Mirror symmetry of Calabi-Yau four-folds with non-trivial cohomology of odd degrees

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Master's Thesis supervised by Thomas Grimm

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Mirror symmetry of CY_4

Introduction

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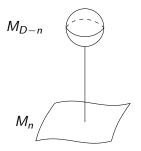
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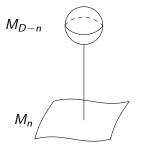
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Of special interest to phenomenology is F-Theory, since it allows insights into strongly-coupled behavior of string-theory. This (in some sense) twelve-dimensional string theory needs to be compactified on a so called (elliptically fibered) Calabi-Yau four-fold. (CY_4)

- Dimensional Reduction
- a Moduli
- Mirror Symmetry of the Torus
- ^(a) Mirror Symmetry on CY_4

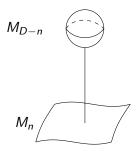
D-dim. gravity theory coupled to matter





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 \Rightarrow *n*-dim. effective theory



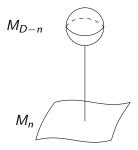
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At every point of the curved spacetime M_n there is a very small deformable internal manifold M_{D-n} whose eigenmodes around a stable ground state correspond to fields on M_n .

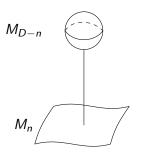
$$\mathsf{mass} \sim \frac{1}{\mathsf{size}~(M_{D-n})}$$

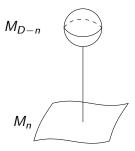
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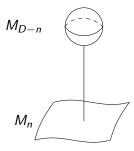




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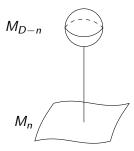


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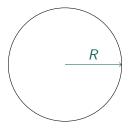
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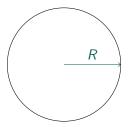
 \Rightarrow Study moduli!

Moduli - Circle



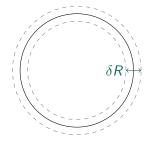
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Moduli - Circle



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Therefore, only the deformations δR are possible.



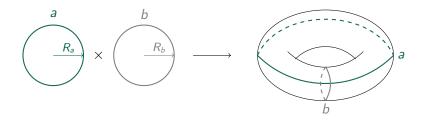
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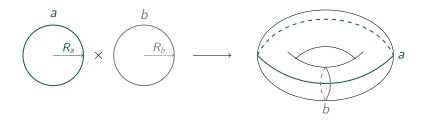
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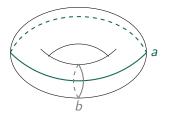


Therefore, we haven now two moduli: R_a, R_b . \Rightarrow Two independent eigenmodes to excite!

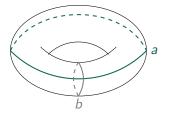
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Mirror symmetry of CY_4

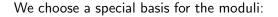
We choose a special basis for the moduli:

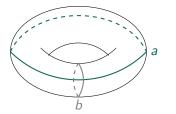


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Volume: $\mathcal{V} = R_a \cdot R_b$



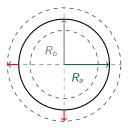


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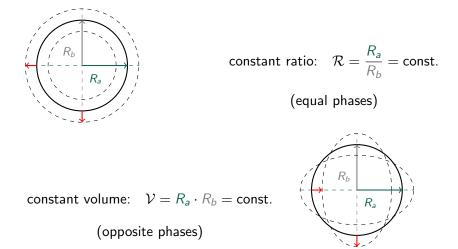
Ratio: $\mathcal{R} = \frac{R_a}{R_b}$

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$$\mathcal{R}=rac{R_a}{R_b}= ext{const.}$$
 (equal phases)

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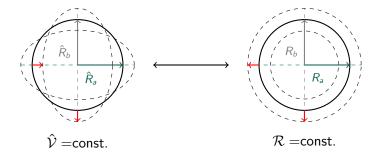
Consider now the torus $\hat{\mathcal{T}}^2$ defined by

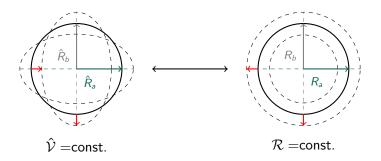
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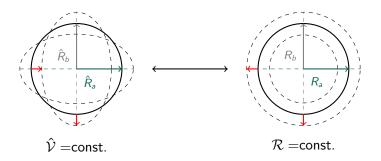
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Therefore, we have the correspondence of eigenmodes:





Due to this correspondence, on both geometries (stable ground states) we have the same content of massless eigenmodes.



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 \Rightarrow We have the same **physics** on both configurations!

This symmetry is called **mirror symmetry** and \hat{T}^2 is the **mirror** of T^2 .

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Mirror symmetry has the advantage that

$$\delta R_b \gg 1, \quad \Rightarrow \quad \delta \hat{R}_b \sim \frac{1}{\delta R_b} \ll 1,$$

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With the same argument, **quantum** effects can be mapped to **classical** effects and are treatable this way.

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Unfortunately, the 8-dimensional case, CY_4 , is hardly studied, beside the analogues of lower dimensional properties.

Mirror Symmetry on CY_4

In our present work, we discuss the so called $h^{2,1}$ -moduli (N'), which where not acessible before, due to the lack of a good base choice for these moduli.

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In particular, we found that the Kähler potential receives corrections for $N' \neq 0$:

$$\mathcal{K} = -\log \mathcal{V} - \log(\int \Omega \wedge \overline{\Omega}) - \operatorname{ReN}_{l}(\int J \wedge \Psi^{l} \wedge \beta^{m}) \operatorname{ReN}_{m}$$

where

$$\Psi' = \frac{1}{2} \operatorname{Re}(f)^{lm} (\alpha_m - i\overline{f}_{mk}\beta^k) \in H^{2,1}(Y_4) \quad [\operatorname{Grimm}]$$

with α_m, β^k a basis of $H^3(Y_4, \mathbb{R})$ and

$$f_{lm} = f_{lm}(z), \quad \frac{df_{lm}}{d\overline{z}_{\overline{K}}} = 0.$$

Mirror symmetry on CY_4

Using mirror symmetry we found that the new correction at large volume/complex structure behaves like

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This implies that the A-model and the B-model of topological string theory on the same CY_4 get **coupled** and can therefore no longer be treated independently.



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On CY_4 **new** moduli arise, that have no lower-dimensional analogue. Mirror symmetry allows us to make first statements about their properties, that are currently not available in the literature. (publication planned)

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Outlook

• Phenomology

- Phenomology
- Geometrical Engineering

- Phenomology
- Geometrical Engineering
- Topological String Theory

- Phenomology
- Geometrical Engineering
- Topological String Theory
- many more!

Thank you for your attention!