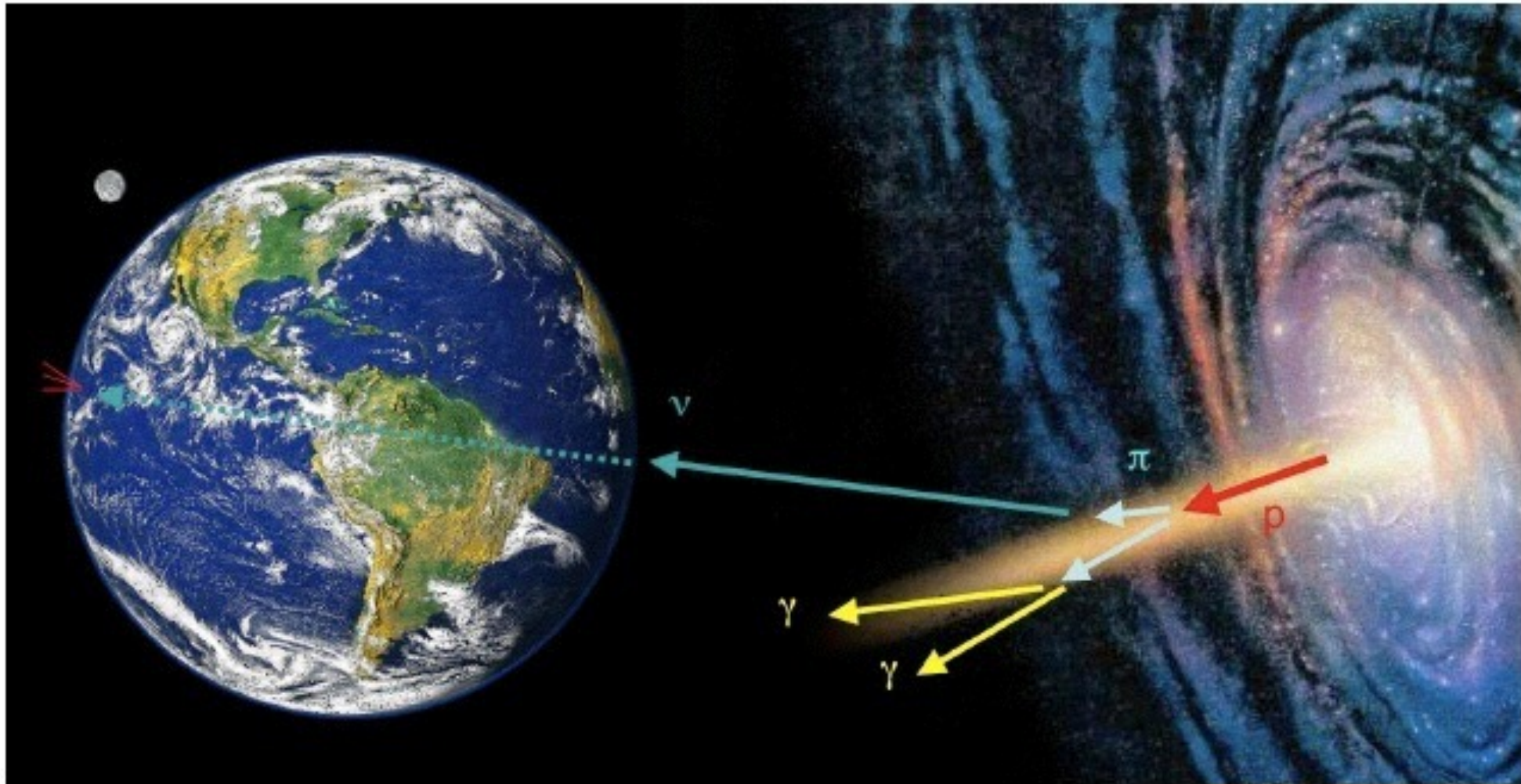


Teilchenphysik mit kosmischen und mit erdgebundenen Beschleunigern



06. Cosmic Accelerators

01.06.2015



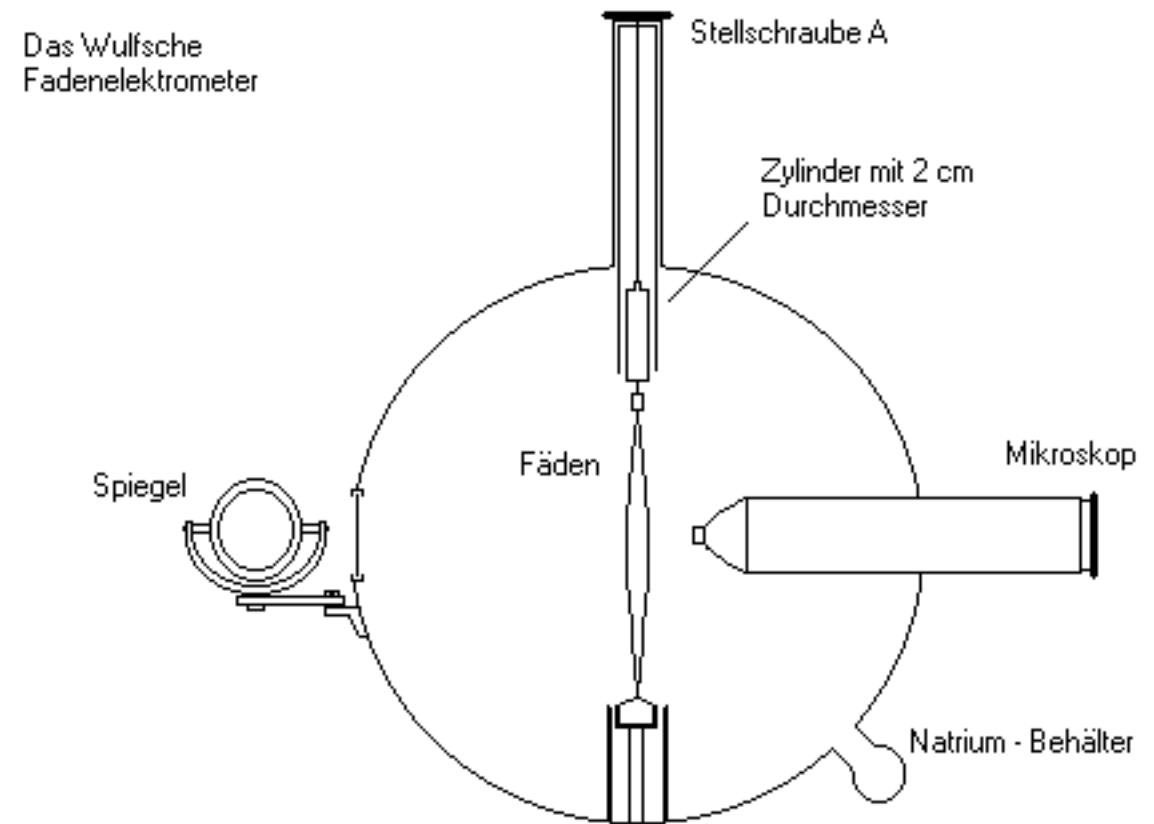
Cosmic Rays: Discovery

- Discovered by Victor Hess 1912
 - ▶ Nobel Prize in physics 1936
- Observation on balloon flights with electrosopes:
 - Rate of discharge reduces with increasing altitude, up to an altitude of 1000 m
 - Above this a strong increase of the discharge rate is observed, at 5000 m it is several times higher than the rate at ground level



Cosmic Rays: Discovery

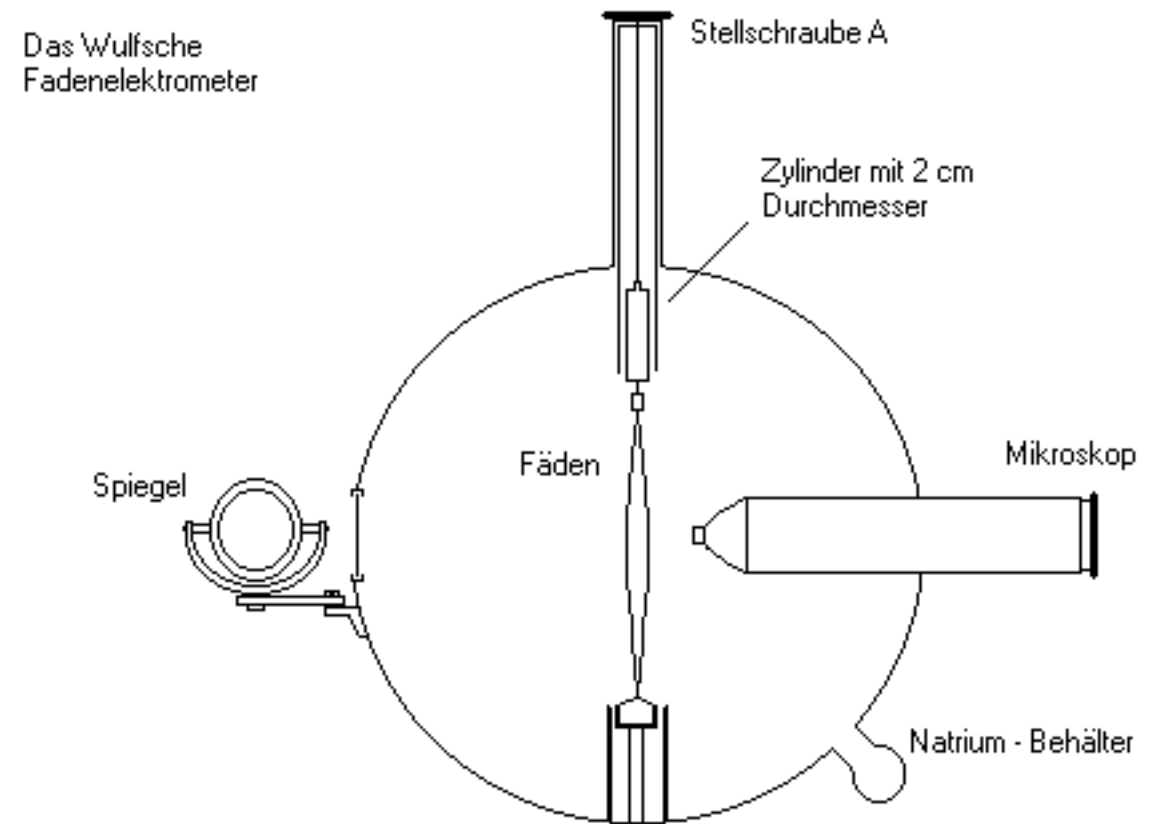
- The experimental method:
 - Electrometer, the distance of (electrostatically charged) strings gives the amount of charge on the strings
 - Discharge via ionising radiation



G. Federmann, Diplomarbeit, U. Wien, 2002

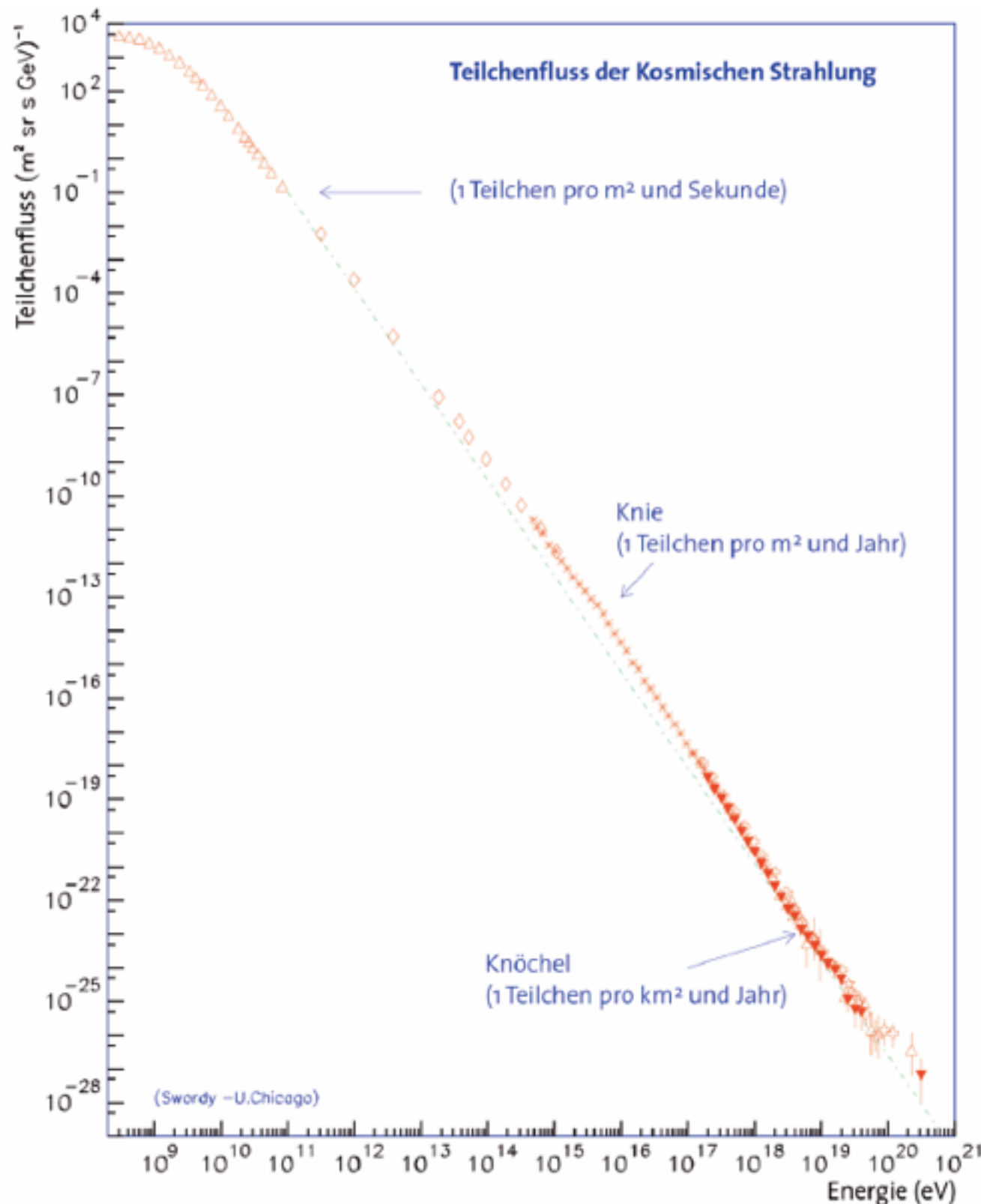
Cosmic Rays: Discovery

- The experimental method:
 - Electrometer, the distance of (electrostatically charged) strings gives the amount of charge on the strings
 - Discharge via ionising radiation
- Interpretation of the observation
 - ▶ Reduction of ambient radioactivity with increasing altitude (less radio nuclides, such as Radon)
 - ▶ The increase of radiation at high altitudes has to be due to extraterrestrial sources
 - ⇒ **“Höhenstrahlung”**



G. Federmann, Diplomarbeit, U. Wien, 2002

Cosmic Rays: Spectrum



- Extends over many orders of magnitude in energy and flux:
 - ▶ GeV (10^9 eV) - ZeV (10^{21})
 - ▶ $>1 \text{ cm}^{-2}\text{s}^{-1}$ - $< 1 \text{ km}^{-2}$ per century

- Follows a power law:

$$\frac{dN}{dE} \propto E^{-\gamma}$$

- $\gamma \sim 2.7$ $E < 10^{15} \text{ eV}$
- $\gamma \sim 3.0$ $10^{15} \text{ eV} < E < 10^{18} \text{ eV}$
- $\gamma \sim 2.7$ $10^{18} \text{ eV} < E$

Energy Density of Cosmic Rays

- Differential flux on earth
(parametrisation valid from $\sim$ GeV to $\sim$ 100 TeV):

$$\frac{dN}{dE} \approx \left(\frac{E}{\text{GeV}} \right)^{-2.7} \frac{\text{particles}}{m^2 \text{ sr s GeV}}$$

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$$\rho_E = \frac{4\pi}{c} \int E \frac{dN}{dE} \approx 1 \frac{\text{eV}}{\text{cm}^3}$$

Cosmic Rays: Power

- Rough estimate of the energy of cosmic rays within the Milky Way

$$V = \pi R^2 d \sim \pi (15 \text{ kpc})^2 (200 \text{ pc}) \sim 4 \times 10^{60} \text{ m}^3$$

$$\Rightarrow E_{\text{Strahlung}} \sim 4 \times 10^{66} \text{ eV} \sim 6.4 \times 10^{47} \text{ J}$$

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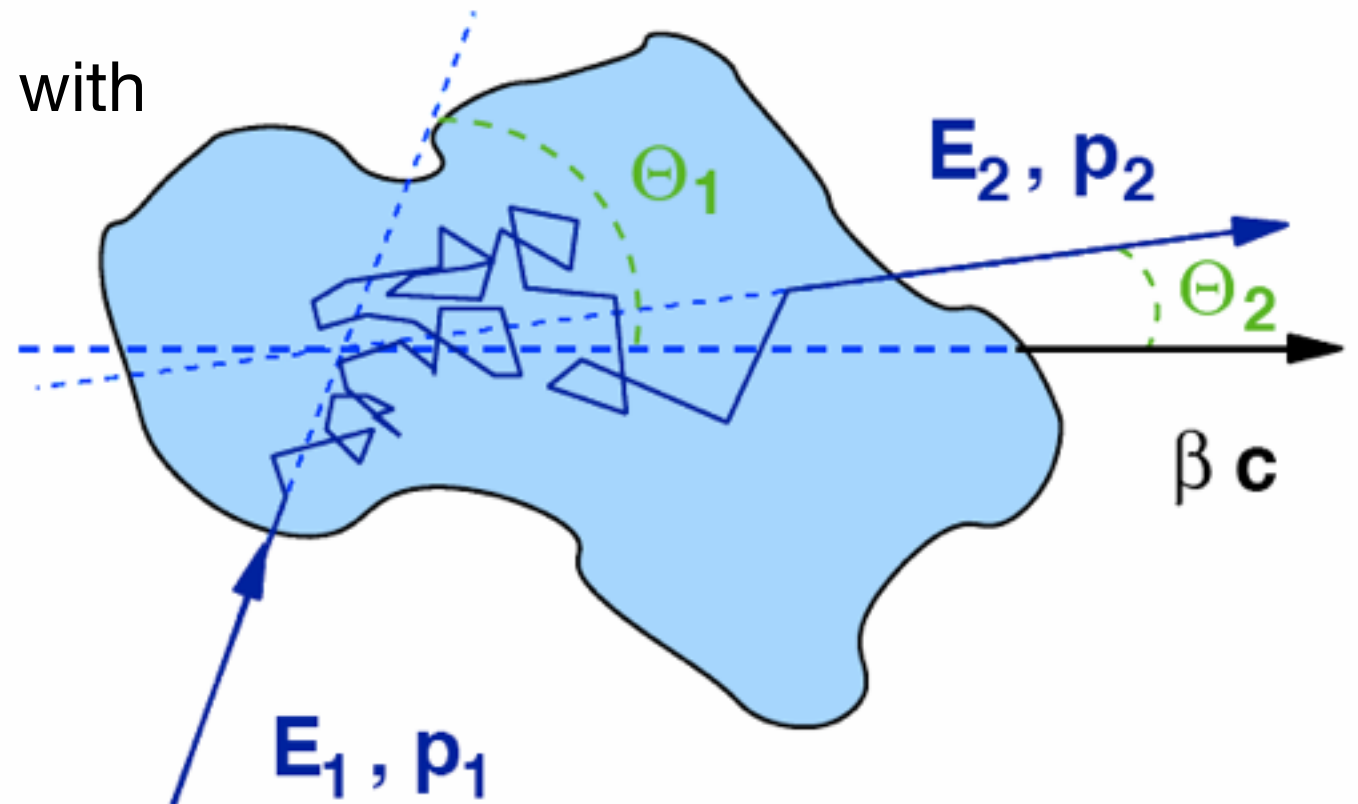
$$P_{\text{SN}} \sim 10^{35} \text{ W}$$

As comparison:
therm. power of the sun
 $\sim 4 \times 10^{26} \text{ W}$,
Milky Way $\sim 2 \times 10^{11}$ stars

- ▶ SNe could be the accelerators, would require $\sim 10\%$ efficiency

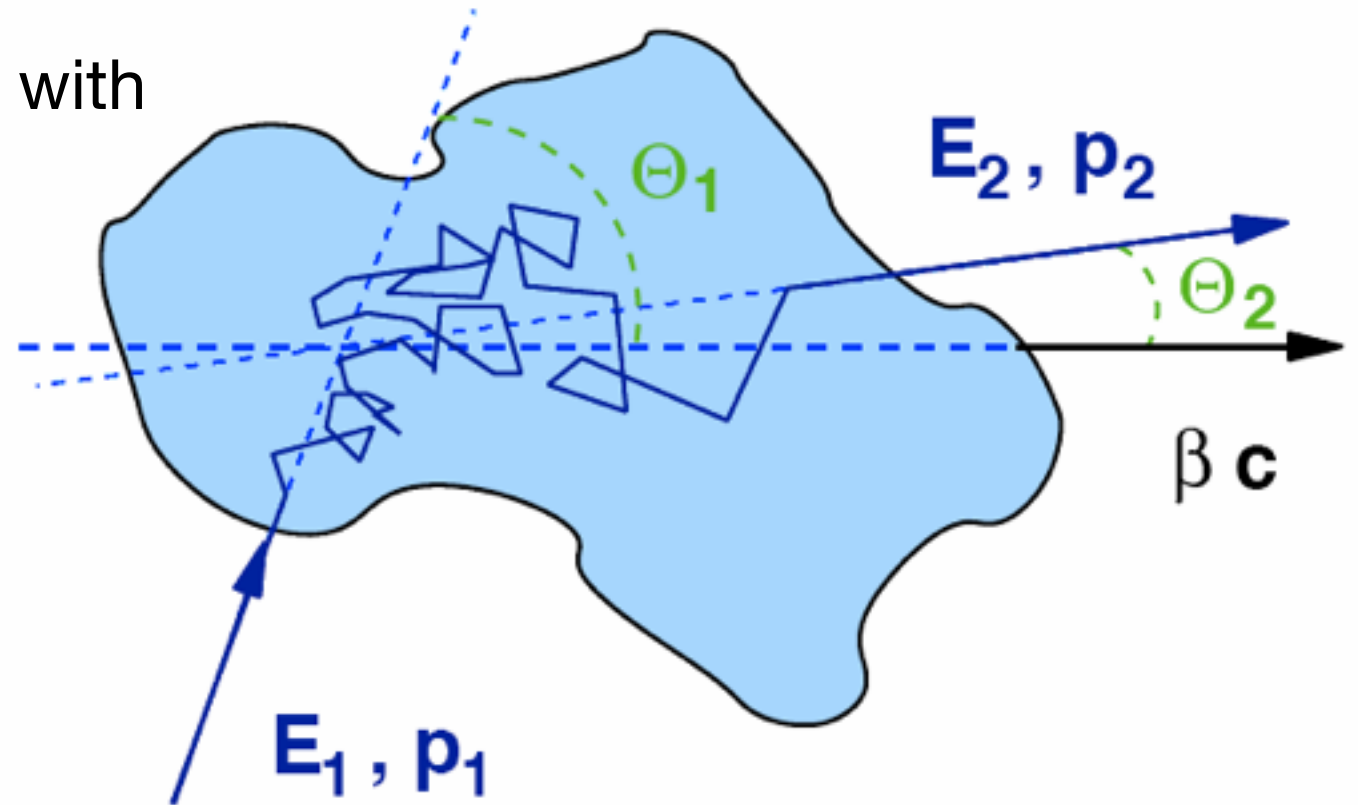
Second Order Fermi Acceleration

- Proposed by Enrico Fermi 1949
- The principle: Collisions of particles with interstellar gas clouds
 - Particle speed $\sim c$
 - Cloud speed βc
 - entrance and exit angle relative to cloud direction Θ_1, Θ_2



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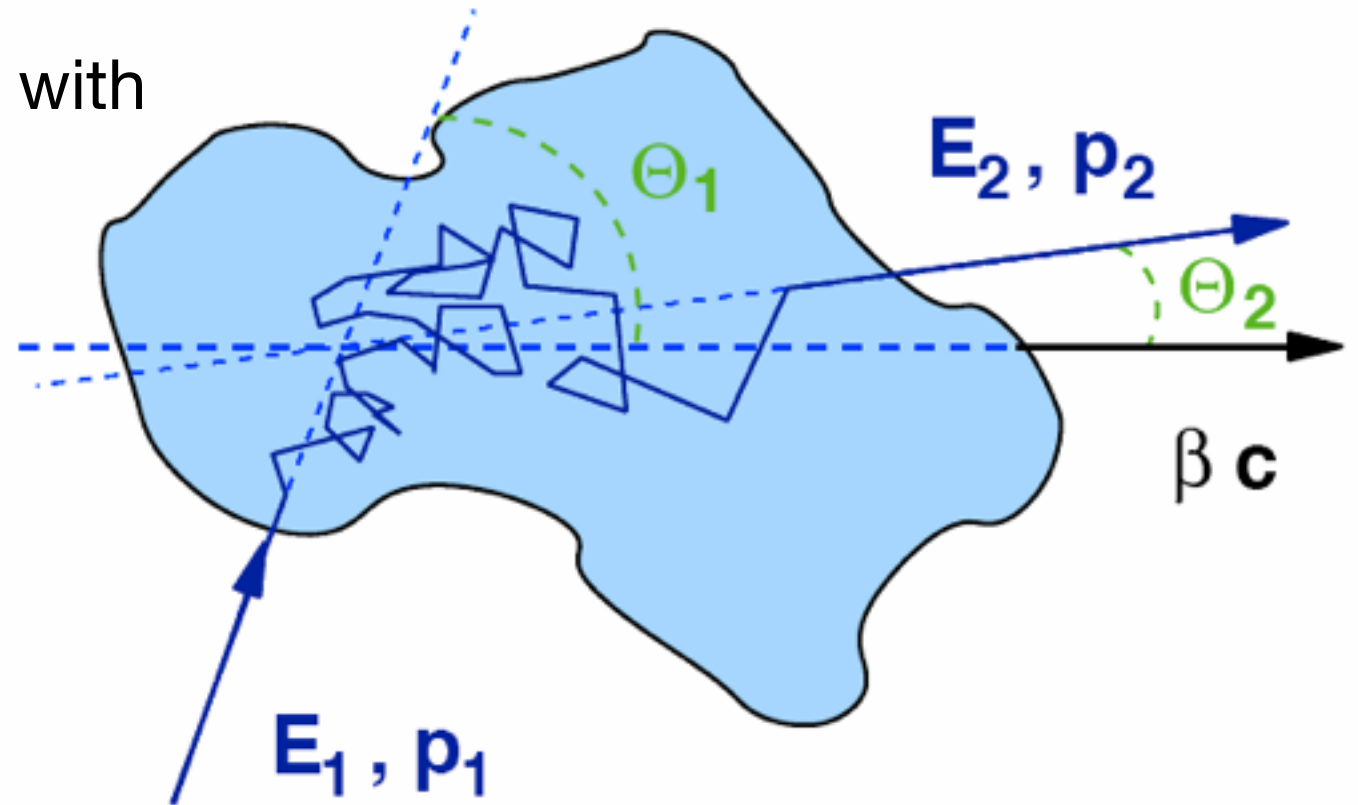


- Boost into cloud rest frame:

$$E'_1 = \gamma E_1 - \beta \gamma p_{||} = \gamma E_1 - \beta \gamma \cos \Theta_1 p \approx \gamma E_1 (1 - \beta \cos \Theta_1)$$

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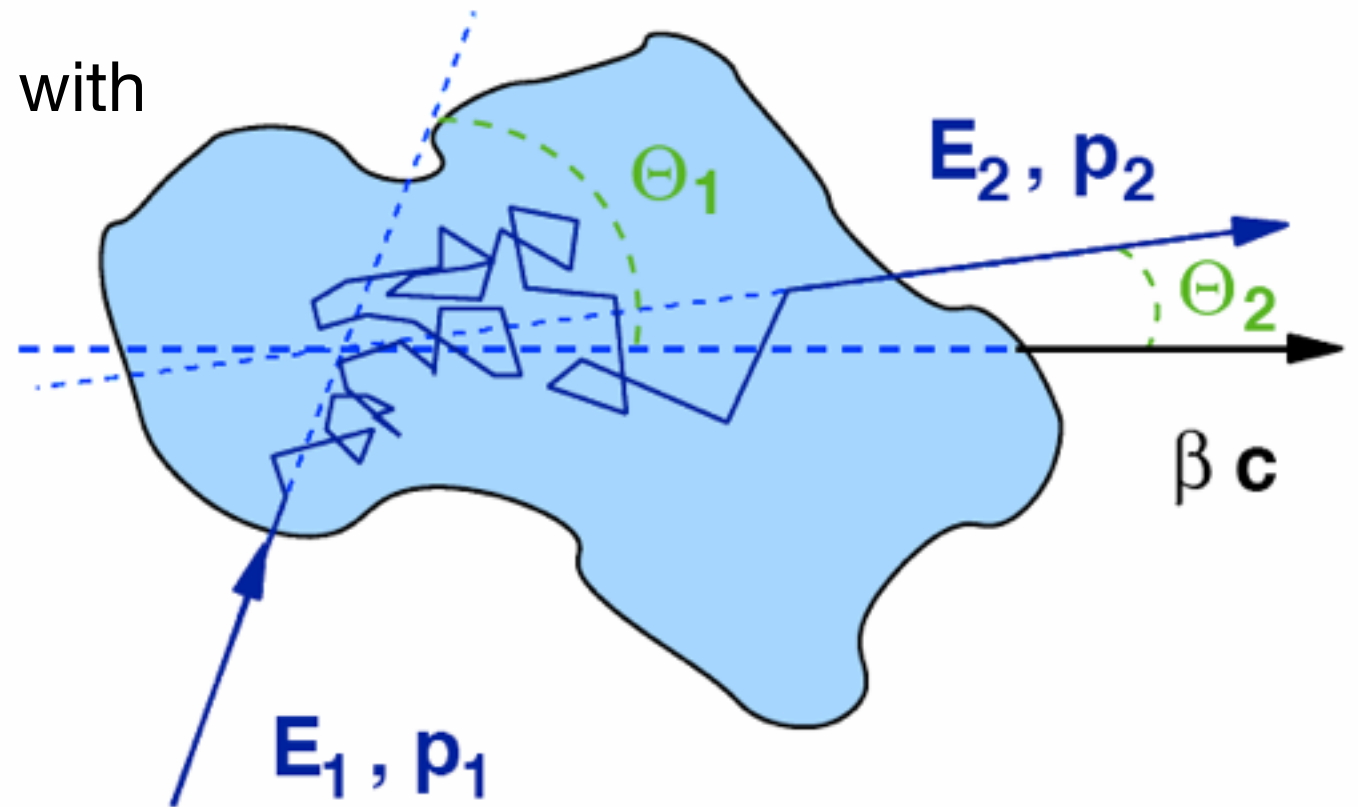
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- Energy difference $\frac{\Delta E}{E} = \frac{E_2 - E_1}{E_1} = \frac{E_2}{E_1} - 1$

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- ▶ Very low acceleration efficiency (proportional to β^2 - second order in β):
 - typical cloud speeds 10^4 m/s $\Rightarrow \beta \sim 3 \times 10^{-5}$
 - mean free path between collisions: ~ 30 pc $\Rightarrow$ every 100 years one collision



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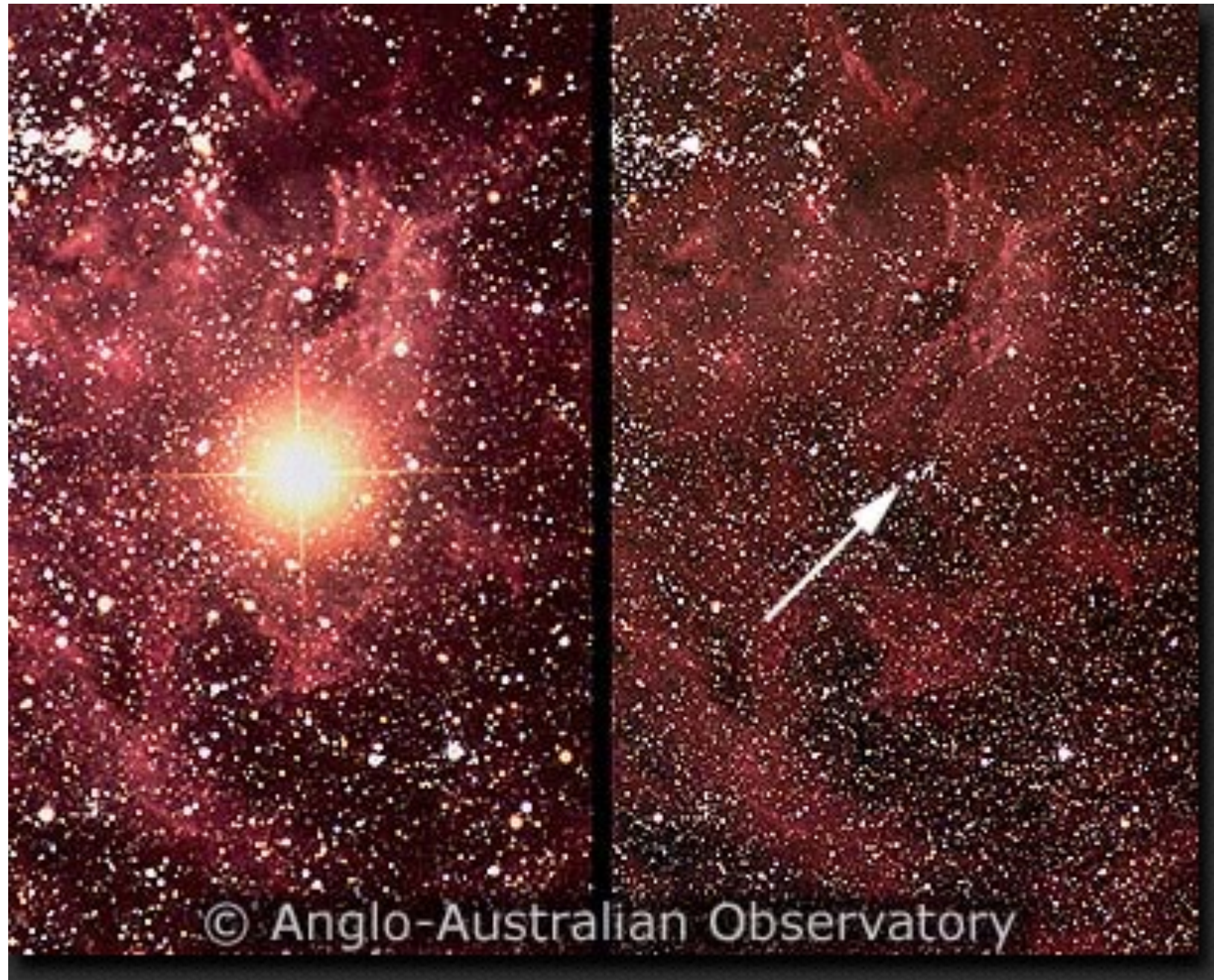
Characteristic time: $\frac{3\tau}{4\beta^2} \approx 6 \times 10^{10} a$

Supernovae

- Classification in to main types
 - SN I: no hydrogen lines in spectrum
 - SN Ia collapse of an accreting white dwarf in a binary star system to a neutron star
 - SN II: hydrogen lines visible
 - Gravitational collapse of a massive star at the end of its life
 - Star burns up to the formation of iron in the core, then no radiation pressure to counter gravitation
 - ▶ Atoms are converted to neutrons via electron capture
 - ▶ Star collapses with a speed of $\sim 0.1 c$
 - ▶ Matter is reflected at the stable neutron star in the core
 - ▶ A shock wave runs outwards
 - ▶ An enormous number of neutrinos is produced ($\sim 10^{58}$), despite their small interaction cross section they further drive the shock wave

Supernova SN1987a

- Supernova explosion 1987 in the great Magellanic Cloud (small partner galaxy of the Milky Way)



Supernova SN1987a

Inner debris of the Supernova 1987A (SN 1987A) ring



Outer bipolar
outflow of
gas and
outer
ring

Inner bipolar
outflow
of debris

Hot fingers
of gas

Ring

Blast
wave

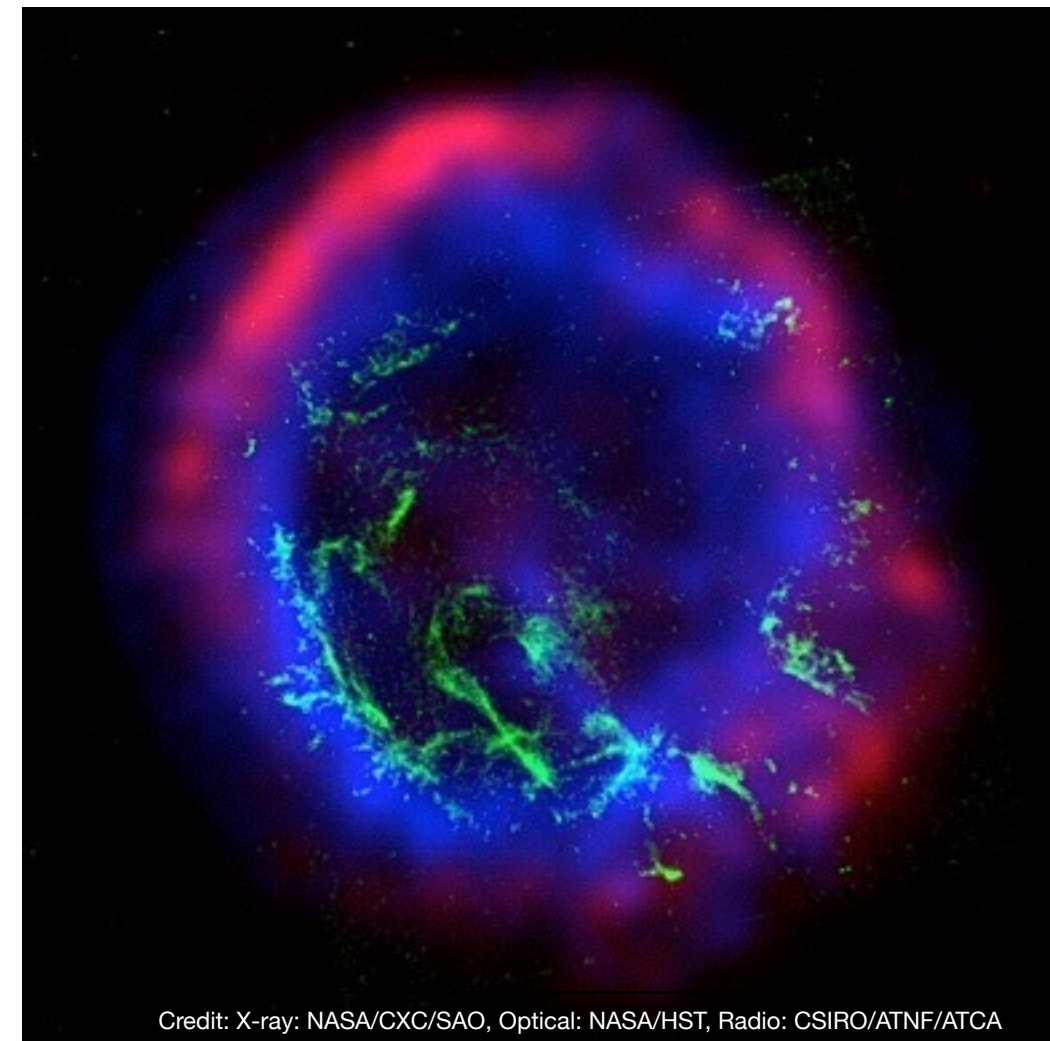
Supernova
debris

Hidden
neutron star
or black hole

Credit: NASA, ESA, and A. Feild (STScI)

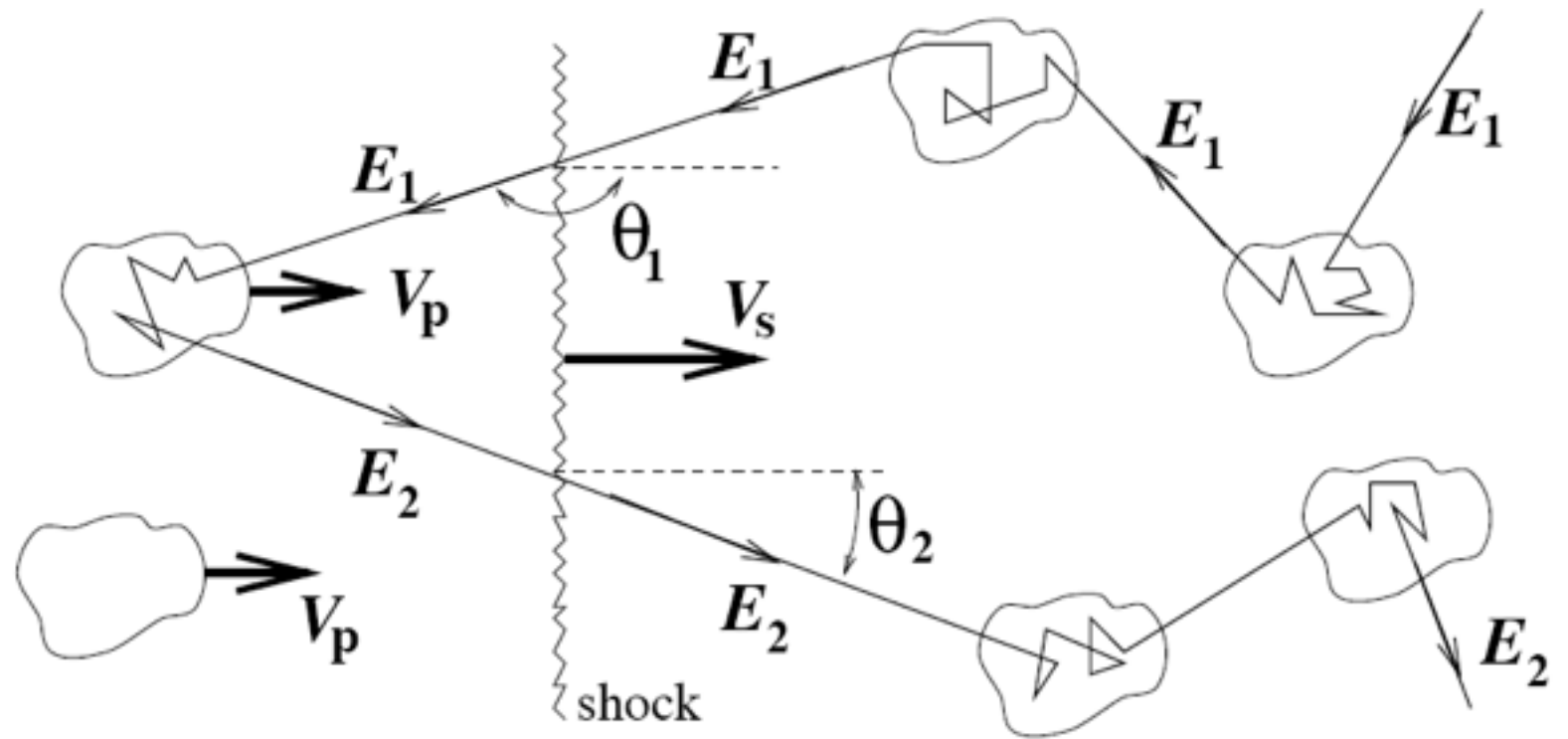
First Order Fermi Acceleration

- Extension of Fermi acceleration concept to supernova shocks
- Formation of the shock wave:
 - SN ejects large amounts of material (several solar masses) with high velocity into the interstellar medium
 - $V_{\text{material}} \gg V_{\text{sound(ISM)}}$, $V_{\text{material}} \sim 10^7 \text{ m/s}$, $V_{\text{sound(ISM)}} \sim 10^4 \text{ m/s}$
 - ▶ Since the matter is much faster than the speed of sound a shock front develops
- The shock wave propagates in the ionized plasma of the ISM (single atoms, specific heat 5/3)
- Hydrodynamics can show that the speed of the shock wave is:
 $V_{\text{shock}}/V_{\text{material}} \sim 4/3$



First Order Fermi Acceleration

- Particle acceleration by repeated crossing of shock fronts
- As for second order Fermi Acceleration the particles are scattered by magnetic field inhomogeneities / turbulences
 - Key “feature”: behind the shock, these turbulences move with the speed of the ejected matter

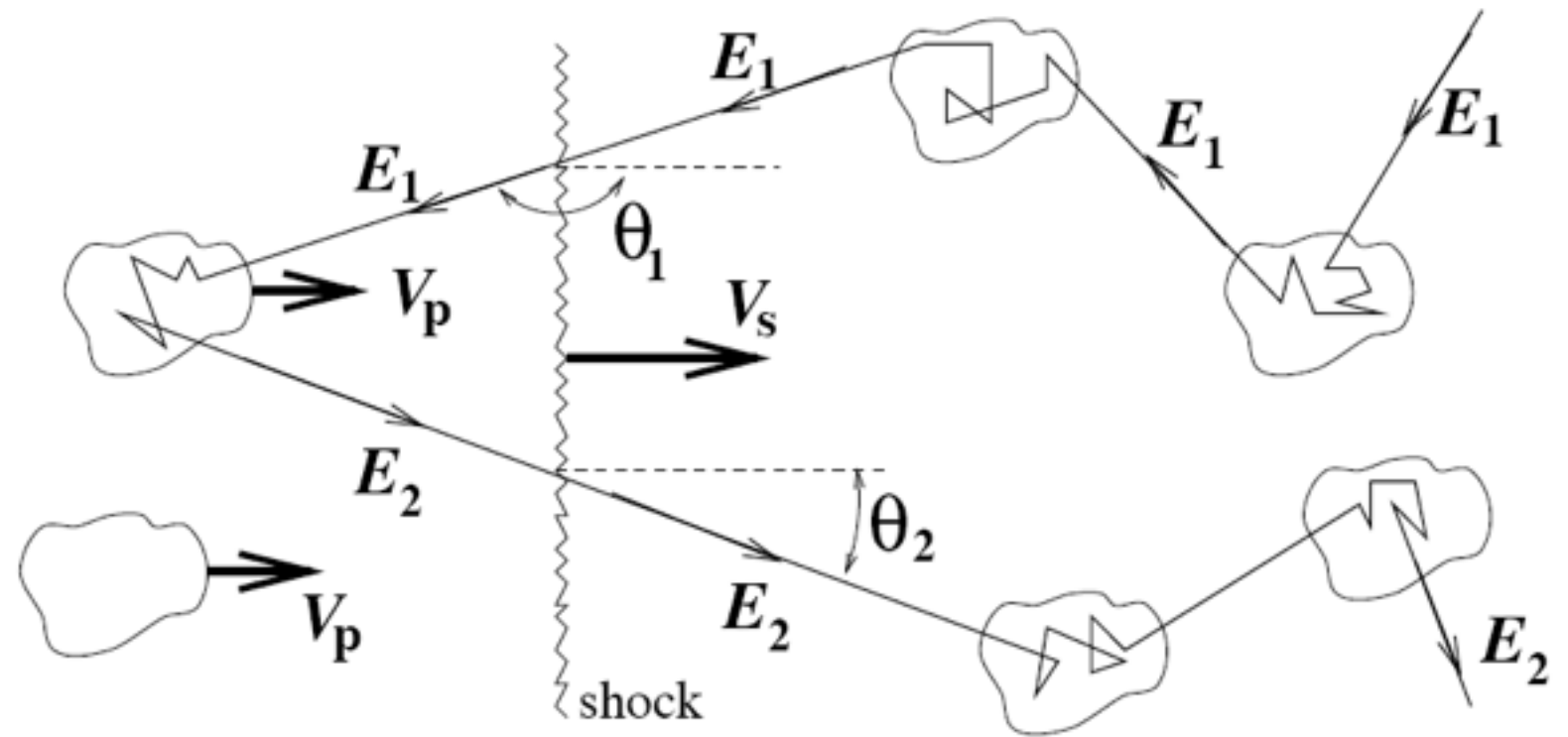


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- Considerations concerning incidence angles of particles:

- Shock front introduces directional motion - projection of flux onto front

Effective area

$$A' = -A \cos \Theta_1$$

Flux crossing shock

$$\frac{dN}{d\cos\Theta_1} \propto -\cos\Theta_1$$

First Order Fermi Acceleration

- Mean value of angles (key point here: Crossing of shock only for $\cos\Theta_1 < 0$, other particles do not contribute!):

$$\begin{aligned}\langle \cos\Theta_1 \rangle &= \frac{\int_{-1}^0 \frac{dN}{d\cos\Theta_1} \cos\Theta_1 d\cos\Theta_1}{\int_{-1}^0 \frac{dN}{d\cos\Theta_1} d\cos\Theta_1} \\ &= \frac{\int_{-1}^0 -\cos^2\Theta_1 d\cos\Theta_1}{\int_{-1}^0 -\cos\Theta_1 d\cos\Theta_1} = -\frac{2}{3}\end{aligned}$$

$$\text{analog: } \langle \cos\Theta_2 \rangle = \frac{2}{3}$$

First Order Fermi Acceleration

- Mean energy change (analogous to second order FA):

$$\frac{\Delta E}{E} = \frac{(1 - \beta \cos \Theta_1)(1 + \beta \cos \Theta'_2)}{1 - \beta^2} - 1$$

β : Speed of matter behind shock wave

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- for $\beta \ll 1$

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- ▶ Substantially more efficient than second order acceleration due to two effects:
 - ▶ large velocity differences of material before and after shock front
 - ▶ directed motion of shock instead of random drifting
- ▶ $\beta \sim 3 \times 10^{-2}$, acceleration linear in β (first order in β)

Energy Spectrum

- Energy gain per cycle

$$\frac{\langle \Delta E \rangle}{E} = \zeta \approx \frac{4}{3} \beta_{Plasma} \approx \beta_{Shock}$$
$$\Rightarrow E = E_0 (1 + \zeta)^k \quad \text{after } k \text{ cycles}$$

- Reminder:

$v_{shock}/v_{material} \sim 4/3$
for strong shocks

Energy Spectrum

- Energy gain per cycle

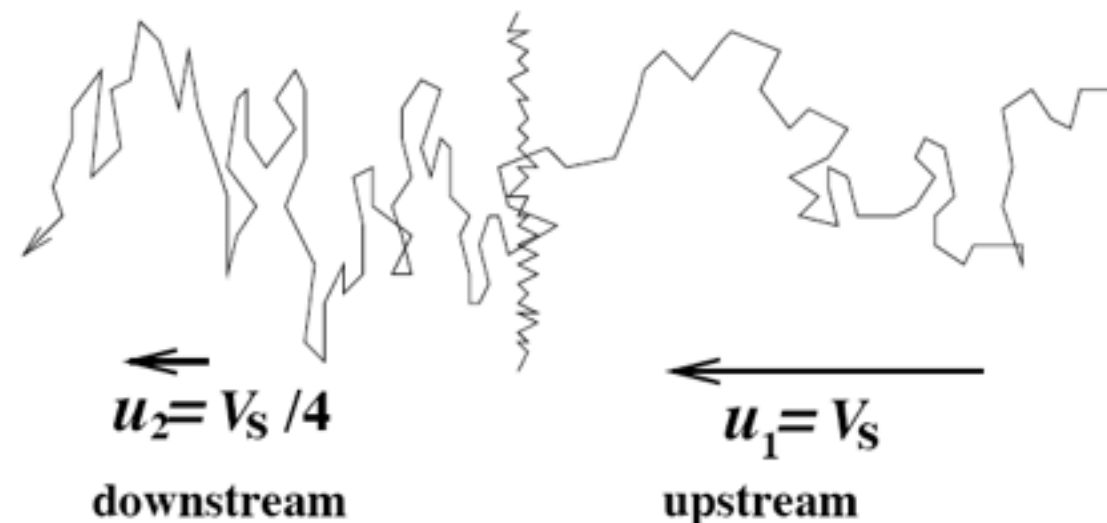
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for strong shocks

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=> A particle can get lost if it falls too far behind the shock



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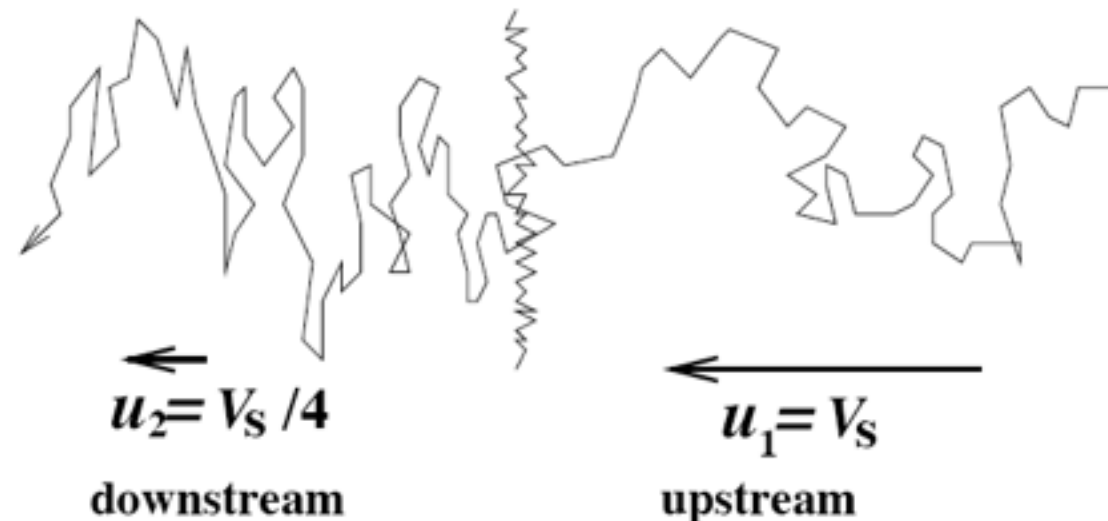
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In general: the material behind the shock is $1/4$ slower than the front itself, the loss rate depends on that difference:

$$R_{loss} = n_{CR} v_s / 4 \quad n_{CR} \text{ is the particle density}$$

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- Particle flux from “upstream” through the shock:

Particle movement relative to shock (particle velocity v_t)

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The probability to cross the shock front at least k times is:

$$P_{\text{cross} > k} = (1 - P_{\text{escape}})^k = \left(1 - \frac{v_s}{v_t}\right)^k \approx (1 - \beta_{\text{Shock}})^k$$

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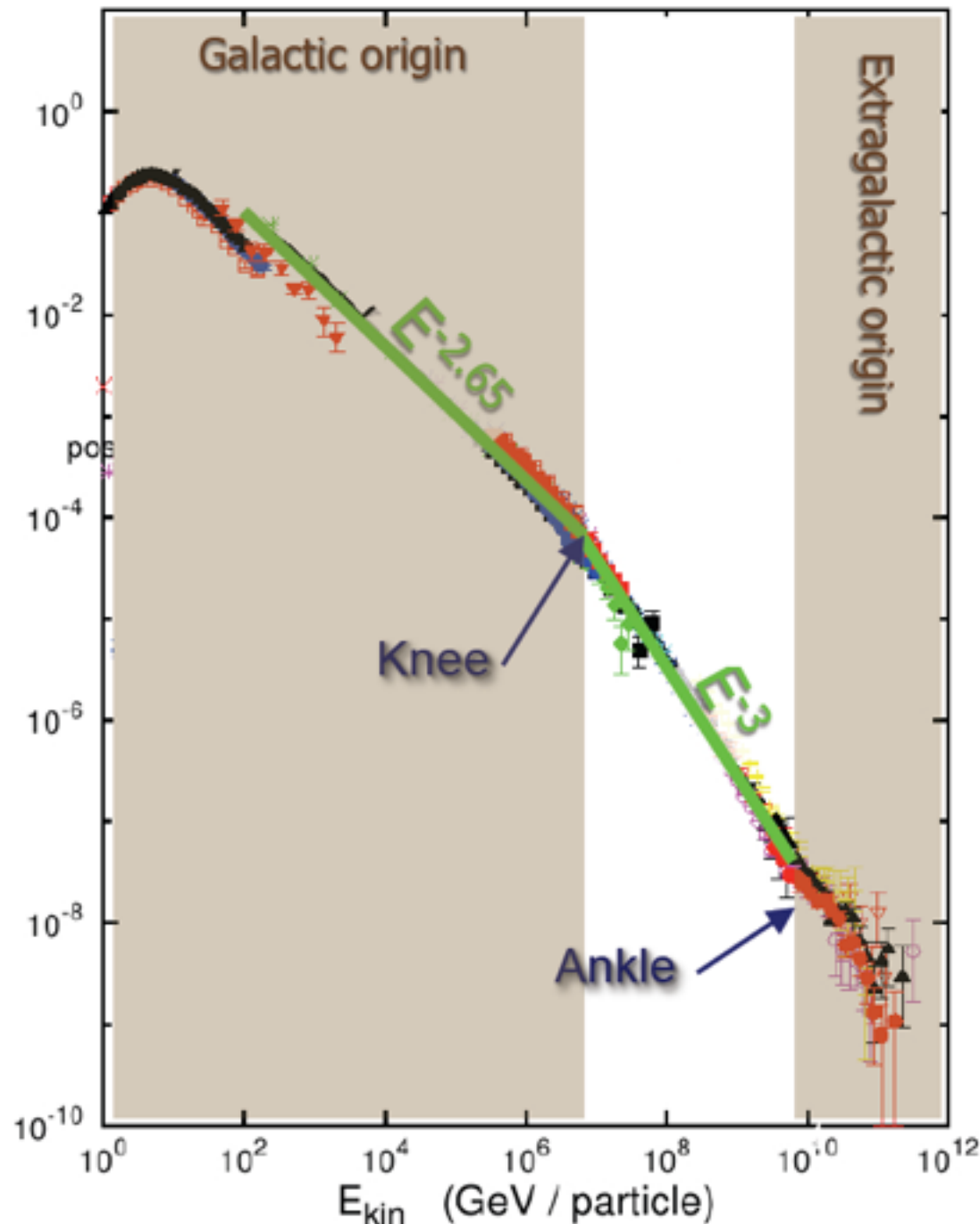
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- Differential particle spectrum

$$\frac{dN}{dE} \propto E^{-2}$$

Maximum Energy

- First order Fermi Acceleration can reach energies up to $\sim 10^{14}$ eV for shock waves originating from supernova explosions (incomplete derivation in Backup)



based on a shock lifetime of ~ 1000 years
a shock speed of $0.03 c$, and a B field in the
nT range

⇒ Extends up to the knee of the cosmic
ray spectrum

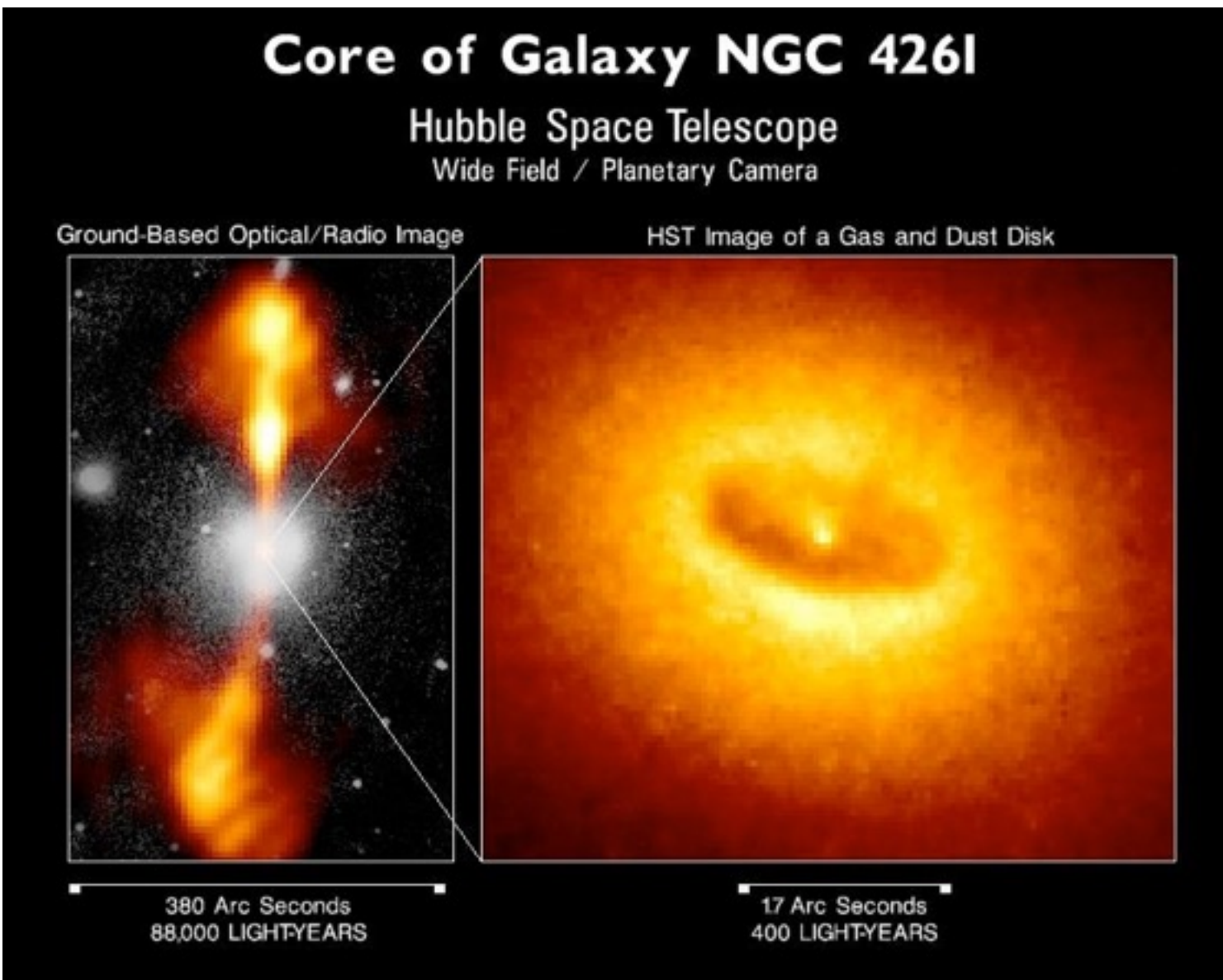
Supernova shock acceleration is well
established as a source for cosmic rays

But: What is the origin of the very highest
energies above 10^{18} GeV?

S. Coutu, TIPP11

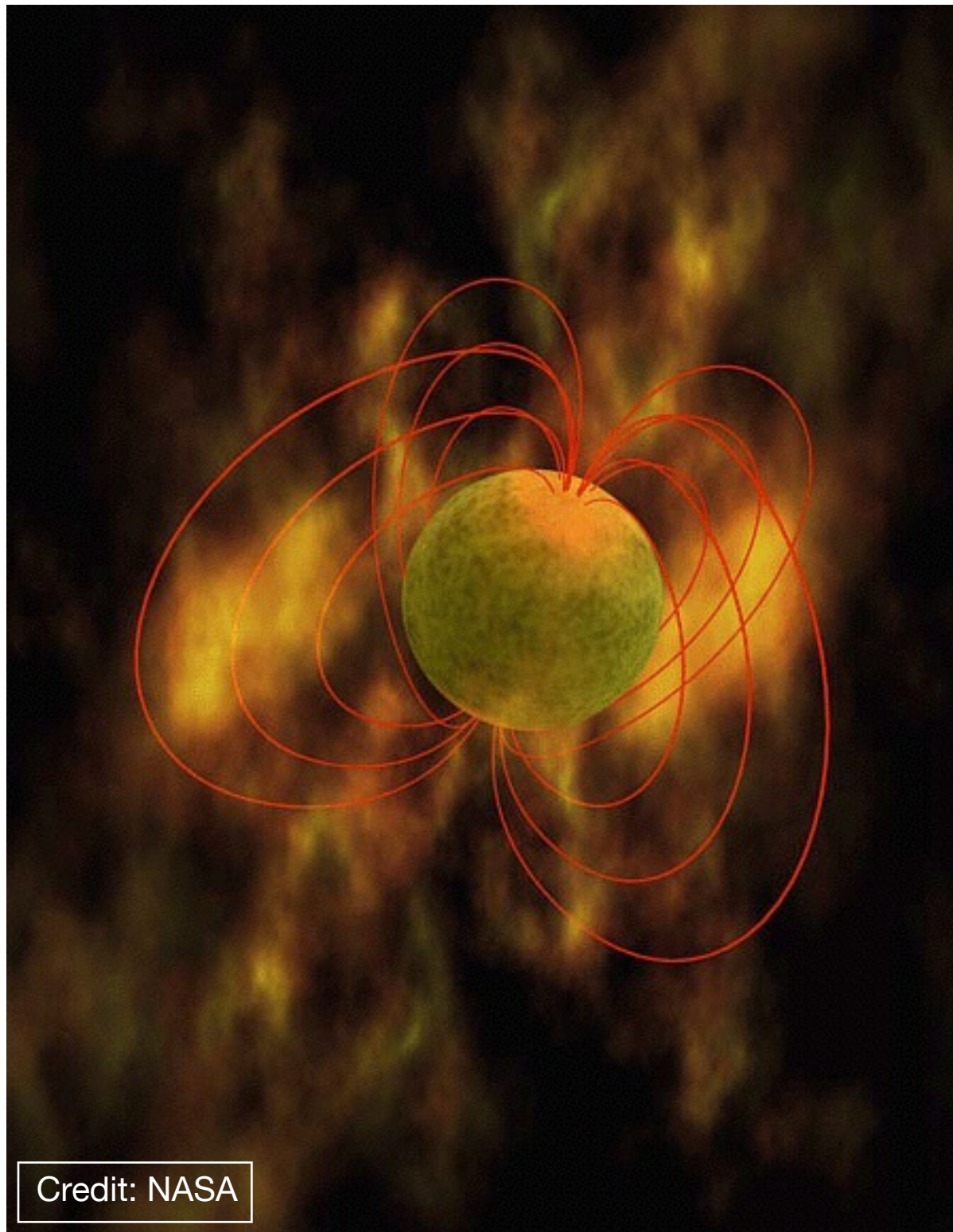
Highest Energies?

- How are energies > 1 PeV reached?



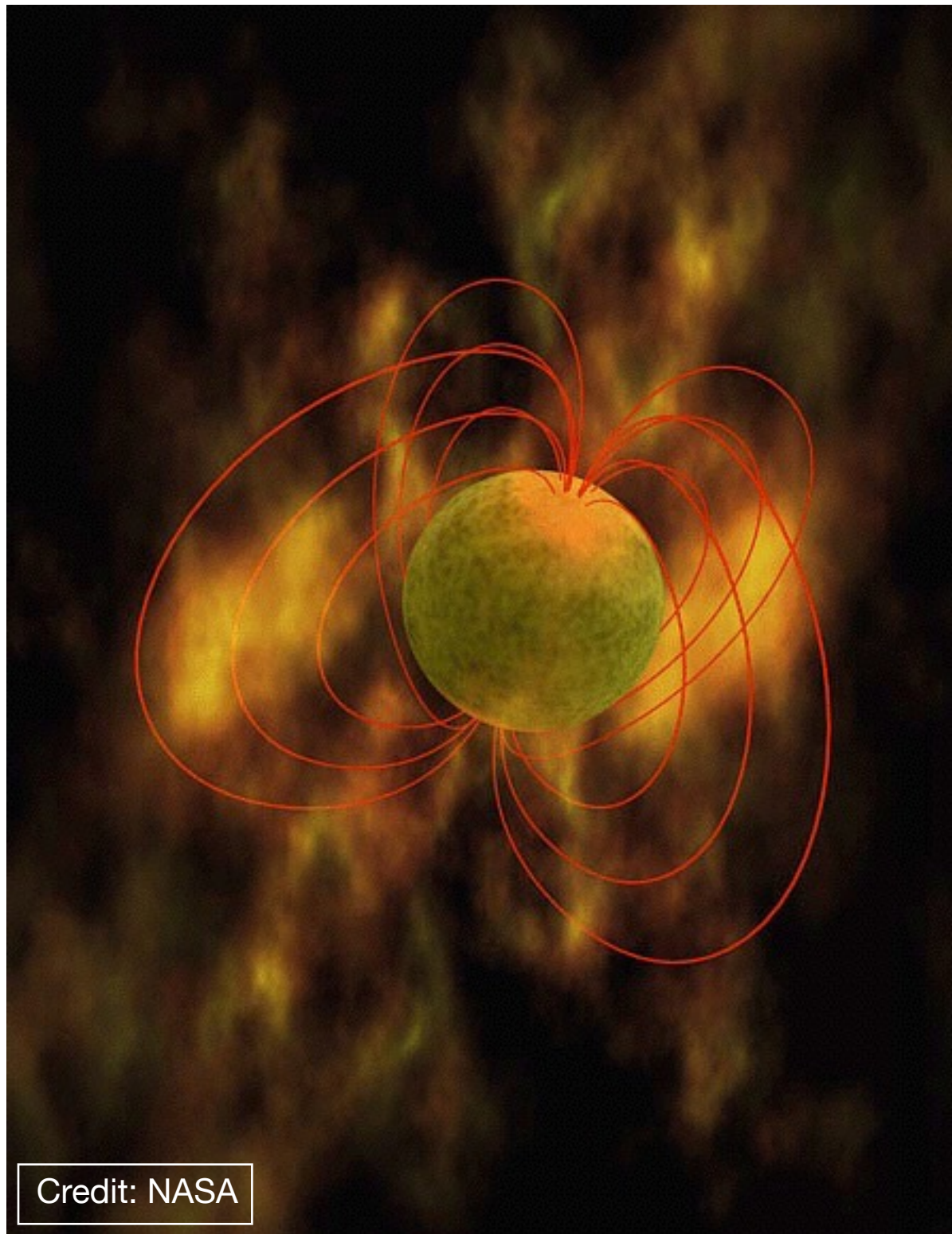
- More energetic events
 - Active galactic nuclei
 - Pulsars (neutron stars)
 - GRB's
- ⇒ Extreme magnetic fields
- ⇒ Shock acceleration in highly relativistic jets: additional γ - Factor

One Example: Neutron Stars



- Neutron stars: compact remnants of supernova explosions
 - radius ~ 10 km
 - extreme rotation: up to $\sim 40\,000$ RPM
 - magnetic fields up to $\sim 10^8$ T
 - mass $\sim 1.4 M_{\text{Sun}}$

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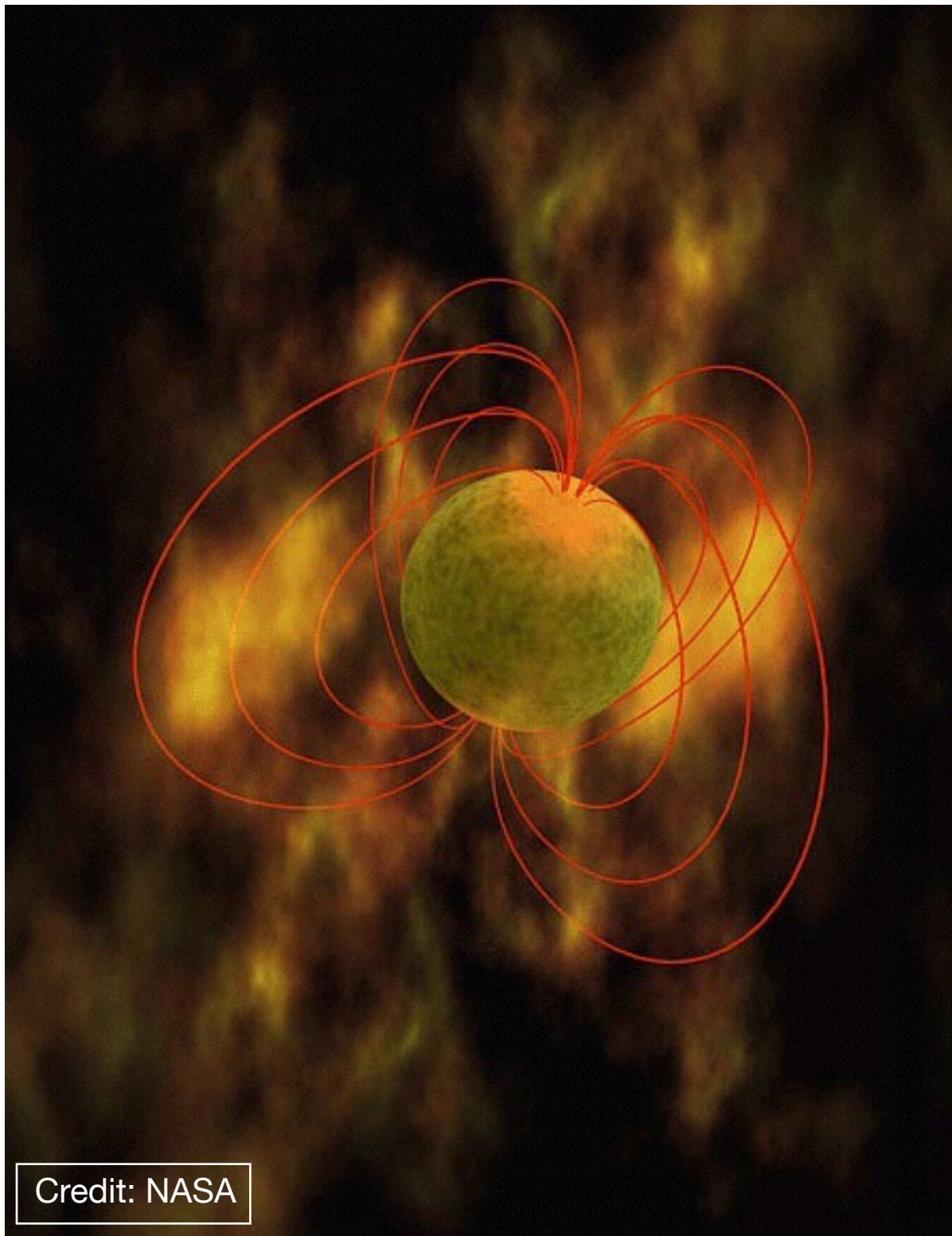


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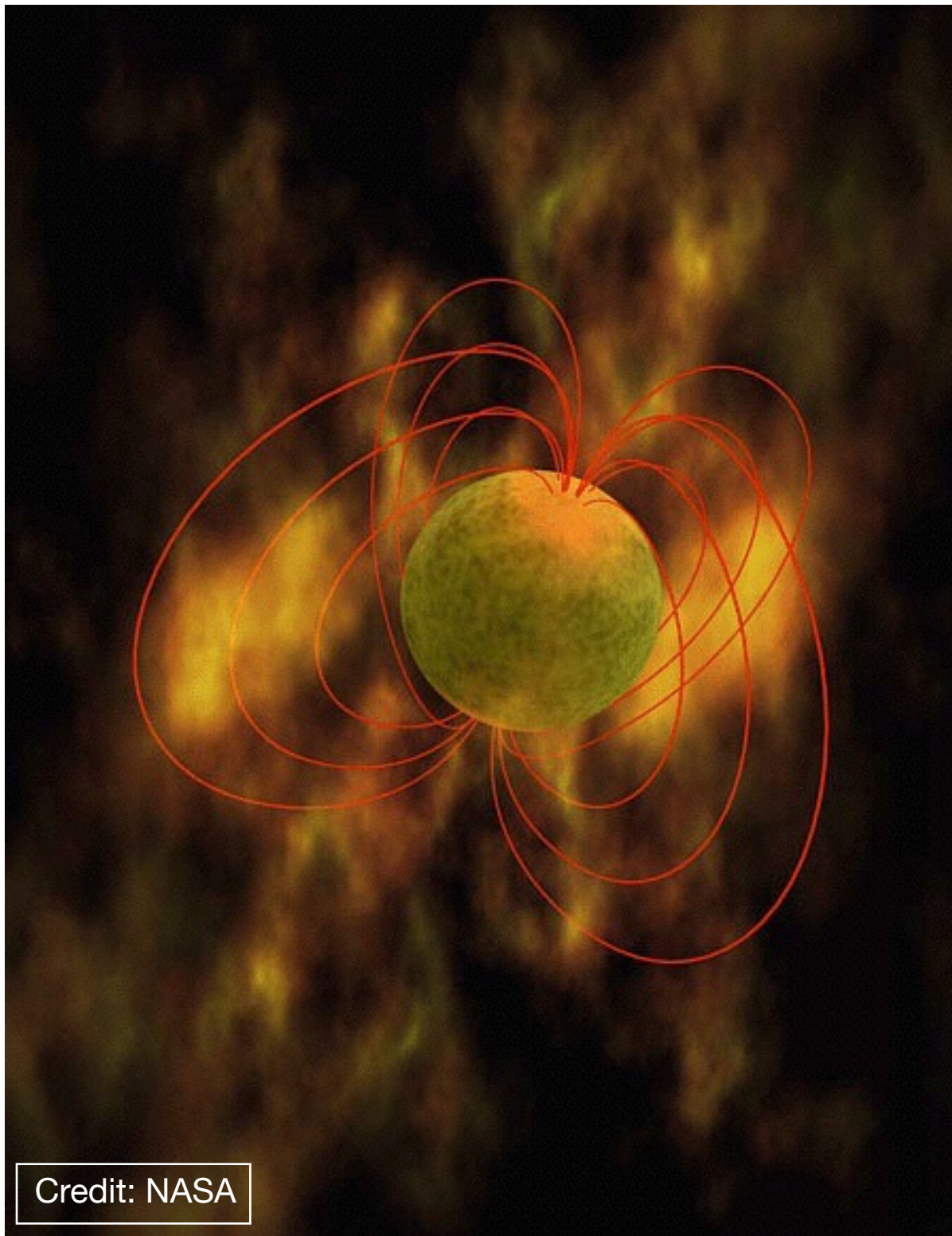
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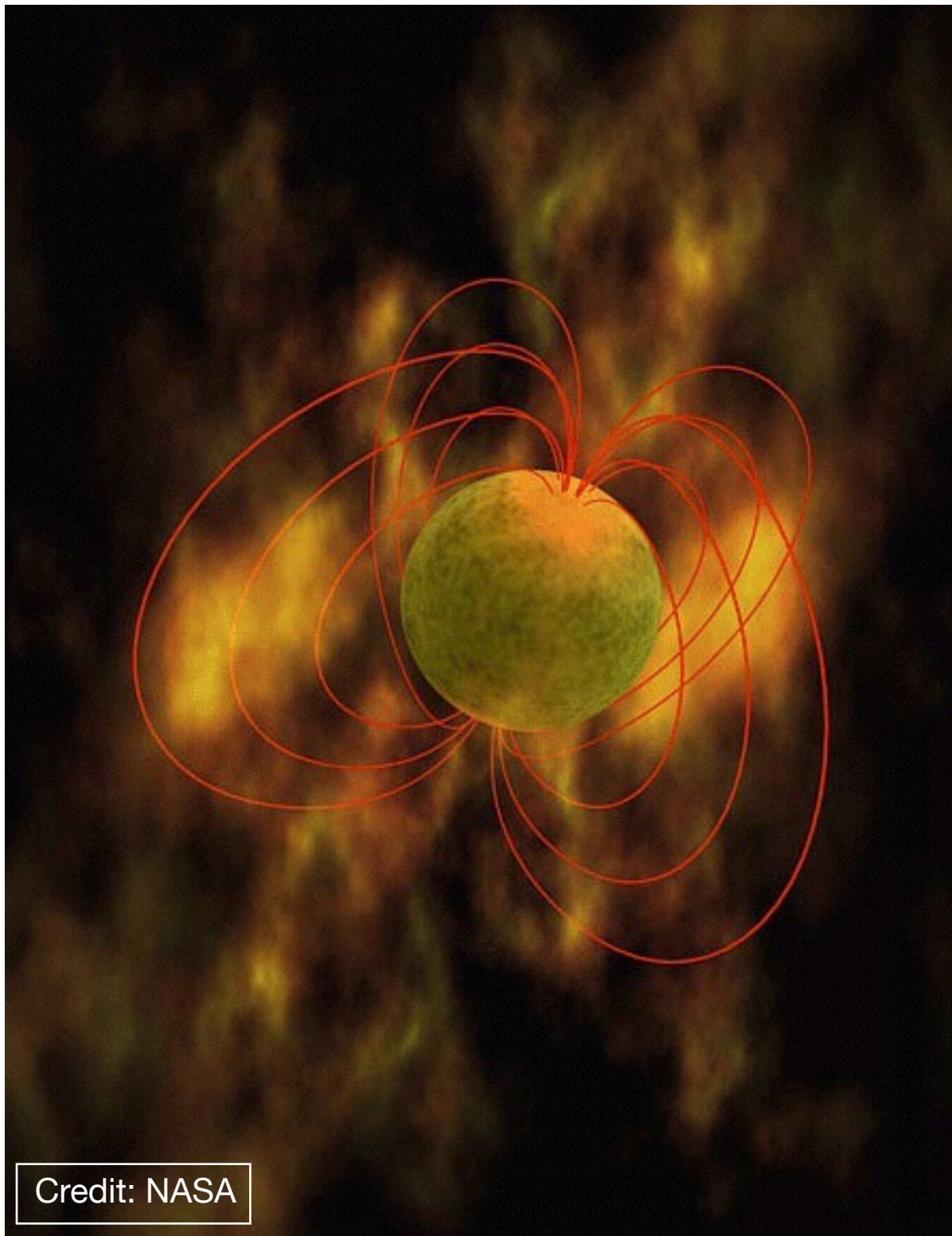
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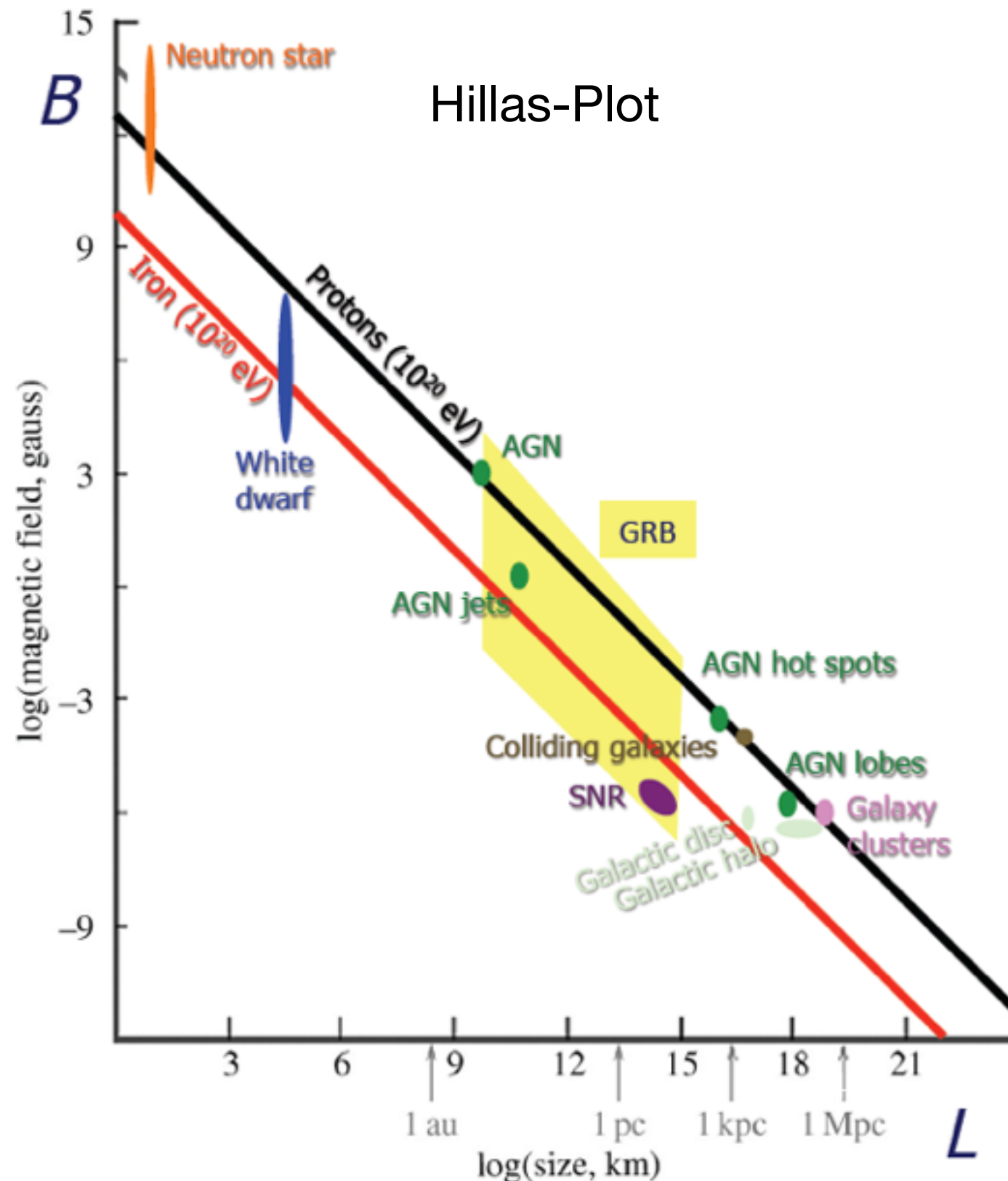
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For $Z = 1$ (protons):

$$E_{\text{max}} \sim 5 \text{ J} = 3 \times 10^{20} \text{ eV}$$

Credit: NASA

Candidates for the Highest Energies



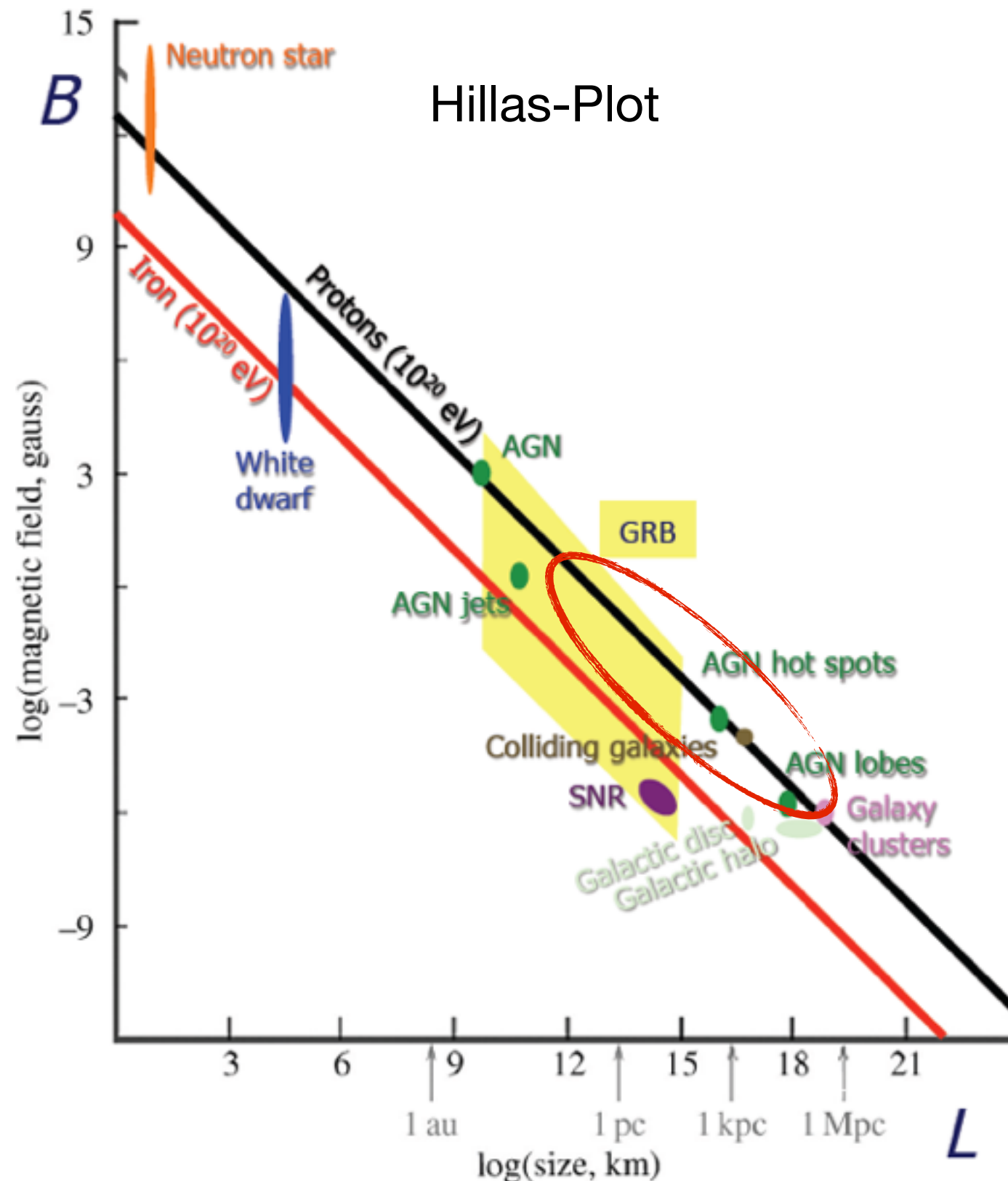
Particles are accelerated as long as they stay in the high field region: $r_L < L$

$$E_{\max} \approx 10^{20} \text{ eV } Z B_{\mu G} L_{100 \text{ kpc}}$$

For Iron nuclei 26 x higher energies are possible relative to protons!

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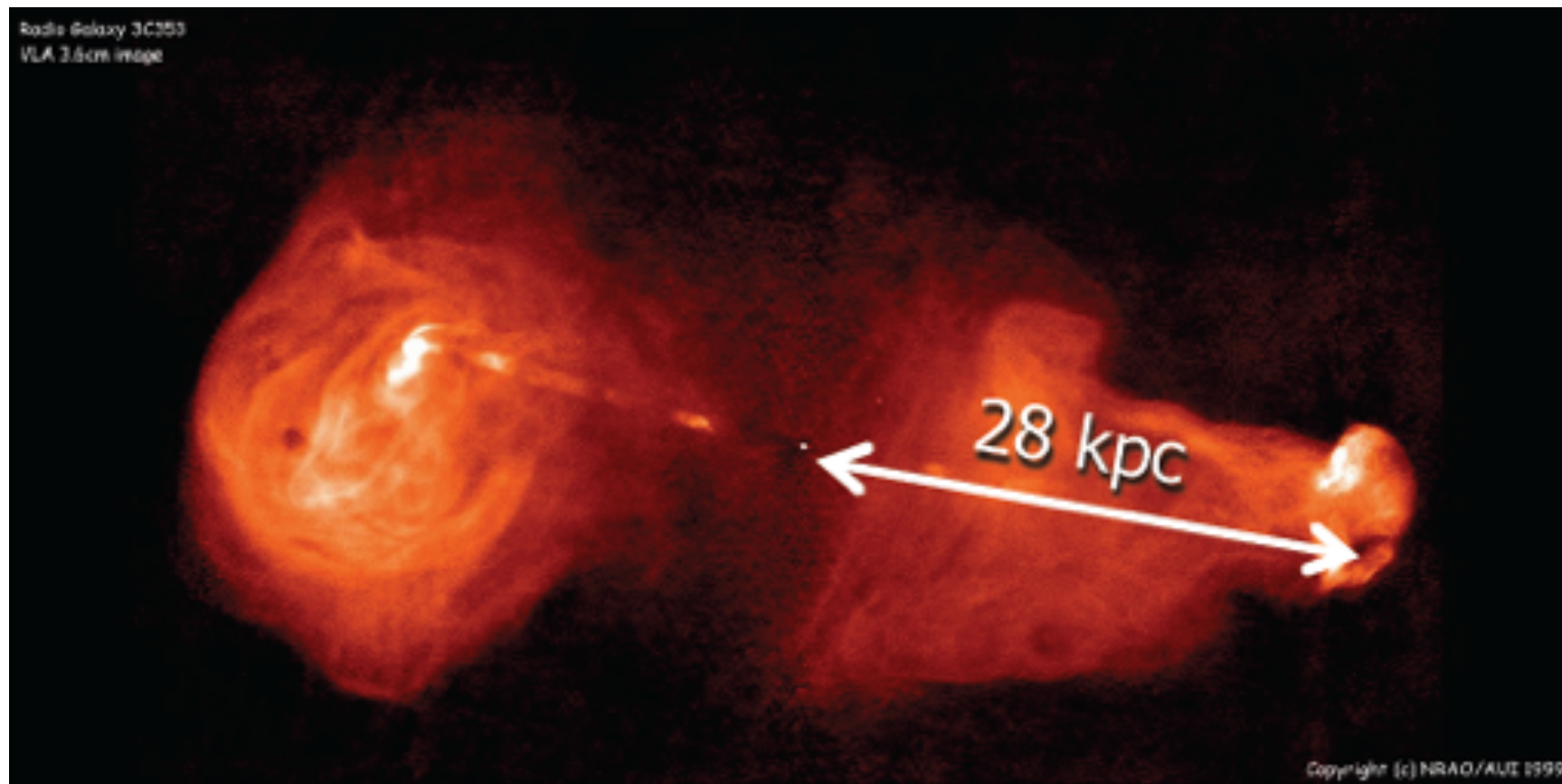
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Beyond this simple consideration:
Radiation losses in the source have to be taken into account:
synchrotron radiation, photo reactions

S. Coutu, TIPP11

One Example

- 3C353 - Active galaxy 130 Mpc away



... more about the highest energies next week!

Propagation of Cosmic Rays

- The source spectrum of shock acceleration follows an E^{-2} distribution, but we observe $E^{-2.7}$, why?
 - ▶ Energy-dependent loss of particles when travelling through the galaxy
- Important contributions:
 - diffusion
 - convection
 - acceleration
- decay of unstable particles and nuclei
- collisions
- cascade production, spallation of heavy nuclei

transport in turbulent
galactic magnetic fields

loss processes

Leaky Box Model

- Very simple model assuming cosmic rays propagate freely in the galaxy with a constant escape / loss probability

$$N(E, t) = N_0(E)e^{-t/\tau_{escape}}$$

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- Adaptation to reality - escape / loss probability is energy dependent:
 - particles with higher energy can more easily leave the “magnetic confinement” in the galaxy and are more likely to participate in inelastic reactions
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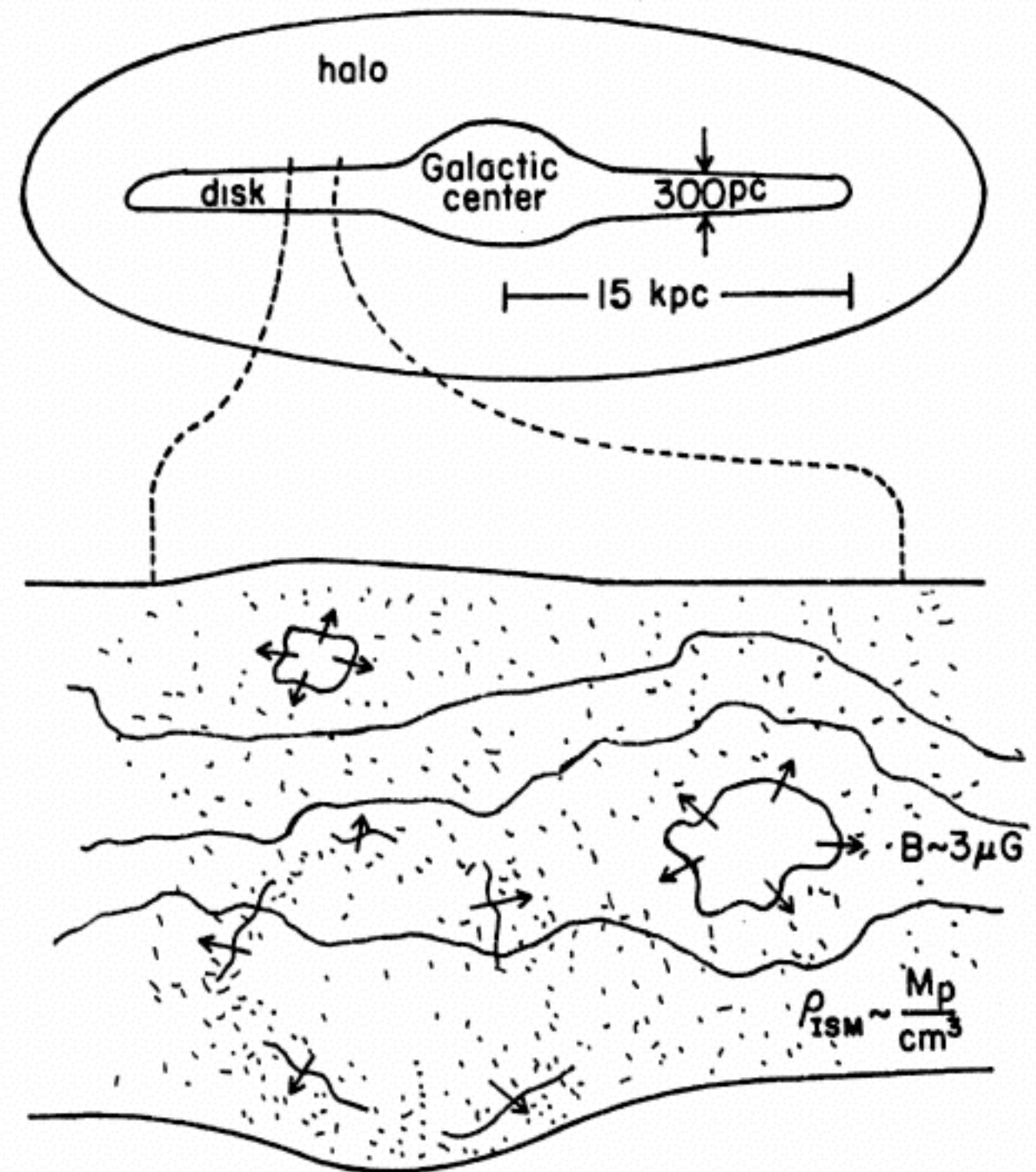
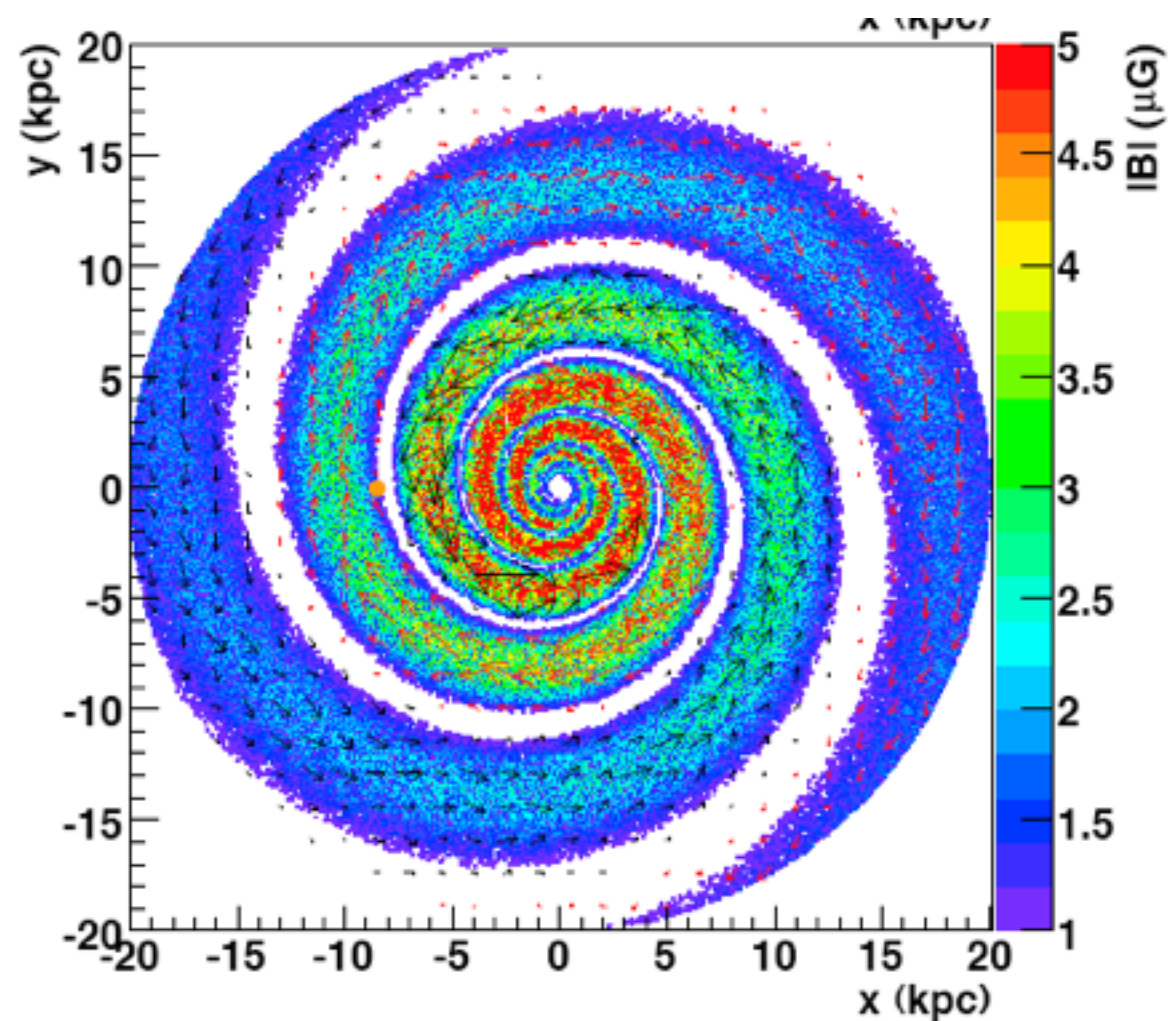
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 - ▶ The observed spectrum gets steeper than the source spectrum: $E^{-2.7}$
- Loss probability due to inelastic reactions depends on amount of traversed matter
 - Density of the ISM in the galaxy: $\sim 1 \text{ proton/cm}^3 \sim 1.7 \times 10^{-24} \text{ g/cm}^3$
 - ▶ per year one particle traverses $\sim 1.5 \times 10^{-6} \text{ g/cm}^2$
 - ▶ loss after traversing $\sim 5 - 10 \text{ g/cm}^2$ (derived from observed composition)
 - ▶ Particles stay in the galaxy for about 5×10^6 years

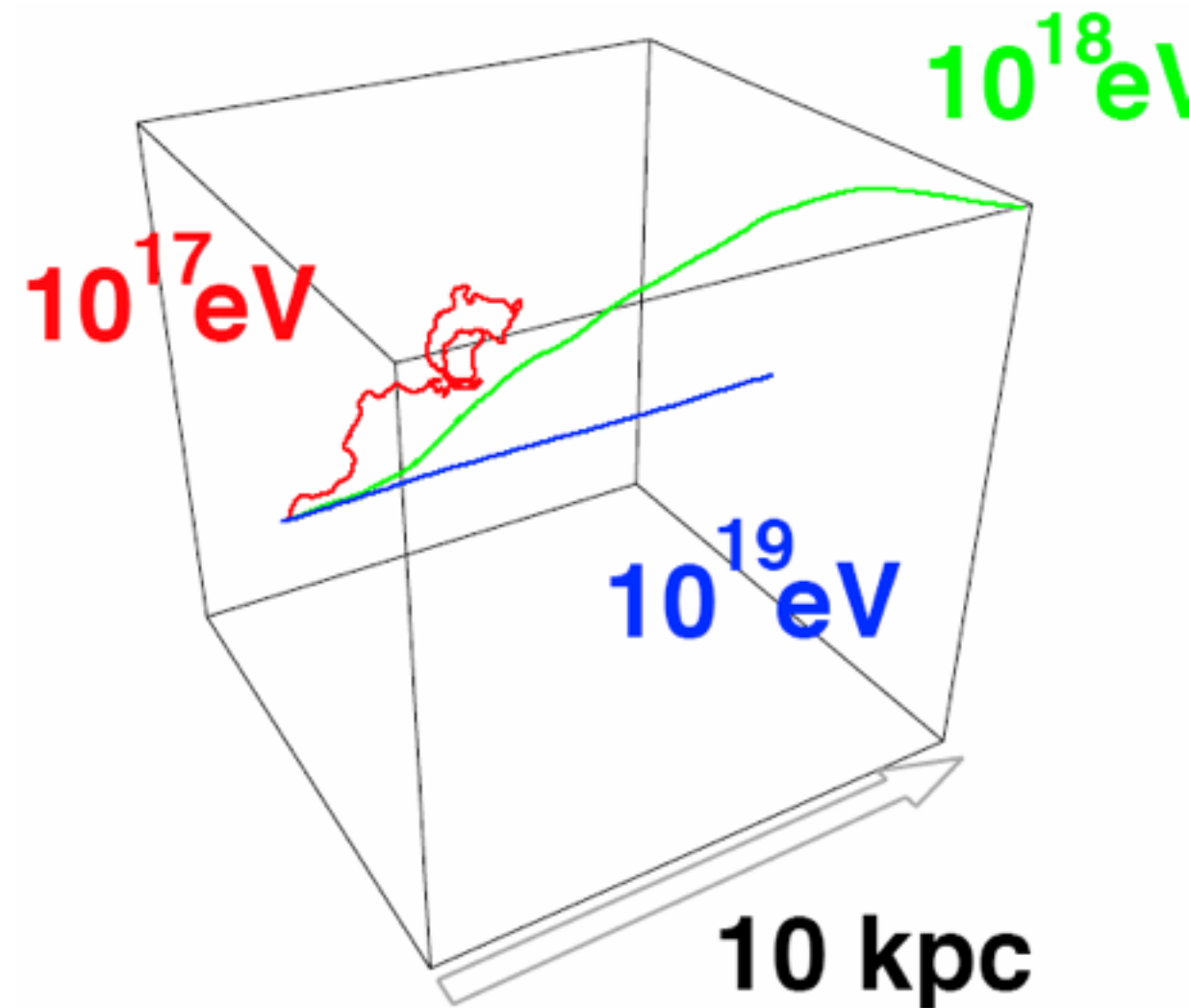
Magnetic Fields in the Galaxy

- Magnetic field in the galaxy along spiral arms, with additional turbulent contributions overlaid
- typical strength ~ 0.1 nT



Propagation of Particles in Magnetic Fields

- Charged particles are deflected by cosmic magnetic fields
- To demonstrate: toy simulation with magnetic fields of ~ 0.1 nT and a coherence length of ~ 100 pc
 - Particles start from the left center with different energies
- Only the very highest energies ($E \sim 10^{19}$ eV) can show the way to their sources - all other particles get substantially deflected and arrive from random directions



Summary

- Cosmic rays are known since 100 years
 - Discovered by Victor Hess on balloon flights
- Acceleration mechanism via scattering on randomly moving cosmic clouds proposed by Fermi 60 years ago (second order Fermi Acceleration)
 - A proof of principle, but insufficient to reach high energies
- Acceleration in shock fronts created by supernovae (first order Fermi Acceleration) can explain energies at least up to $\sim 10^{14}$ eV
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Next Lecture: 08.06., “Cosmic Rays I”, F. Simon

Topics - Overview

13.04.	Einführung / Introduction
20.04.	Achtung - keine Vorlesung! No Lecture!
27.04.	Erdegebundene Beschleuniger / Accelerators
04.05.	Detektoren in der Nicht-Beschleuniger-Physik / Detectors
11.05.	Das Standardmodell / The Standard Model
18.05.	QCD und Jet Physik an Lepton Beschleunigern
25.05.	Pfingsten - Keine Vorlesung! No Lecture
01.06.	Kosmische Beschleuniger / Cosmic Accelerators
08.06.	Kosmische Strahlung I / Cosmic Rays I
15.06.	Kosmische Strahlung II / Cosmic Rays II
22.06.	Präzisionsexperimente (g-2) / Precision Experiments
29.06.	Neutrinos I
06.07.	Neutrinos II
13.07.	Dunkle Materie & Dunkle Energie / Dark Matter & Dark Energy



Backup

Erreichbare Energie

- Die Rate des Energiezuwachses ist gegeben durch die Dauer eines Zyklus und durch den Zuwachs pro Zyklus:

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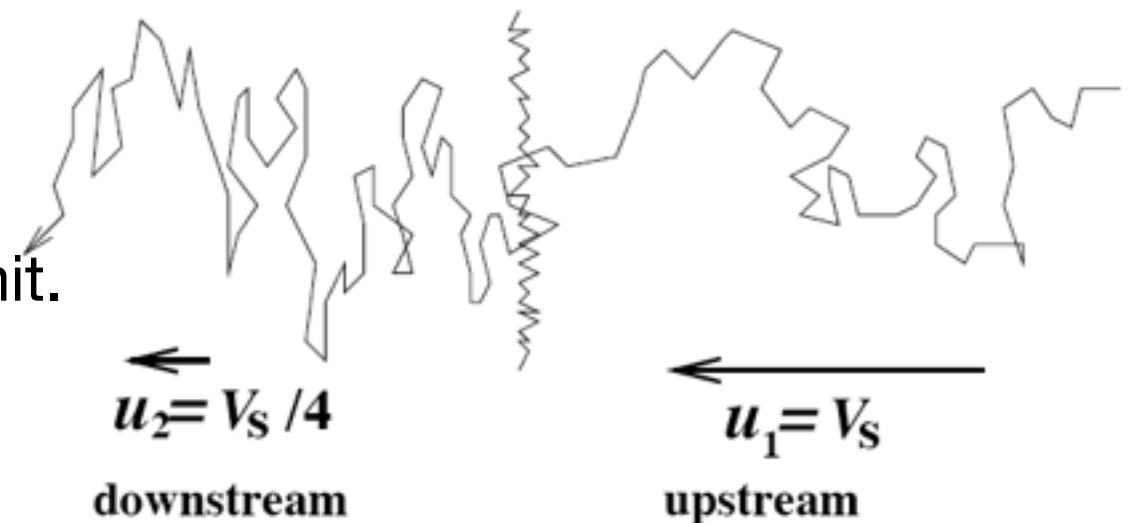
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gerichtete Bewegung: $u_2 t$



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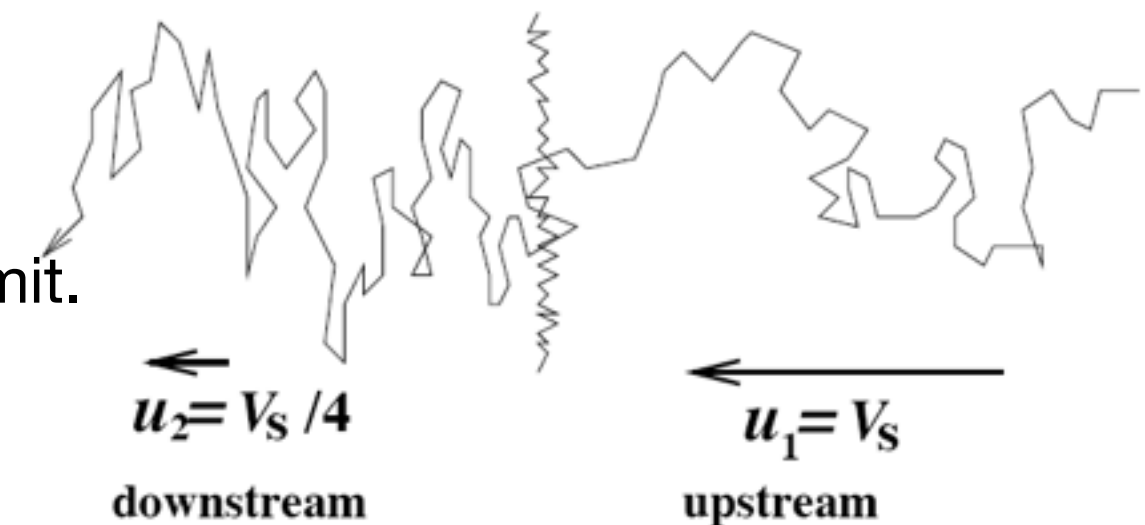
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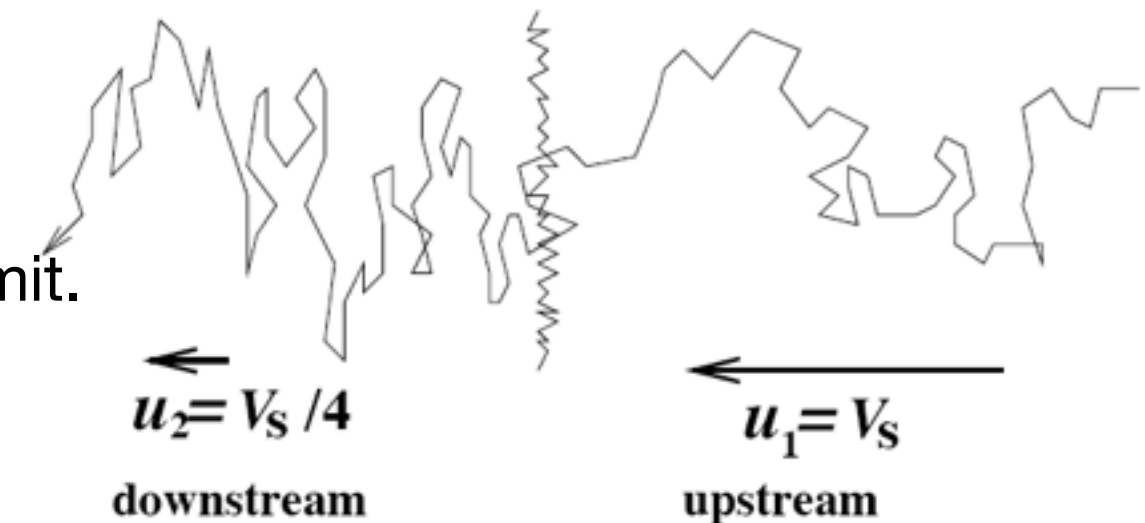
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Hohe Wahrscheinlichkeit, den Schock für immer zu verlassen: $\sqrt{k_2 t} \ll u_2 t$

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- “Grenze”, die entscheidet, ob ein Teilchen “verloren” ist: $\sim k_2/u_2$
- Verweildauer hinter dem Schock aus Teilchendichte und Übergangsrate:

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- Damit ergibt sich die Zyklus-Dauer:

$$t_{cycle} = t_1 + t_2 \approx \frac{4}{c} \left(\frac{k_1}{u_1} + \frac{k_2}{u_2} \right) = \frac{4}{\beta_{Schock} c^2} (k_1 + 4k_2)$$

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- Die Diffusionskonstante hängt vom Magnetfeld ab (Diffusionslänge muss mindestens so groß sein wie der Larmor-Radius, damit Streuung an Magnetfeldänderungen funktioniert) - Bohm-Diffusions-Koeffizient:

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- Erreichbare Energie:

$$E_{max} = \int_0^{t_{acc}} \frac{dE}{dt} dt = \int_0^{t_{acc}} \frac{E \beta_{Schock}}{\frac{20E}{3c^2 \beta_{Schock} ZeB}} dt = \frac{3}{20} \beta_{Schock}^2 c^2 ZeB t_{acc}$$

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$$k = \frac{1}{3} r_L c, \quad r_L = \frac{p}{ZeB} \sim \frac{E}{cZeB}$$

- Es folgt für die Zyklus-Dauer für $k_1 = k_2$

$$t_{cycle} = \frac{20}{3} \frac{E}{c^2 \beta_{Schock} ZeB} \propto E$$

- Erreichbare Energie:

$$E_{max} = \int_0^{t_{acc}} \frac{dE}{dt} dt = \int_0^{t_{acc}} \frac{E \beta_{Schock}}{\frac{20E}{3c^2 \beta_{Schock} ZeB}} dt = \frac{3}{20} \beta_{Schock}^2 c^2 ZeB t_{acc}$$

- Für typische Werte ($\beta_{Schock} \sim 0.03$, $B \sim 0.3$ nT, $t_{acc} \sim 1000$ Jahre)
 $E_{max} \sim 10^{14}$ eV (für Protonen)
► bis zum Knie der Verteilung