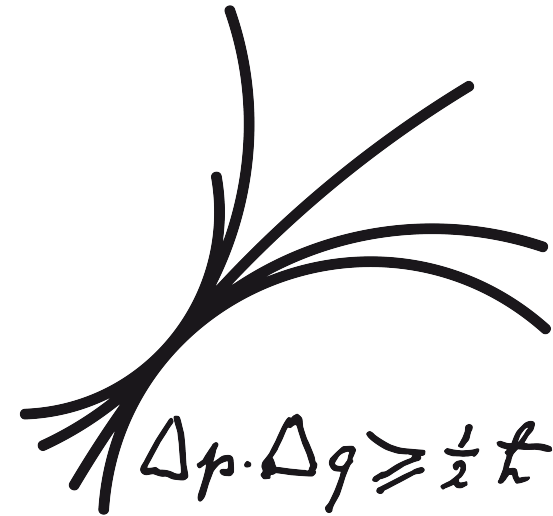


Sterile Neutrinos as Dark Matter



Alexander Merle
MPP Munich
Germany



MAX-PLANCK-GESELLSCHAFT

Based on:

JCAP **1506** (2015) 011, Phys. Lett. **B749** (2015) 283, Int. J. Mod. Phys. **D22** (2013) 1330020, JCAP **1403** (2014) 028, JCAP **1107** (2011) 023, Phys. Rev. **D88** (2013) 113004,...

Prospects in Low Mass DM, MPP Munich, 30-11-2015

Contents:

1. To WIMP or not to WIMP
2. Sterile neutrinos
3. Non-thermal DM production
4. Active-sterile transitions
5. Decay production
6. To line or no to line
7. Conclusions

Contents:

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2. Sterile neutrinos
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4. Active-sterile neutrinos
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6. To line or no to line
7. Conclusions

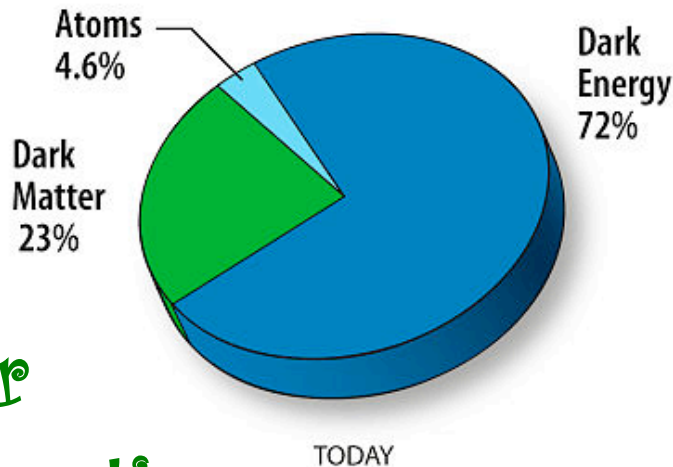
My agenda looks longer than it is!! ;-)

1. To WIMP or not to WIMP

What do we actually know about Dark Matter?!?

KNOWN 😊

- abundance
- rough distribution
- important for structure formation



UNKNOWN ☹️

- identity
- production mechanism
- exact velocity spectrum

THUS: We should be careful not to overlook possibilities just because they are called "non-standard"!!!

1. To WIMP or not to WIMP

Maybe our most natural guess for the identity of DM is a yet unknown elementary particle:

- historically most natural possibility: **WIMP**
 - ↳ because:
 - weak interaction known → “NATURAL”
 - stable WIMPs are predicted in particular by SUSY and by extra dimensions → “THEORY MOTIVATION”
 - comparatively good detection prospects → “EXPERIMENTAL INTEREST”
- **BUT**: unfortunately no clear detection so far...
 - *we should seriously think about alternatives, which may become “standard” at some point*

2. Sterile neutrinos

What is a **sterile neutrino** and why could it be a good Dark Matter candidate?!?

- ordinary (“active”) neutrino ν_a : known elementary particle with very small mass and only weak interactions
- sterile neutrino ν_s : may have a larger mass (value theoretically not predicted) and does not at all participate in standard interactions (**BUT**: small mixing with ν_a)
- **thus**: if produced in the right amounts and with a suitable velocity spectrum, ν_s could act as DM if they are sufficiently stable

NB: NO constraint from oscillations!!!

2. Sterile neutrinos

Indeed, a sterile neutrino with a (typical) mass of a few keV may act as DM, but...

- needs **non-standard production mechanism** (ordinary thermal freeze-out does not work due to tiny coupling)
 - **warm/cold/non-thermal** (interesting for structure and/or galaxy formation)
- typically, this is decaying Dark Matter: $N_1 \rightarrow \nu + \gamma$
 - **monoenergetic X-ray signal** e.g. from galaxies
- strong connection to ordinary neutrinos
 - **concrete models can be tested using light neutrinos**

3. Non-thermal production mechanisms

3. Non-thermal production mechanisms

Dark Matter:

THE STANDARD LORE

Dark Matter is in equilibrium with everything in the early Universe and then “freezes out”.

→ THERMAL SPECTRUM

NON-THERMAL?!?

Non-standard, exotic, unnecessary,...

3. Non-thermal production mechanisms

Dark Matter:

THE STANDARD LORE

Dark Matter

BUT: In reality we don't know the spectrum of Dark Matter very well at all!!

WHY NOT GOING NON-THERMAL?!?

NON-THERMAL?!?

Non-standard, exotic, unnecessary,...

3. Non-thermal production mechanisms

Dark Matter:

What does this mean?!?

THERMAL

$$f(p) = \frac{1}{\exp\left(\frac{\sqrt{p^2+m^2}}{T}\right) \pm 1}$$

HOT

$$T \gg m$$

WARM/COOL

$$T \sim m$$

COLD

$$T \ll m$$

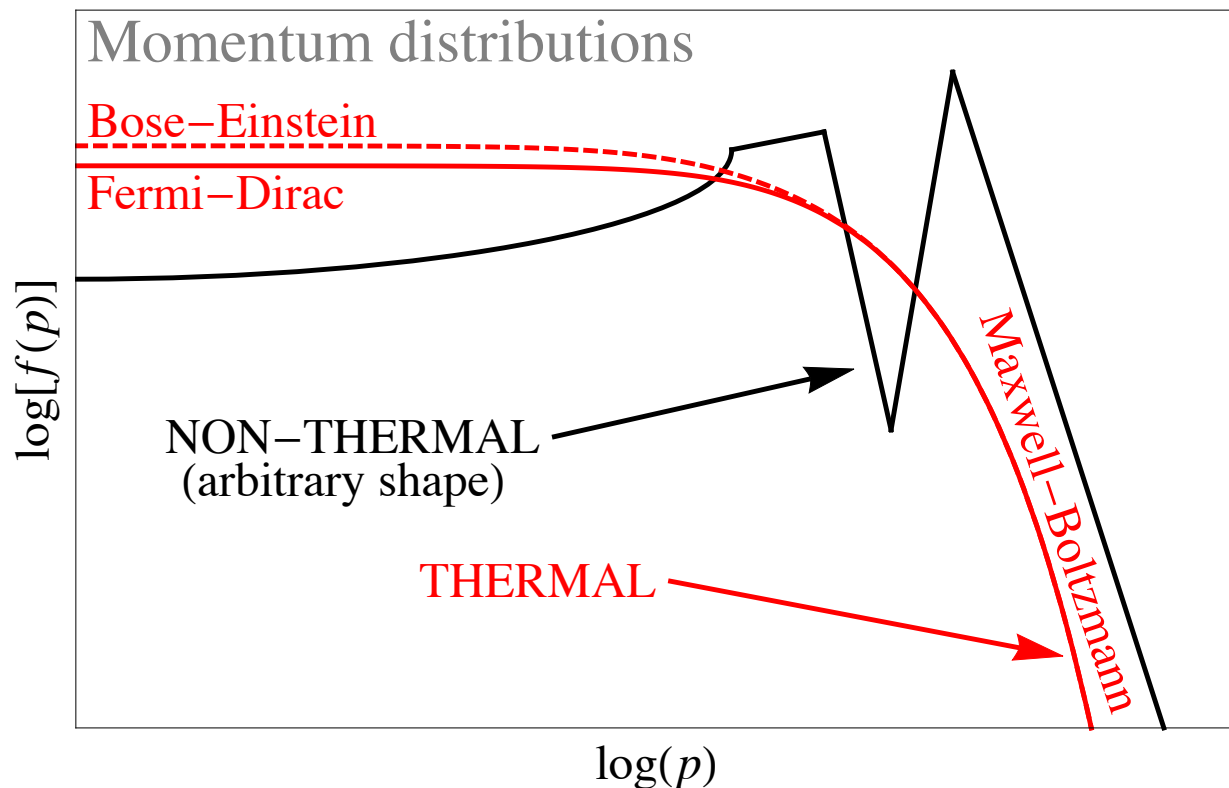
NON-THERMAL

$f(p)$ arbitrary

with:

$$\int_{p=0}^{\infty} p^2 f(p) dp < \infty$$

T not defined!!!



3. Non-thermal production mechanisms

For sterile neutrinos, all but the last mechanism produce a non-thermal shape:

ν_a - ν_s transitions [Dodelson, Widrow: Phys. Rev. Lett. **72** (1994) 17]

Resonant transitions [Shi, Fuller: Phys. Rev. Lett. **82** (1999) 2832], [Canetti *et al.*: Phys. Rev. **D87** (2013) 093006], [Venumadhav *et al.*: 1507.06655], ...

Particle decays [Asaka *et al.*: Phys. Lett. **B638** (2006) 401], [Anisimov *et al.*: Phys. Lett. **B671** (2009) 211], [Bezrukov, Gorbunov: JHEP **1005** (2010) 010], [Kusenko, Petraki: Phys. Rev. **D77** (2008) 065014], [AM, Niro, Schmidt: JCAP **1403** (2013) 028], [Frigerio, Yaguna: Eur. Phys. J. **C75** (2015) 31], [Adulpravitchai, Schmidt: JHEP **1501** (2015) 006], [AM, Totzauer: JCAP **1506** (2015) 011], [AM, Schneider: Phys. Lett. **B749** (2015) 283], [Lello, Boyanovsky: Phys. Rev. **D91** (2015) 063502], [Matsui, Nojiri: Phys. Rev. **D92** (2015) 025045], [Lello, Boyanovsky: 1508.04077], ...

Diluted thermal overproduction [Bezrukov *et al.*: Phys. Rev. **D81** (2010) 085032], [King, AM: JCAP **1208** (2012) 016], [Nemevsek *et al.*: JCAP **1207** (2012) 006], [Patwardhan *et al.*: 1507.01977], ...

3. Non-thermal production mechanisms

For sterile neutrinos, all but the last mechanism produce a non-thermal shape:

PRACTICALLY EXCLUDED!!!

ν_a - ν_s transitions

[Dodelson, Widrow: Phys. Rev. Lett. **72** (1994) 17]

Resonant transitions

[Shi, Fuller: Phys. Rev. Lett. **82** (1999) 2832], [Canetti *et al.*: Phys. Rev. **D87** (2013) 093006], [Venumadhav *et al.*: 1507.06655], ...

Particle decays

[Asaka *et al.*: Phys. Lett. **B671** (2009) 211], [Bezruko *et al.*: Phys. Lett. **B671** (2009) 211], [Bezruko *et al.*: Phys. Lett. **B671** (2009) 211], [Kusenko, Petraki: Phys. Rev. **D77** (2008) 065014], [Frigerio, Yaguna: Eur. Phys. J. **C75** (2015) 31], [Adulpravitchai, Schmidt: JCAP **1403** (2013) 028], [Frigerio, Yaguna: Eur. Phys. J. **C75** (2015) 31], [AM, Totzauer: JCAP **1506** (2015) 011], [AM, Schneider: Phys. Lett. **B749** (2015) 283], [Lello, Boyanovsky: Phys. Rev. **D91** (2015) 063502], [Matsui, Nojiri: Phys. Rev. **D92** (2015) 025045], [Lello, Boyanovsky: 1508.04077], ...

OKAY!!!

Diluted thermal overproduction

[Bezrukov *et al.*: Phys. Rev. **D81** (2010) 085032], [King, AM: JCAP **1208** (2012) 016], [Nemevsek *et al.*: JCAP **1207** (2012) 006], [Patwardhan *et al.*: 1507.01077]

BARELY ALLOWED...

ADVERTISEMENT

Planning Page for the White Paper on keV Sterile Neutrino Dark Matter

This initiative was launched at the NIAPP Workshop "Neutrinos in Astro and Particle Physics" held at TU Munich on July 14-25 2014. Part of this workshop was devoted to review the evidence for and against - keV scale - sterile neutrino as a possible warm/cold Dark Matter candidate. At the cross road of particle physics, astrophysics, and cosmology, the observational constraints, the production mechanisms in the early Universe, and the experimental perspectives have been discussed. As an outcome of the workshop the participants proposed to gather the current status of this field in a white paper.

Outline and Section Editors

Editorial Committee: Marco Drewes, Thierry Lasserre, Alexander Merle, Susanne Mertens

I - Neutrinos in the Standard Model of Particle Physics and Beyond

(Section Editors: Carlo Giunti and André de Gouvea)

1. Current status of Neutrino Masses and Oscillations
2. Open questions in Neutrino Physics
 - 2.1 Neutrino Masses
 - 2.2 Neutrino Nature, Dirac or Majorana
 - 2.3 Neutrino Mass Hierarchy
 - 2.4 Neutrino CP violation
 - 2.5 Additional neutrino states
3. Sterile Neutrinos
 - 3.1 eV-scale
 - 3.2 keV-scale
 - 3.3 GeV, TeV, and \gg TeV scales

II - Neutrinos in The Standard Model of Cosmology

(Section Editors: Julien Lesgourgues and Alessandro Mirizzi)

1. Cosmological Concordance Model (J. Hamann, G....
2. Active neutrinos in Cosmology (J. Lesgourgues, S. ...)
3. Sterile neutrinos in Cosmology
 - 3.1 eV-scale (M. Archidiacono, N. Saviano)
 - 3.2 KeV-scale (A. Boiarskiy, O. Ruchayskiy)
 - 3.3 MeV-scale (S. Pascoli)
 - 3.4 GeV-TeV (A. Ibarra)
 - 3.5 Leptogenesis (P. Di Bari)

The topic starts to be recognised by a larger community... a dedicated White Paper is on the way!

4. Active-Sterile Transitions

Simplest idea:

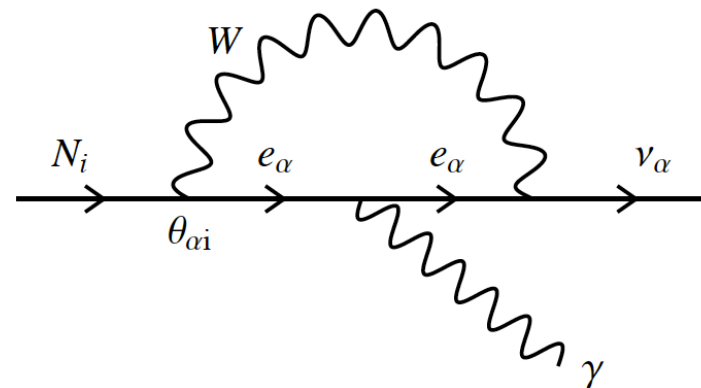
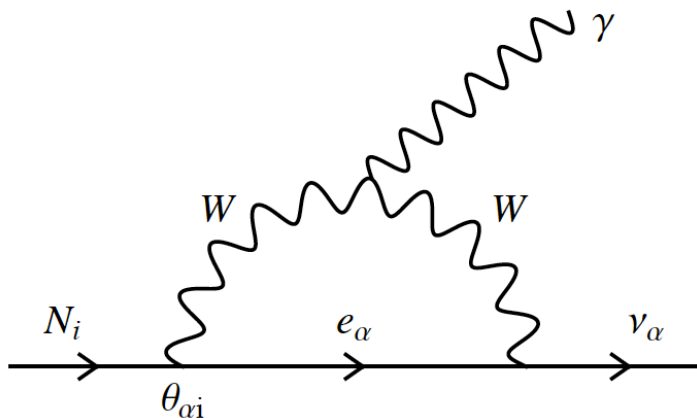
- sterile neutrinos are not entirely sterile
 - *their small mixing with active neutrinos connect them to the Standard Model*
 - **Why not using this to efficiently produce them?!?**
 - *non-resonant active-sterile transitions*
("Dodelson-Widrow")
 - *resonant active-sterile transitions*
("Shi-Fuller")
- ... *these are the "standard" options*

The Dodelson-Widrow mechanism

It could all be sooo simple...

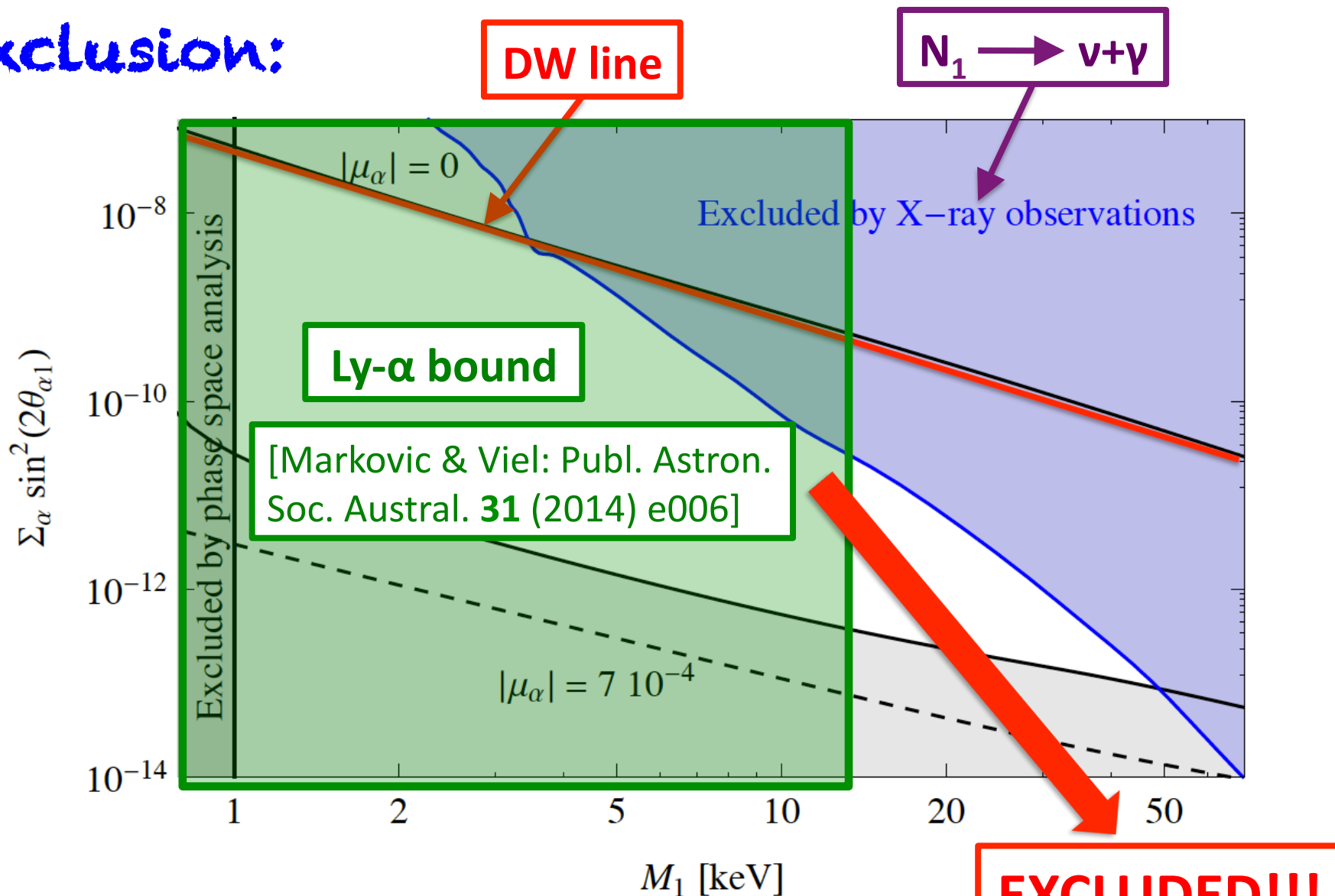
[Dodelson, Widrow: Phys. Rev. Lett. 72 (1994) 17]

- slow non-resonant “oscillations” of active into sterile neutrinos can gradually produce the DM from the thermal plasma (just like “freeze-in”) \rightarrow nice & simple
- this mechanism produces relatively hot DM \rightarrow large mass M_1 needed, BUT decay into X-rays scales like M_1^5 :



The Dodelson-Widrow mechanism

Exclusion:



The Shi-Fuller mechanism

Maybe there is a good way out!

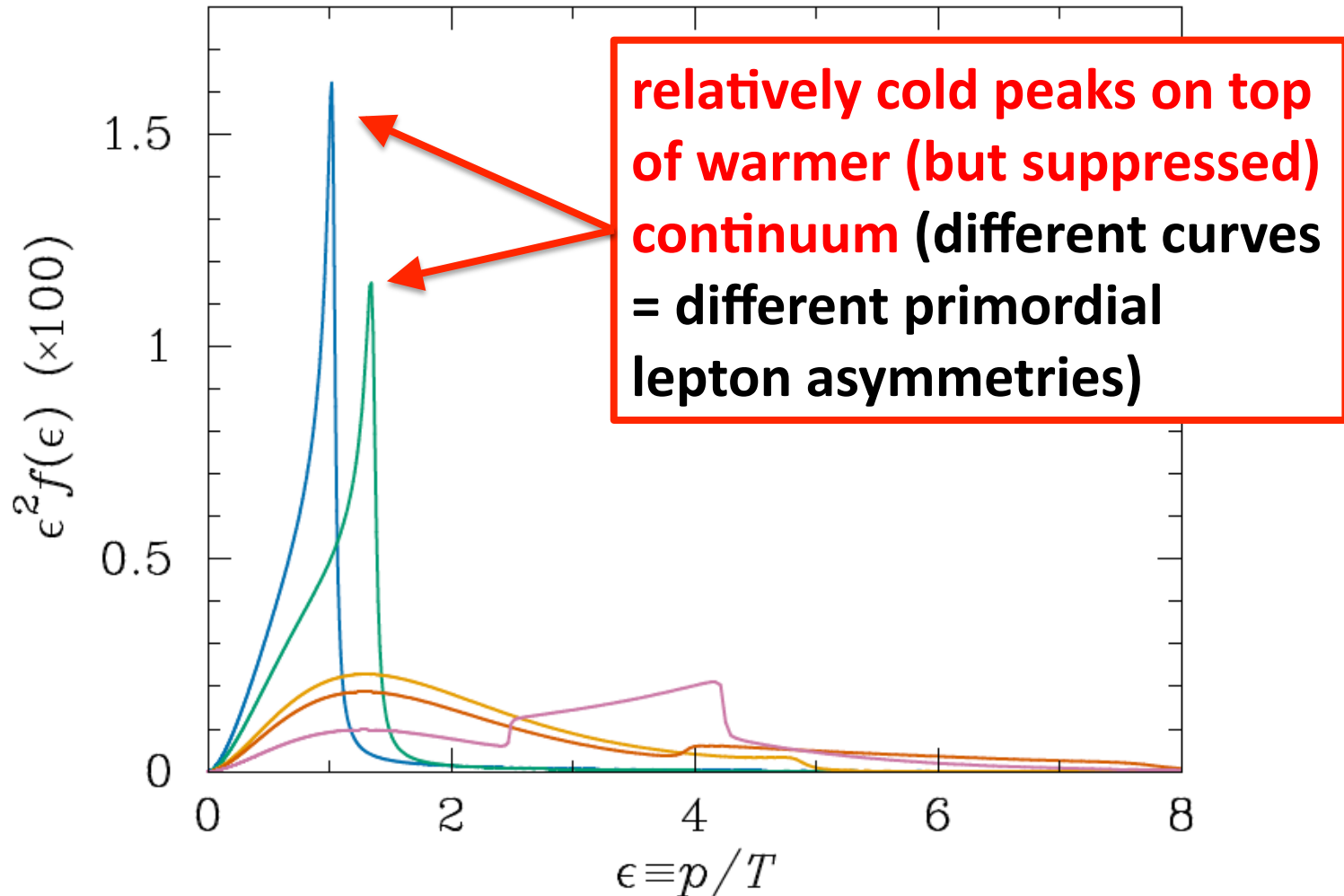
[Shi, Fuller: Phys. Rev. Lett. 82 (1999) 2832]

- just like for ordinary neutrinos in the Sun, active-sterile neutrino **transitions could be resonantly enhanced by a sizeable lepton number asymmetry $|\mu_\alpha|$** present in the early Universe
- this would produce a large amount of ν_s at a specific (momentum-dependent) resonance temperature
→ *cooler spectrum*
- **BUT: the origin of such a primordial lepton number asymmetry is unclear...**

The Shi-Fuller mechanism

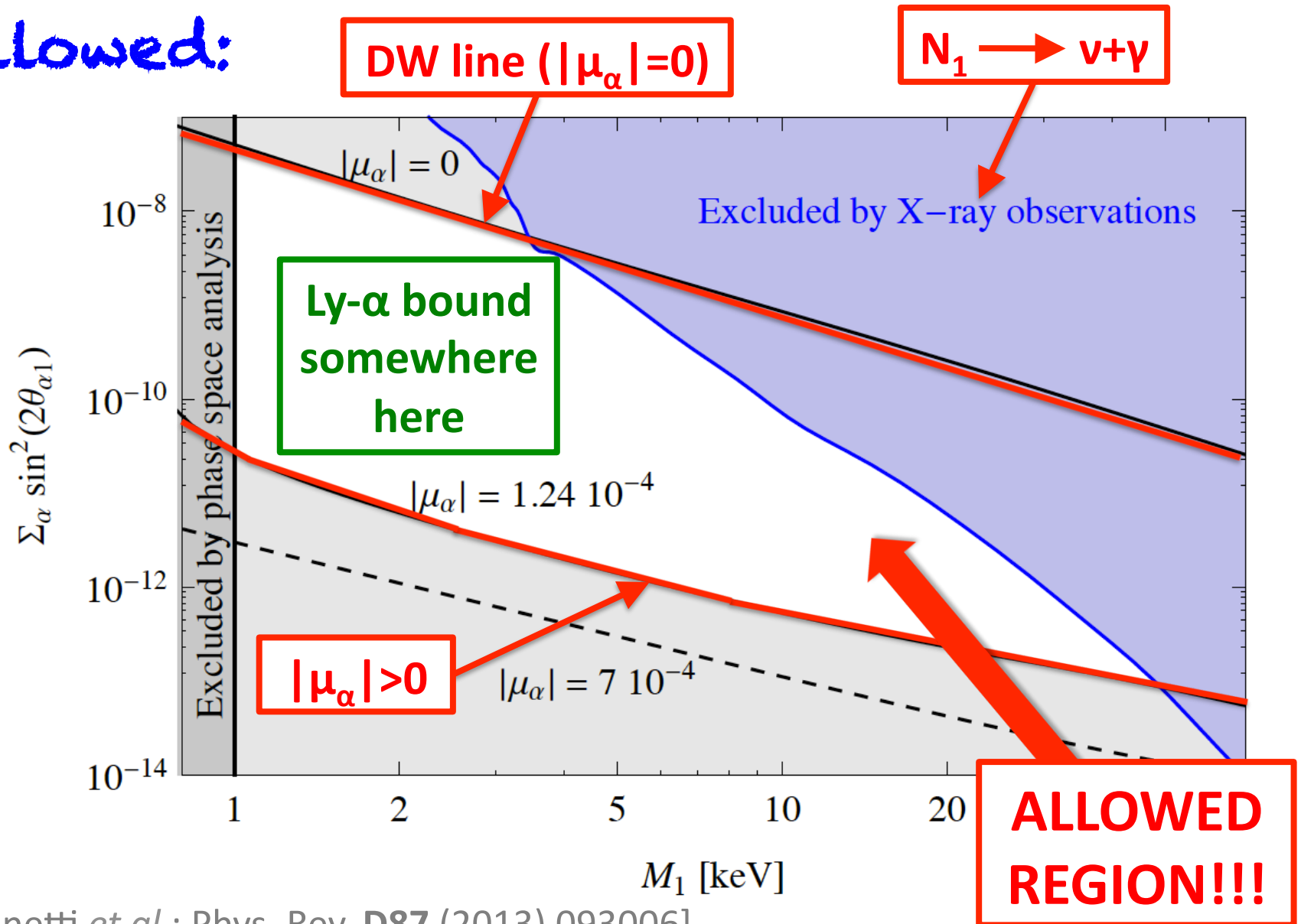
This produces to non-trivial spectra:

[Abazajian: Phys. Rev. Lett. 112 (2014) 161303]



The Shi-Fuller mechanism

Allowed:



5. Decay production

5. Decay production

E.g. scalar decays: e.g. $S \rightarrow N_1 N_1$

- **decaying inflaton**

[Asaka *et al.*: Phys. Lett. **B638** (2006) 401]

[Anisimov *et al.*: Phys. Lett. **B671** (2009) 211]

[Bezrukov, Gorbunov: JHEP **1005** (2010) 010]

- **singlet scalar that freezes out**

[Kusenko: Phys. Rev. Lett. **97** (2006) 241301; Kusenko, Petraki: Phys. Rev. **D77** (2008) 065014]

[Frigerio, Yaguna: Eur. Phys. J. **C75** (2015), 1]

[AM, Schneider: Phys. Lett. **B749** (2015) 283; AM, Totzauer: JCAP **1506** (2015) 011]

- **singlet scalar that freezes in**

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[Adulpravitchai, Schmidt: JHEP **1501** (2015) 006]

[AM, Schneider: Phys. Lett. **B749** (2015) 283; AM, Totzauer: JCAP **1506** (2015) 011]

[Klasen, Yaguna: JCAP **1311** (2013) 039]

- **other particle that decays**

[Lello, Boyanovsky: Phys. Rev. **D91** (2015) 063502]

[Lello, Boyanovsky: 1508.04077]

5. Decay production

- **EXAMPLE: SINGLET SCALAR “S” FREEZES IN OR OUT**

[Kusenko: Phys. Rev. Lett. **97** (2006) 241301; Kusenko, Petraki: Phys. Rev. **D77** (2008) 065014]

[AM, Niro, Schmidt: JCAP **1403** (2013) 028; AM, Totzauer: JCAP **1506** (2015) 011]

Two-step process: scalar S must be produced before decaying \rightarrow many possibilities, e.g.:

5. Decay production

• EXAMPLE: SINGLET SCALAR "S" FREEZES IN OR OUT

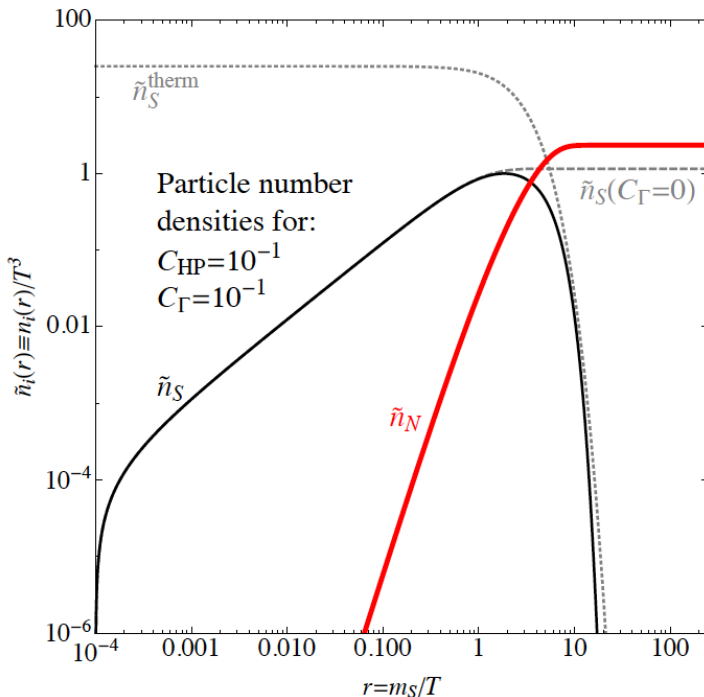
[Kusenko: Phys. Rev. Lett. **97** (2006) 241301; Kusenko, Petraki: Phys. Rev. **D77** (2008) 065014]

[AM, Niro, Schmidt: JCAP **1403** (2013) 028; AM, Totzauer: JCAP **1506** (2015) 011]

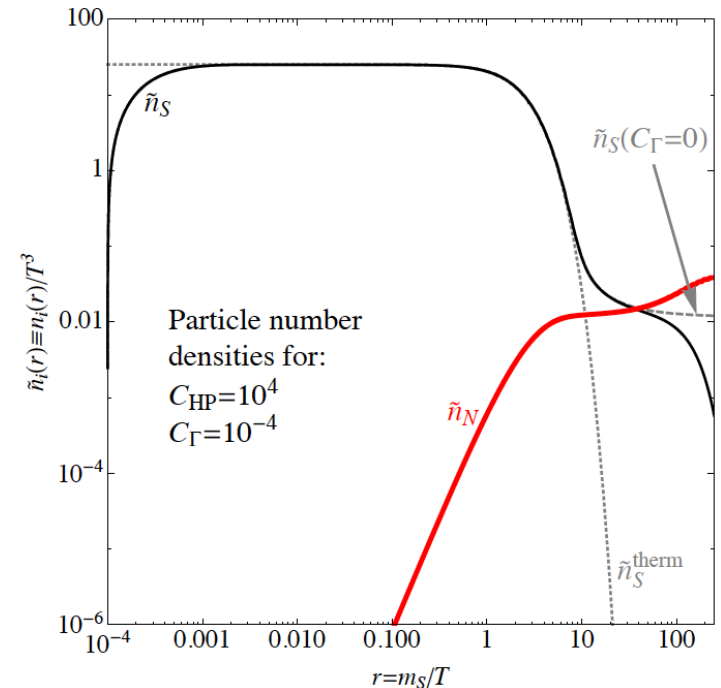
Two-step process: scalar S must be produced before decaying → many possibilities, e.g.:

S freezes-in and decays afterwards

S freezes-out and decays both in and out of equilibrium



Abundance



5. Decay production

• EXAMPLE: SINGLET SCALAR “S” FREEZES IN OR OUT

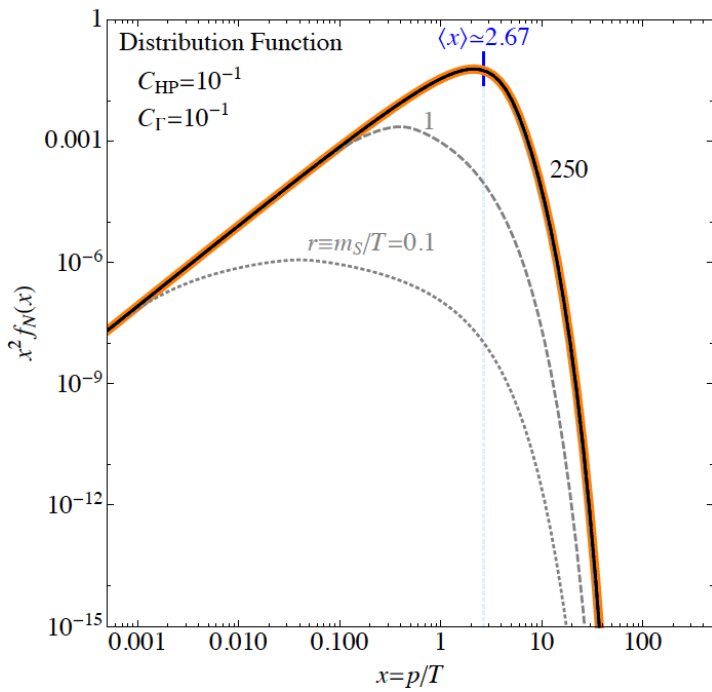
[Kusenko: Phys. Rev. Lett. **97** (2006) 241301; Kusenko, Petraki: Phys. Rev. **D77** (2008) 065014]

[AM, Niro, Schmidt: JCAP **1403** (2013) 028; AM, Totzauer: JCAP **1506** (2015) 011]

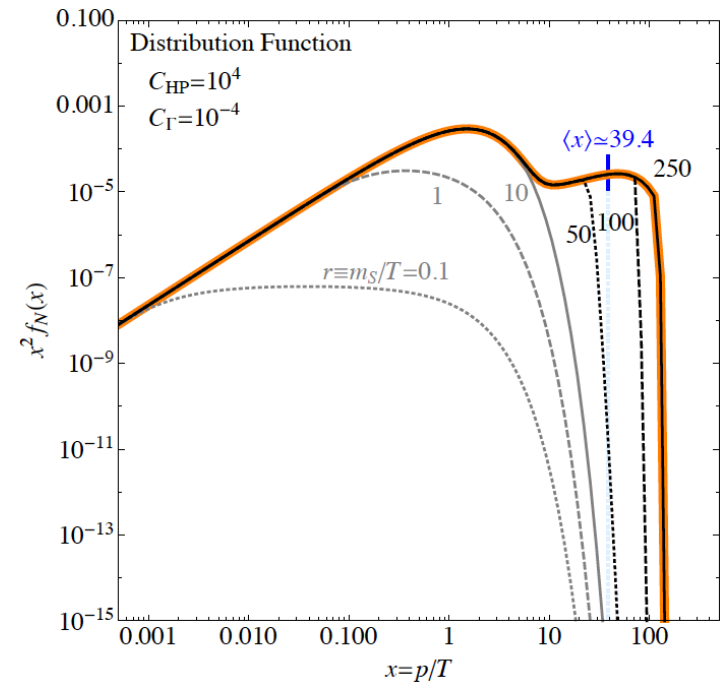
Two-step process: scalar S must be produced before decaying → many possibilities, e.g.:

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S freezes-out and decays both in and out of equilibrium



Distribution function



5. Decay production

• EXAMPLE: SINGLET SCALAR "S" FREEZES IN OR OUT

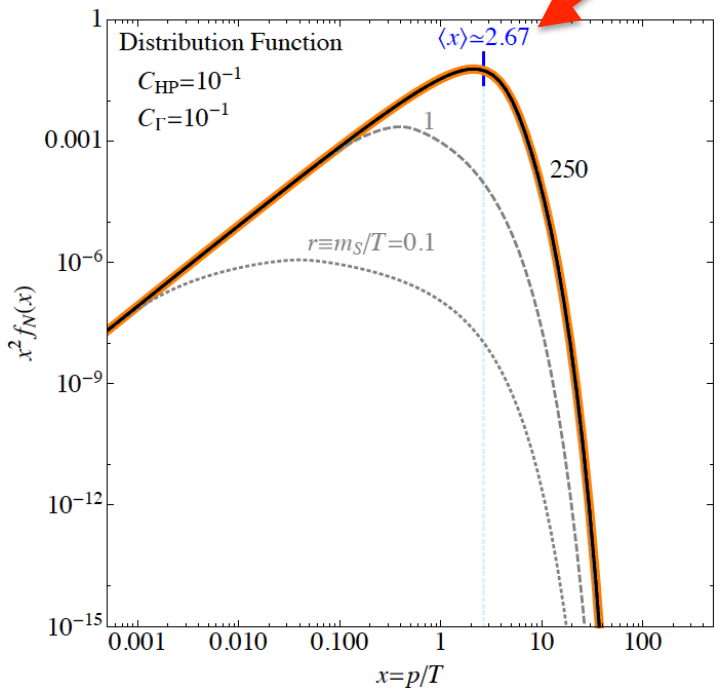
[Kusenko: Phys. Rev. Lett. **97** (2006) 241301; Kusenko, Petraki: Phys. Rev. **D77** (2008) 065014]

[AM, Niro, Schmidt: JCAP **1403** (2013) 028; AM, Totzauer: JCAP **1506** (2015) 011]

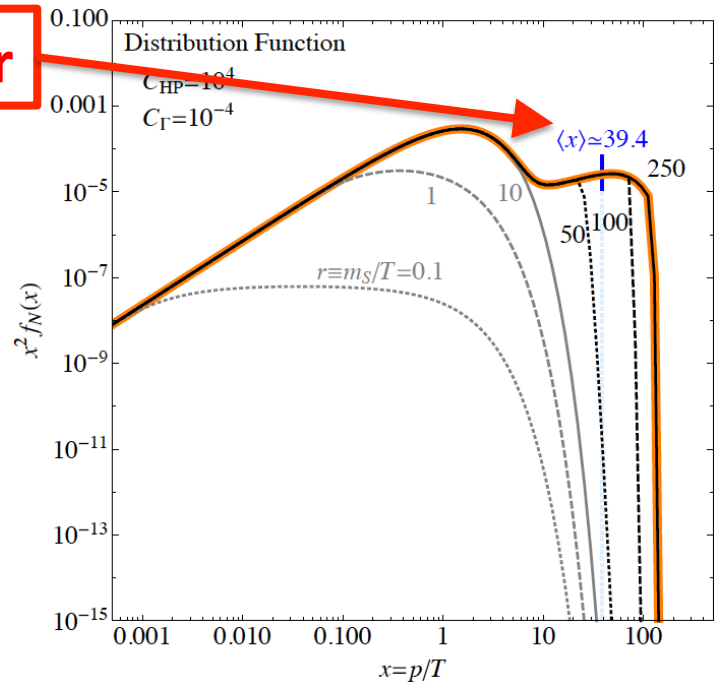
Two-step process: scalar S must be produced before decaying → many possibilities, e.g.:

S freezes-in and decays afterwards **colder**

S freezes-out and decays both in and out of equilibrium **warmer**



Distribution function



5. Decay production

• EXAMPLE: SINGLET SCALAR “S” FREEZES IN OR OUT

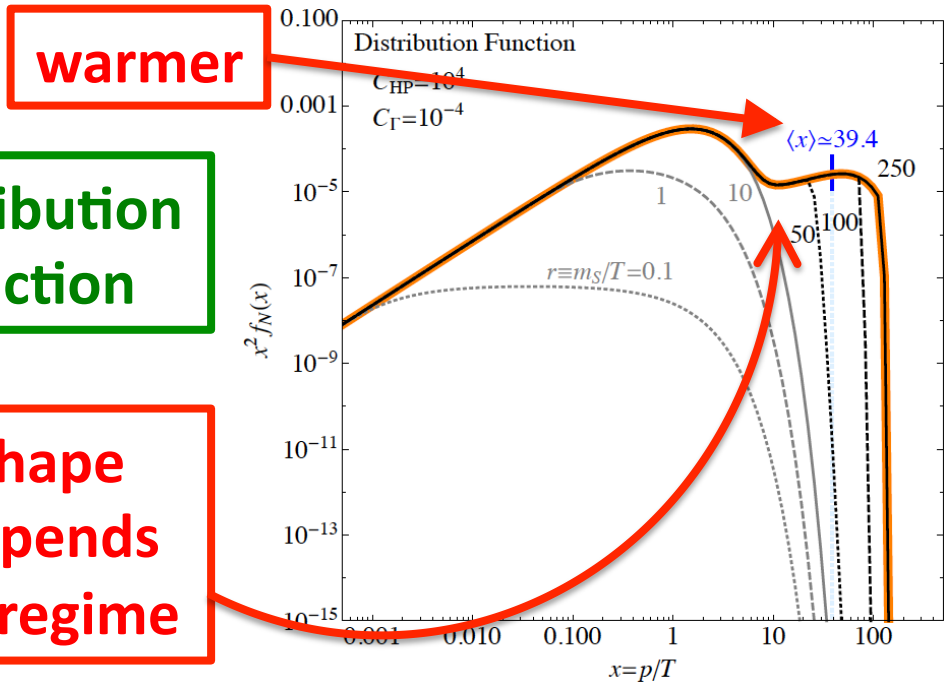
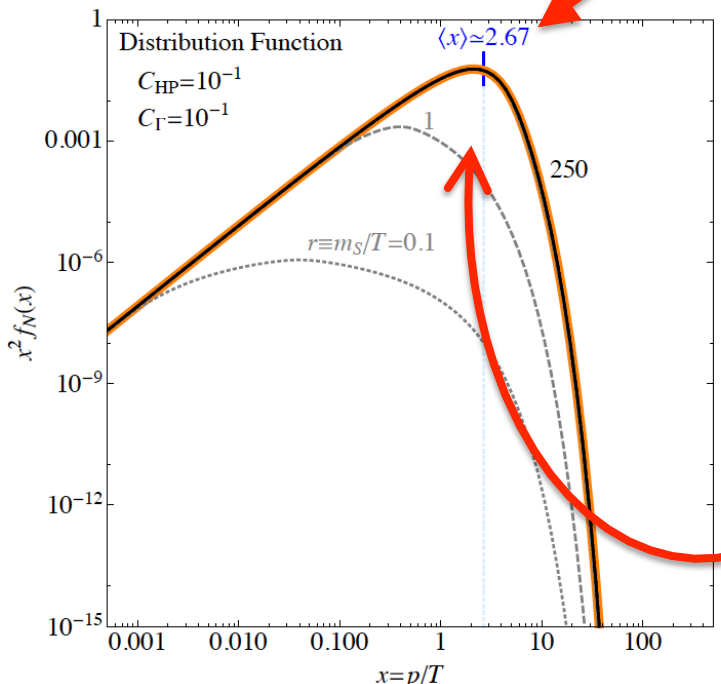
[Kusenko: Phys. Rev. Lett. **97** (2006) 241301; Kusenko, Petraki: Phys. Rev. **D77** (2008) 065014]

[AM, Niro, Schmidt: JCAP **1403** (2013) 028; AM, Totzauer: JCAP **1506** (2015) 011]

Two-step process: scalar S must be produced before decaying → many possibilities, e.g.:

S freezes-in and decays afterwards **colder**

S freezes-out and decays both in and out of equilibrium



Distribution function

shape depends on regime

warmer

5. Decay production

• EXAMPLE: SINGLET SCALAR “S” FREEZES IN OR OUT

[Kusenko: Phys. Rev. Lett. 97 (2006) 241301; Kusenko, Petraki: Phys. Rev. D77 (2008) 065014]

[AM, Niro, Schmidt: JCAP 1403 (2013) 028; AM, Totzauer: JCAP 1506 (2015) 011]

Two-step process: scalar S must be produced before decaying → many possibilities, e.g.:

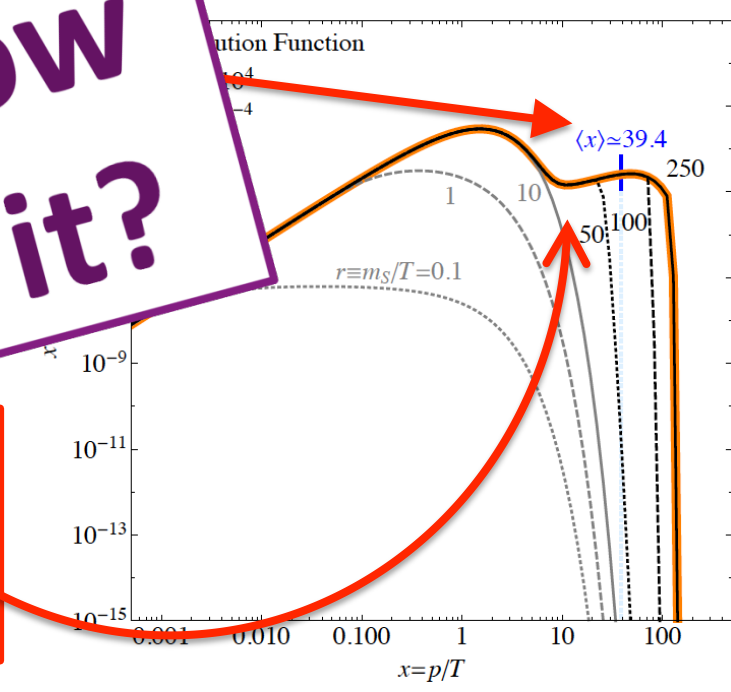
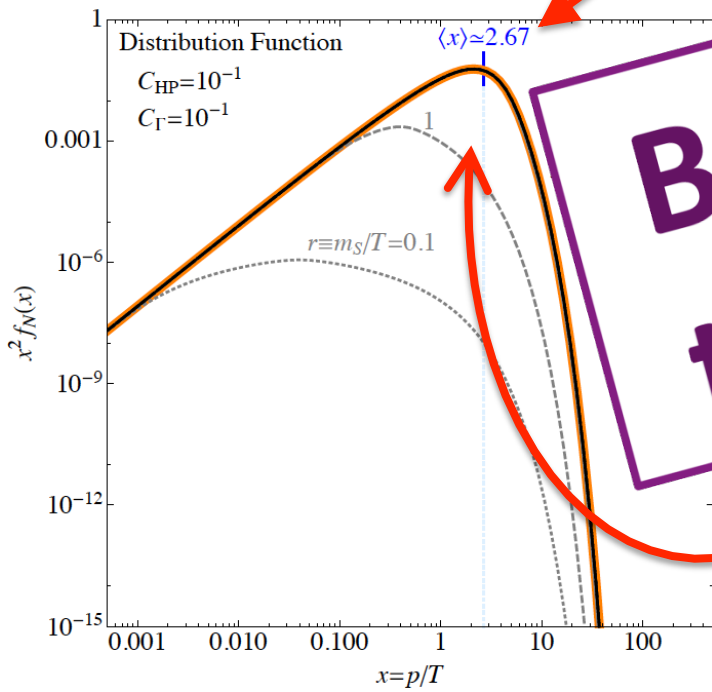
S freezes-in and decays afterwards

colder

S freezes-out and decays both in and out of equilibrium

BUT: HOW to test it?

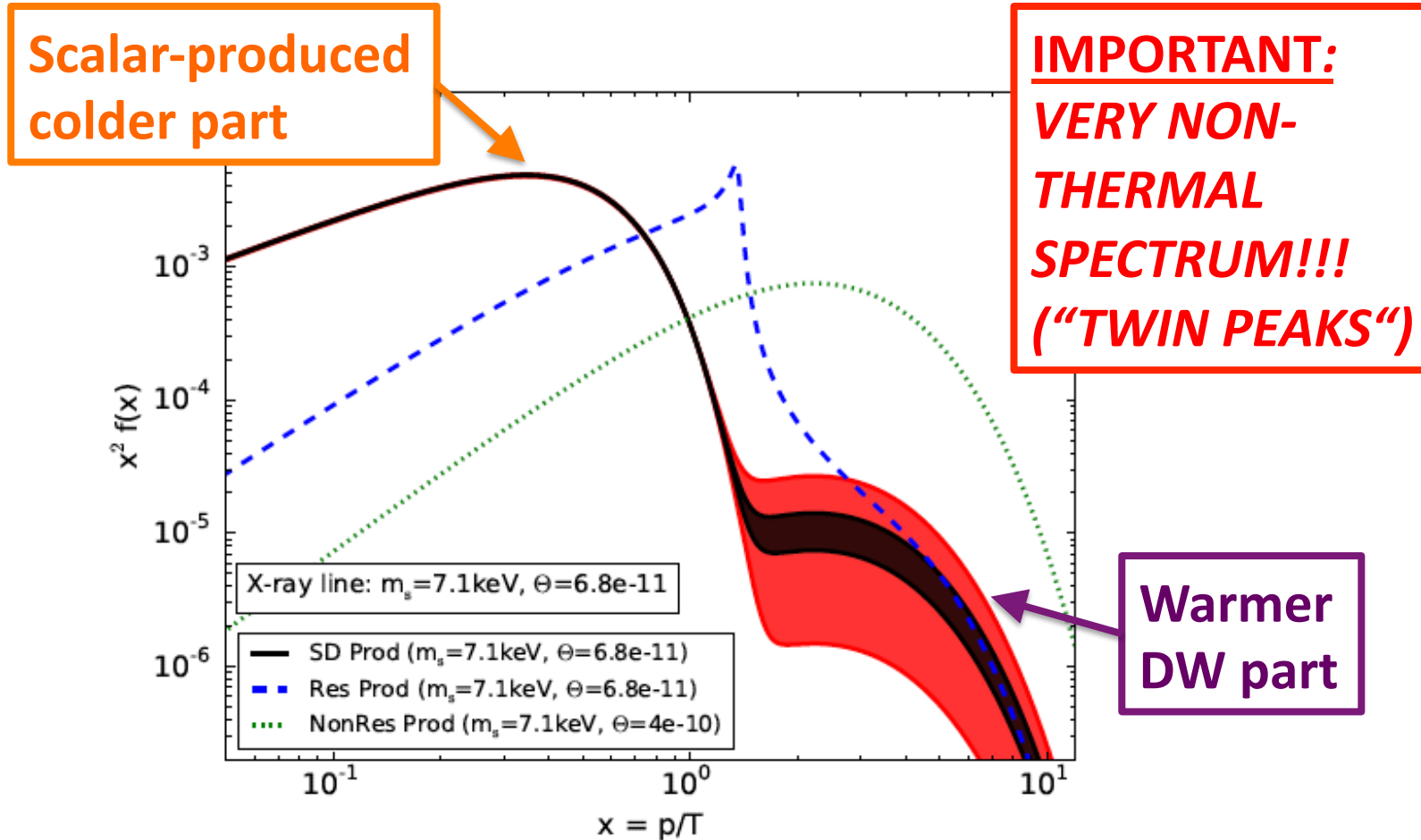
shape depends on regime



5. Decay production

- Cold Dark Matter is not perfect: **NON-THERMAL DM CAN BE VERY INTERESTING FOR STRUCTURE FORMATION**

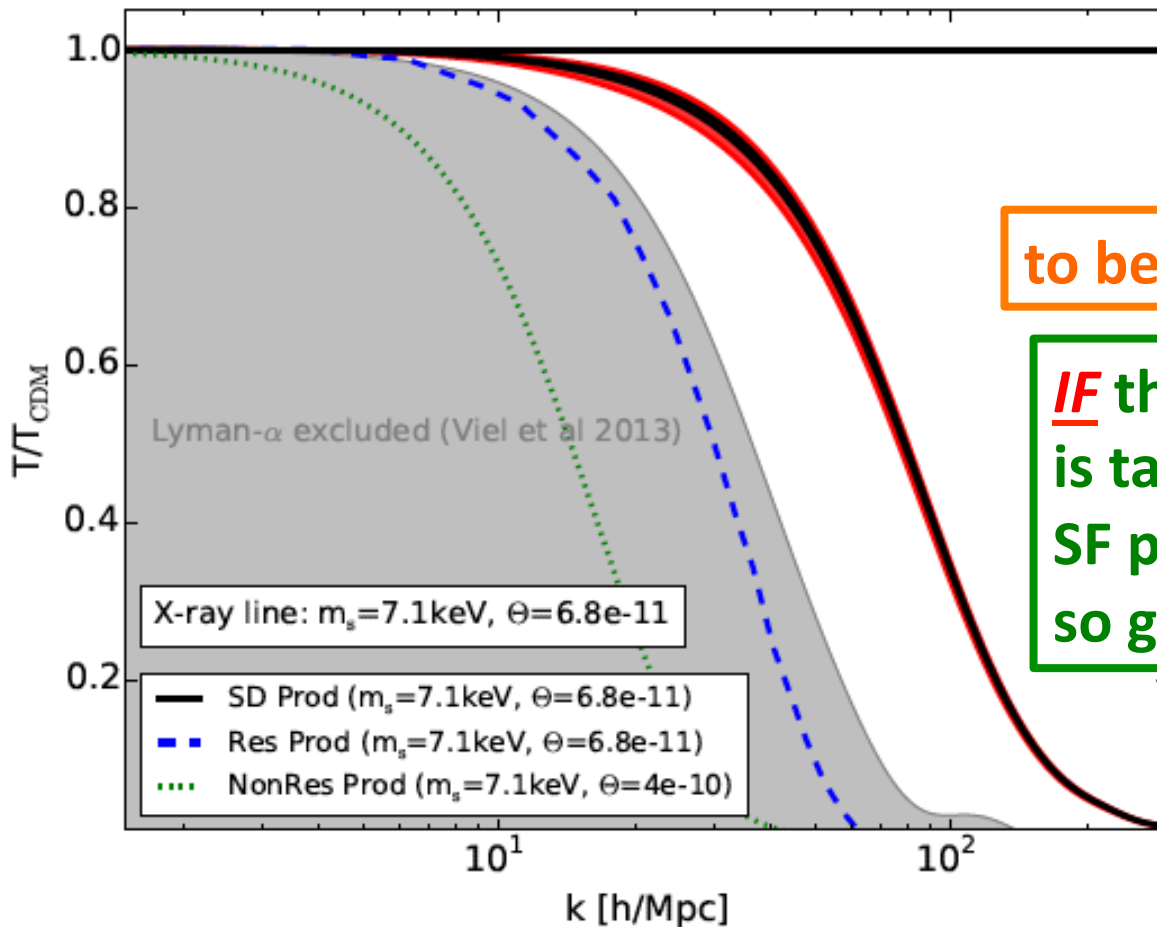
[AM, Schneider: Phys. Lett. B749 (2015) 283]



5. Decay production

- Cold Dark Matter is not perfect: **NON-THERMAL DM CAN BE VERY INTERESTING FOR STRUCTURE FORMATION**

[AM, Schneider: Phys. Lett. B749 (2015) 283]



to be discussed in a minute

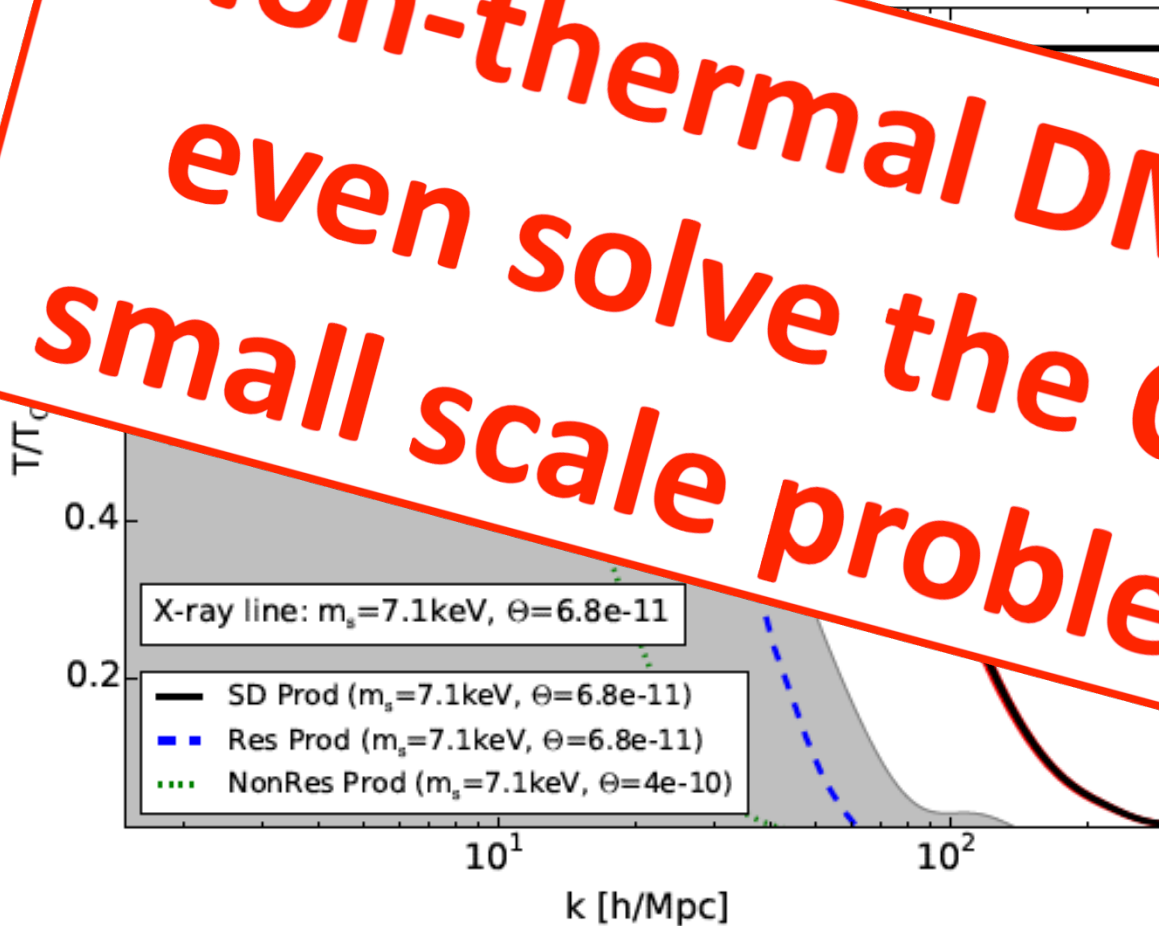
IF the 3.5 keV line signal is taken seriously, then SF production is not in so good shape anymore

5. Decay production

- Cold Dark Matter is not perfect: **NON-THERMAL DM CAN BE VERY INTERESTING FOR STRUCTURE FORMATION**

...ider: Phys. Lett. B749 (2015) 283]

Non-thermal DM may even solve the CDM small scale problems!!



al
n
in
more

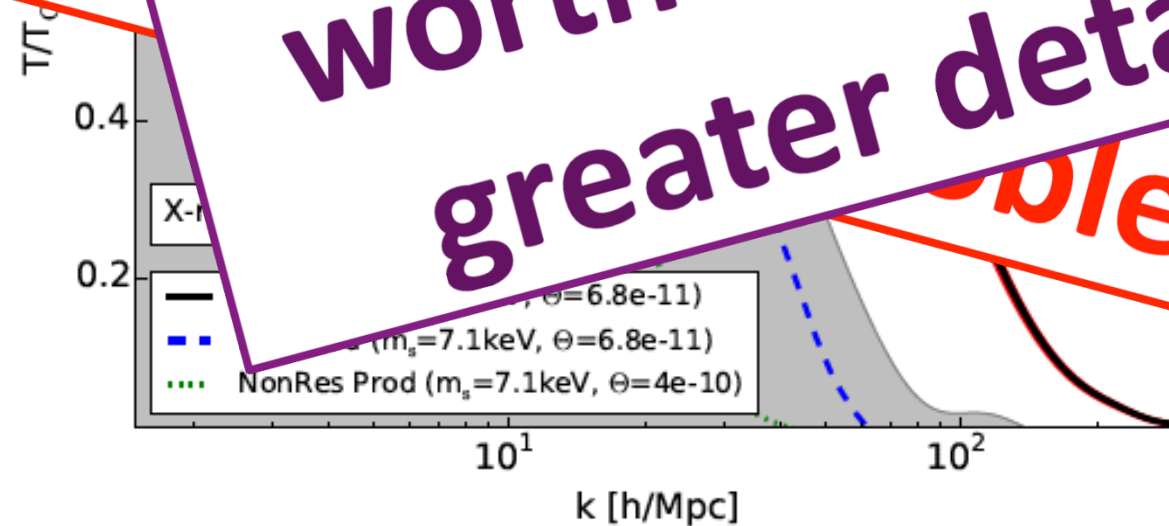
5. Decay production

- Cold Dark Matter is not perfect: **NON-THERMAL DM CAN BE VERY INTERESTING FOR STRUCTURE FORMATION**

Journal: Phys. Lett. B749 (2015) 2831

Non-th

Decay production
worth to study in
greater detail!



problems!!

al
n
in
more

5. Decay production

• EXAMPLE: SINGLET SCALAR “S” FREEZES IN OR OUT

[Kusenko: Phys. Rev. Lett. **97** (2006) 241301; Kusenko, Petraki: Phys. Rev. **D77** (2008) 065014]

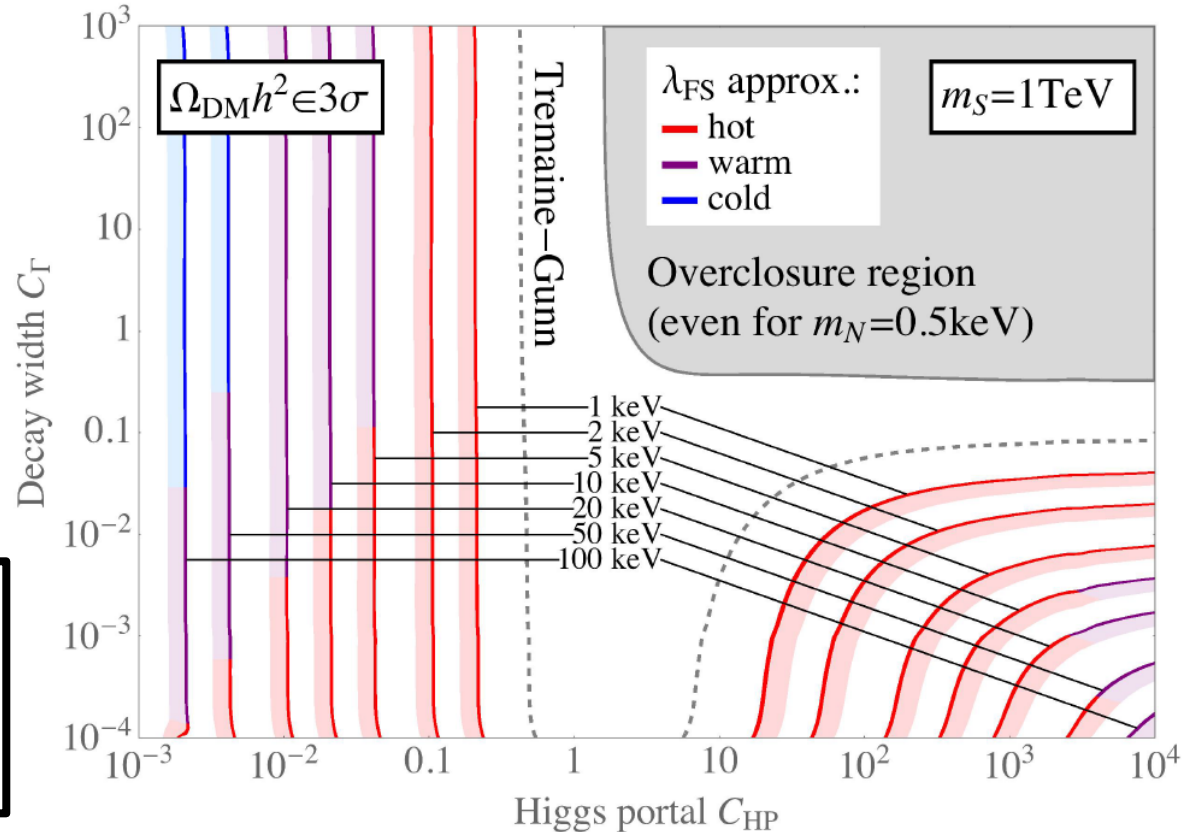
[AM, Niro, Schmidt: JCAP **1403** (2013) 028; AM, Totzauer: JCAP **1506** (2015) 011]

Free-streaming horizon

$$r_{\text{FS}} = \int_{t_{\text{in}}}^{t_0} \frac{\langle v(t) \rangle}{a(t)} dt$$



Decides about whether the keV sterile neutrinos are **HOT**, **WARM**, or **COLD**



→ ALL possible depending on effective parameters (C_Γ, C_{HP})

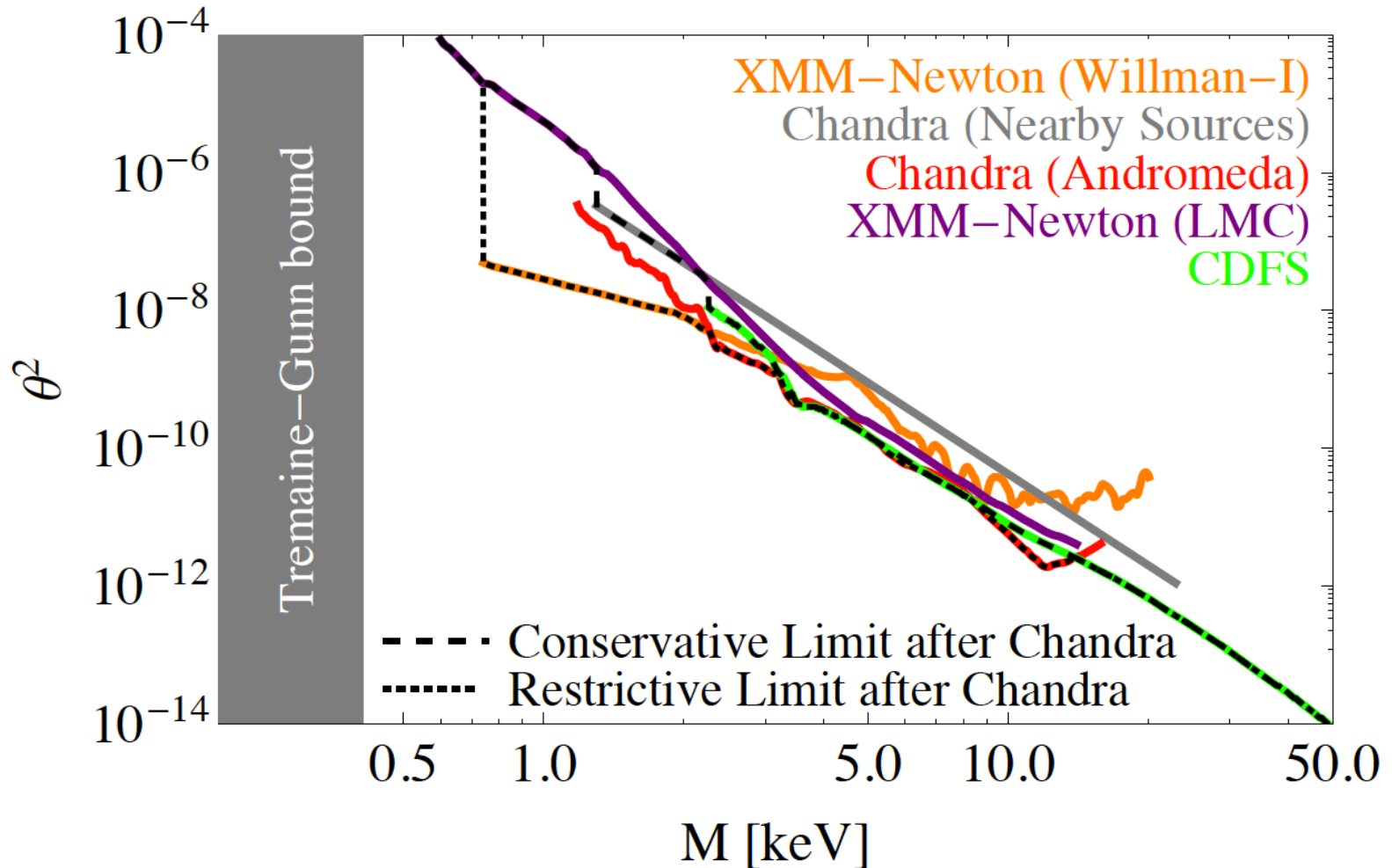
→ **STRUCTURE FORMATION STUDY OUT SOON!**

[AM, Schneider, Totzauer: 1512.XXXXX]

6. *To Line Or Not To Line*

6. To Line Or Not To Line

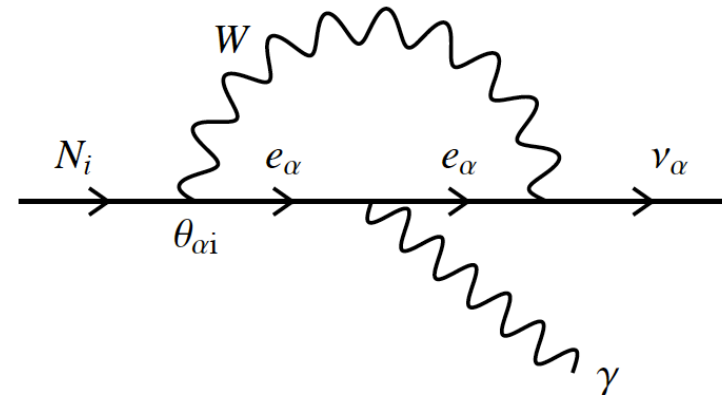
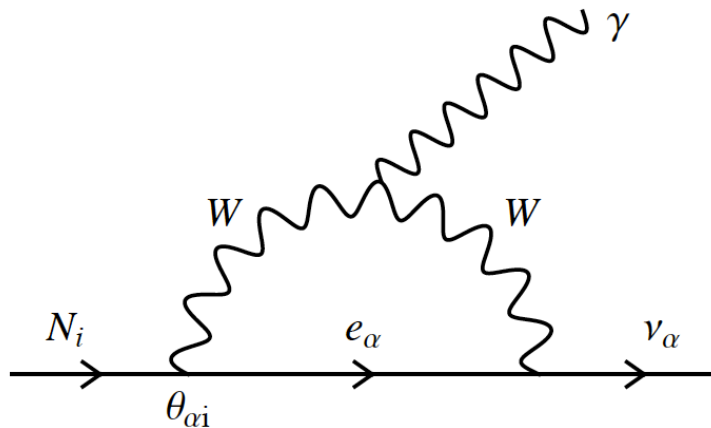
Up to February 2014 we only had bounds:



6. To Line Or Not To Line

Do we now have a signal?!?

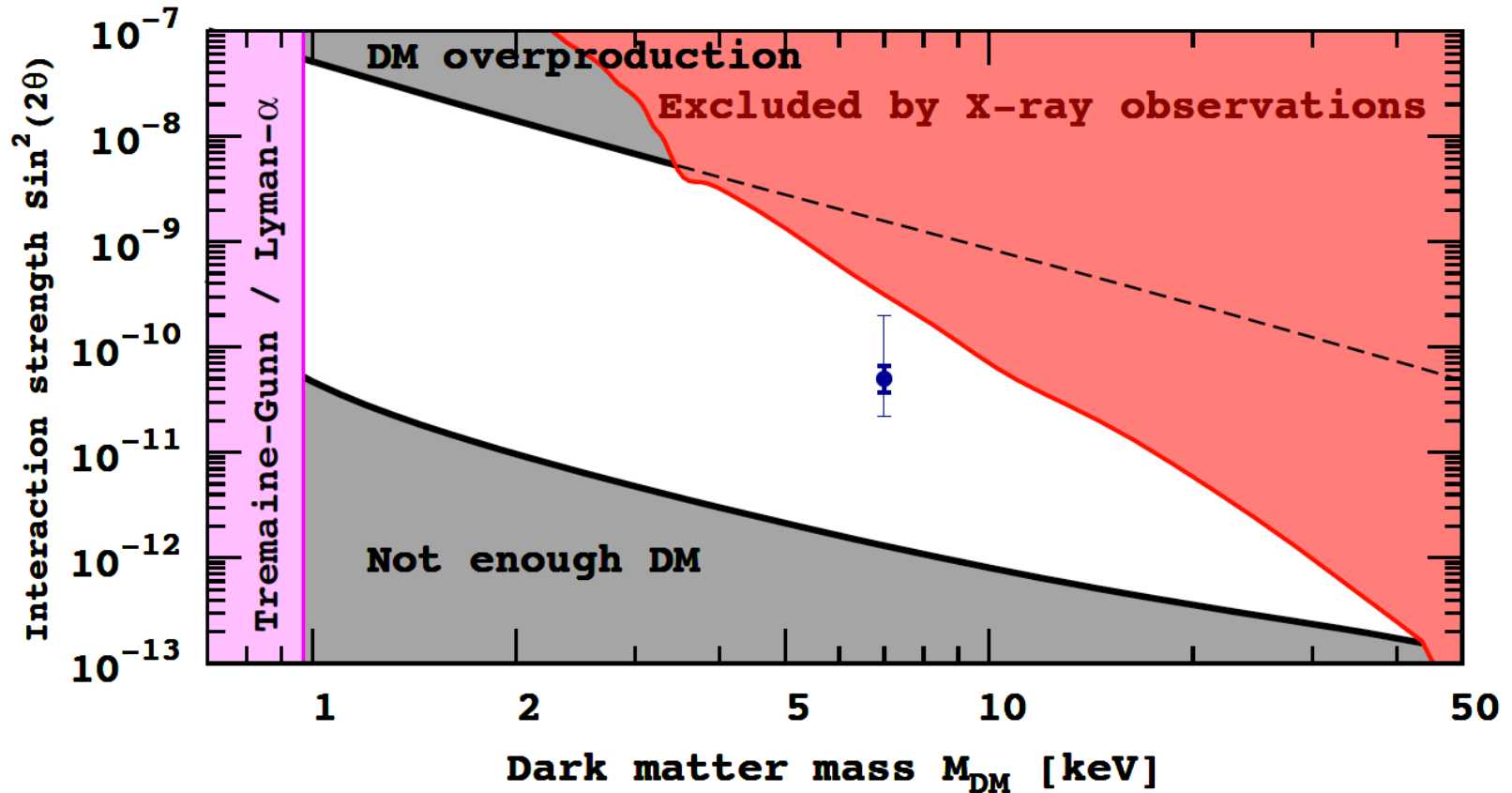
- a possible hint: the “3.5” keV line (**ACTUALLY**: 3.6 keV, at least according to the first paper)
 - stacked XMM-Newton spectrum of 73 galaxy clusters + division into 3 subsamples (e.g. Perseus cluster only):
 $m_s \approx 7.1 \text{ keV}$ & $\sin^2(2\theta) \approx 7 \times 10^{-11}$ [Bulbul *et al.*: *Astrophys. J.* **789** (2014) 13]
 - XMM-Newton spectrum of Perseus cluster & Andromeda galaxy:
 $m_s = 7.06 \pm 0.05 \text{ keV}$ & $\sin^2(2\theta) = (2.2-20) \times 10^{-11}$ [Boyarsky *et al.*: *Phys. Rev. Lett.* **113** (2014) 251301]



6. To Line Or Not To Line

Do we now have a signal?!?

- resulting region:



[Boyarisky *et al.*: 1402.4119]

6. *To Line Or Not To Line*

BUT: Not everybody agrees...

... Little "war" on the arXiv:

Main arguments on the "con"-side:

- **signal not seen in all data sets** [e.g. Jeltema, Profumo: Mon. Not. Roy. Astron. Soc. 450 (2015) 2143; Malyshev, Neronov, Eckert: Phys. Rev. D90 (2014) 103506; Riemer-Sorensen: 1405.7943]
- **wrong use of atomic data** [e.g. Jeltema, Profumo: Mon. Not. Roy. Astron. Soc. 450 (2015) 2143]
- **signal does not correctly follow DW-profile** [e.g. Carlson, Jeltema, Profumo: JCAP 1502 (2015) 009]

ON THE OTHER HAND:

COUNTER-ARGUMENTS EXIST FOR ALL OF THEM...

6. To Line Or Not To Line

My personal take on this:

THE STATISTICS OF SIGNAL IS BAD!!!

→ CLEAR:

will depend on analysis method

weak position from the very beginning

→ ALSO CLEAR:

no need to argue at this stage...

... let's just wait for Astro-H or LOFT
(beginning of 2016 is not too far)!!!

7. Conclusions

7. Conclusions

- **Dark Matter...** may be very different from WIMPs, only time will tell!
- **Production Mechanisms for sterile neutrinos...** are non-standard and need to be studied!
- ***back-up slides available:***

TERRESTRIAL EXPERIMENTS

MODEL BUILDING

DETAILS ON DM PRODUCTION



**THANK
YOU!!!**

ADVERTISEMENT



Workshop on keV sterile neutrino Dark Matter:

- at IAS Garching
- 7 – 9 December 2015

v-Dark

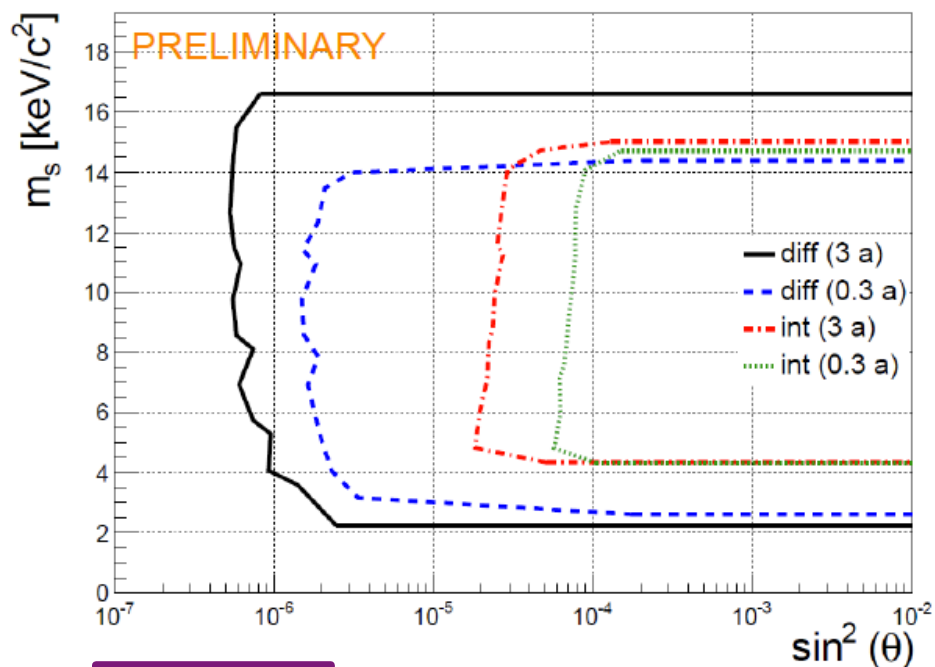


**BACK-UP
SLIDES**

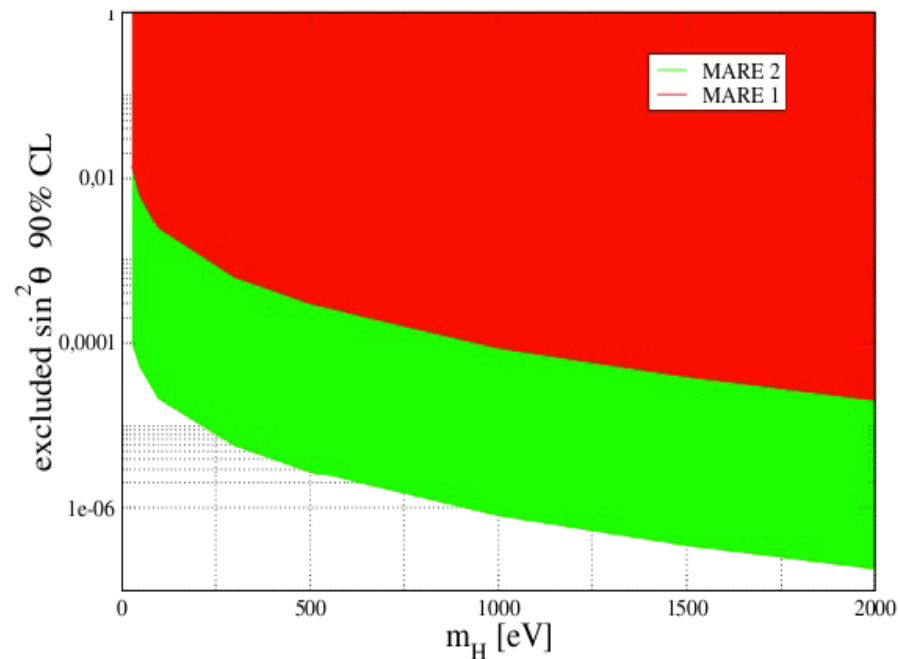
8. BACK-UP: Terrestrial Experiments

There are several Earth-based experiments attempting to measure something...

- β decays: can produce ν_s if mass $<$ Q-value



KATRIN [Korzeczek: Meudon `14]



MARE [Ferri: Meudon `13]

→ upgrade to TRISTAN

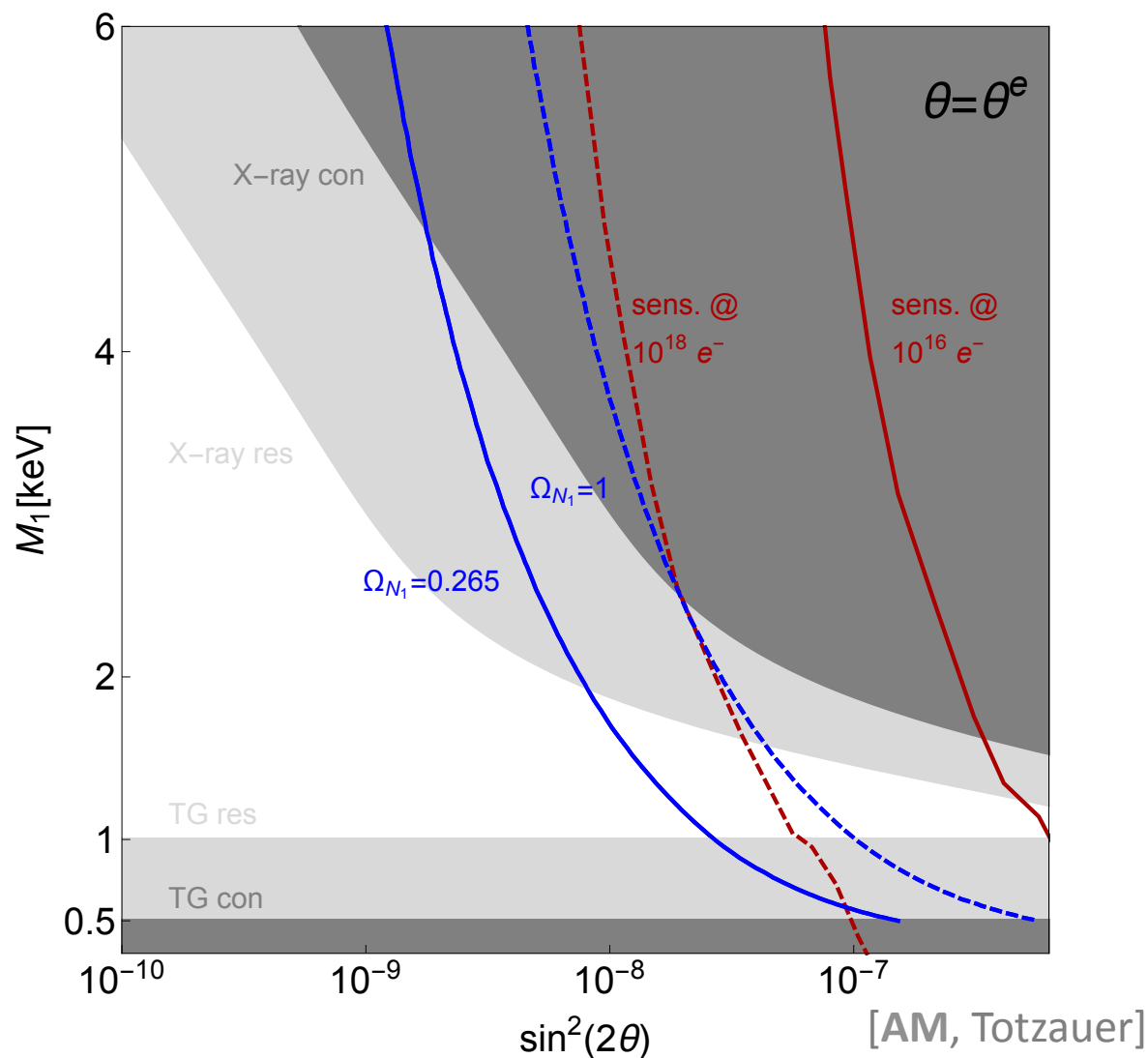
8. BACK-UP: Terrestrial Experiments

BUT: KATRIN/TRISTAN IS ONLY STICKING IN THE DARK!!!

→ *only sensitive if sterile neutrinos mix with electrons*

→ *even then, there is ZERO parameter space left which would still be allowed*

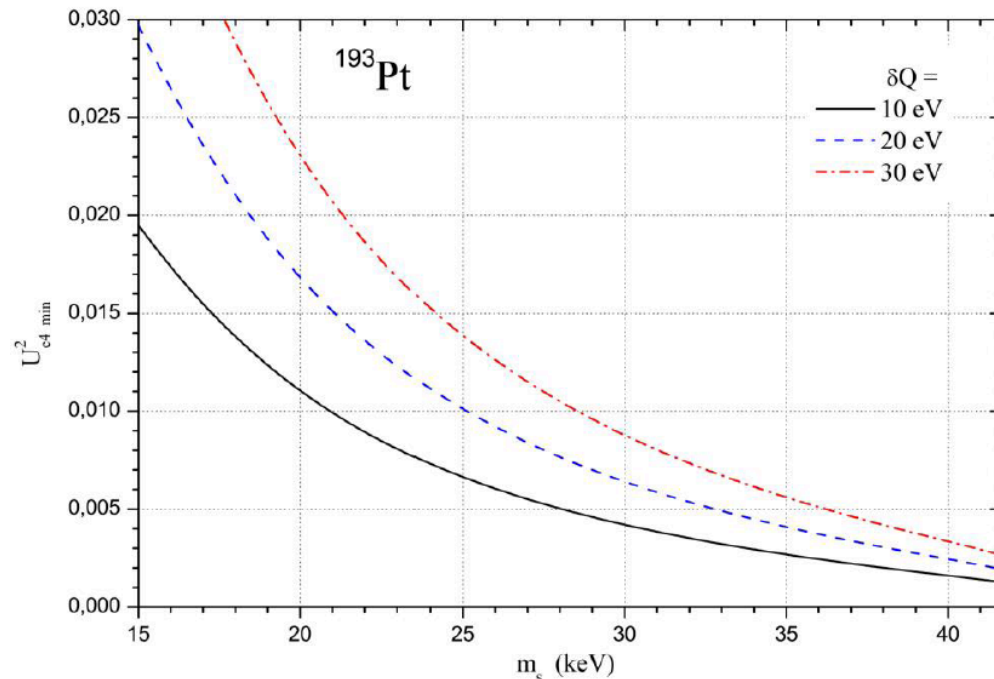
KATRIN reach for keV steriles



8. BACK-UP: Terrestrial Experiments

Other experiments unfortunately can't do any better...

- electron capture: alternative process



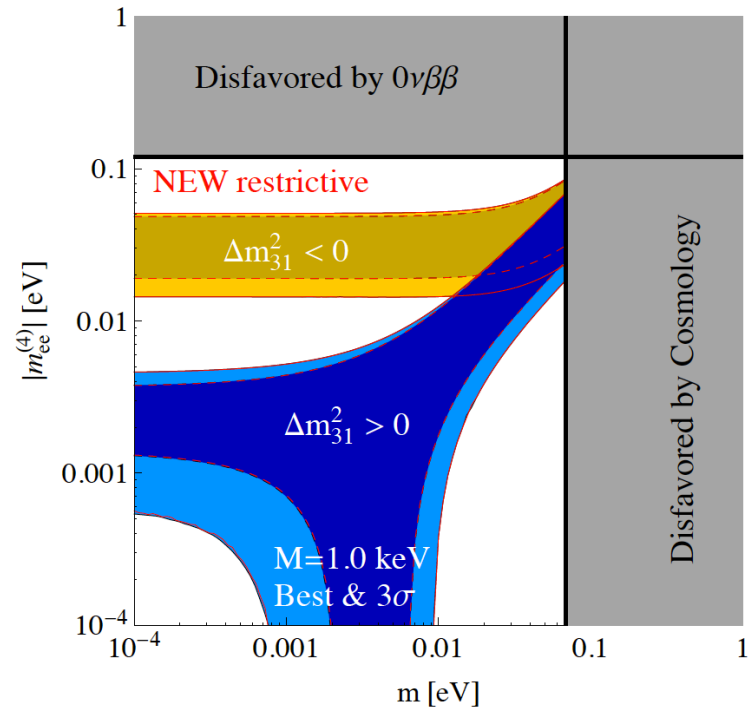
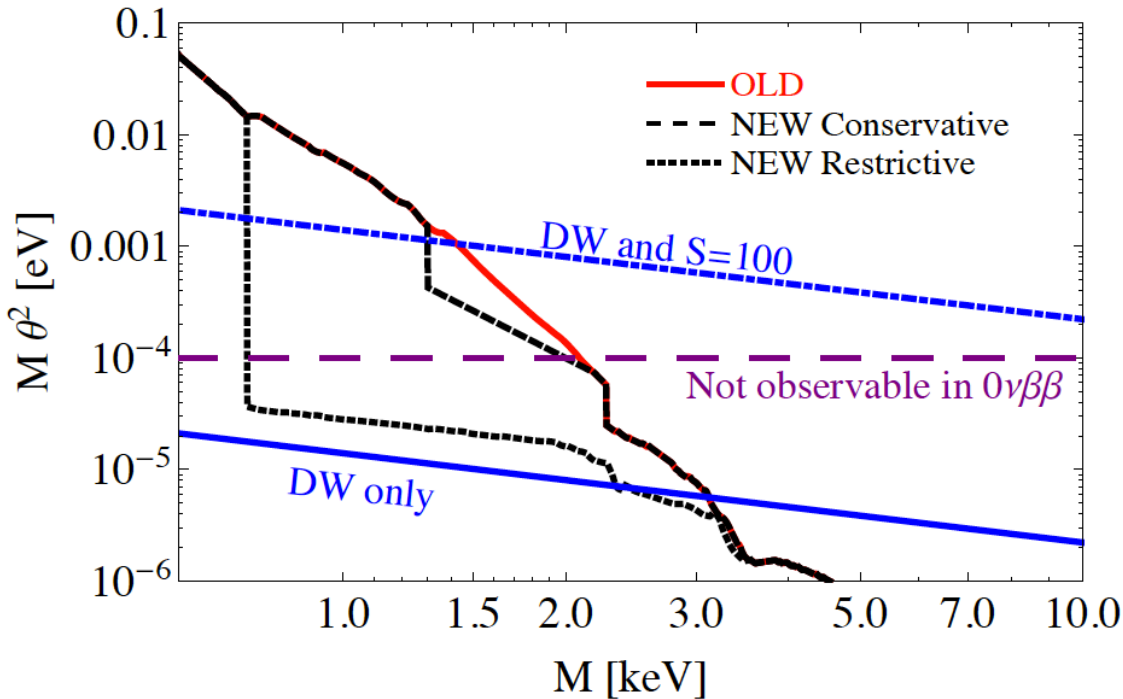
[Filianin *et al.*: 1402.4400]

- *combine techniques to go beyond rate measurements*
- *e.g.: cryogenic micro-calorimetry \rightarrow de-excitation spectrum (perturbed by ν_s) & Penning traps \rightarrow Q-value*
- *other proposals: Holmes, ECHO, ...*

\rightarrow complementary approach, but does it help?

8. BACK-UP: Terrestrial Experiments

But many do not seem ideal for this...



[Bezrukov: Phys. Rev. **D72** (2005) 071303]

[Asaka *et al.*: JHEP **1104** (2011) 011]

[AM, Niro: Phys. Rev. **D88** (2013) 113004]

→ essentially no effect due to strong X-ray bound ☹️

8. BACK-UP: Model Building Aspects

8. BACK-UP: Model Building Aspects

Simple framework: ν MSM

[Asaka, Blanchet, Shaposhnikov: Phys. Lett. B631 (2005) 151]

- **SM + 3 RH neutrinos** at (**keV, GeV- ϵ , GeV+ ϵ**)
 - can accommodate for ν -oscillations, BAU, and WDM
- provides fundamental connections between two clear signs for BSM physics: **neutrinos & Dark Matter**
- very **minimalistic extension of the SM**: only singlet (RH) neutrinos and lepton number violation
- **BUT: keV mass not explained**
GeV-degeneracy not explained
 ν -mixing pattern not explained
hardly testable

→ **MODEL BUILDING NEEDED!!!**

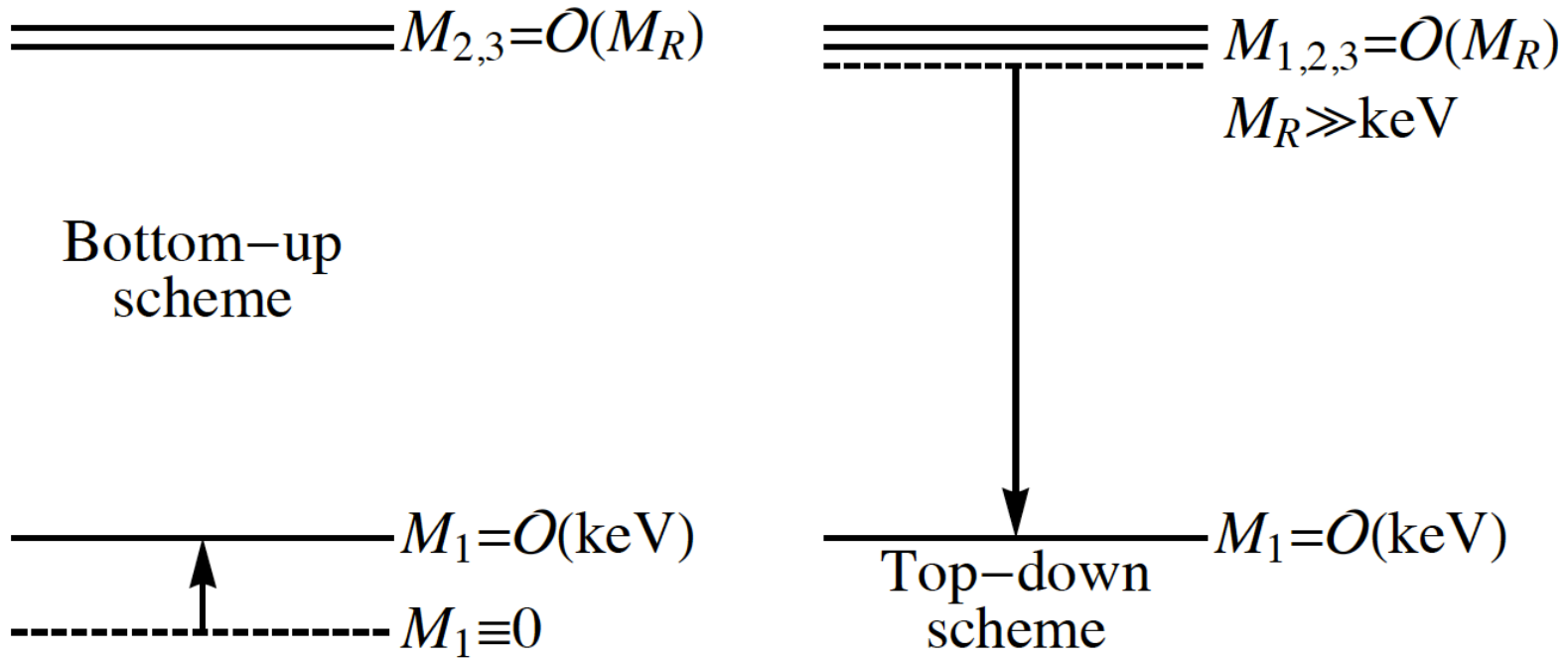
8. BACK-UP: Model Building Aspects

- **Differences to “ordinary” model building:**
 - we need an explanation for the keV scale:
 - not considered to be “fundamental”
 - need some mechanism

8. BACK-UP: Model Building Aspects

- Differences to “ordinary” model building:

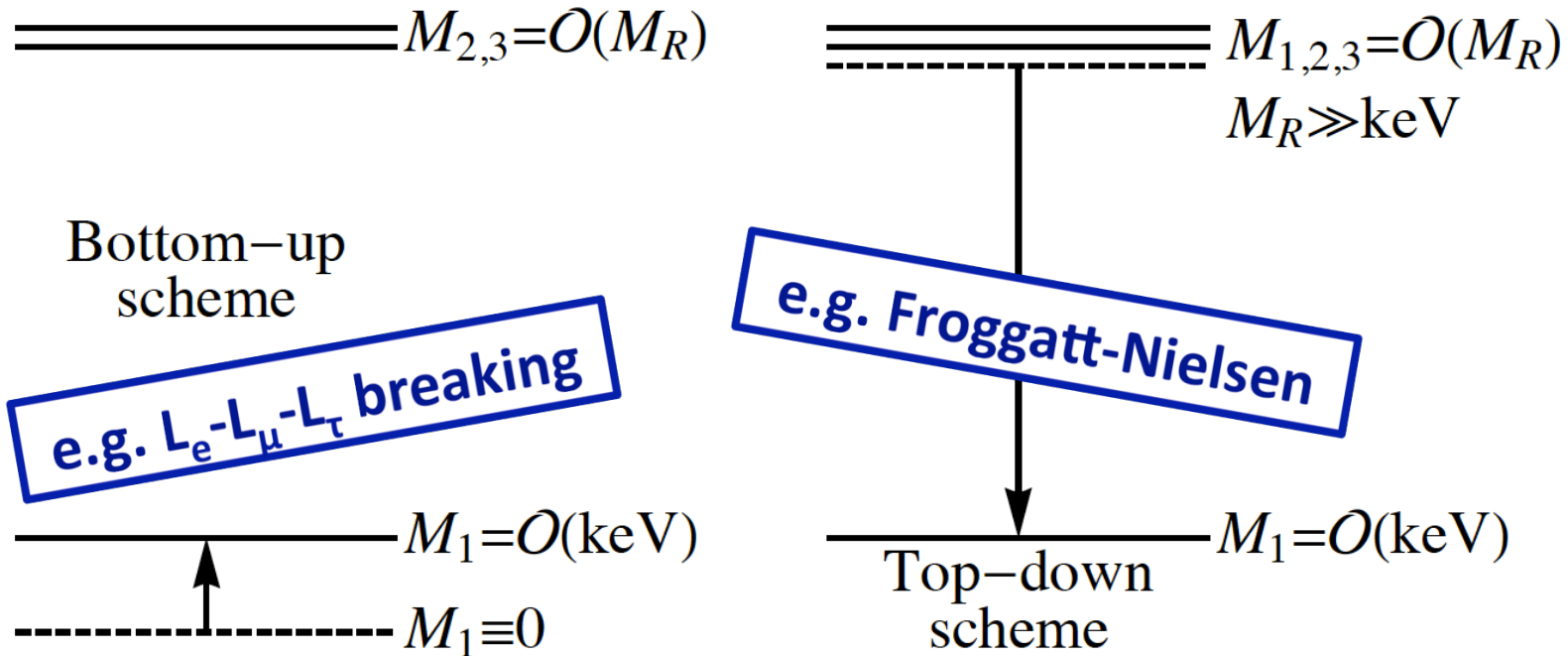
- we need an explanation for the keV scale:
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8. BACK-UP: Model Building Aspects

- Differences to “ordinary” model building:

- we need an explanation for the keV scale:
 - not considered to be “fundamental”
 - need some mechanism → two generic schemes:



[AM: Int. J. Mod. Phys. D22 (2013) 1330020]

→ Most models fall into one or the other category!

8. BACK-UP: Model Building Aspects

Example

- Model based on a flavour symmetry: $\mathcal{F} = L_e - L_\mu - L_\tau$

- **original**: [Petcov: Phys. Lett. **B110** (1982) 245]

- **2 RH neutrinos**: [Grimus, Lavoura: JHEP **0009** (2000) 007]

- **3 RH neutrinos**:

 - [Barbieri, Hall, Tucker-Smith, Strumia, Weiner: JHEP **9812** (1998) 017]

 - [Mohapatra: Phys. Rev. **D64** (2001) 091301]

- ***application to keV sterile neutrinos***:

 - [Shaposhnikov: Nucl. Phys. **B763** (2007) 49]

 - [Lindner, AM, Niro: JCAP **1101** (2011) 034]

- **general features**:

 - ***symmetry*** → patterns: $(0, m, m)$ & $(0, M, M)$

 - ***broken*** → **small mass, degeneracy lifted**

8. BACK-UP: Model Building Aspects

Example

- Model based on a flavour symmetry: $\mathcal{F} = L_e - L_\mu - L_\tau$
 - charge assignment under global U(1) [or: Z_4]:

	L_{eL}	$L_{\mu L}$	$L_{\tau L}$	e_R	μ_R	τ_R	N_{1R}	N_{2R}	N_{3R}	ϕ	Δ
\mathcal{F}	1	-1	-1	1	-1	-1	1	-1	-1	0	0

- then, only symmetry preserving terms are allowed:

$$\mathcal{L}_{\text{mass}} = -\frac{1}{2} \overline{\Psi^C} \mathcal{M}_\nu \Psi + h.c.$$

with: $\Psi \equiv ((\nu_{eL})^C, (\nu_{\mu L})^C, (\nu_{\tau L})^C, N_{1R}, N_{2R}, N_{3R})^T$

→ mass matrix:

$$\mathcal{M}_\nu = \begin{pmatrix} 0 & m_L^{e\mu} & m_L^{e\tau} & m_D^{e1} & 0 & 0 \\ m_L^{e\mu} & 0 & 0 & 0 & m_D^{\mu2} & m_D^{\mu3} \\ m_L^{e\tau} & 0 & 0 & 0 & m_D^{\tau2} & m_D^{\tau3} \\ \hline m_D^{e1} & 0 & 0 & 0 & M_R^{12} & M_R^{13} \\ 0 & m_D^{\mu2} & m_D^{\tau2} & M_R^{12} & 0 & 0 \\ 0 & m_D^{\mu3} & m_D^{\tau3} & M_R^{13} & 0 & 0 \end{pmatrix}$$

8. BACK-UP: Model Building Aspects

Example

- Model based on a flavour symmetry: $\mathcal{F} = L_e - L_\mu - L_\tau$
 - eigenvalues of \mathcal{M}_ν (**with μ - τ symmetry!!**):
 - light neutrinos: $(\lambda_+, \lambda_-, 0)$
 - heavy neutrinos: $(\Lambda_+, \Lambda_-, 0)$
 - with: $\lambda_\pm = \pm \sqrt{2} \left[m_L - \frac{m_D^2}{M_R} \right]$ $\Lambda_\pm = \pm \sqrt{2} M_R$
 - mass patterns:
 - light ν 's: $(0, \lambda_+, \lambda_-) \rightarrow$ okay up to degeneracy
 - heavy N 's: $(0, \Lambda_+, \Lambda_-) \rightarrow 0 \ll M$, but still $0 \neq \text{keV}$
 - WAY OUT: **broken symmetry**
 - \rightarrow will remedy the above issues
 - \rightarrow important: no matter how the breaking is achieved, the results will always look similar

8. BACK-UP: Model Building Aspects

Example

- Model based on a flavour symmetry: $\mathcal{F} = L_e - L_\mu - L_\tau$
 - pragmatic: *soft breaking* [Lindner, AM, Niro: JCAP **1101** (2011) 034]
 - we assumed small breaking terms and worked out their consequences:

→ new mass matrix:

$$\begin{pmatrix} s_L^{ee} & m_L^{e\mu} & m_L^{e\tau} & m_D^{e1} & 0 & 0 \\ m_L^{e\mu} & s_L^{\mu\mu} & 0 & 0 & m_D^{\mu2} & m_D^{\mu3} \\ m_L^{e\tau} & 0 & s_L^{\tau\tau} & 0 & m_D^{\tau2} & m_D^{\tau3} \\ \hline m_D^{e1} & 0 & 0 & S_R^{11} & M_R^{12} & M_R^{13} \\ 0 & m_D^{\mu2} & m_D^{\tau2} & M_R^{12} & S_R^{22} & 0 \\ 0 & m_D^{\mu3} & m_D^{\tau3} & M_R^{13} & 0 & S_R^{33} \end{pmatrix}$$

→ new eigenvalues: $\Lambda_s = S$, $\Lambda'_\pm = S \pm \sqrt{2}M_R$

$$\lambda_s = s \quad \lambda'_\pm = s \pm \sqrt{2} \left[m_L - \frac{m_D^2}{M_R} \right] + \frac{5m_D^2 S}{4M_R^2}$$

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natural assumption: like p-n isospin symmetry

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$$\begin{pmatrix} s_L^e & m_L^e & m_L^\tau & m_D^e & 0 & 0 \\ m_L^{e\mu} & s_L^{\mu\mu} & 0 & 0 & m_D^{\mu 2} & m_D^{\mu 3} \\ m_L^{e\tau} & 0 & s_L^{\tau\tau} & 0 & m_D^{\tau 2} & m_D^{\tau 3} \\ \hline m_D^{e1} & 0 & 0 & S_R^{11} & M_R^{12} & M_R^{13} \\ 0 & m_D^{\mu 2} & m_D^{\tau 2} & M_R^{12} & S_R^{22} & 0 \\ 0 & m_D^{\mu 3} & m_D^{\tau 3} & M_R^{13} & 0 & S_R^{33} \end{pmatrix}$$

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keV neutrino

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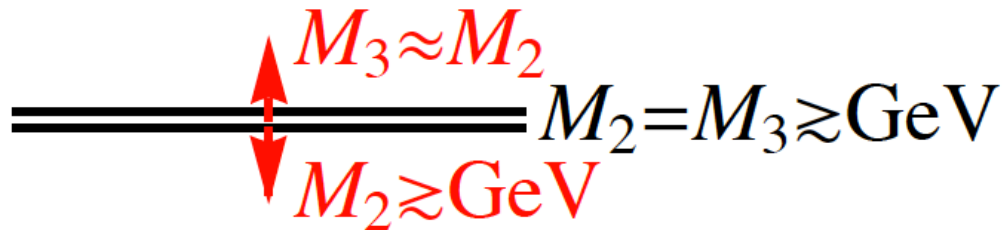
$$\lambda_s = s \quad \lambda'_\pm = s \pm \sqrt{2} \left[m_L - \frac{S}{M_R} \right]$$

small splitting

8. BACK-UP: Model Building Aspects

Example

- Model based on a flavour symmetry: $\mathcal{F} = L_e - L_\mu - L_\tau$
 - mass shifting scheme:



$$L_e - L_\mu - L_\tau \& \mu - \tau$$

→ clear bottom-up type scheme

~~$$L_e - L_\mu - L_\tau \& \mu - \tau$$~~



8. BACK-UP: Model Building Aspects

Example

- Model based on a flavour symmetry: $\mathcal{F} = \mathbf{L}_e - \mathbf{L}_\mu - \mathbf{L}_\tau$
 - mixings also require soft breaking:

$$\mathcal{M}_l \mathcal{M}_l^\dagger \simeq \begin{pmatrix} m_e^2 + m_\mu^2 \lambda^2 & m_\mu^2 \lambda & 0 \\ m_\mu^2 \lambda & m_\mu^2 & 0 \\ 0 & 0 & m_\tau^2 \end{pmatrix}$$

$$\tan^2 \theta_{12} \simeq 1 - 2\sqrt{2}\lambda + 4\lambda^4 - 2\sqrt{2}\lambda^3 \rightarrow \theta_{12} \simeq 33.4^\circ$$

$$\boxed{\lambda = \theta_{12} - \pi/4} \quad |U_{e3}| \simeq \frac{\lambda}{\sqrt{2}} \rightarrow \theta_{13} \simeq 8^\circ,$$

$$\sin^2 2\theta_{23} \simeq 1 - 4\lambda^4 \rightarrow \theta_{23} \simeq 45^\circ.$$

- prediction for the masses (under assumptions):

$$|m_1| = 0.0486 \text{ eV}, |m_2| = 0.0494 \text{ eV}, \text{ and } |m_3| = 0.0004 \text{ eV}$$

The Dodelson-Widrow mechanism

BUT: Is this really the end of the story?!?

NO!!! *We have disproved two common prejudices...*

[AM, Schneider, Totzauer: 1511.XXXXX]

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1st prejudice:

DW produces spectrum with thermal shape

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DW produces spectrum with thermal shape

- appears as approximation in the original papers:

[Dodelson, Widrow: Phys. Rev. Lett. **72** (1994) 17; DW & Colombi: Astrophys. J. **458** (1996) 1]

$$f(q) = \frac{\beta_{\text{DW}}}{\exp(q/T_{\text{DW}}) + 1}$$

- used in structure/galaxy formation computations:

[Viel *et al.*: Phys. Rev. **D71** (2005) 063534; Herpich *et al.*: Mon. Not. Roy. Astron. Soc. **442** (2014) 176; Menci, Fiore, Lemastra: Mon. Not. Roy. Astron. Soc. **421** (2012) 2384; Lovell *et al.*: Mon. Not. Roy. Astron. Soc. **439** (2014) 300,...]

The Dodelson-Widrow mechanism

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1st prejudice:

DW produces spectrum with thermal shape
- in reality, even if g_S is taken constant, which it is NOT
during DW production:

$$f_N^{\text{DW}}(T_f, p) \approx \frac{1}{\exp\left(\frac{p}{T_f} \left(\frac{\langle g_S \rangle}{g_S(T_f)}\right)^{1/3}\right) + 1} \int_{T_{\text{ini}}}^{T_f} dT_2 h\left(T_2, \frac{T_2}{T_f} \left(\frac{g_S(T_2)}{g_S(T_f)}\right)^{1/3}, p\right)$$

This function h (contains active-sterile mixing, mass difference, etc.) and the whole integral needs to vary **SLOWLY** with the momentum p !

→ not the case in reality... → NO THERMAL SHAPE!!!

The Dodelson-Widrow mechanism

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2nd prejudice:

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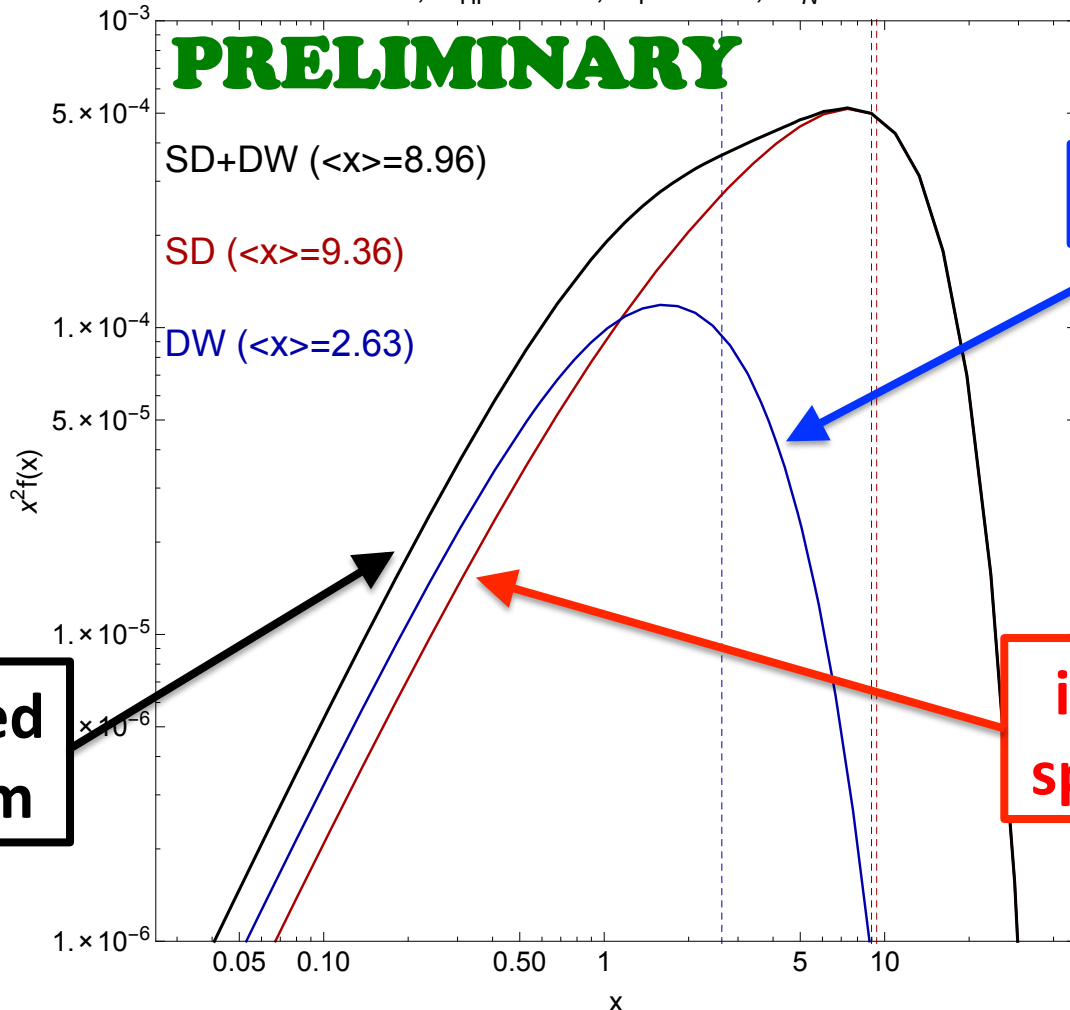
2nd prejudice:

DW always produces a too hot spectrum

- different for *non-zero initial abundance*: if DM is already present before DW sets in, the spectrum may experience non-trivial modifications
- example: scalar decay production + SUBSEQUENT modification by the Dodelson-Widrow mechanism

The Dodelson-Widrow mechanism

$T=10$ MeV, $C_{HP}=10^{-1.5}$, $C_{\Gamma}=10^{-3.0}$, $m_N=6.35$ keV



Combined spectrum

DW-modification

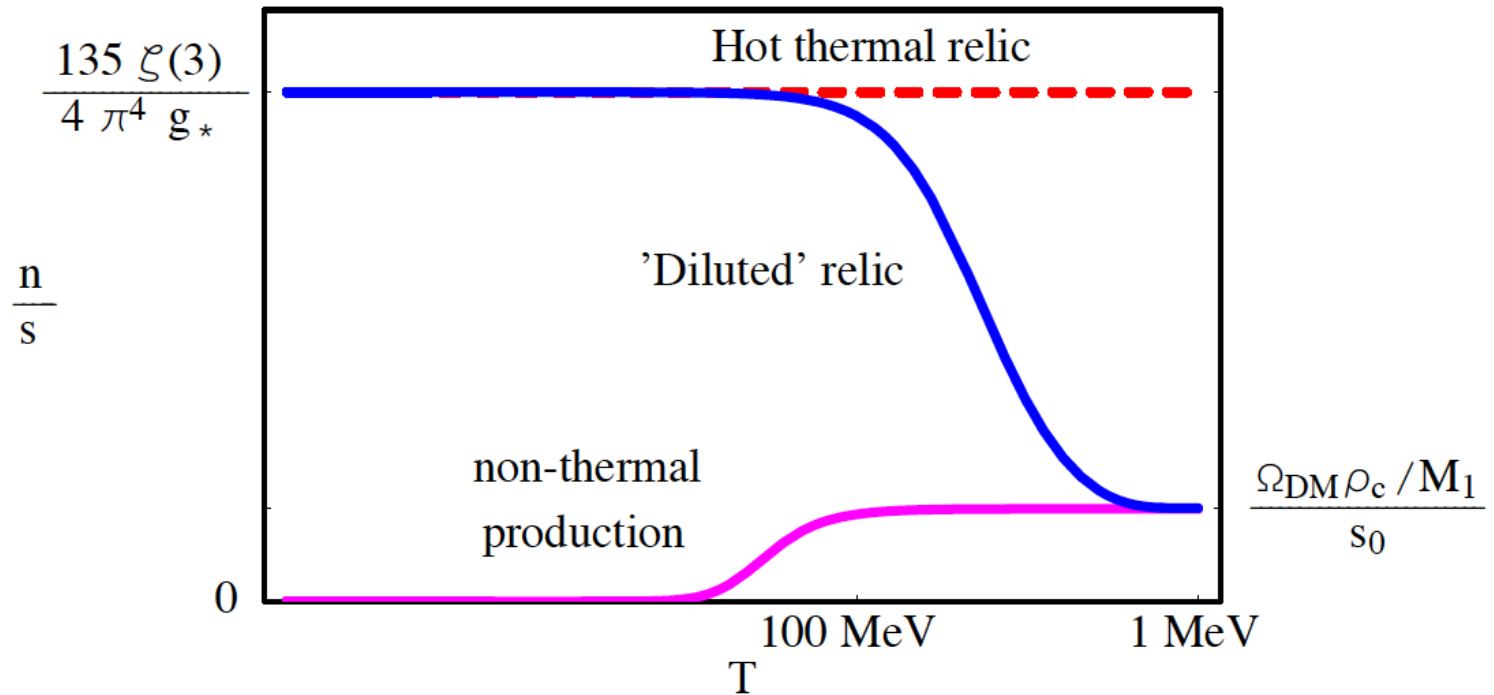
initial scalar decay spectrum (see later)

may shift the average momentum to LOWER values (“DW-cooling”)

BUT: SMALL EFFECT!!! (DW-part can at most be a $\sim 25\%$ modification)

Diluted Thermal Production

- Beyond the SM: sterile neutrino may be charged or have other interactions sufficiently strong to equilibrate it
→ *overproduction unless diluted by entropy production*



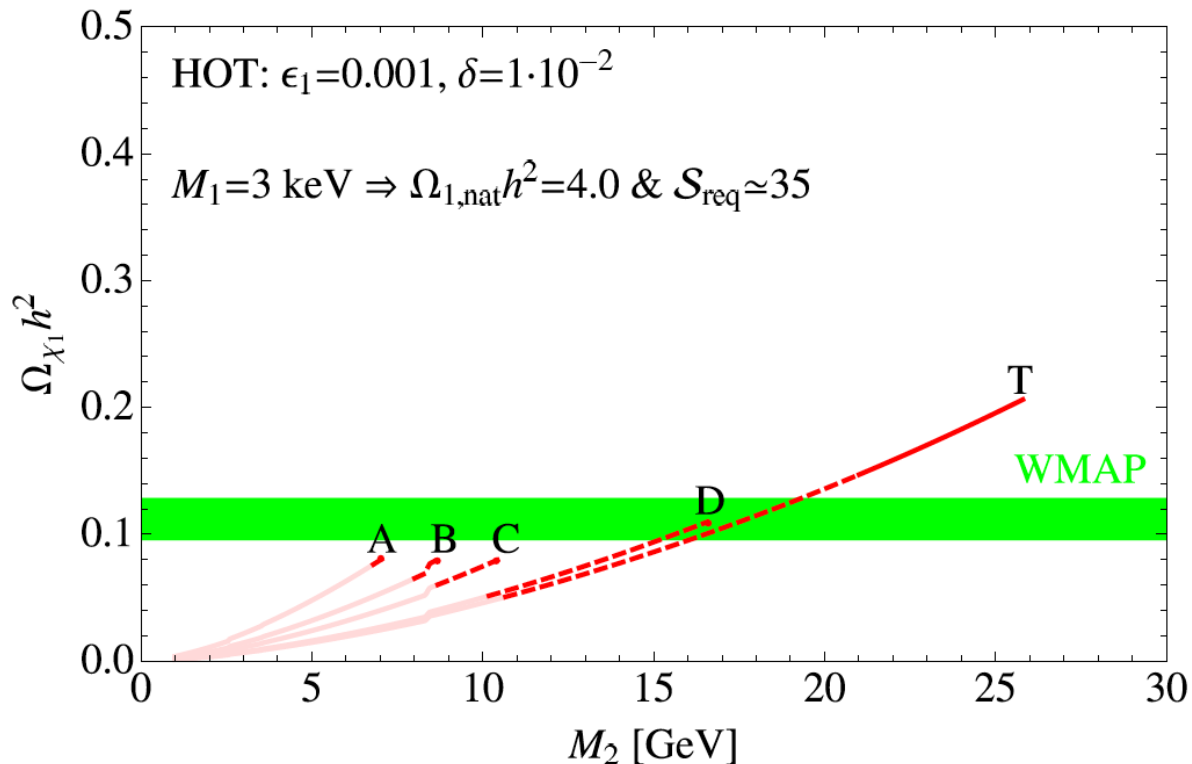
[Bezrukov, Hettmansperger, Lindner: Phys. Rev. **D81** (2010) 085032]

[King, **AM**: JCAP **1208** (2012) 016]

[Nemevsek, Senjanovic, Zhang: JCAP **1207** (2012) 006]

Diluted Thermal Production

- requirement 1: the diluter has to decay **slowly enough** to maximise entropy production
 - requirement 2: the diluter has to decay **fast enough** not to disturb big bang nucleosynthesis
- ➔ *hardly any viable parameter space left* ☹️



[King, AM:
JCAP 1208 (2012) 016]