Beyond-the-standard-model contributions to rare B decays analyzed with variational-Bayes enhanced adaptive importance sampling

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Bayes' formula:

$$P(\boldsymbol{\theta}|\mathcal{D}, \mathbf{M}) = \frac{P(\mathcal{D}|\boldsymbol{\theta}, \mathbf{M})P(\boldsymbol{\theta}|\mathbf{M})}{P(\mathcal{D}|\mathbf{M})} = \frac{P(\mathcal{D}|\boldsymbol{\theta}, \mathbf{M})P(\boldsymbol{\theta}|\mathbf{M})}{\int P(\mathcal{D}|\boldsymbol{\theta}, \mathbf{M})P(\boldsymbol{\theta}|\mathbf{M})d\boldsymbol{\theta}}$$

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model independent search for new physics (effective theory):

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$$\boldsymbol{\theta} = \text{effective couplings } \mathcal{C}_i, \dots$$



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## model independent search for new physics (effective theory):



 $\boldsymbol{\theta} = \text{effective couplings } \mathcal{C}_i, \dots$  $\mathcal{D} = \text{detector events}$  $M = EFT, SM, \dots$ 



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no standard algorithm so far



- Adaptive importance sampling with the variational-Bayes approach
- 3 Model independent search for new physics



# Adaptive importance sampling with the variational-Bayes approach

## 1 Overview

- 2 Adaptive importance sampling with the variational-Bayes approach
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## 4 Summary

## Adaptive importance sampling

$$\int P(x)dx = \int \frac{P(x)}{q(x)}q(x)dx \approx \frac{1}{N}\sum_{n=1}^{N}\frac{P(x_n)}{q(x_n)} \equiv \hat{\mu}^N \text{ where } x_n \sim q$$

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squared uncertainty (variance):

$$var(\hat{\mu}^N) = \frac{1}{N} \left[ \int \frac{P(x)}{q(x)} P(x) dx - \left( \int P(x) dx \right)^2 \right]$$

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minimize  $var(\hat{\mu}^N)$  with respect to the *proposal q* 

# Adaptive importance sampling with the variational-Bayes approach



### pypmc

#### '''This example illustrates how to run a Markov Chain using pypmc'''

import numpy as np
import pypmc

```
# define a proposal
```

```
prop_dof = 50.
prop_sigma = np.array([[0.1 , 0. ]
,[0. , 0.02]])
prop = pvpmc.densitv.student t.LocalStud
```

# define the target; i.e., the function
# In this case, it is a Gaussian with me to
# covariance "target\_sigma".
#

def unnormalized\_log\_pdf\_gauss(x, mu, inv\_sigma): diff = x - mu return -0.5 \* diff.dot(inv sigma).dot(diff)

```
log target = lambda x: unnormalized log pdf gauss(x, target mean, i
```

#### # choose a bad initialization start = np.array([-2., 10.])

#### # define the markov chain object

mc = pypmc.sampler.markov chain.AdaptiveMarkovChain(log target, prc

# run burn-in
mc.run(10\*\*4)

# delete burn-in from history
mc.history.clear()



(effective samples).

#### rel\_tol -

Relative tolerance  $\epsilon$ . If two consecutive values of the log likelihood bound,  $L_t, L_{t-1}$ , are close, declare convergence. More precisely, check that

$$\left\|\frac{L_t-L_{t-1}}{L_t}\right\|<\epsilon.$$

#### abs\_tol -

Absolute tolerance  $\epsilon_a.$  If the current bound  $L_t$  is close to zero,  $(L_t < \epsilon_a),$  declare convergence if

$$||L_t - L_{t-1}|| < \epsilon_a$$

verbose - Output status information after each update.

#### set\_variational\_parameters()

Reset the parameters to the submitted values or default.

Use this function to set the prior value (indicated by the subscript  $\theta$  as in  $\alpha_0$ ) or the initial value (e.g.,  $\alpha$ ) used in the iterative procedure to find the posterior value of the variational distribution.

Every parameter can be set in two ways:

1. It is specified for only one component, then it is copied to all other components.

2. It is specified separately for each component as a K vector.

The prior and posterior variational distributions of  $\mu$  and  $\Lambda$  for each component are given by

$$q(\boldsymbol{\mu}, \boldsymbol{\Lambda}) = q(\boldsymbol{\mu}|\boldsymbol{\Lambda})q(\boldsymbol{\Lambda}) = \prod_{k=1}^{K} \mathcal{N}(\boldsymbol{\mu}_{k}|\boldsymbol{m}_{k}, (\beta_{k}\boldsymbol{\Lambda}_{k})^{-1})\mathcal{W}(\boldsymbol{\Lambda}_{k}|\boldsymbol{W}_{k}, \nu_{k}),$$

where  ${\cal N}$  denotes a Gaussian and  ${\cal W}$  a Wishart distribution. The weights  $\pi$  follow a Dirichlet distribution

$$q(\pi) = Dir(\pi | \alpha).$$

This function may delete results obtained by update().

### https://pypi.python.org/pypi/pypmc

Warning

## Model independent search for new physics

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## 4 Summary

## The standard model (SM) of particle physics cannot explain:

- dark matter
- neutrino masses
- hierarchy problem
- strong CP problem
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- dark matter
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## new physics (NP) required exact structure unknown $\Rightarrow$ model independent analysis

## Effective theory

effective Lagrangian for  $b \to s \ell^+ \ell^-$  (SM):

$$\mathcal{L}_{int} = \frac{4G_F}{\sqrt{2}} \frac{\alpha_e}{4\pi} V_{tb} V_{ts}^* \sum_i C_i \mathcal{O}_i + ... + \text{h.c.}$$

 $\mathcal{O}_{9} = \begin{bmatrix} \bar{s} \gamma_{\mu} P_{L} & b \end{bmatrix} \begin{bmatrix} \bar{\ell} \gamma^{\mu} \ell \end{bmatrix} \qquad \qquad \mathcal{O}_{10} = \begin{bmatrix} \bar{s} \gamma_{\mu} P_{L} & b \end{bmatrix} \begin{bmatrix} \bar{\ell} \gamma^{\mu} \gamma_{5} \ell \end{bmatrix}$ 



## Effective theory

effective Lagrangian for  $b \rightarrow s \ell^+ \ell^-$  (beyond-SM):

$$\mathcal{L}_{int} = \frac{4G_F}{\sqrt{2}} \frac{\alpha_e}{4\pi} V_{tb} V_{ts}^* \sum_i C_i \mathcal{O}_i + \dots + \text{h.c.}$$

 $\mathcal{O}_{9}^{(\prime)} = \left[\bar{s}\gamma_{\mu}P_{L(R)}b\right]\left[\bar{\ell}\gamma^{\mu}\ell\right] \qquad \qquad \mathcal{O}_{10}^{(\prime)} = \left[\bar{s}\gamma_{\mu}P_{L(R)}b\right]\left[\bar{\ell}\gamma^{\mu}\gamma_{5}\ell\right]$ 

 $\mathcal{O}_{S}^{(\prime)} = \begin{bmatrix} \bar{s} P_{R(L)} b \end{bmatrix} \begin{bmatrix} \bar{\ell} \ell \end{bmatrix} \qquad \mathcal{O}_{P}^{(\prime)} = \begin{bmatrix} \bar{s} P_{R(L)} b \end{bmatrix} \begin{bmatrix} \bar{\ell} \gamma_{5} \ell \end{bmatrix}$  $\mathcal{O}_{T} = \begin{bmatrix} \bar{s} \sigma_{\mu\nu} b \end{bmatrix} \begin{bmatrix} \bar{\ell} \sigma^{\mu\nu} \ell \end{bmatrix} \qquad \mathcal{O}_{T5} = \begin{bmatrix} \bar{s} \sigma_{\mu\nu} b \end{bmatrix} \begin{bmatrix} \bar{\ell} \sigma^{\mu\nu} \gamma_{5} \ell \end{bmatrix}$ 

## Experimental constraints

•  $B \to K\mu^+\mu^-$ :  $\mathcal{B}, A_{FB}, F_H$ 

- LHCb 2014 (arXiv:1403.8044, arXiv:1403.8045)
- CDF 2012

(http://www-cdf.fnal.gov/physics/new/bottom/120628.blessed-b2smumu\_96)

• 
$$B_s \rightarrow \mu^+ \mu^-$$
:  $\mathcal{B}$ 

• LHCb+CMS 2014 (arXiv:1411.4413)

• 
$$B \to K^* \mu^+ \mu^-$$
:  $\mathcal{B}$ 

- LHCb 2013 (arXiv:1304.6325)
- CMS 2013 (arXiv:1308.3409)
- CDF 2012

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## Joint fit of $\mathcal{C}_{10}^{(\prime)}, \mathcal{C}_{S}^{(\prime)}, \mathcal{C}_{P}^{(\prime)}, \mathcal{C}_{T}, \mathcal{C}_{T5}$ , and 29 nuisance parameters



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## • first *simultaneous* fit

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- first *simultaneous* fit
- interference  $\mathcal{C}_{10}^{(\prime)} \leftrightarrow \mathcal{C}_{S,P}^{(\prime)}$  in  $\mathcal{B}(B_s \to \mu^+ \mu^-)$

## Joint fit of $\mathcal{C}_{10}^{(\prime)}, \mathcal{C}_S^{(\prime)}, \mathcal{C}_P^{(\prime)}, \mathcal{C}_T, \mathcal{C}_{T5}$ , and 29 nuisance parameters



- first *simultaneous* fit
- interference  $\mathcal{C}_{10}^{(\prime)} \leftrightarrow \mathcal{C}_{S,P}^{(\prime)} \text{ in }$  $\mathcal{B}(B_s \to \mu^+ \mu^-)$  $\Rightarrow$  larger uncertainty than obtained for fixed  $\mathcal{C}_{10}^{(\prime)} = \mathcal{C}_{10}^{(\prime)\mathrm{SM}}$ arXiv:1205.5811. arXiv:1206.0273. arXiv:1407.7044

## Summary



## model-independent search for new physics



## model-independent search for new physics

• simultaneous fit of  $C_{10}^{(\prime)}, C_{S}^{(\prime)}, C_{P}^{(\prime)}, C_{T}, \text{ and } C_{T5}$  $\Rightarrow$  updated constraints





## model-independent search for new physics

• simultaneous fit of  $C_{10}^{(\prime)}, C_5^{(\prime)}, C_P^{(\prime)}, C_T$ , and  $C_{T5}$  $\Rightarrow$  updated constraints

• no significant deviation from the SM

## Nuisance parameters



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