

Damping rate
of a scalar
field at high
temperatures

Philipp H.

Motivation

Which damping
rate?

An application:
inflation

Formula for
the damping
rate

Linear response

Application of
linear response

High
temperature
regime

A simple model

The self energy

Another model

Kubo formula

Summary

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Philipp Henkenjohann

26.10.2015

Motivation

Which damping rate?

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Summary

- real scalar field performing small oscillations around its potential minimum
- interactions with other fields lead to damping of oscillations \rightarrow decay
- produced particles constitute a plasma which has effects on further decay

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Summary

- homogeneous scalar field drives inflation
- slowly rolls down its potential
- after reaching minimum: oscillation
- → production of particles

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Summary

- influence of plasma makes damping rate temperature dependent
- this effects the temperature evolution in the early universe
- maximal temperature can have implications for production of dark matter candidates

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Summary

- density matrix ρ determines thermal and quantum mechanical average of operators: $\langle B \rangle = \text{tr}(\rho B)$
- von Neumann equation $\frac{\partial \rho}{\partial t} = [\rho, H]$
- consider free Hamiltonian + small perturbation $H = H_0 + V(t)$, where $V(t) = \int d^3\mathbf{x} J(t, \mathbf{x}) A(\mathbf{x})$
- integration and expanding in J gives to first order:

$$\langle B \rangle = \langle B \rangle_{\text{eq.}} - \int d^4x' J(x') \Delta_{BA}^R(x - x')$$

- $\langle \dots \rangle_{\text{eq.}}$ can be calculated in thermal field theory

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- use this formula for the field operator ϕ ($A = B = \phi$)
- "kick" the field homogeneously: $J(x) = J_0\delta(t)$

$$\Rightarrow \delta\phi = -\frac{J_0}{2M_{\text{th}}}\sin(M_{\text{th}}t)e^{-\gamma t}$$

with

$$\gamma \approx -\frac{\text{Im}\Pi(M_{\text{th}} + i\epsilon, \mathbf{0})}{2M_{\text{th}}}$$

- thermal mass squared $M_{\text{th}}^2 = m^2 + \text{corrections}$
- $\Pi = \text{self energy calculated from thermal field theory}$

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Summary

Consider the following model:

$$\mathcal{L} = \mathcal{L}_{\text{KG},\phi} + \mathcal{L}_{\text{KG},\chi} - \frac{\lambda}{4!}\chi^4 - g\phi\chi^2$$

- thermal masses correspond to frequencies of oscillation
- at high temperatures thermal masses are dominated by corrections
- assume self coupling of plasma to be much larger than coupling of the decay-channel
- from the point of view of the plasma: ϕ -field barely oscillates, therefore:

$$\rightarrow \gamma \approx \lim_{\omega \rightarrow 0} -\frac{\text{Im}\Pi(\omega + i\epsilon, \mathbf{0})}{2\omega}$$

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- What does Π look like?
- to leading order in g and "all orders" in λ one has

$$\text{Im}\Pi(\omega + i\epsilon, \mathbf{k}) \propto g^2 \Delta_{\chi^2 \chi^2}^R(\omega, \mathbf{k})$$

$$\rightarrow \gamma \propto \lim_{\omega \rightarrow 0} -\frac{\Delta_{\chi^2 \chi^2}^R(\omega, \mathbf{0})}{2\omega}$$

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In a slightly different model one has

$$\gamma \propto \lim_{\omega \rightarrow 0} - \frac{\Delta_{(\partial\chi)^2(\partial\chi)^2}^R(\omega, \mathbf{0})}{2\omega}.$$

and

$$\partial_\mu \chi \partial^\mu \chi \propto T^\mu_\mu$$

$$\rightarrow \gamma \propto \lim_{\omega \rightarrow 0} - \frac{\Delta_{T^\mu_\mu T^\nu_\nu}^R(\omega, \mathbf{0})}{2\omega}$$

[D. Boedeker, JCAP 0606 (2006) 027]

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Use Kubo formula from hydrodynamics involving bulk viscosity

ζ

$$\zeta = \frac{1}{9} \lim_{\omega \rightarrow 0} \frac{1}{\omega} \Delta_{T^\mu_\mu T^\nu_\nu}^R(\omega, \mathbf{0})$$

$$\rightarrow \gamma \propto \zeta$$

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Summary

- real scalar field performs damped oscillations around potential minimum
- damping rate has influence on temperature evolution in the early Universe for example
- formula for damping rate can be obtained from linear response
- damping rate can be related to a hydrodynamic quantity (up to now only in one model)