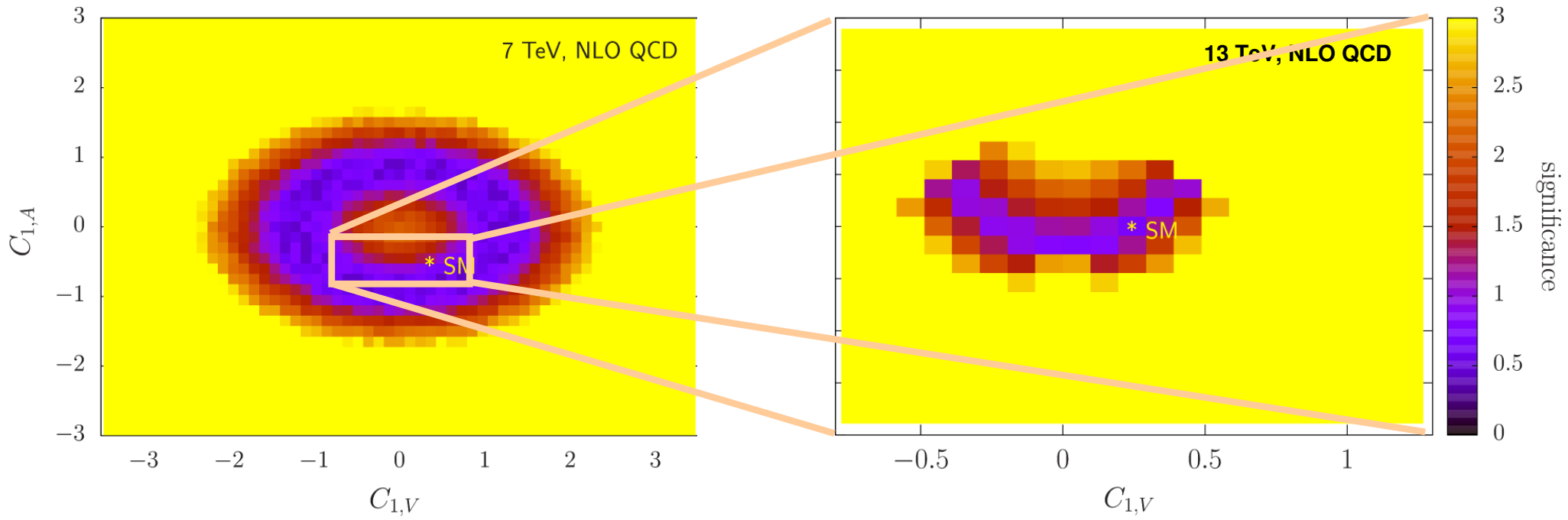


ttbar + Z

Future projection of sensitivities at the 13 TeV LHC

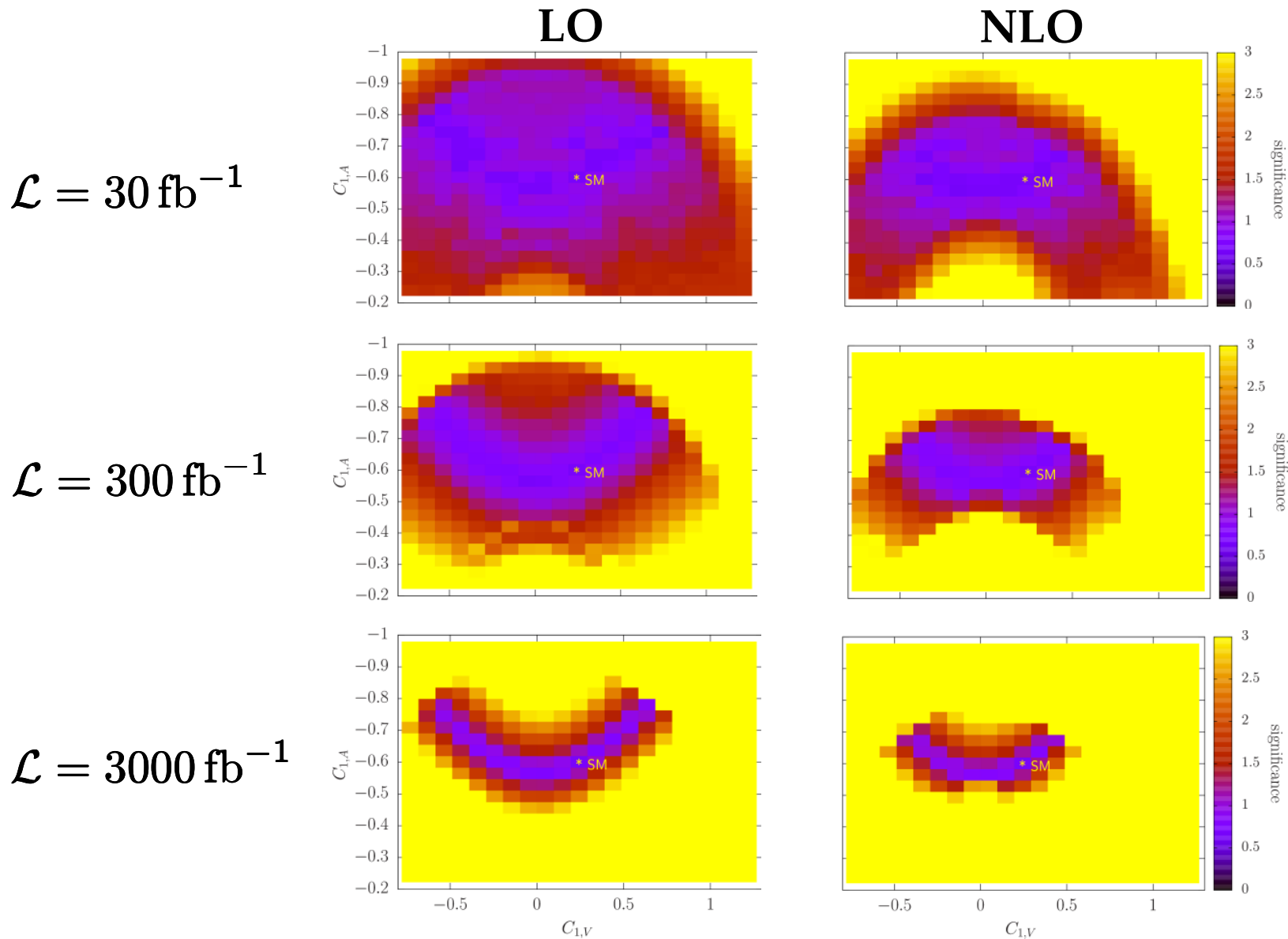


$t\bar{t} + Z$

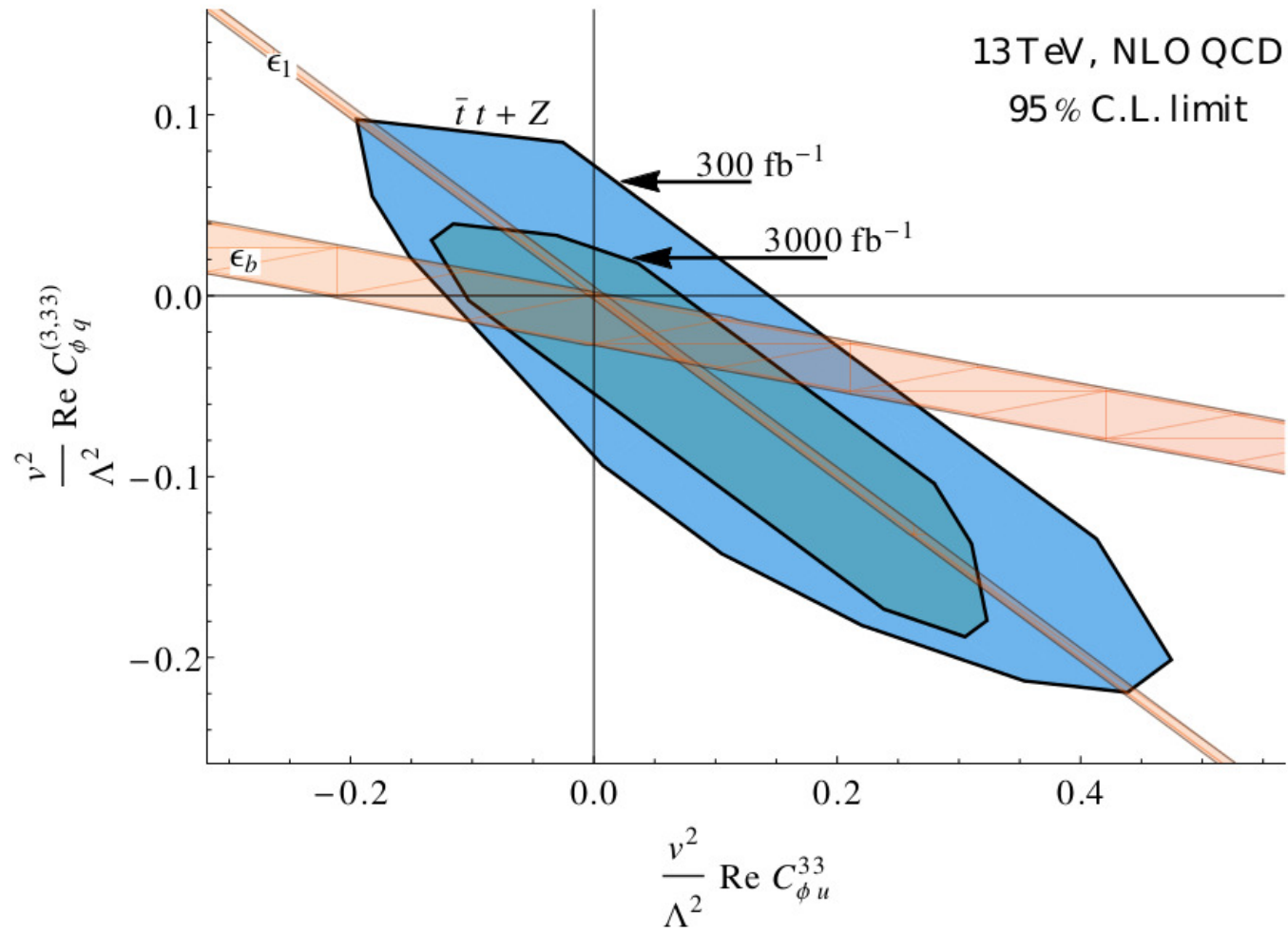


ttbar + Z

Future projection of sensitivities at the 13 TeV LHC



$t\bar{t} + Z$



SUMMARY

- LHC will provide large event samples of top quark pairs in association with electroweak bosons. The study of these processes will mark a new era in top quark physics.
- Direct sensitivity to top electroweak couplings
- Photon radiation off the decay products is crucial for accurate description
- Analyses at NLO significantly improve the sensitivity to couplings