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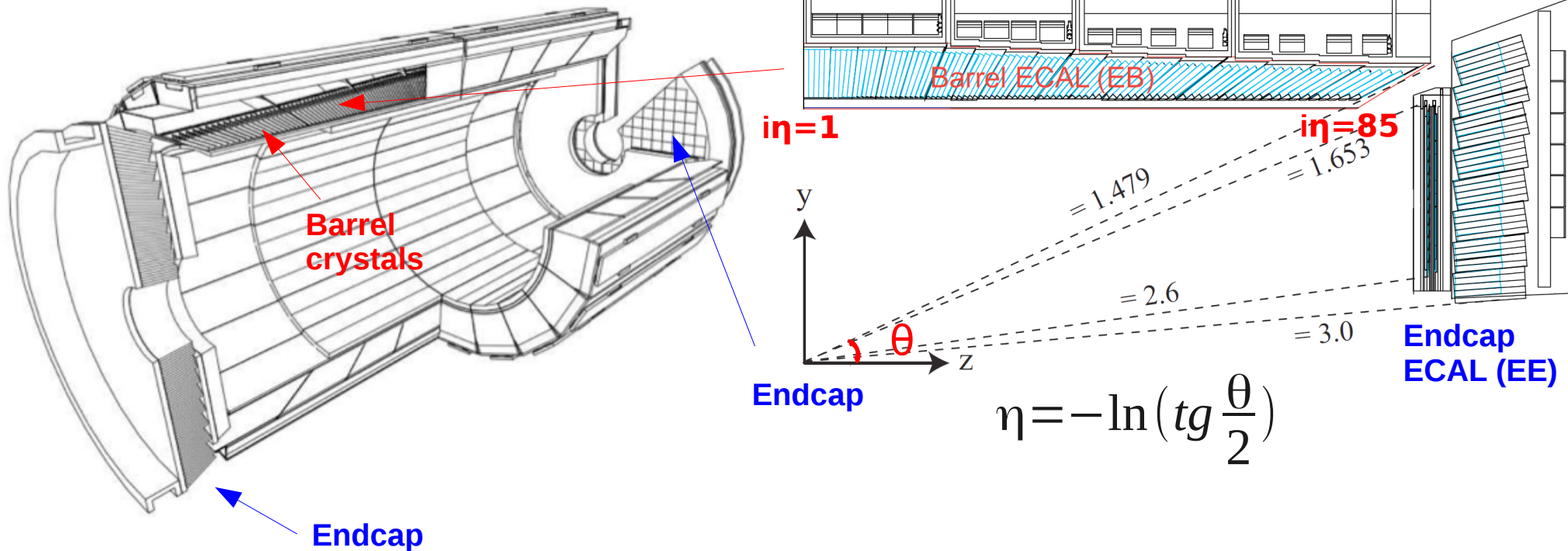


Tracker material study with the energy flow through the CMS electromagnetic calorimeter

Riccardo Paramatti , [Ambra Provenza](#)

The electromagnetic calorimeter (ECAL)

To detect photons and electrons



- Hermetic, homogeneous made of lead tungstate crystals (75848)
- **Barrel (EB)** $|\eta| < 1.479$
- **Endcaps (EE)** $1.479 < |\eta| < 3$
- Good energy resolution

Crystal identified with $(i\eta, i\phi)$

$i\eta \in [1, 85] \rightarrow \text{EB+}$

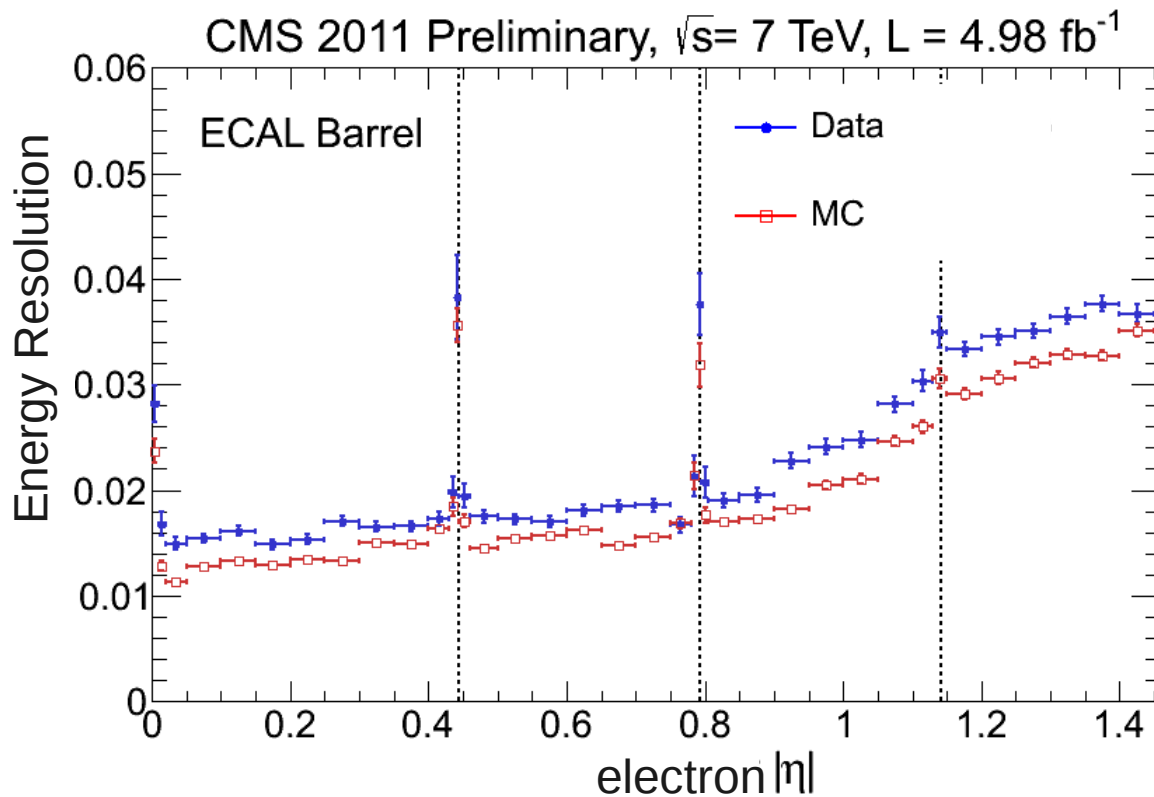
$i\eta \in [-85, -1] \rightarrow \text{EB-}$

360 crystals in $i\phi$ for each value of $i\eta$

Problem

ECAL energy resolution is a crucial parameter in a lot of analysis

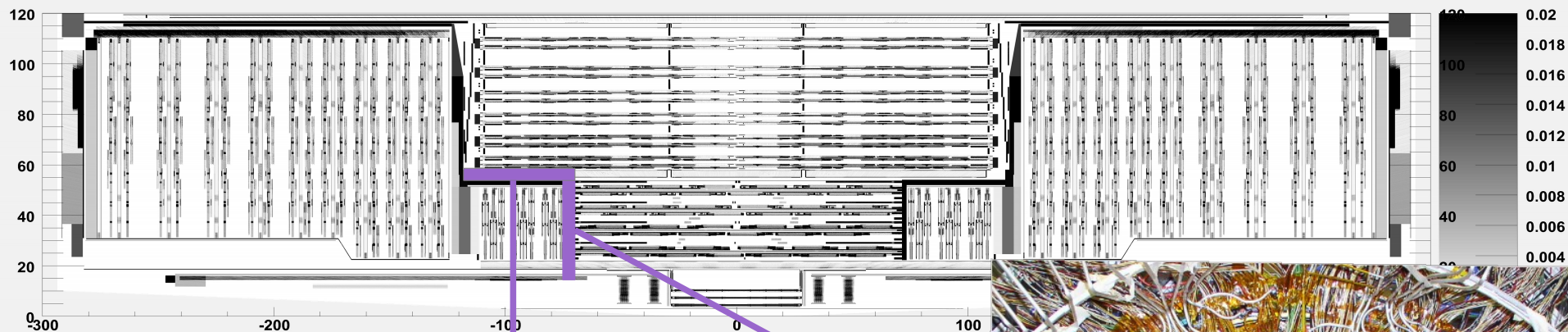
Problem: ECAL energy resolution obtained in $Z \rightarrow e^+e^-$ events \neq ECAL energy resolution expected by MonteCarlo simulation



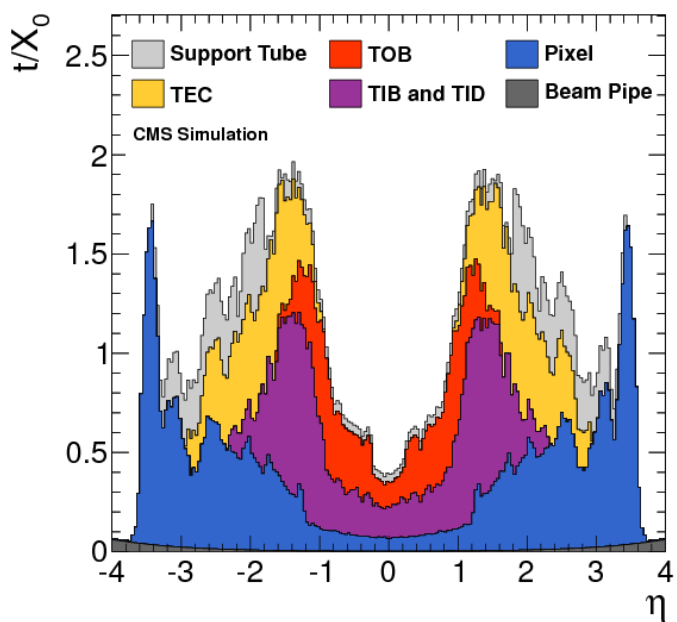
Cause: Inaccurate description of the tracker material (services)

Problem

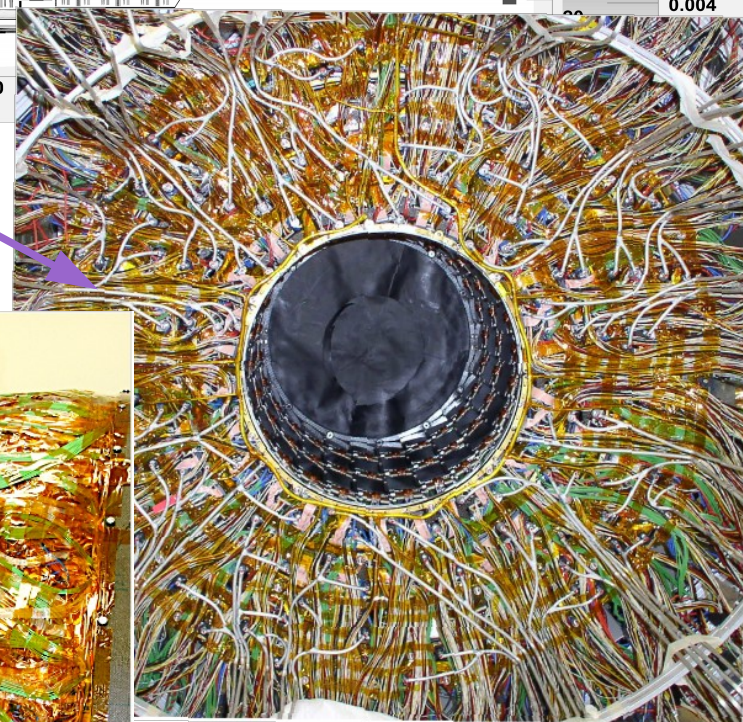
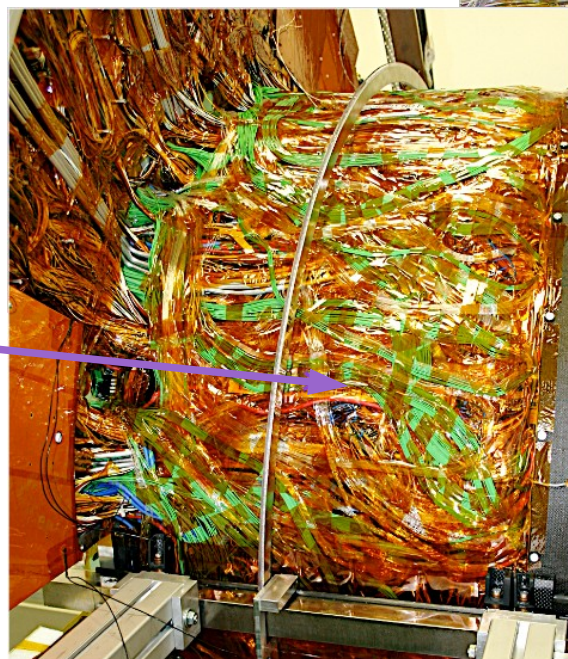
Tracker Structure



Distribution of the tracker material as a function of η



SERVICES



Problem

Combined action of the tracker **material** and the **magnetic field (B)**

For the **material**
 $\gamma \rightarrow e^+e^-$

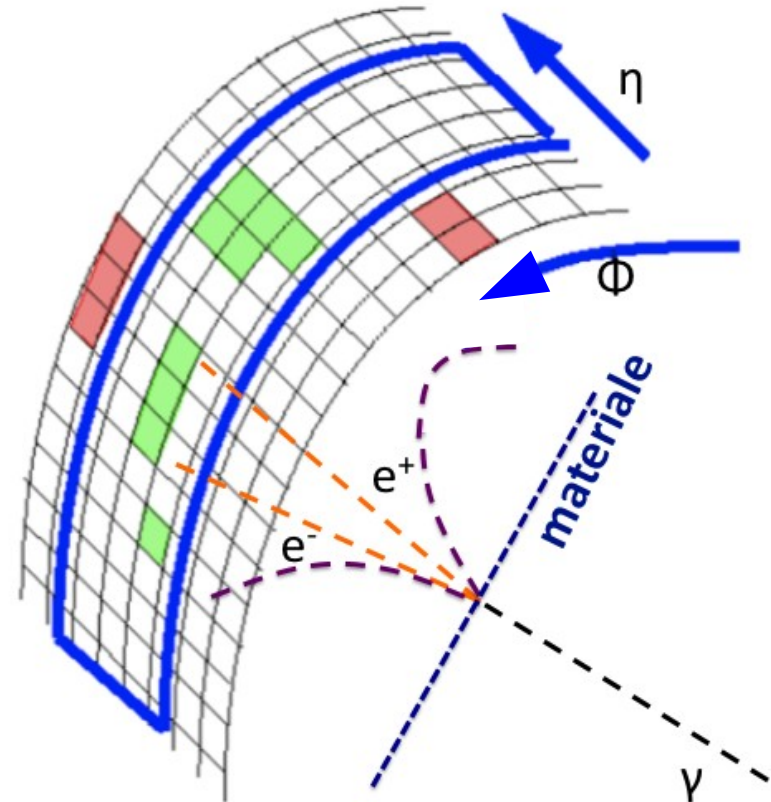
B changes the trajectory of e^+ and e^- .



some e^+ or e^- can not reach ECAL

Goal of the study:

Obtain a better description of the tracker material than the estimate made during the detector construction



e^+e^- without **B**
 e^+e^- with **B**

Methods to estimate the tracker material

Study the momentum lost:

Momentum variation \propto material

- Use of high energy electrons that radiate photons for bremsstrahlung

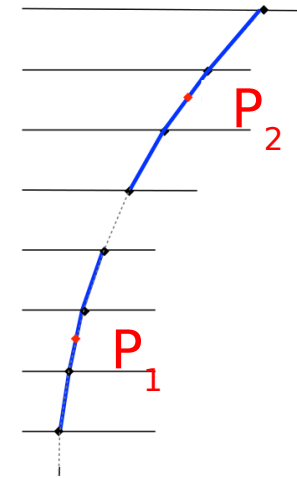
$$f_{brem} = \frac{p_i - p_f}{p_i}$$

- Use of charged pion with $P_t \sim 1$ GeV that do multiple scattering: study of the difference between P_2 e P_1

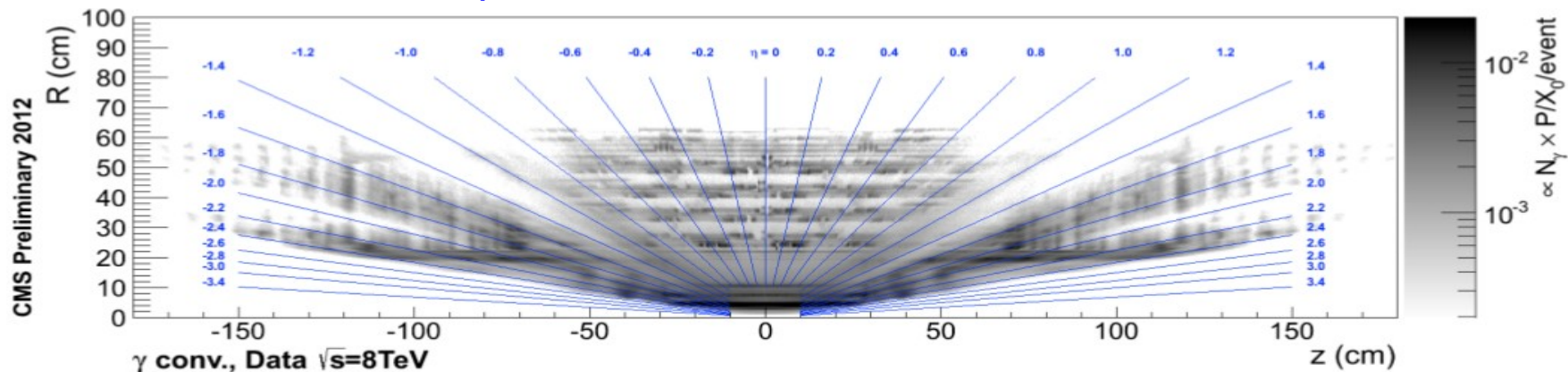
Conversion method:

In presence of material $\gamma \rightarrow e^+e^-$

Pairs produced \propto material



Conversions map



Methods non sensitive to the outer layers of the tracker material

The Energy Flow method

Energy Flow through the ECAL crystals:

$$S_{xtal} = \sum_i (E_t^i)_{xtal}$$

Transverse energy sum for an high number of Minimum Bias events

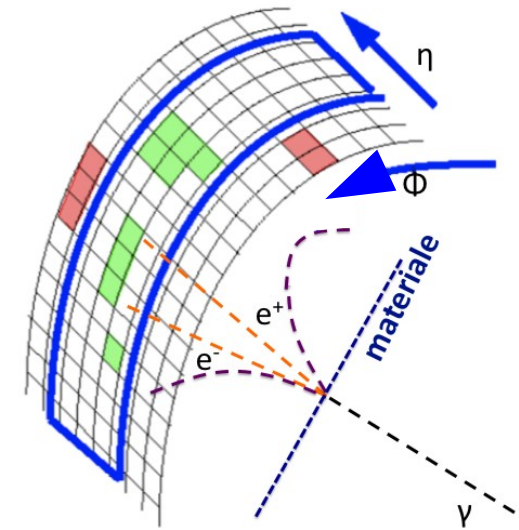
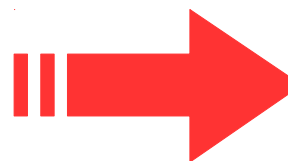
Energy Flow ratio in ECAL crystals:

$$R_{xtal} = \frac{(S_{xtal})_{Boff}}{(S_{xtal})_{Bon}}$$



Gives a measure of the amount of tracker material because of the combined action of the **tracker material** and the **magnetic field**.

In fact



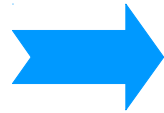
$$(S_{xtal})_{Bon} < (S_{xtal})_{Boff}$$

Data

Run of Minimum Bias events with **magnet off (Boff)** ($\sim 2.6 \cdot 10^8$ events)
Run of Minimum Bias events with **magnet on (Bon)** ($\sim 2.0 \cdot 10^8$ events)
Taken in the same period

In the analysis only crystals with energy deposits between a **lower threshold** and an **upper threshold** have been considered

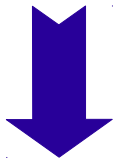
$$E_{\min} = 400 \text{ MeV}$$



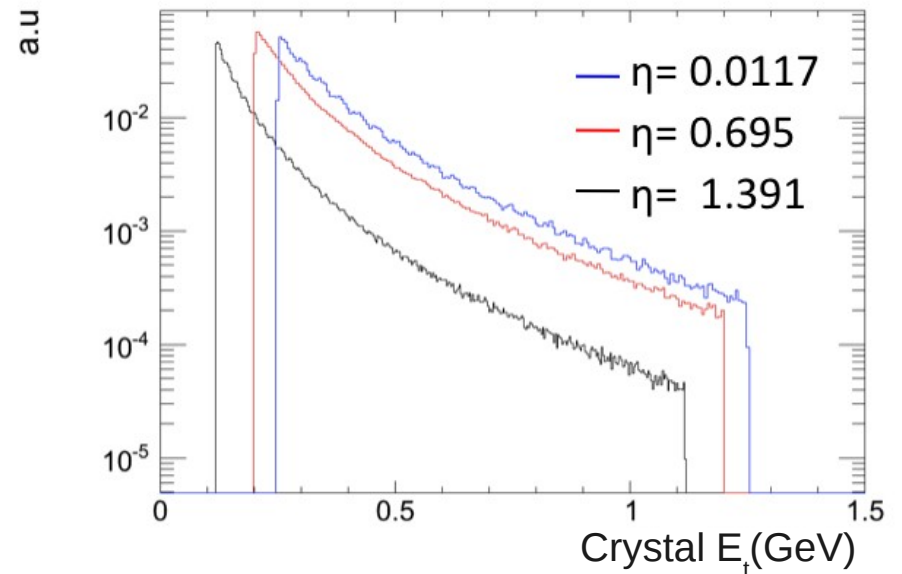
To cut the noise

Maximum value of transverse energy

$$(E_t)_{\max} = (E_t)_{\min} + 1 \text{ GeV}$$



To reduce fluctuations
caused by rare high energy
deposits



Average number of energy deposits
between the cuts $\sim 3.7 \cdot 10^5$

Analysis Flow

- Obtain energy flow for magnet off and magnet on events separately
- Determine the energy flow ratio for each ECAL crystal
- Compute the corrections for the effect that influence the energy flow:
 - ◆ Beam spot position effect
 - ◆ Border effect
- Comparison between data and MonteCarlo (MC):

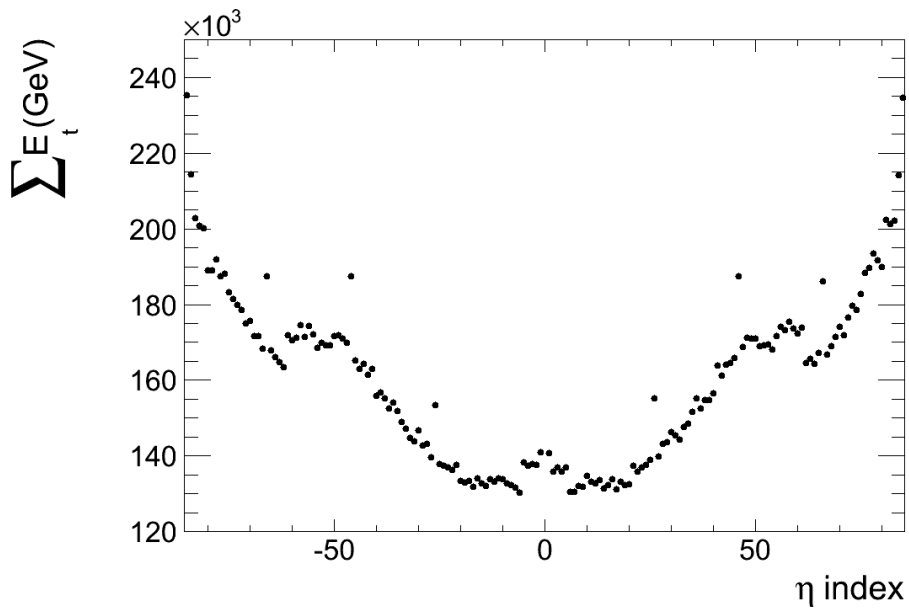
$$\frac{R_{data}}{R_{MC}} \neq 1 \longrightarrow$$

Tracker material not well implemented in MC

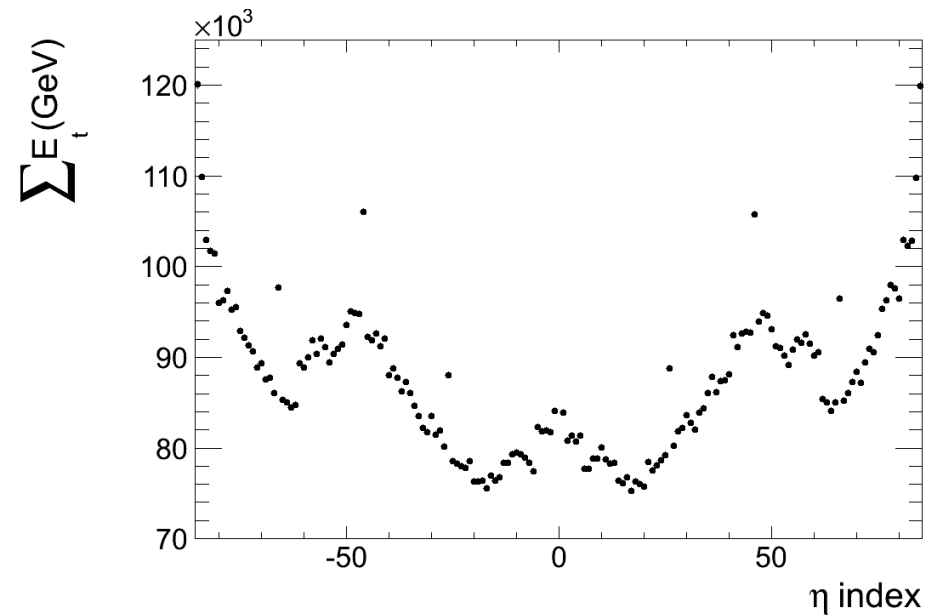
Energy Flow through ECAL

$S_{xtal} = \sum_i (E_t^i)_{xtal}$ as a function of η index for magnet off and magnet on data

Magnet off data



Magnet on data



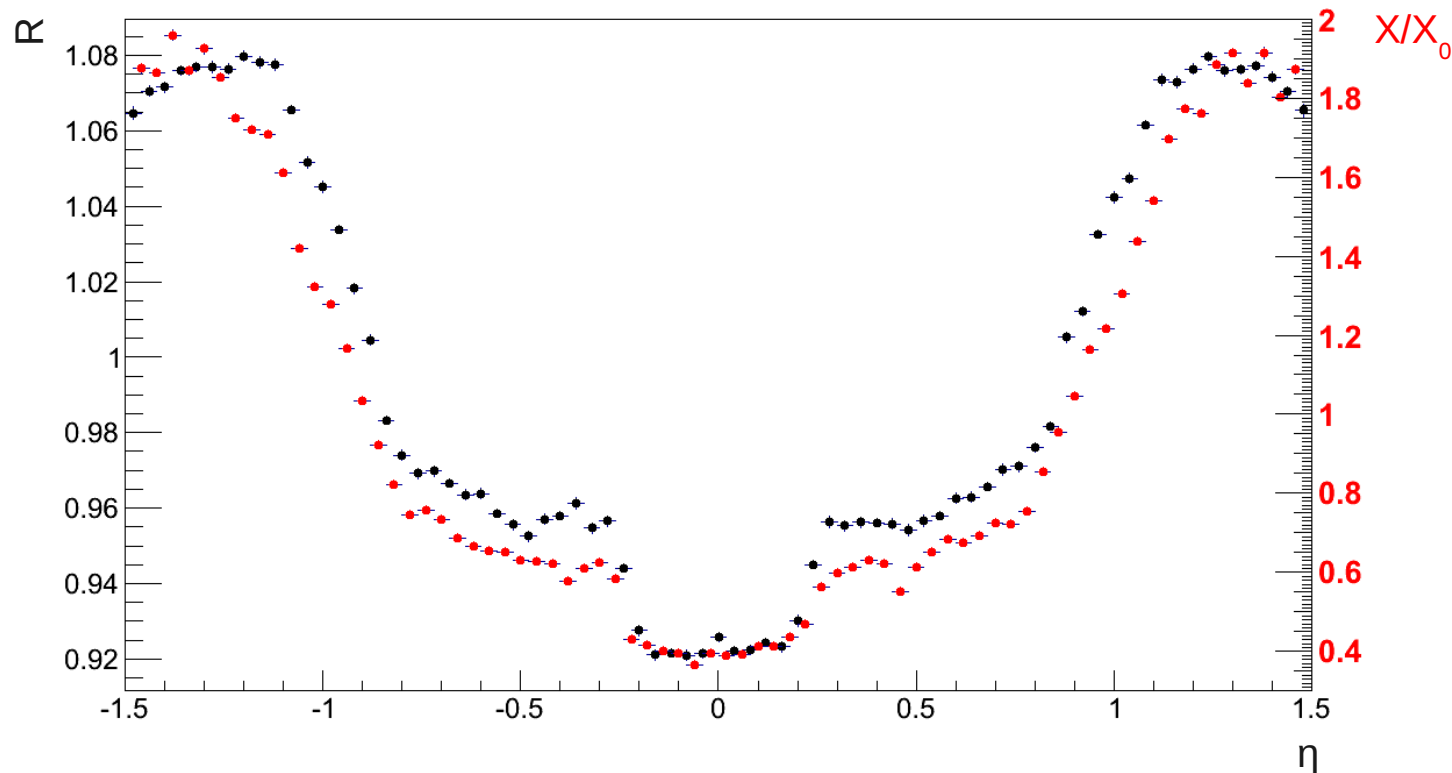
Energy Flow Ratio

Energy flow ratio

Mean of S_{xtal} in the
ECAL barrel

$$R_{xtal} = \frac{(S_{xtal} / \langle S \rangle)_{Boff}}{(S_{xtal} / \langle S \rangle)_{Bon}}$$

$$S_{xtal} = \sum_i (E_t^i)_{xtal}$$



Correlation
between R and
the **tracker**
material (X/X_0)

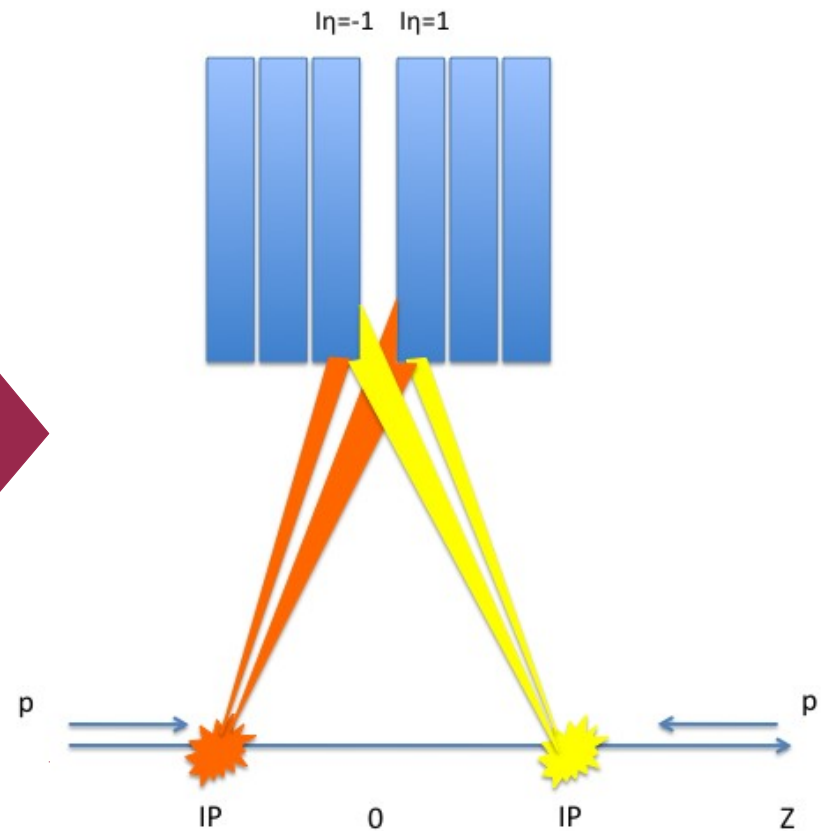
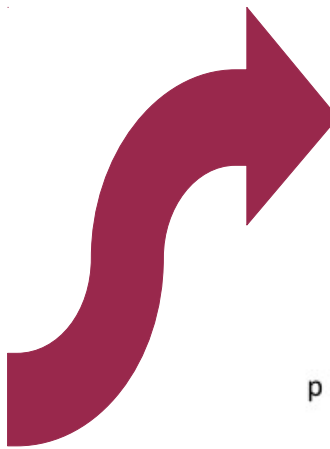
Beam spot position effect

The interaction point of the two beams doesn't match always with the center of the detector ($z=0$). It can take place at a few cm from $z=0$.

IP= Interaction Point Position

$z=0$ center of the detector

**Beam spot position
effect on the
energy flow**

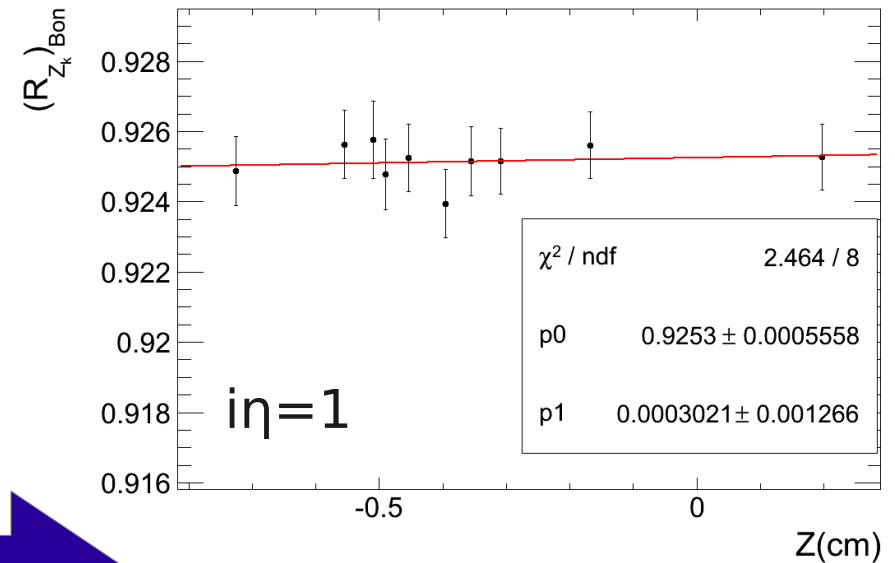
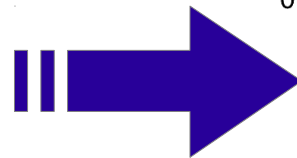
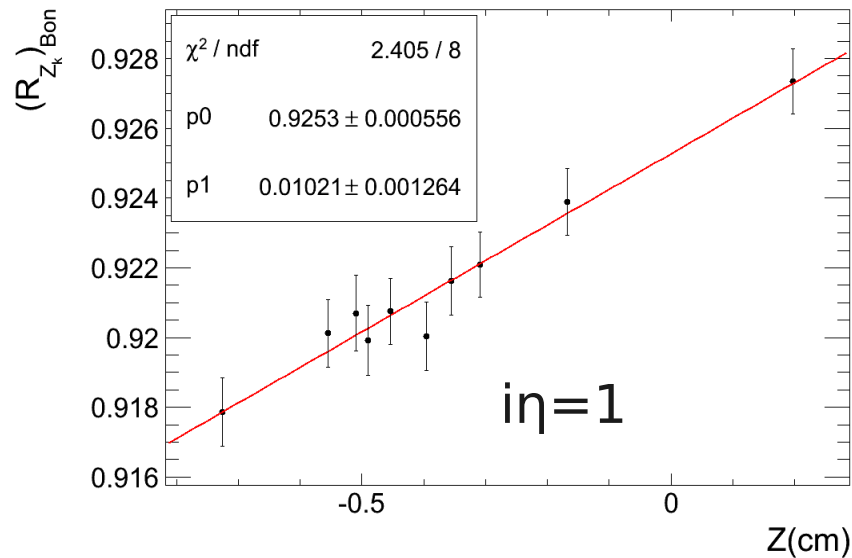


Beam spot position effect

Computing of the corrections:

R has been calculated in Z_k intervals and it has been studied as a function of Z

Bon

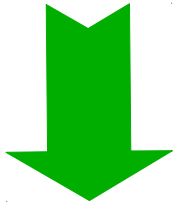


The corrections lead E_t and R to the value that they would have if IP is equal to $z=0$

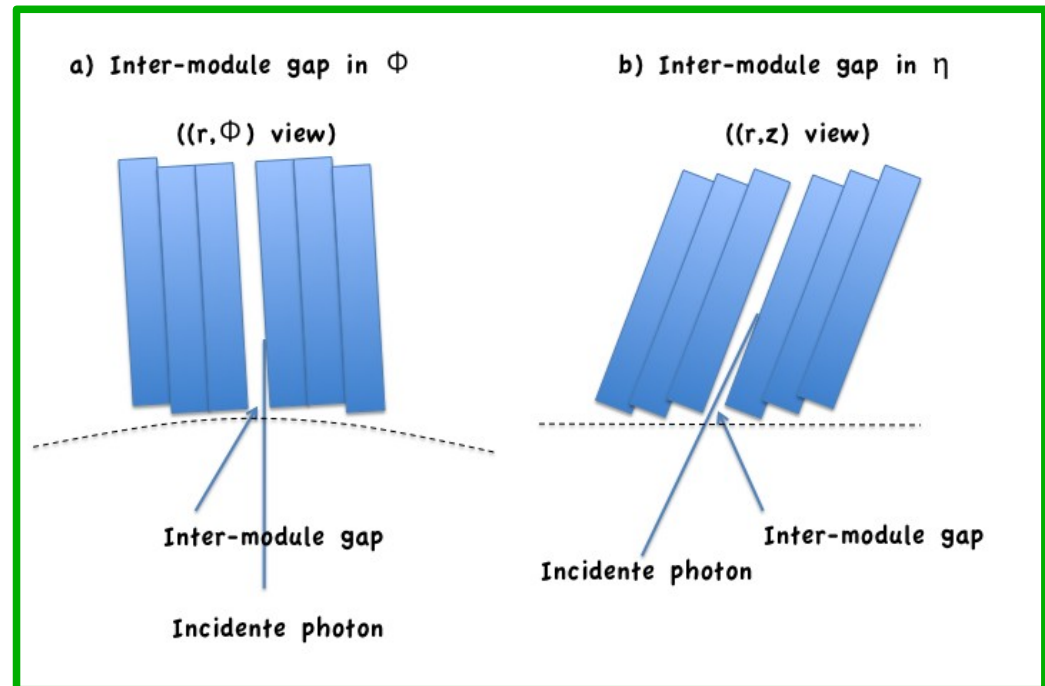
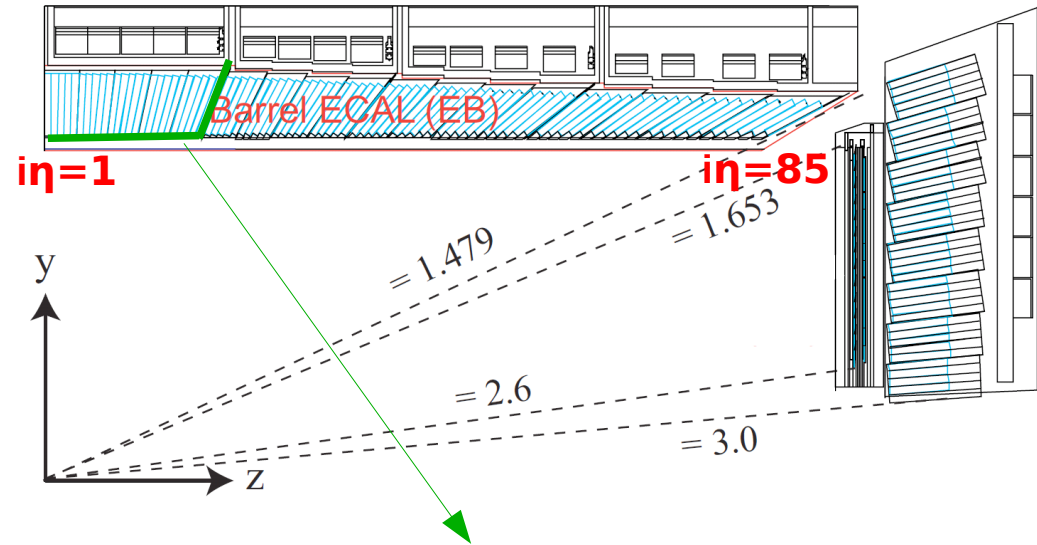
Border effect

The crystal axis is not pointing to the center of CMS. They are tilted by 3° with respect to the center of the detector, to maximize the detector acceptance.

The particles entering in the module gap hit the lateral face of the crystal at the border

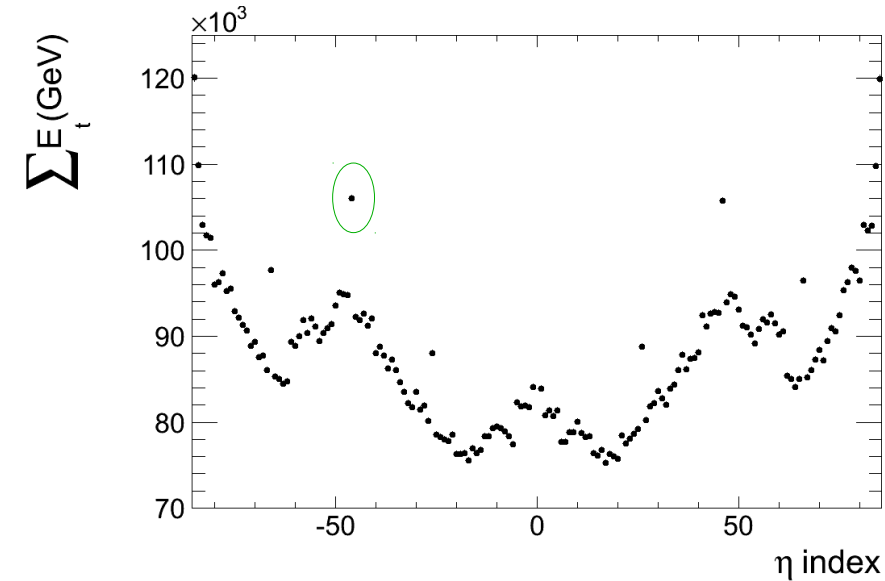


Border crystal collects more energy than the other



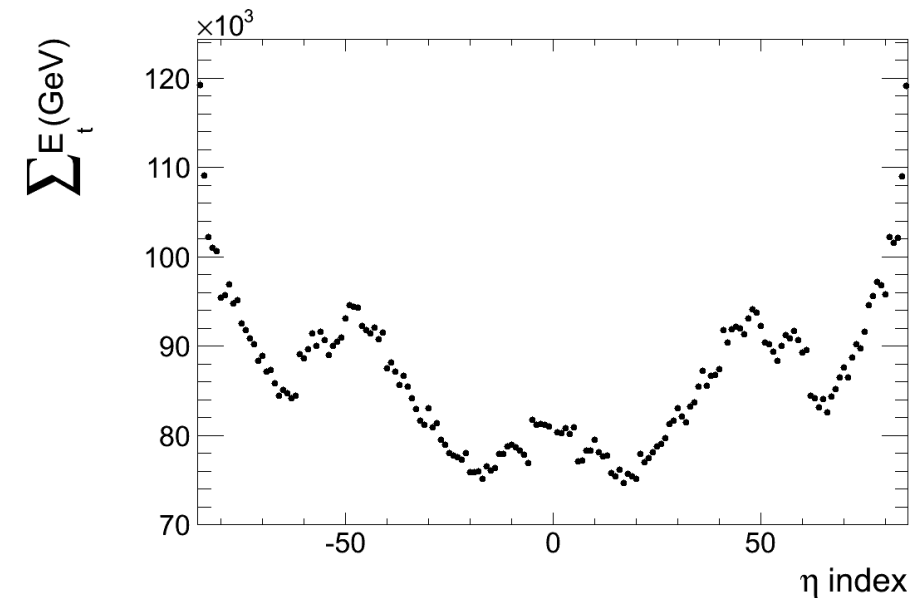
Border effect

Energy flow for magnet off data



The corrections have been determined by normalizing $\langle S_{\text{crystal}} \rangle$ to the average in the adjacent rings

Energy flow for magnet off data, corrected for the border effect

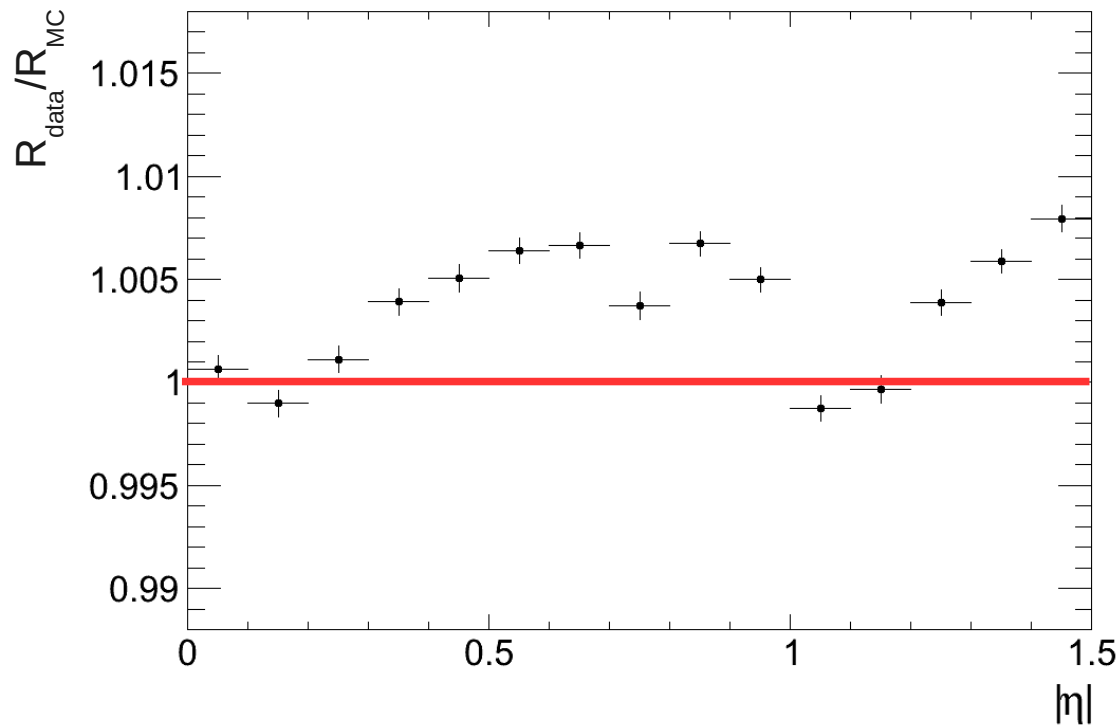


Corrections take values from $\sim 3\% \rightarrow \sim 15\%$

Comparison between data & MonteCarlo

The ratio $R_{\text{data}}/R_{\text{MC}}$ gives informations about the inaccuracy on the description of the tracker material.

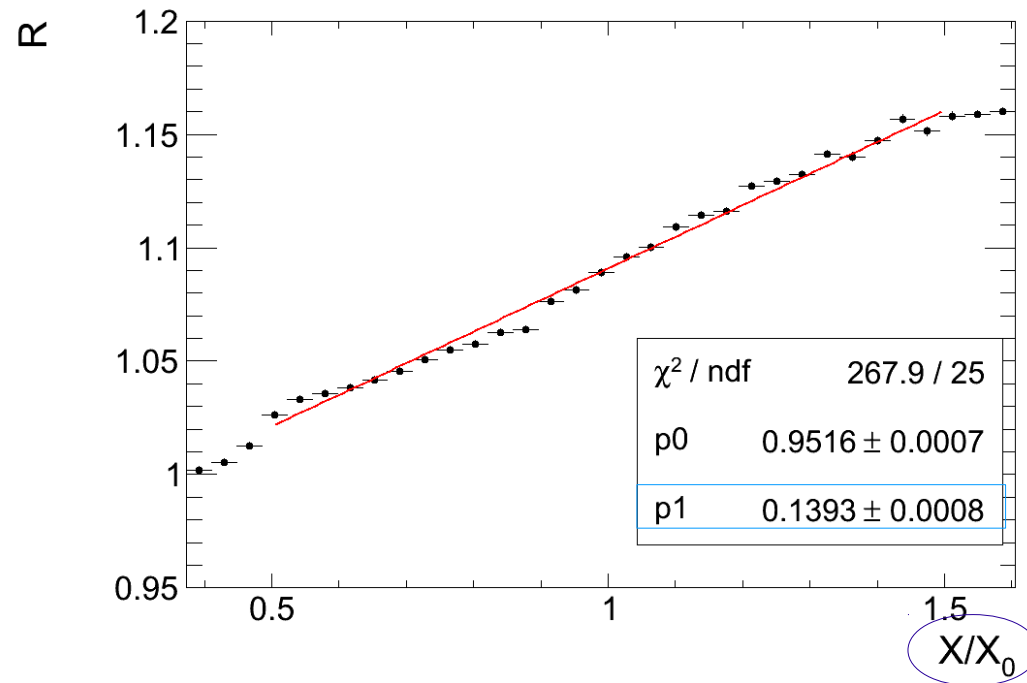
$R_{\text{data}}/R_{\text{MC}} \neq 1 \rightarrow$ tracker material not well implemented in the MC



These measures have to be calibrated to have a direct information on the needed additional material (in radiation length)

Comparison between data & MonteCarlo

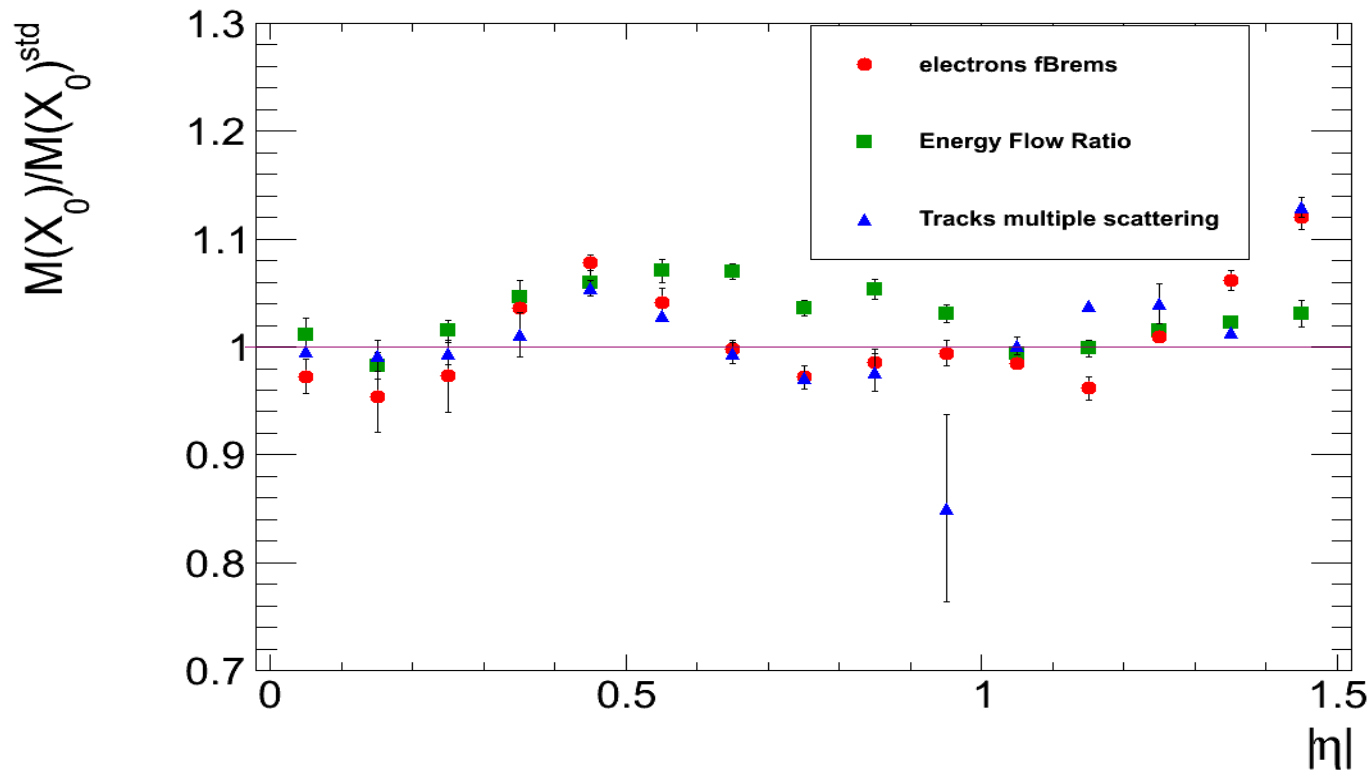
Mean of R as a function of X/X_0



$1/p1$ Factor used to convert $R_{\text{data}}/R_{\text{MC}}$ in the amount of material to be added to the material in the MC simulation ($\Delta M(X_0)$)

$$\frac{R_{\text{data}}}{R_{\text{MC}}} = 0.01 \rightarrow \Delta M(X_0) = 0.07 X_0$$

Results



$$M(X_0) = \boxed{M(X_0)^{std}} + \Delta M(X_0)$$

Measure of the material made during the detector construction

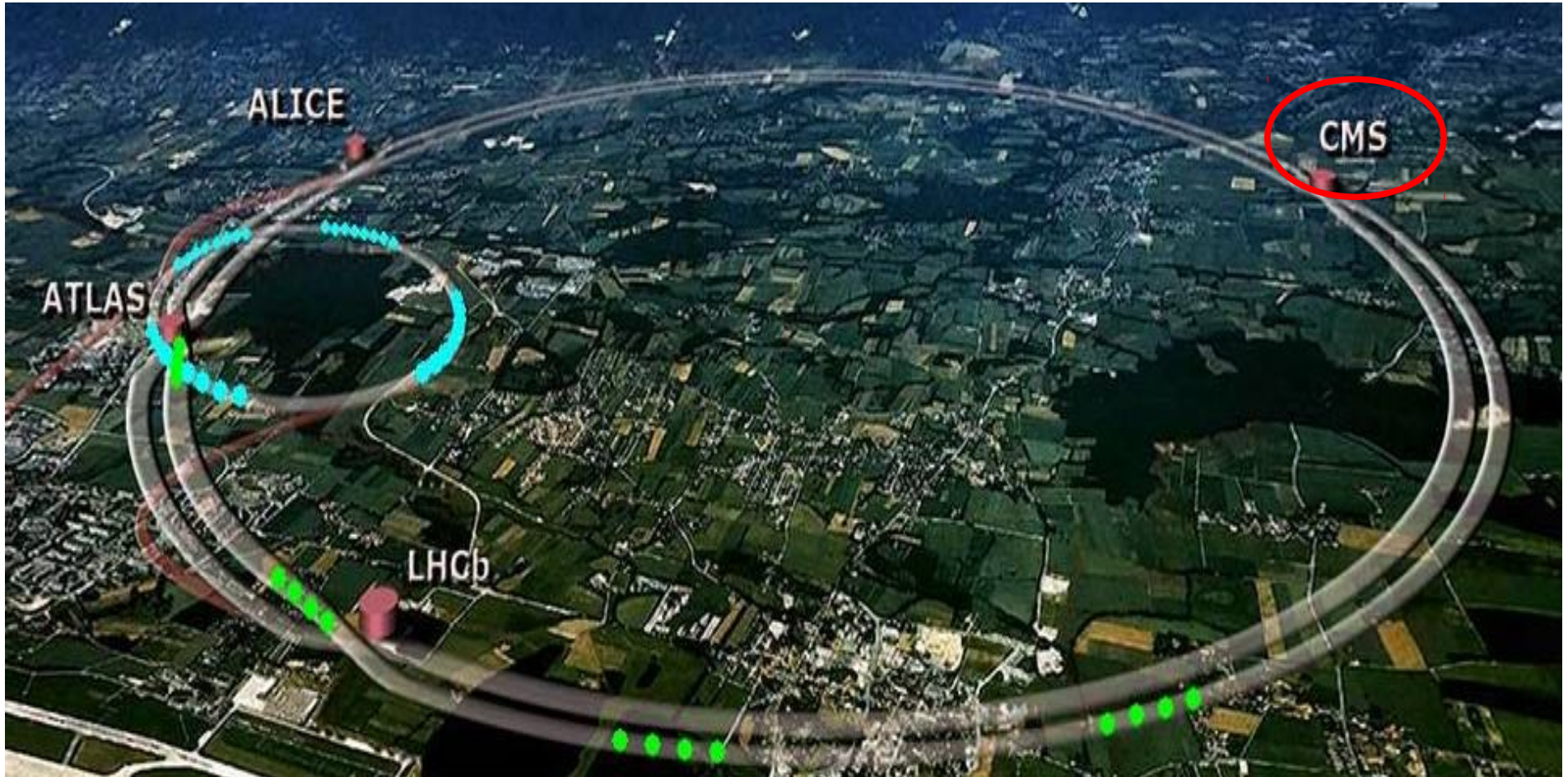
- Quite good agreement with other methods
- Larger additional material for $0.5 < |\eta| < 1$
- The **energy flow method** is the only one that take into account the outer layers of the tracker material

Conclusions

- ECAL energy resolution is a key parameter in a lot of analysis: Its knowledge is crucial
- A new method to study the amount of tracker material in front of ECAL has been proposed
- This method uses only calorimetric quantity:
It is the only one that take into account the outer layers of the tracker material
- Good agreement with the other methods
- These results are used in the new MC production aimed to the validation of the new measure of the tracker material through the data-MC compatibility in the energy resolution.

BACKUP

CMS experiment at LHC



CMS: Compact Muon Solenoid, one of the two multi-purpose experiment at **LHC**

LHC: Large Hadron Collider, p-p accelerator at CERN, Geneva
Run 1: 2009 → 2012 $\sqrt{s}=7-8$ TeV
Run 2: from 2015 $\sqrt{s}=13-14$ TeV

The electromagnetic calorimeter (ECAL)

	<i>PbWO₄</i>	NaI(Tl)	BGO
Density [g/cm ³]	8.28	3.67	7.13
Radiation length [cm]	0.89	2.59	1.12
Molière radius [cm]	2.2	4.5	2.4
Peak emission [nm]	425	410	480
LY(related to NaI(Tl)[%])	1.3	100	15
Time emission [ns]	5-15	250	300

Longitudinal distance for which an electron traversing the material loses on average 1/e of its energy through diffusion processes. For E ~ TeV the 98% of the longitudinal development is contained in 25X₀.

$$R_M = \frac{21.2 \text{ MeV} * X_0}{E_C [\text{MeV}]}$$

describe the transversal development of an electromagnetic shower. The 90% of the shower is contained in a cylinder with a radius equal to 3.5 R_M

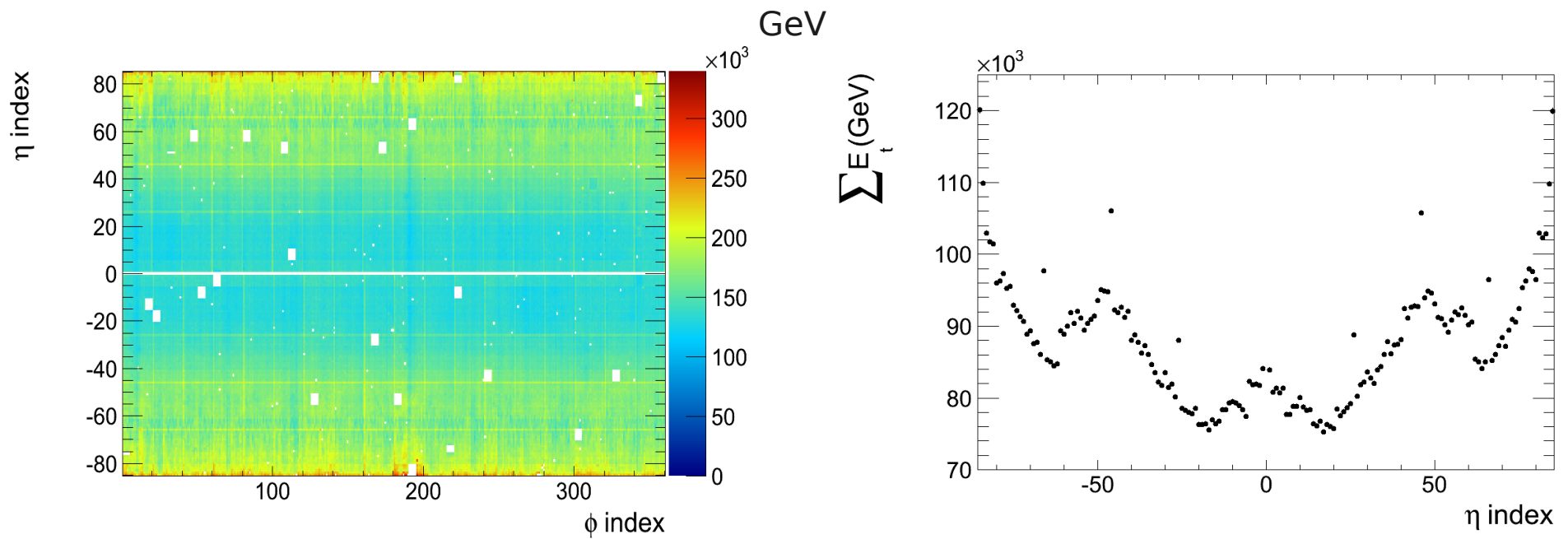
For e $E_C = \frac{610 \text{ MeV}}{Z}$

Energy resolution

$$\frac{\sigma(E)}{E} = \frac{S}{\sqrt{E}} \oplus \frac{N}{E} \oplus C$$

Energy Flow through ECAL

$$S_{xtal} = \sum_i (E_t^i)_{xtal} \quad \text{For magnet off data}$$

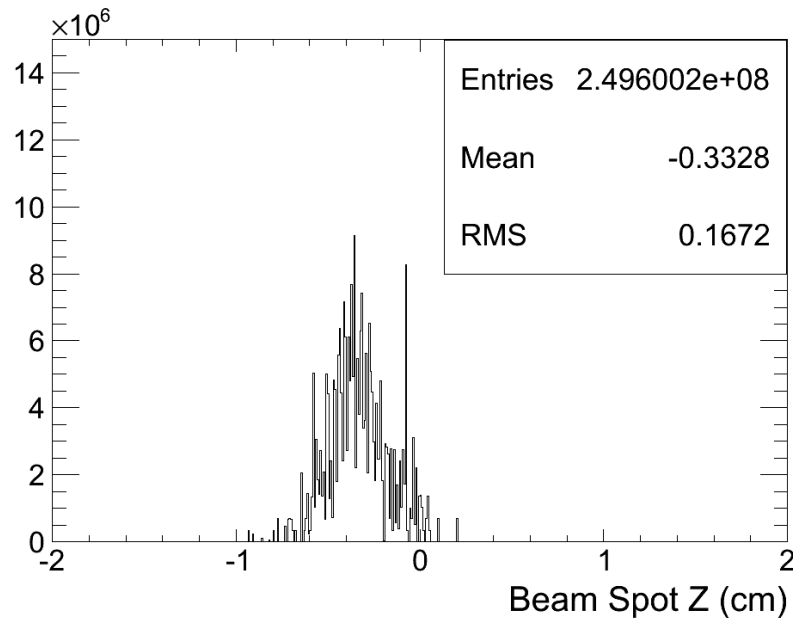


Beam spot position

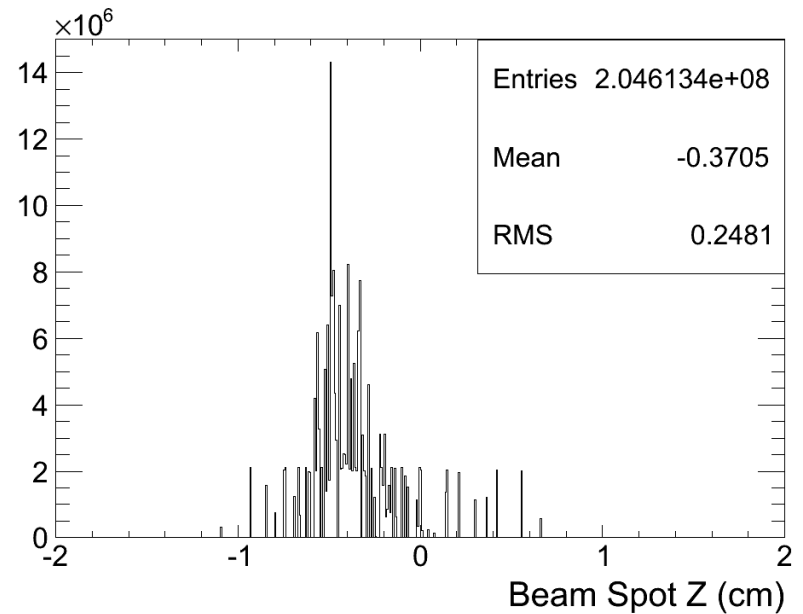
The interaction point of the two beams doesn't match always with the center of the detector ($z=0$). It can take place at a few cm from $z=0$.

Beam spot position distribution for the two groups of data

Magnet off data



Magnet on data



Beam spot position

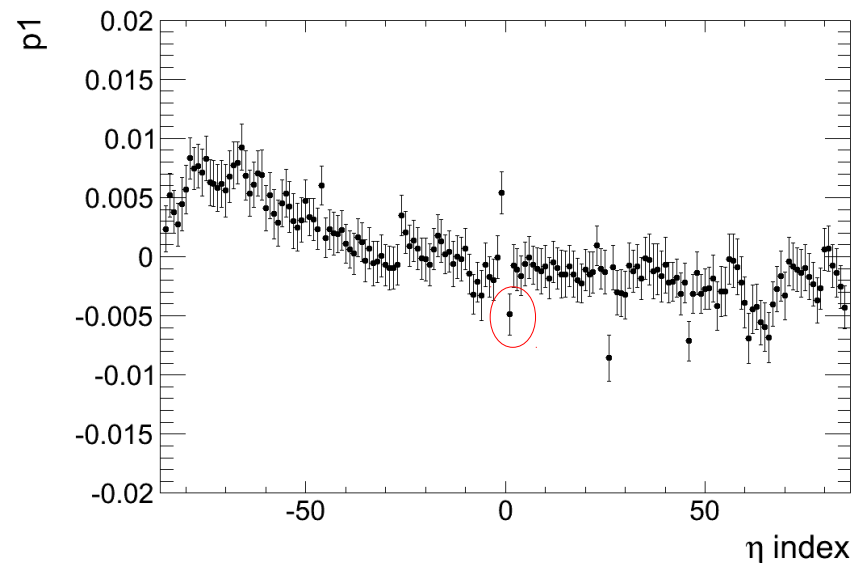
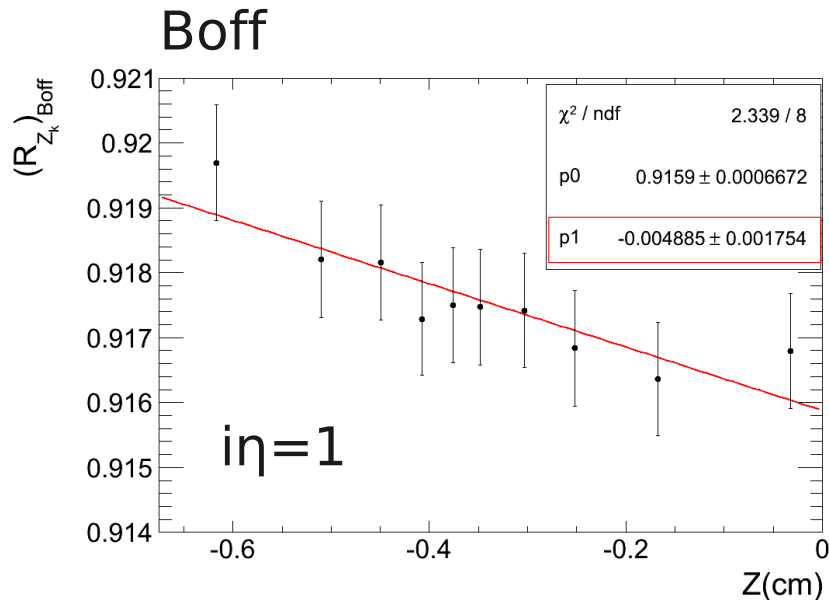
Computing of the corrections for Boff data:

The beam spot position distribution has been divided into 10 intervals (Z_k):

This quantity has been defined

$$(R_{Z_k})_{Boff} = \frac{((\sum E_t^i / \langle \sum E_t \rangle)_{Z_k})_{Boff}}{(\sum E_t^i / \langle E_t \rangle)_{Bon}}$$

It has been studied as a function of the beam spot



Beam spot position

Computing of the corrections for Boff data:

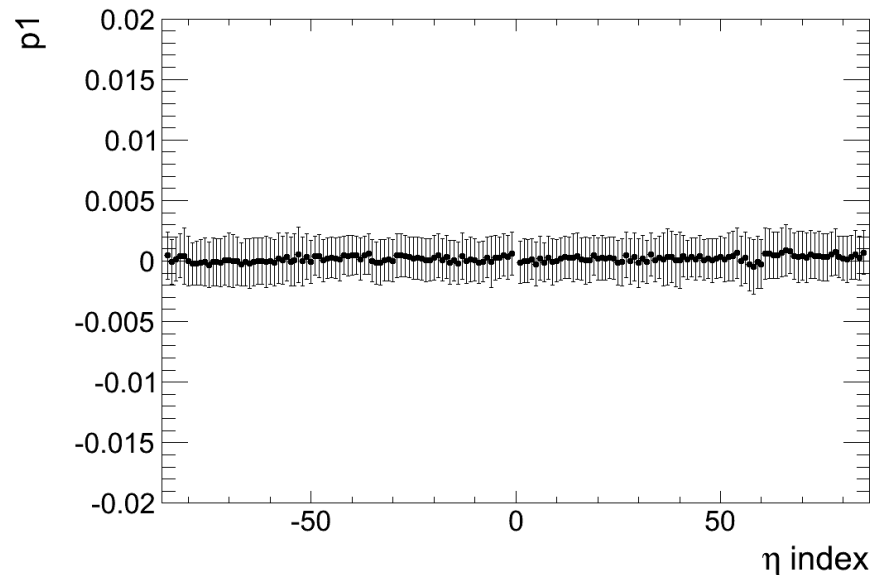
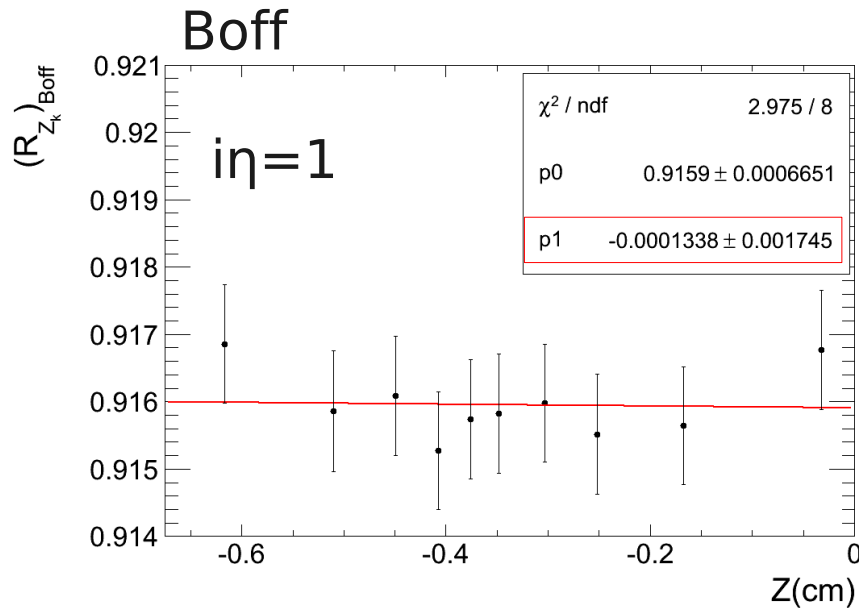
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After the correction

$$(E_t^{corr})_{Boff} = \frac{E_t}{p1 * Z + 1}$$



Beam spot position

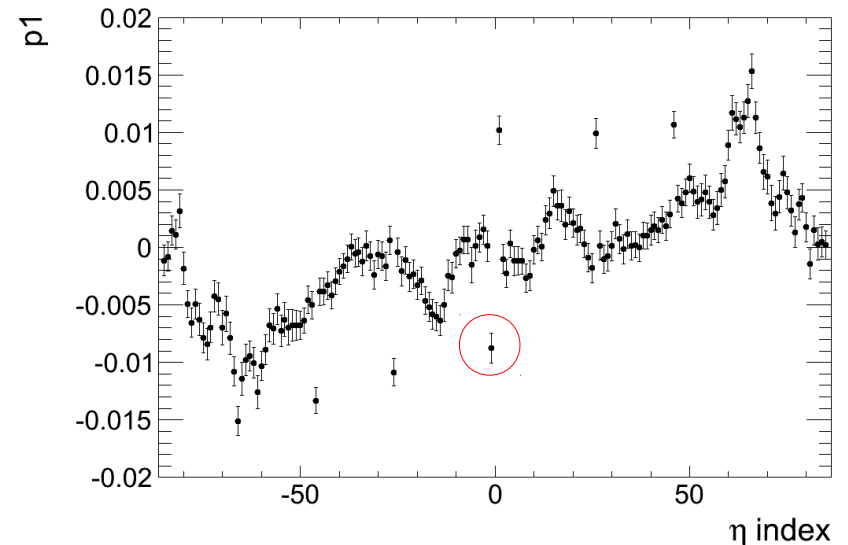
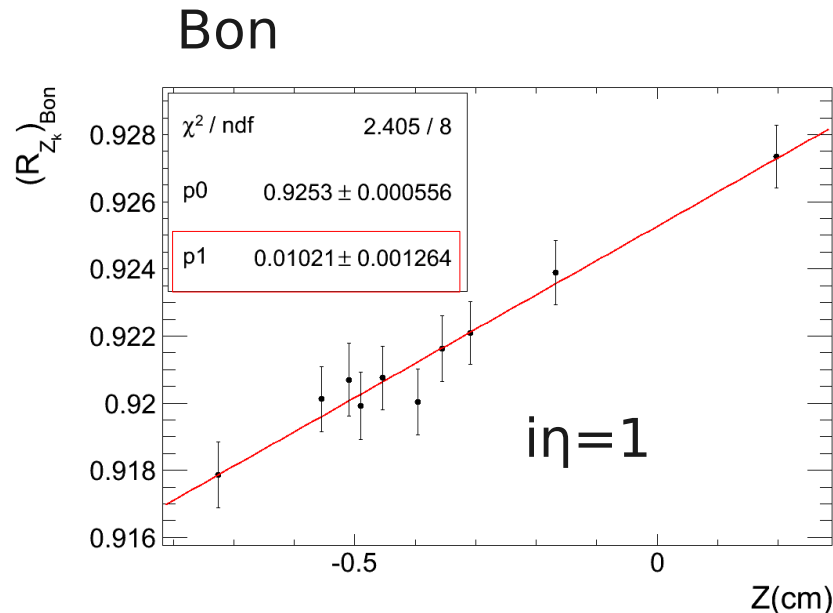
Computing of the corrections for Bon data:

The beam spot distribution has been divided into 10 intervals (Z_k):

This quantity has been defined

$$(R_{Z_k})_{Bon} = \frac{(\sum E_t^i / \langle \sum E_t \rangle)_{Boff}}{((\sum E_t^i / \langle E_t \rangle)_{Z_k})_{Bon}}$$

They have been studied as a function of the beam spot



Beam spot position

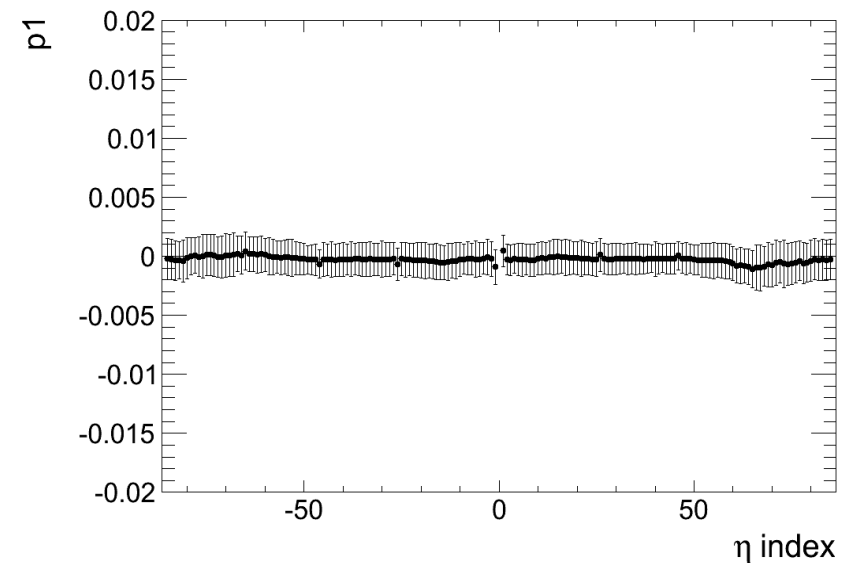
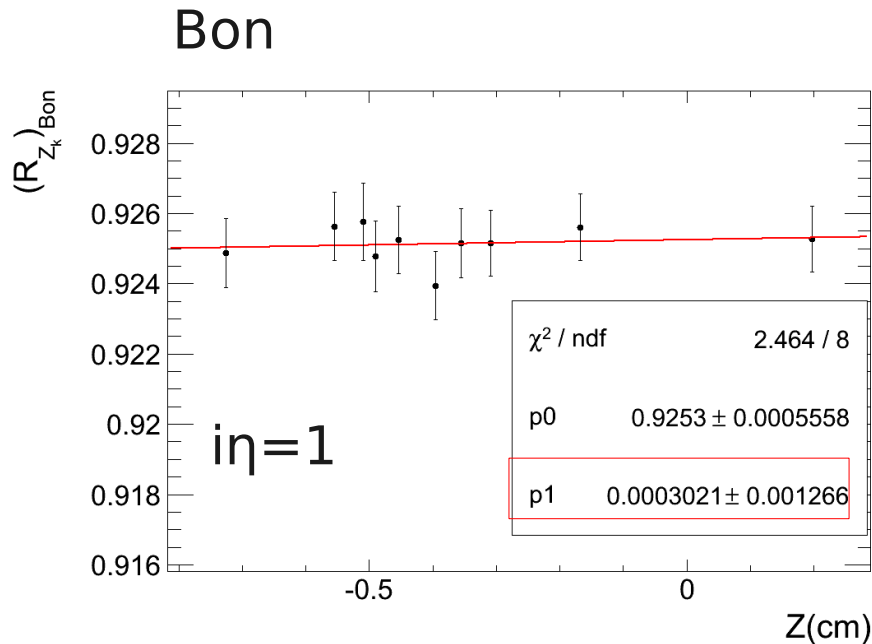
Computing of the corrections for Bon data:

The beam spot position distribution has been divided into 10 intervals (Z_k):

This quantity has been defined

$$(R_{Z_k})_{Bon} = \frac{(\sum E_t^i / \langle \sum E_t \rangle)_{Boff}}{((\sum E_t^i / \langle E_t \rangle)_{Z_k})_{Bon}}$$

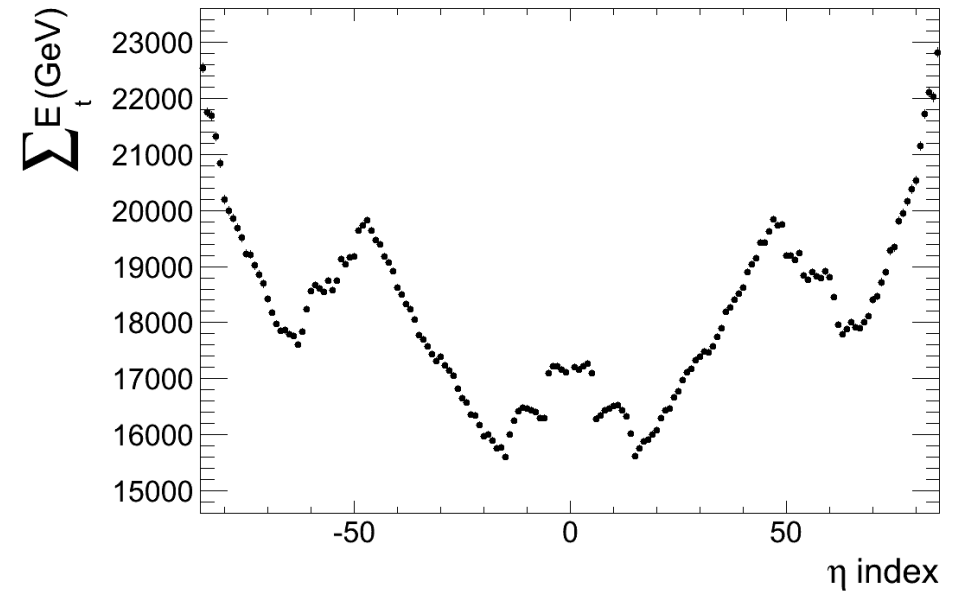
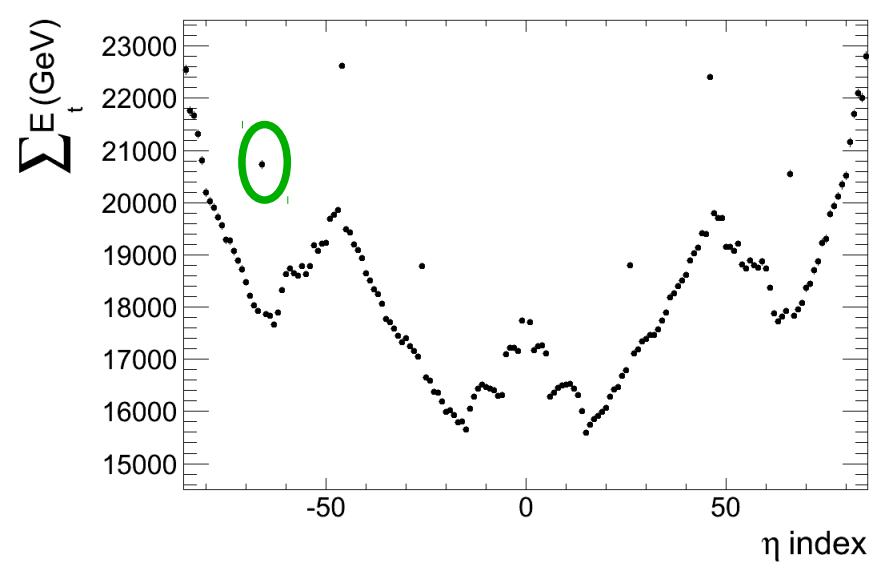
After the correction $(E_t^{corr})_{Bon} = E_t * (p1 * Z + 1)$



Border Effect

Corrections derived from MC and for η and ϕ ECAL coordinate

η border correction



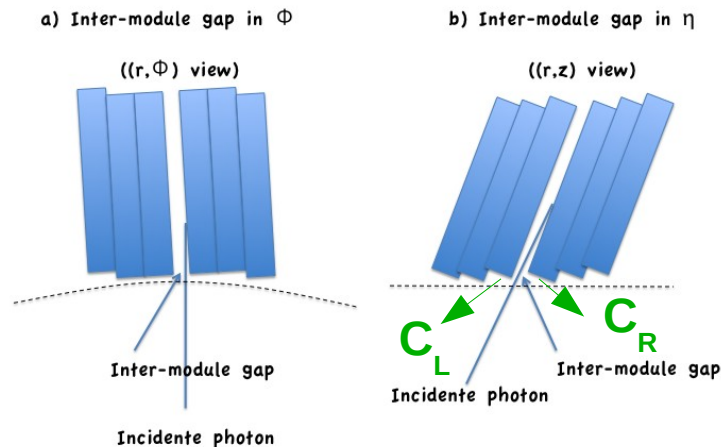
$$C_{\eta_i}^{\pm} = \frac{\langle E_t(\eta_i) \rangle - (\langle E_t(\eta_{i+1}) \rangle + \langle E_t(\eta_{i-1}) \rangle) / 2}{(\langle E_t(\eta_{i+1}) \rangle + \langle E_t(\eta_{i-1}) \rangle) / 2}$$

Total correction

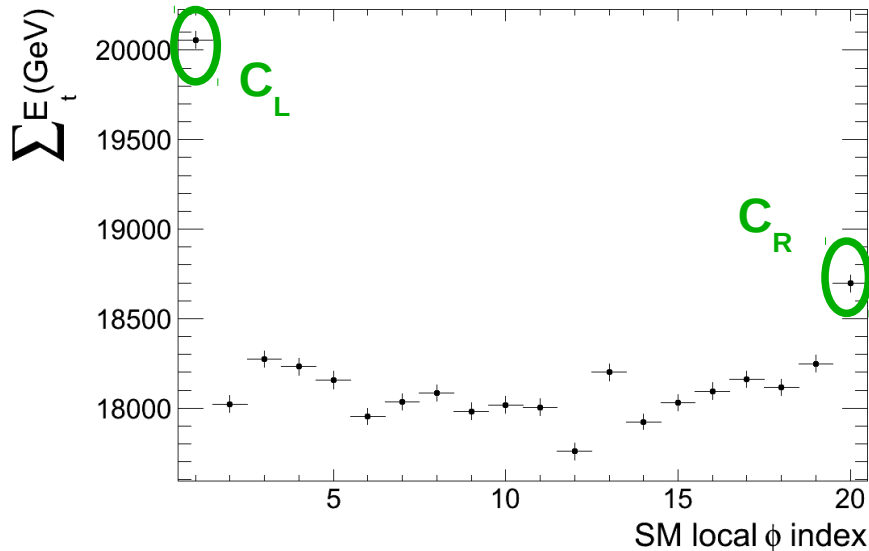
$$C_{\eta_i} = \frac{C_{\eta_i}^+ + C_{\eta_i}^-}{2}$$

Border Effect

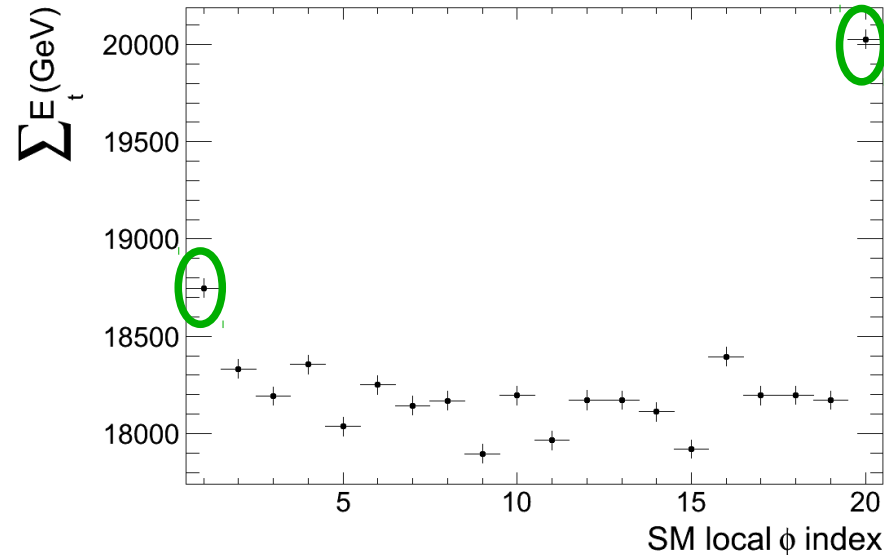
$i\phi$ border correction for magnet on MC, derived for EB+ and EB- separately due to the crystals staggering



EB -



EB +



$$(C_{\Phi}^{\pm})_{R(L)} = \frac{\langle E_{t_B} \rangle - \langle E_t \rangle}{\langle E_t \rangle}$$

Total correction

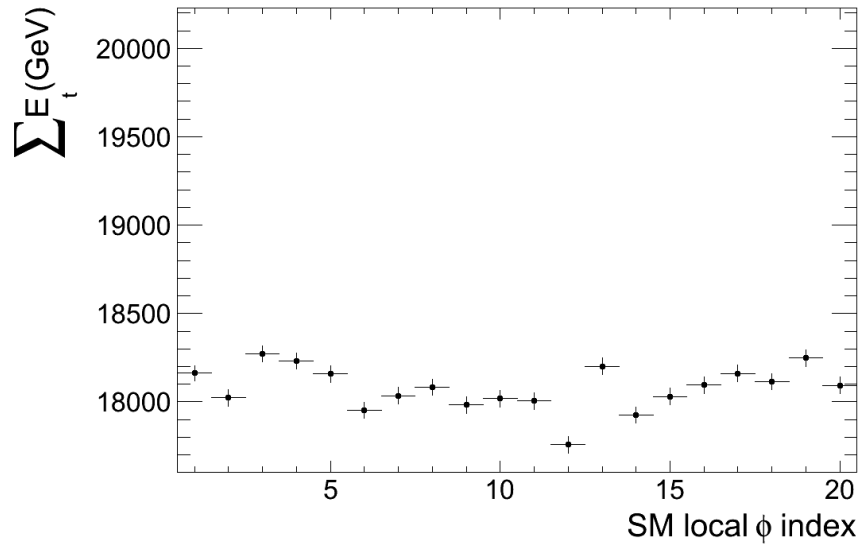
$$(C_{\Phi})_{R(L)} = \frac{(C_{\Phi}^+)_{R(L)} + (C_{\Phi}^-)_{R(L)}}{2}$$

Border Effect

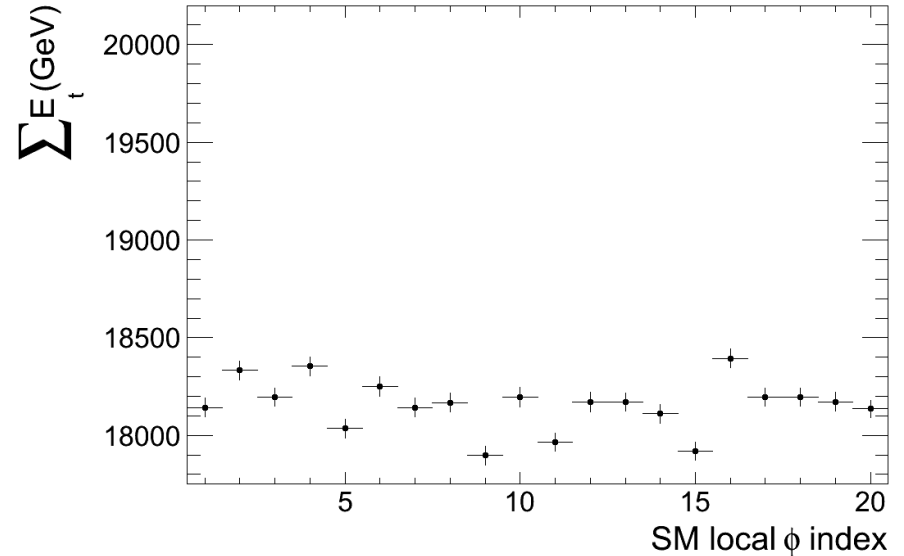
$i\phi$ border correction for magnet on MC, derived for EB+ and EB- separately due to the crystals staggering

After corrections

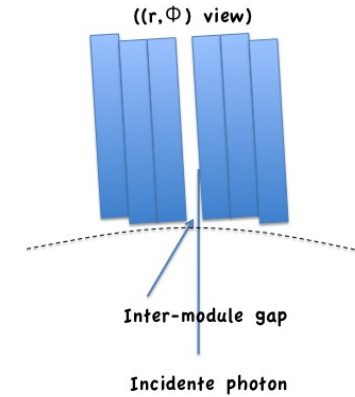
EB -



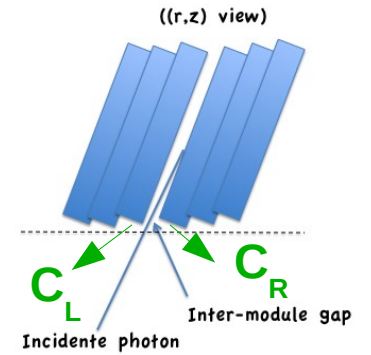
EB +



a) Inter-module gap in Φ



b) Inter-module gap in η



Border Effect

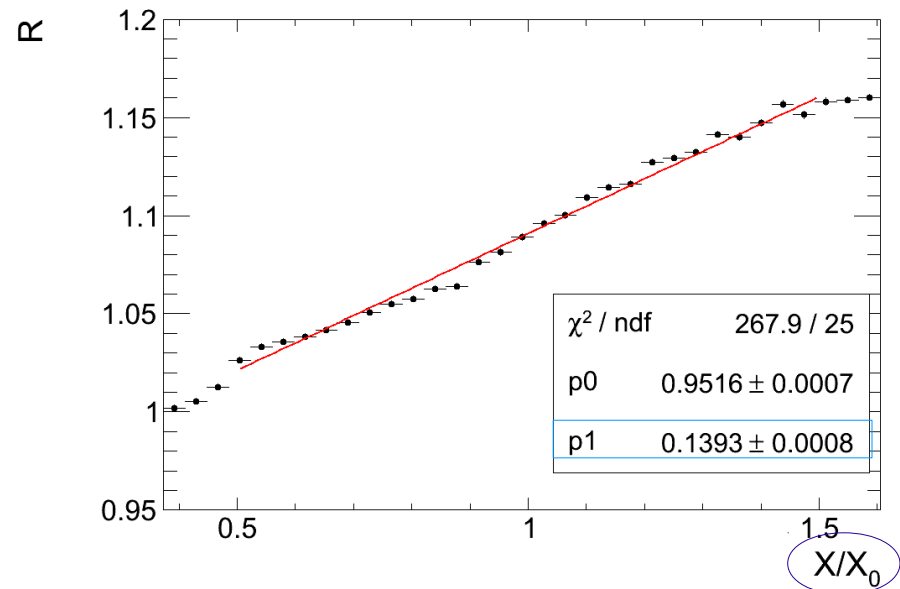
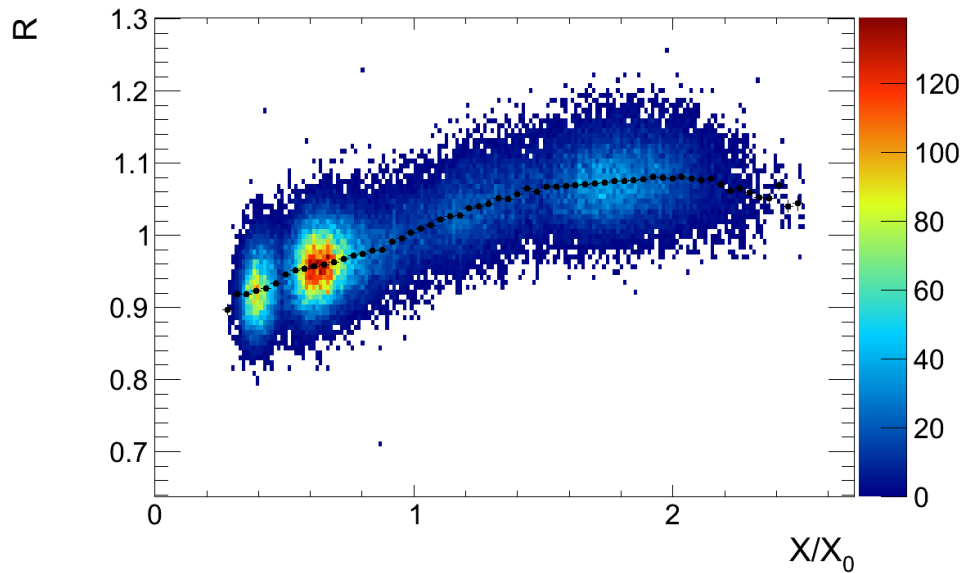
Table of border corrections for magnet off and magnet on MC

	B_{off}	B_{on}
$C_{\eta_i=1}$	(3.5 ± 0.1)%	(3.3 ± 0.2)%
$C_{\eta_i=26}$	(11.0 ± 0.3)%	(11.2 ± 0.3)%
$C_{\eta_i=46}$	(13.6 ± 0.3)%	(14.6 ± 0.3)%
$C_{\eta_i=66}$	(14.0 ± 0.4)%	(15.4 ± 0.4)%
$(C_{\phi}^+)_{R}$	(9.8 ± 0.8)%	(10.1 ± 0.7)%
$(C_{\phi}^-)_{R}$	(10.3 ± 0.7)%	(10.8 ± 0.8)%
$(C_{\phi}^+)_{L}$	-	(3.2 ± 1.0)%
$(C_{\phi}^-)_{L}$	-	(3.4 ± 1.0)%
$(C_{\phi})_{R}$	(10.0 ± 1.1)%	(10.4 ± 1.1)%
$(C_{\phi})_{L}$	-	(3.3 ± 1.0)%

Boff and Bon correction are not the same so they have been derived separately


Correlation between R and tracker material

Correlation between the energy flow ratio and the tracker material (X/X_0)



$$\Delta M(X_0) = \frac{1}{p1} \left(\frac{R_{\text{dati}}}{R_{\text{MC}}} - 1 \right)$$

$$M(X_0) = M(X_0)^{\text{std}} + \Delta M(X_0)$$


Measure of the material made during the detector construction