

keV Sterile Neutrino Dark Matter from Singlet Scalar Decays

IMPRS-EPP student seminar

based on 1502.01011 in collaboration with Alexander Merle

Maximilian Totzauer

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- 1 What is a sterile neutrino in the first place?
- 2 Why spend time on thinking about sterile neutrinos?
- 3 keV sterile neutrinos as Dark Matter
 - How to constrain Dark Matter models in general
 - Production mechanisms for keV steriles as Dark Matter
- 4 Production from scalar decays – The details
 - The model
 - What about the constraints?
- 5 Conclusion & Outlook

└ What is a sterile neutrino in the first place?

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	I	II	III	
mass	2.4 MeV	1.27 GeV	171.2 GeV	0
charge	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0
name	u	c	t	g
	up	charm	top	gluon
Quarks	d	s	b	γ
	down	strange	bottom	photon
Leptons	ν_e	ν_μ	ν_τ	Z
	electron neutrino	muon neutrino	tau neutrino	weak force
	e	μ	τ	H
	electron	muon	tau	Higgs boson
	0.511 MeV	105.7 MeV	1.777 GeV	91.2 GeV
	-1	-1	-1	0
	Left	Left	Left	spin 1
	Right	Right	Right	spin 0
				95.4 GeV
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	0 ν_e N_1 electron (neutrino) sterile neutrino	0 ν_μ N_2 muon (neutrino) sterile neutrino	0 ν_τ N_3 tau (neutrino) sterile neutrino	91.2 GeV 0	>114 GeV
				Z weak force	H Higgs boson
				spin 1	spin 0
Leptons	0.511 MeV -1	105.7 MeV -1	1.777 GeV -1	80.4 GeV -1	
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- **Mixing** of heavy mass eigenstates into the eigenstates of the weak interaction (**mixing angles**, usually denoted θ).

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- Neutrinos in the SM are **massless**, but **neutrino oscillations** suggest something different! \Rightarrow (Heavy) sterile neutrinos naturally allow for small masses of the active neutrinos (**seesaw mechanism**).

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- **Leptogenesis**
- A massive, electrically neutral particle is a natural **candidate for Dark Matter**.
- Recent excitement about **tentative signal** ($N \rightarrow \nu\gamma$) at $E_\gamma = 3.55 \text{ keV} \Rightarrow M_R = 7.1 \text{ keV} ?$

Which experimental insights constrain DM models?

- **Relic abundance** constraint: ($\Omega_{\text{DM}} h^2 = 0.1188$).
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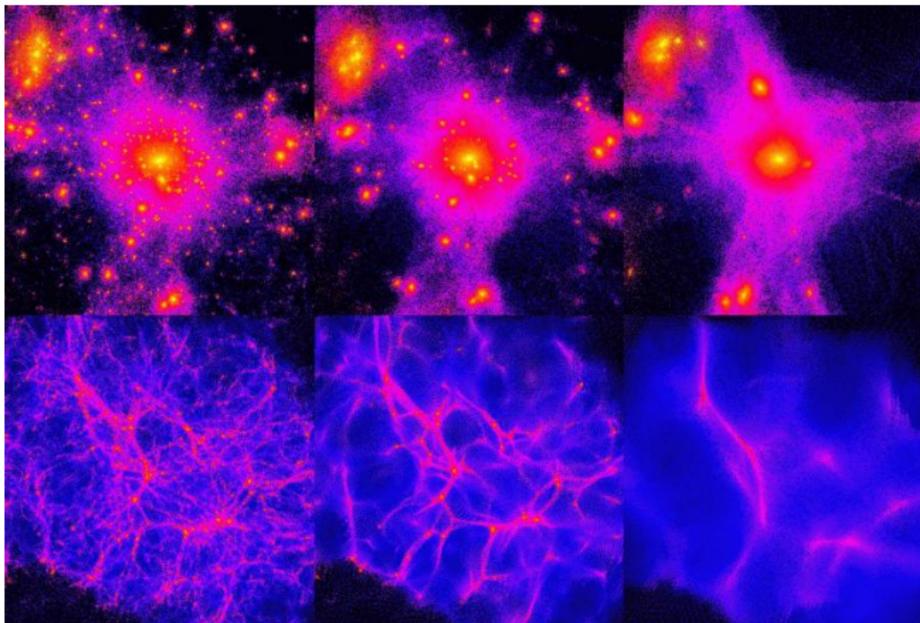
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- **Constraints on N_{eff}** : Additional degrees of freedom in radiation heavily constrained by the CMB and BBN.
- **Structure formation**: Cold Dark Matter fits the observed large-scale structure almost perfectly (e.g. **missing satellite problem**). If Dark Matter was too light (and thereby **most likely** too fast), there would be less structure today!

Hot, warm or cold Dark Matter? – Structure formation



Cold (left), warm (centre) and hot (right) Dark Matter

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Alternative: **freeze-in**: Interaction (e.g. $\chi\chi \leftrightarrow X_{\text{SM}}X_{\text{SM}}$) too feeble to obtain equilibrium. Occasional production of DM from plasma until $T \leq m_{\text{DM}}$.

- Relic abundance $\propto \sigma_{\text{ann}}$
- **Non-thermal spectrum**

Possible production mechanisms for keV steriles as Dark Matter

- **Dodelson-Widrow-mechanism** (hep-ph/9303287) Production from the plasma via active-sterile-mixing; equilibrium never achieved; (**freeze-in type**)
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⇒ **numerical treatment** on the level of **distribution functions!**

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where

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- Mixing θ switched off in this model.

Sterile neutrinos from scalar decays – particle physics

The relevant parameters of the setup:

- Yukawa coupling y ($-\mathcal{L} \supset \frac{y}{2} S \bar{N}^c N$).
- Higgs portal λ ($-\mathcal{L} \supset 2\lambda (H^\dagger H) S^2$)
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Convenient parametrisation for our analysis:

- $\mathcal{C}_\Gamma \equiv \frac{M_0}{m_S} \frac{y^2}{16\pi} = \frac{M_0}{m_S} \frac{\Gamma}{m_S}$ (effective decay width)
- $\mathcal{C}_{\text{HP}} \equiv \frac{M_0}{m_S} \frac{\lambda^2}{16\pi^3}$ (effective Higgs portal)
- with $M_0 \equiv \left(\frac{45 M_{\text{Pl}}^2}{4\pi^3 g_*} \right)^{1/2} = 7.35 g_*^{-1/2} \times 10^{18} \text{ GeV}$

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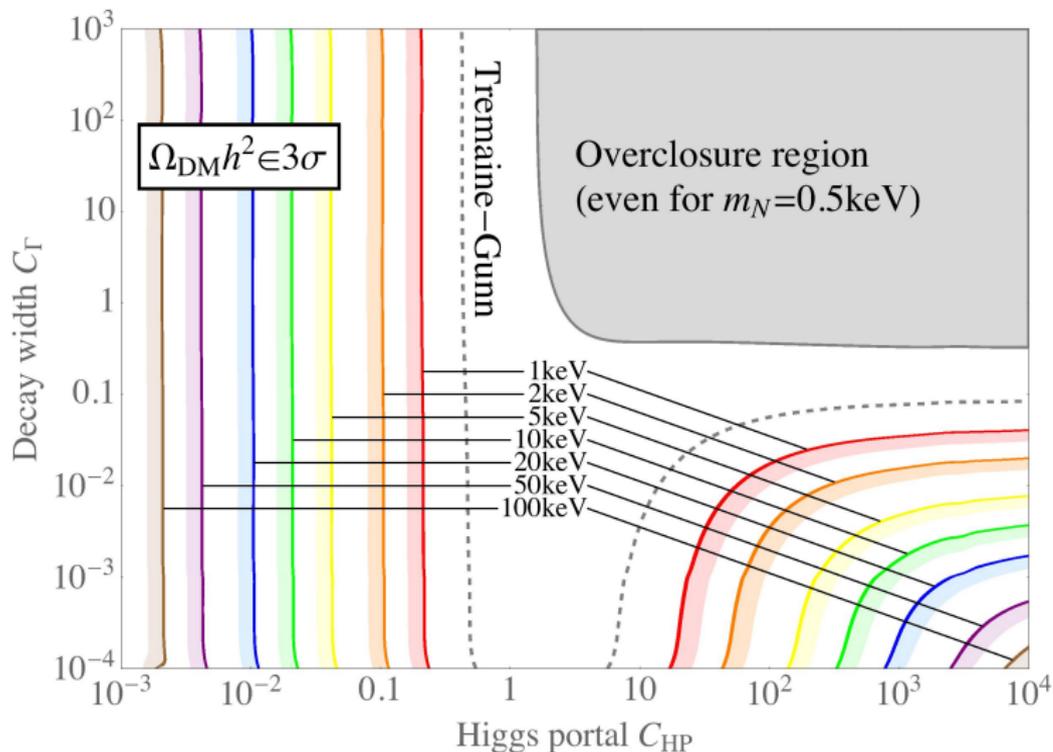
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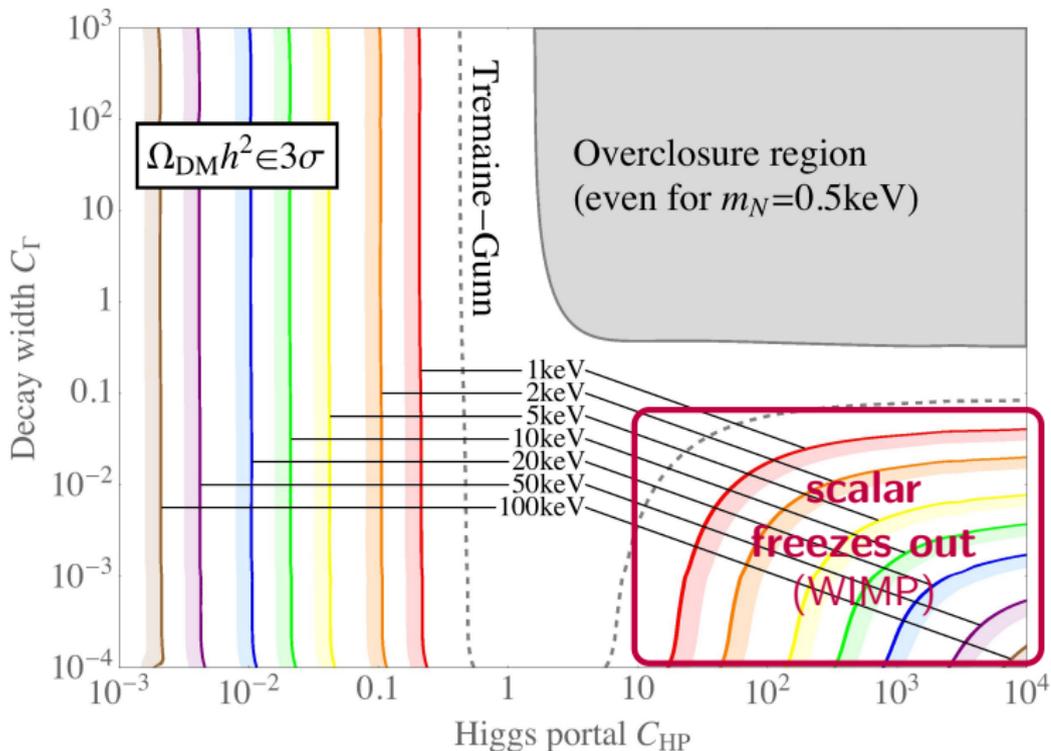
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⇒ Numerically difficult partial integro-differential equations.

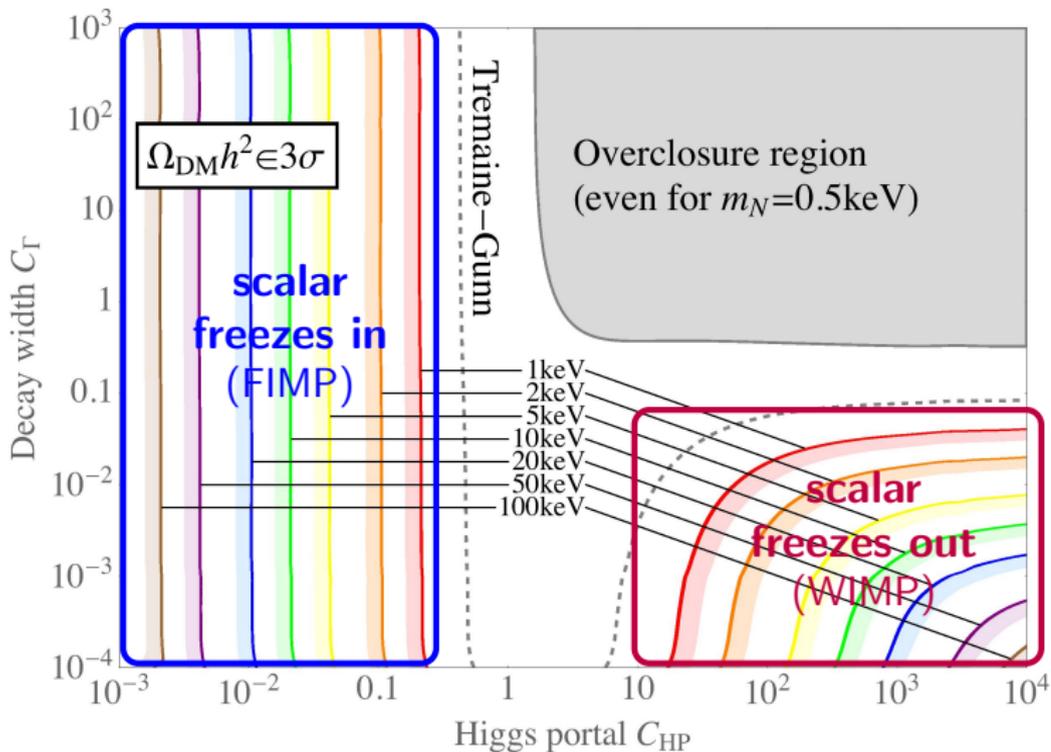
Confronting the model with reality – relic abundance



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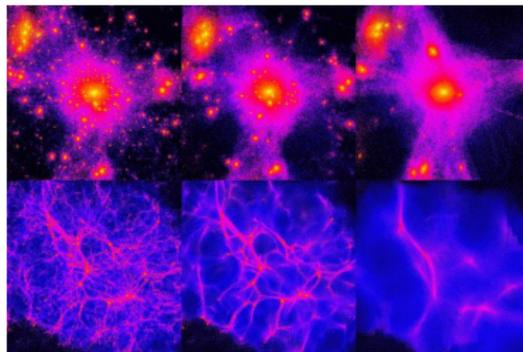


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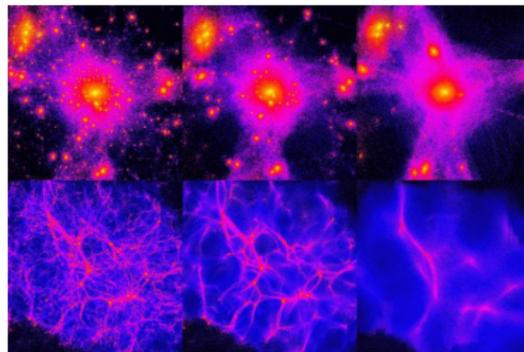
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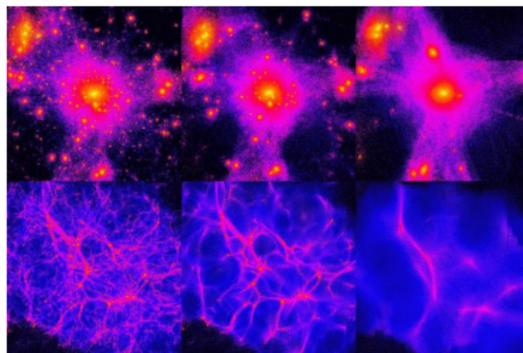
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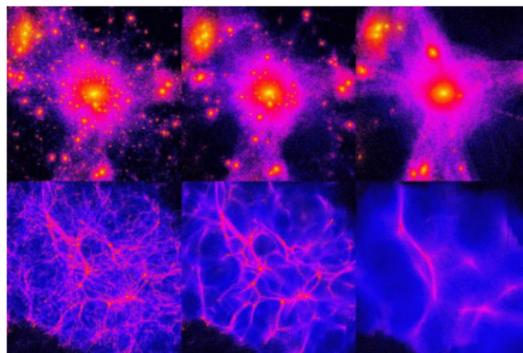
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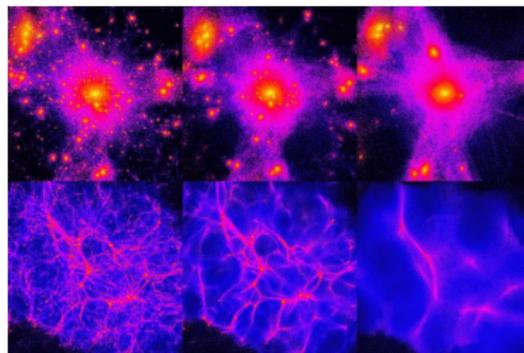
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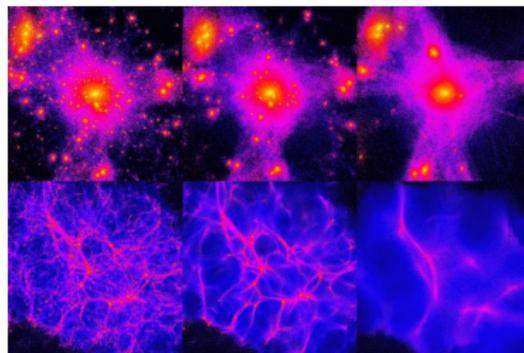
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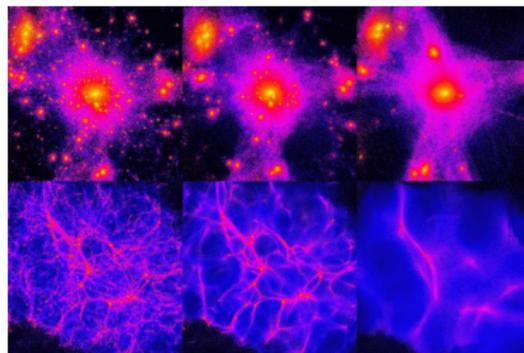
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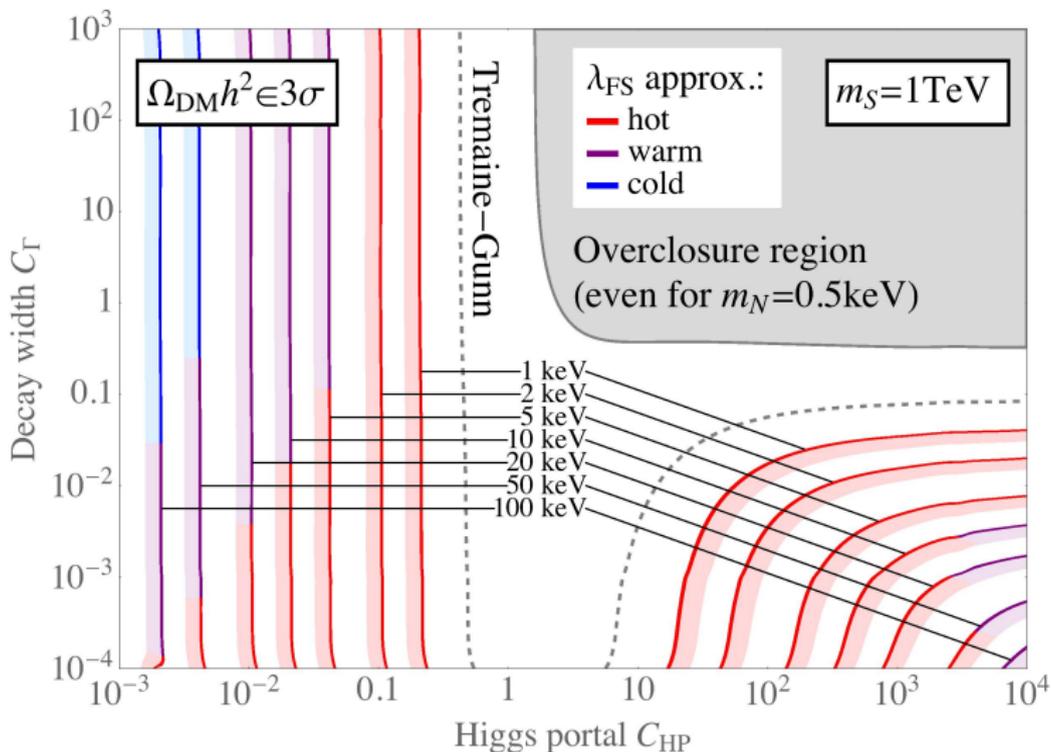
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- Substantial uncertainties in time-temperature relation dt/dT .

\Rightarrow Use λ_{FS} for the time being, more detailed analyses ongoing.

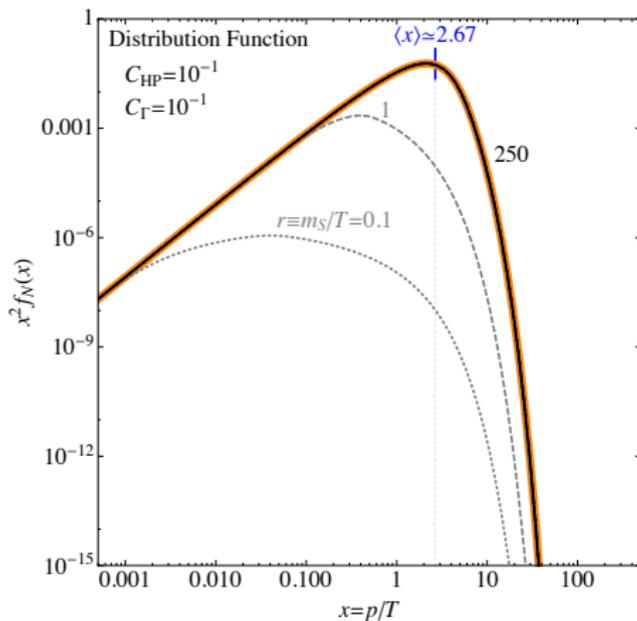


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Why might λ_{FS} be particularly bad in our model?

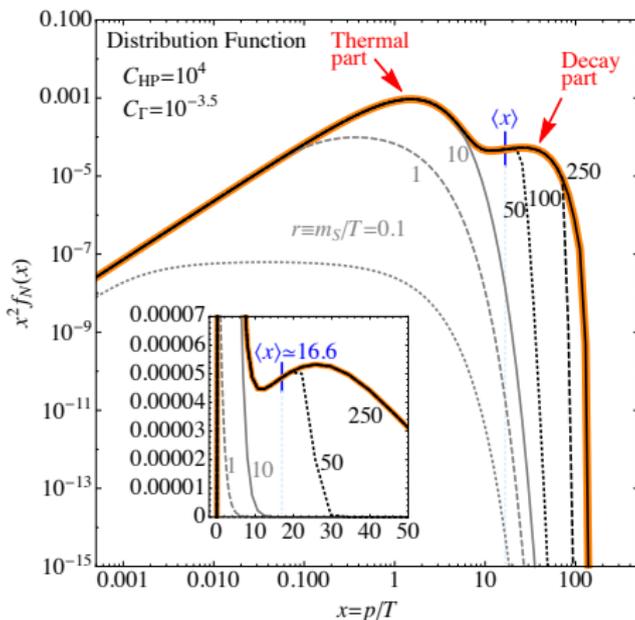
FIMP case: SN spectrum cooler than thermal one ($\langle x \rangle_{\text{th}} = 3.15$)



$\Rightarrow \lambda_{\text{FS}}$ probably not so bad...

Why might λ_{FS} be particularly bad in our model?

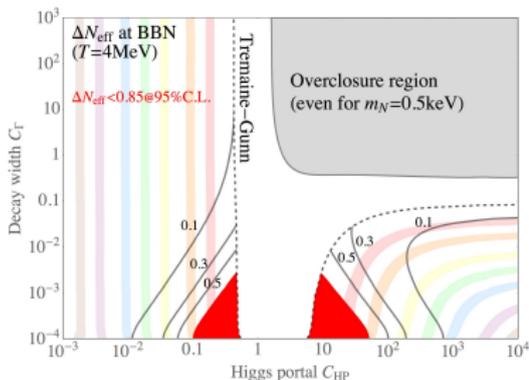
WIMP case: SN from scalars in equilibrium and frozen-out scalars!



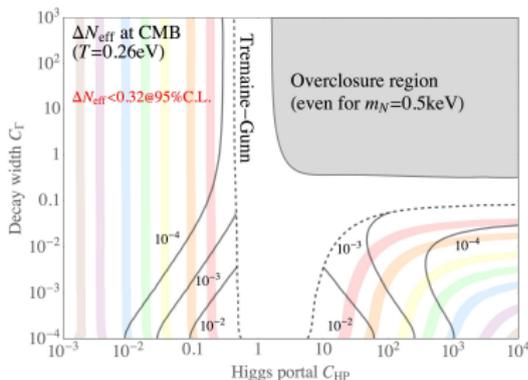
$\Rightarrow \lambda_{\text{FS}}$ probably not so good...

Confronting the model with reality – CMB and BBN

Extra radiation and the time of CMB decoupling / BBN would leave imprints in power spectrum / primordial He-abundance:



(a) Constraints from Big Bang Nucleosynthesis (BBN)
(red patches excluded)



(b) Constraints from the Cosmic Microwave Background (CMB)
(nothing excluded)

Conclusion

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- The framework offers **many viable scenarios for the scalar to be produced** (FIMP / WIMP).
- Implications for structure formation need more scrutiny to take into account the full information of the spectrum (\rightarrow **small scale problems**).

Outlook

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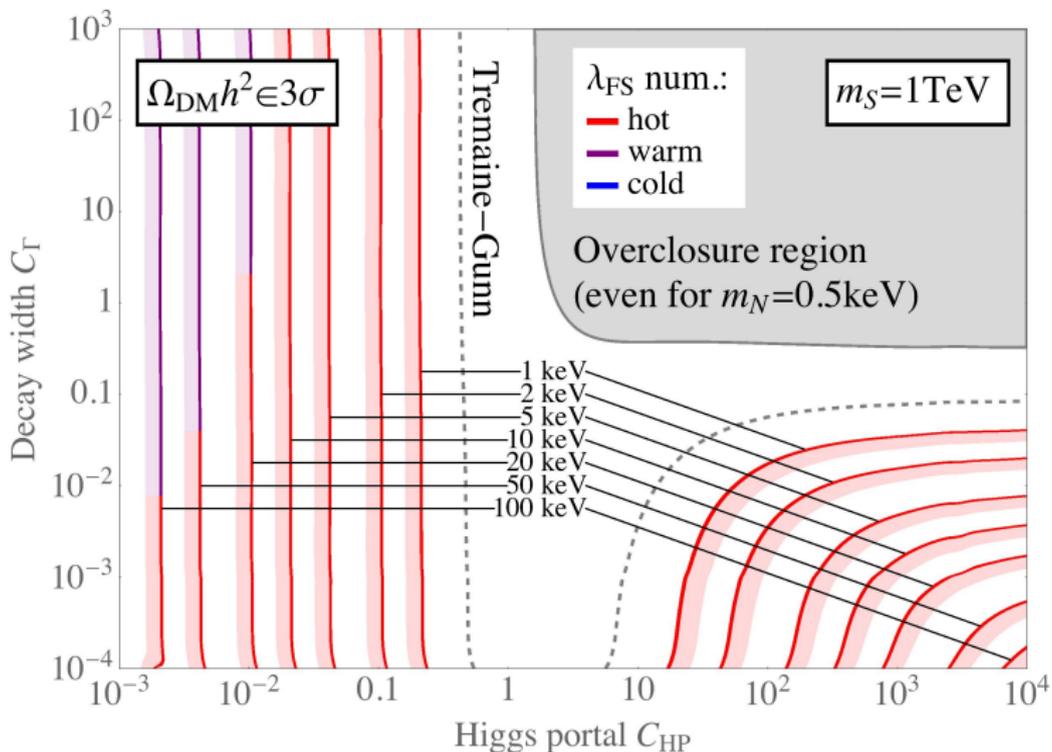
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Outlook

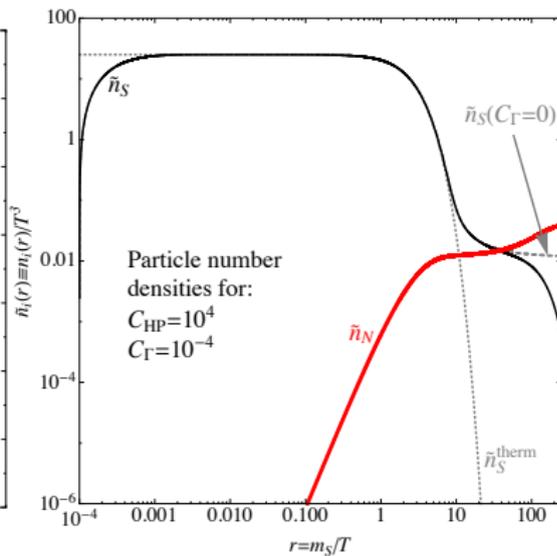
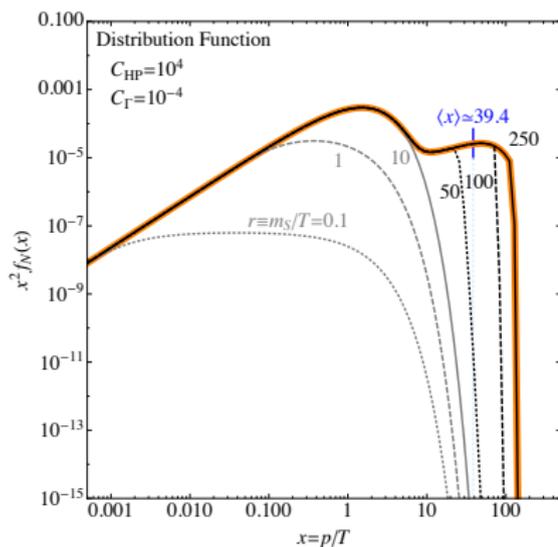
- Fully numerical implementation of the **evolution of the d.o.f.(time-temperature relation)** (work in progress).
- Extend analysis to **scalar masses below the electroweak scale**, which opens up many more interaction processes.
- Analysis of **transfer functions** of structure formation using **full spectral information** (work in progress).
- **Include DW contribution** (non-zero mixing θ , work in progress).

Thank you for your attention!

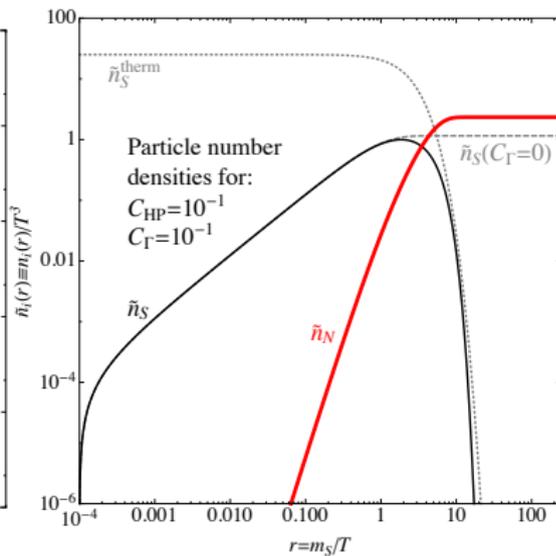
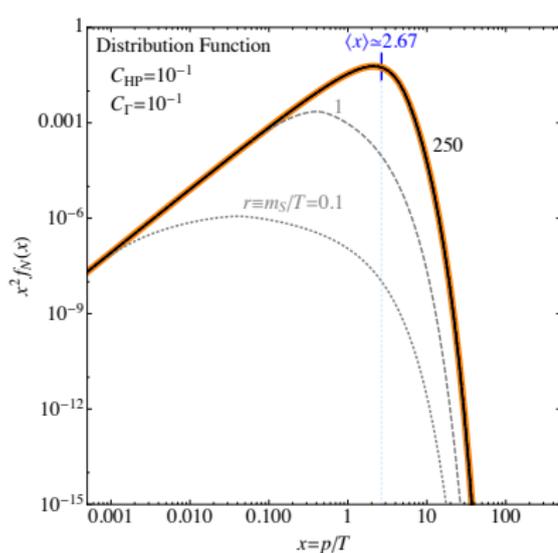
Backup I – Numerical free-streaming horizon



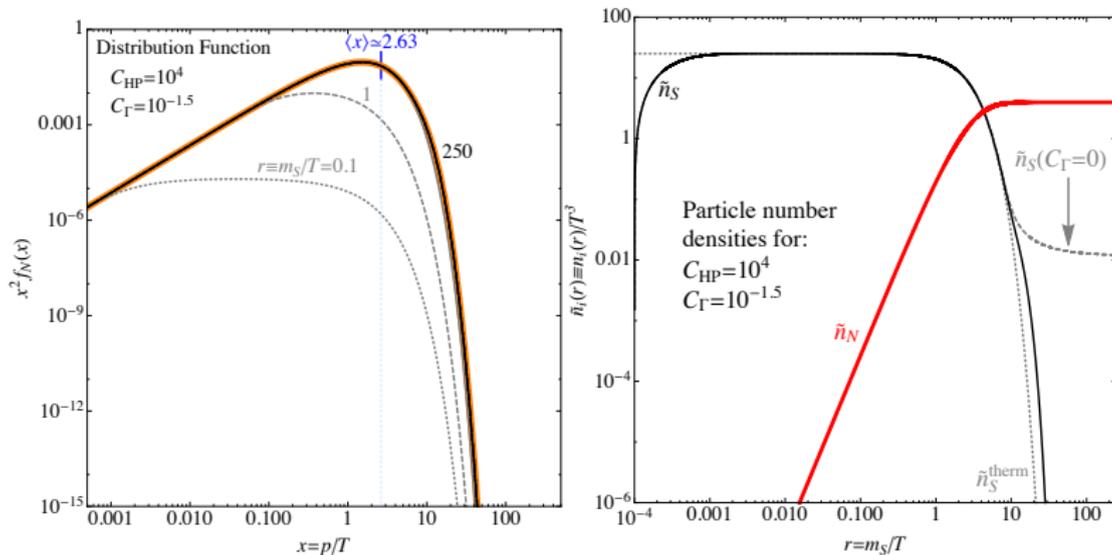
Backup II – Evolution of abundances vs. evolution of distribution function



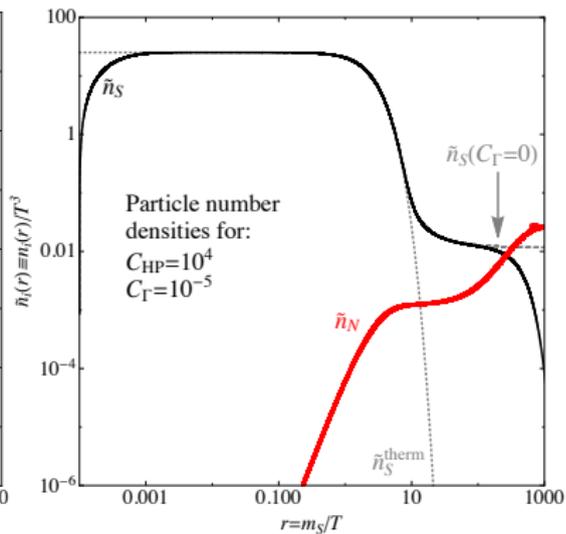
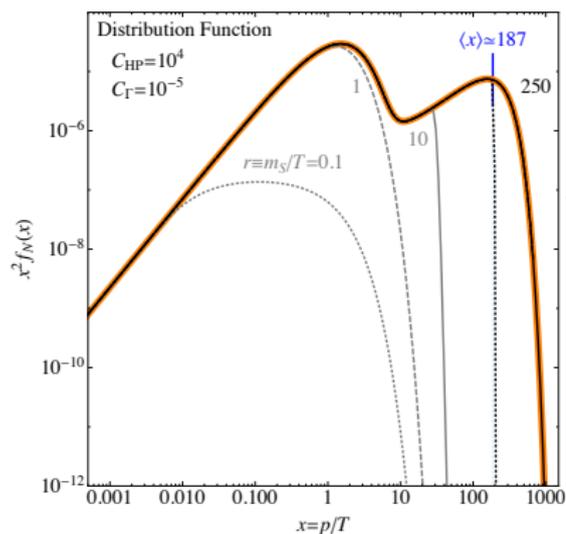
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Backup III – Assessing the approximation of $g_* = \text{const.}$ 