

# Non-associative Deformations of Geometry in Double Field Theory

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based on JHEP 04(2014)141 by R. Blumenhagen, MF, F. Haßler, D. Lüst, R. Sun

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# Outline

- Deformation Quantization
- T-Duality and non-associativity in String Theory
- Non-associative Deformations of Geometry in DFT

# Canonical Quantization

Replace Poisson-bracket by commutator:

$$\{x, p\}_{PB} = 1 \quad \rightarrow \quad [x, p] = i\hbar$$

Fulfilled for instance by operators:  $\hat{p} = -i\hbar \frac{\partial}{\partial x}$

## Deformation Quantization:

No operators, instead change multiplication law: Replace  $f \cdot g$  by

$$f \star g := f \cdot g + \frac{i\hbar}{2} \begin{pmatrix} \partial_x f \\ \partial_p f \end{pmatrix}^T \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \begin{pmatrix} \partial_x g \\ \partial_p g \end{pmatrix}$$

Insert coordinate and momentum:

$$\left. \begin{aligned} x \star p &= x \cdot p + \frac{i\hbar}{2} \\ p \star x &= p \cdot x - \frac{i\hbar}{2} \end{aligned} \right\} [x, p] = x \star p - p \star x = i\hbar$$

## Take home message:

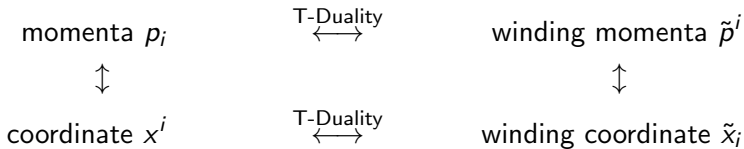
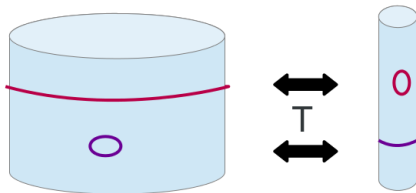
Commutation relations realized by deformed  
product 9709040

$$f \star g := f \cdot g + \frac{i\hbar}{2} \omega^{ij} \partial_i f \partial_j g + \mathcal{O}(\hbar^2)$$



# T-Duality

Closed strings wind around compactified dimensions:



# Non-geometric Fluxes

T-Duality mixes  $G$  and  $B \Rightarrow$  change of geometry



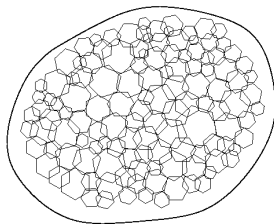
$$H_{ijk} \quad \longleftrightarrow^{T^k} \quad F_{ij}{}^k \quad \longleftrightarrow^{T^j} \quad Q_i{}^{jk} \quad \longleftrightarrow^{T^i} \quad R^{ijk}$$



# Non-associative Geometry

Blumenhagen, Lüst, Plauschinn et alii:

$$[x^a, x^b] \cong R^{abc} p_c$$



“fuzzy” geometry due to Heisenberg uncertainty:

$$\Delta x^a \Delta x^b \cong \langle [x^a, x^b] \rangle \neq 0$$

non-vanishing Jacobi identity!  $\hat{=}$  non-associative operators!

Not possible in ordinary quantum mechanics!

Deformed product vanishes for observables by momentum  
conservation! 1106.0316 by Blumenhagen, Deser, Lüst, Plauschinn, Rennecke

**Our work: Investigate in double field theory how  
non-associativity vanishes!**

## Double Field Theory

Combine  $\begin{pmatrix} \text{normal} \\ \text{winding} \end{pmatrix}$  in 2D vector

$$P_M = \begin{pmatrix} p_i \\ \tilde{p}^i \end{pmatrix} \quad \partial_M = \begin{pmatrix} \partial_i \\ \tilde{\partial}^i \end{pmatrix} \quad X_M = \begin{pmatrix} x_i \\ \tilde{x}^i \end{pmatrix}$$

⇒ Coordinates and winding on equal footage!

**BUT: Constraints needed for consistency!**

# Non-associative Deformations of Geometry in DFT

Translate deformed product into DFT:

$$f \Delta g \Delta h = f \cdot g \cdot h + \underbrace{\mathcal{F}_{ABC}}_{\text{contains H,f,Q,R}} \partial^A f \partial^B g \partial^C h$$

We found: Deformation vanishes by consistency constraints!

Deformation in physical situations (action)  $\hat{=}$  integration:

$$\int_{DFT} \mathcal{F}^{ABC} \mathcal{D}_A f \mathcal{D}_B g \mathcal{D}_C h \stackrel{\text{PI}}{=} - \int_{DFT} \underbrace{\mathcal{Z}^{AB}}_{\text{Bianchi } \mathcal{Z}^{AB}=0!} f \mathcal{D}_A g \mathcal{D}_B h$$

## Another deformation:

DFT allows for another deformation:

$$f \Delta g \Delta h = f \cdot g \cdot h + \check{\mathcal{F}}_{ABC} \partial^A f \partial^B g \partial^C h$$

⇒ Generalization of open strings in B-field background 9812219

No reason to vanish! Integral:

$$\int_{DFT} \check{\mathcal{F}}^{ABC} \mathcal{D}_A f \mathcal{D}_B g \mathcal{D}_C h \stackrel{\text{PI}}{=} - \int_{DFT} \underbrace{\mathcal{G}^{AB}}_{\text{eom: } \mathcal{G}^{AB}=0!} f \mathcal{D}_A g \mathcal{D}_B h$$

## Conclusion

**associativity of observables preserved by crucial ingredients  
of double field theory**

$$\check{\mathcal{F}}_{ABC} \partial^A f \partial^B g \partial^C h$$

equation of motion

continuity equation

$$\mathcal{F}_{ABC} \partial^A f \partial^B g \partial^C h$$

Bianchi identity

closure of algebra

# Outlook

Future research directions:

- Derive higher orders of the product (ongoing)
- Non-associativity in Hamiltonian formalism