

# ***CP* Violation in $K$ and $B_d$ mesons: Exp vs Standard Model**

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## **Summary**

- ✓ On the (less than perfect) consistency between the amount of *CP* violation in the  $K$  and  $B_d$  systems within the Standard Model
- ✓ How Minimal Flavour Violation can help this consistency test
- ✓ What would make the MFV corrections really distinguishable from the plain SM

**Based on:**

Buras, DG (PRD 08 & PRD 09)

## Nécessaire on $K$ mesons and their $CP$ violation

The  $d$  and  $s$  quarks can form, through strong interactions, the following bound states

$$|K^0\rangle \sim \begin{pmatrix} d \\ \bar{s} \end{pmatrix}$$

$$|\bar{K}^0\rangle \sim \begin{pmatrix} \bar{d} \\ s \end{pmatrix}$$

and phase conventions  
can be defined so that



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 \end{array} \right.$$

$K^0$  and  $\bar{K}^0$  mix into each other because of weak interactions.

However, if  $CP$  were a good symmetry, one would end up with the physical states:

$$\begin{array}{l}
 |K\rangle_{\text{even}} = \frac{|K^0\rangle - |\bar{K}^0\rangle}{\sqrt{2}} \quad \left( \begin{array}{l} CP = +1 \text{ admixture:} \\ \text{decays into } |\pi\pi\rangle \end{array} \right) \\
 |K\rangle_{\text{odd}} = \frac{|K^0\rangle + |\bar{K}^0\rangle}{\sqrt{2}} \quad \left( \begin{array}{l} CP = -1 \text{ admixture:} \\ \text{has to decay into } |\pi\pi\pi\rangle \end{array} \right)
 \end{array}$$

## Nécessaire 2: what is $\epsilon_K$

However, the actual physical admixtures are (slightly) different:

$$|K_S\rangle \propto |K\rangle_{\text{even}} + \bar{\epsilon} |K\rangle_{\text{odd}}$$

$$|K_L\rangle \propto |K\rangle_{\text{odd}} + \epsilon |K\rangle_{\text{even}}$$

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The magnitude of this CP violation is accessed experimentally by measuring the amplitude ratios:

$$\eta_{+-} = \frac{\langle \pi^+ \pi^- | K_L \rangle}{\langle \pi^+ \pi^- | K_S \rangle}$$

$$\eta_{00} = \frac{\langle \pi^0 \pi^0 | K_L \rangle}{\langle \pi^0 \pi^0 | K_S \rangle}$$

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It turns out that the corresponding types of  $CP$  violation can be disentangled by the following quantities:

$$\epsilon_K = \frac{1}{3} (\eta_{00} + 2\eta_{+-})$$

**“Indirect” CP violation  
(through mixing)**

$$\epsilon' = \frac{1}{3} (\eta_{+-} - \eta_{00})$$

**“Direct” CP violation  
(directly in the decay)**

**General formula  
for  $\epsilon_K$**

$$\epsilon_K = e^{i\phi_\epsilon} \sin \phi_\epsilon \left( \frac{\text{Im}(M_{12}^K)}{\Delta M_K} + \xi \right)$$

The quantities relevant to this formula are



$$\Delta M_K \equiv m_{K_L} - m_{K_S} \simeq 3.5 \times 10^{-15} \text{GeV}$$

$$\Delta \Gamma_K \equiv \Gamma_{K_L} - \Gamma_{K_S} \simeq -7.4 \times 10^{-15} \text{GeV}$$

**Note:**

$$\Delta \Gamma_K \approx -2\Delta M_K$$

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  $M_{12}^K \equiv \langle K^0 | \mathcal{H}_{\Delta S=2} | \bar{K}^0 \rangle$

**Amplitude for K-mixing:**  
sensitive to non-SM contributions

  $\xi \equiv \frac{\text{Im}A_0}{\text{Re}A_0}$

with  $A_0$  the amplitude for the decay  
 $K^0 \rightarrow \pi\pi$  (0-isospin)



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**Note**

The formula typically adopted in phenomenology takes

- $\xi \rightarrow 0$
- $\phi_\epsilon = 45^\circ$



Since both the deviations of  $\phi_\epsilon$  from  $45^\circ$  and  $\xi$  from zero are corrections, one can rewrite the general formula for  $\epsilon_K$  as

$$\epsilon_K = \kappa_\epsilon \times \epsilon_K(\xi = 0, \phi_\epsilon = 45^\circ)$$

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### How close is $\kappa_\epsilon$ to unity?

In the Standard Model, we estimated  
[ See: Buras, DG, PRD08 ]

$$\kappa_\epsilon = 0.92 \pm 0.02$$

Note: the corrections from  $\xi \neq 0$  AND  $\phi_\epsilon \neq 45^\circ$  have like sign.

This accident builds up a  
– **8 % total correction!**

## How to estimate $\kappa_\epsilon$

As we saw before,  
 $\kappa_\epsilon$  is defined by the relation

$$\epsilon_K = \kappa_\epsilon \times \epsilon_K(\xi = 0, \phi_\epsilon = 45^\circ) \quad \Rightarrow$$

$$\kappa_\epsilon = \frac{\sin \phi_\epsilon}{1/\sqrt{2}} \times (1 + \Delta_\epsilon)$$

- Parameterizes the effect of  $\xi \neq 0$ .
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However,  $\kappa_\epsilon$  can be estimated indirectly, through  $\epsilon'/\epsilon$ , using the relation

$$\frac{\epsilon'}{\epsilon} = -\omega \Delta_\epsilon (1 - \Omega)$$


- $\omega = \text{Re}(A_2)/\text{Re}(A_0) = 0.045$  is known very precisely (“ $\Delta I = 1/2$  rule”)
- $\Omega$  represents (basically) the ratio between EW-penguin and QCD-penguin contributions to  $\epsilon'/\epsilon$
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$$\Delta_\epsilon = -\frac{1}{\omega(1 - \Omega)} \left[ \frac{\epsilon'}{\epsilon} \right]_{\text{exp}}$$


$$= -0.061 \pm 0.014$$

Using

- $\epsilon'/\epsilon = (1.66 \pm 0.26) \times 10^{-3}$
- $\Omega = (0.4 \pm 0.1)$  [within the SM]

See the analysis by:  
**Buras-Jamin,**  
**JHEP04**

## Consequences of the $\epsilon_K^{SM}$ suppression factor

-  Once the rest of the input is fixed, the formula for  $\epsilon_K^{SM}$  allows to predict  $\sin 2\beta$ , that measures the amount of  $CP$  violation in the  $B_d$ -system.

In particular, the constraint from  $\epsilon_K$  works as follows

$$|\epsilon_K^{\text{exp}}| = |\epsilon_K^{\text{th,SM}}| \propto \kappa_\epsilon \times \sin 2\beta$$



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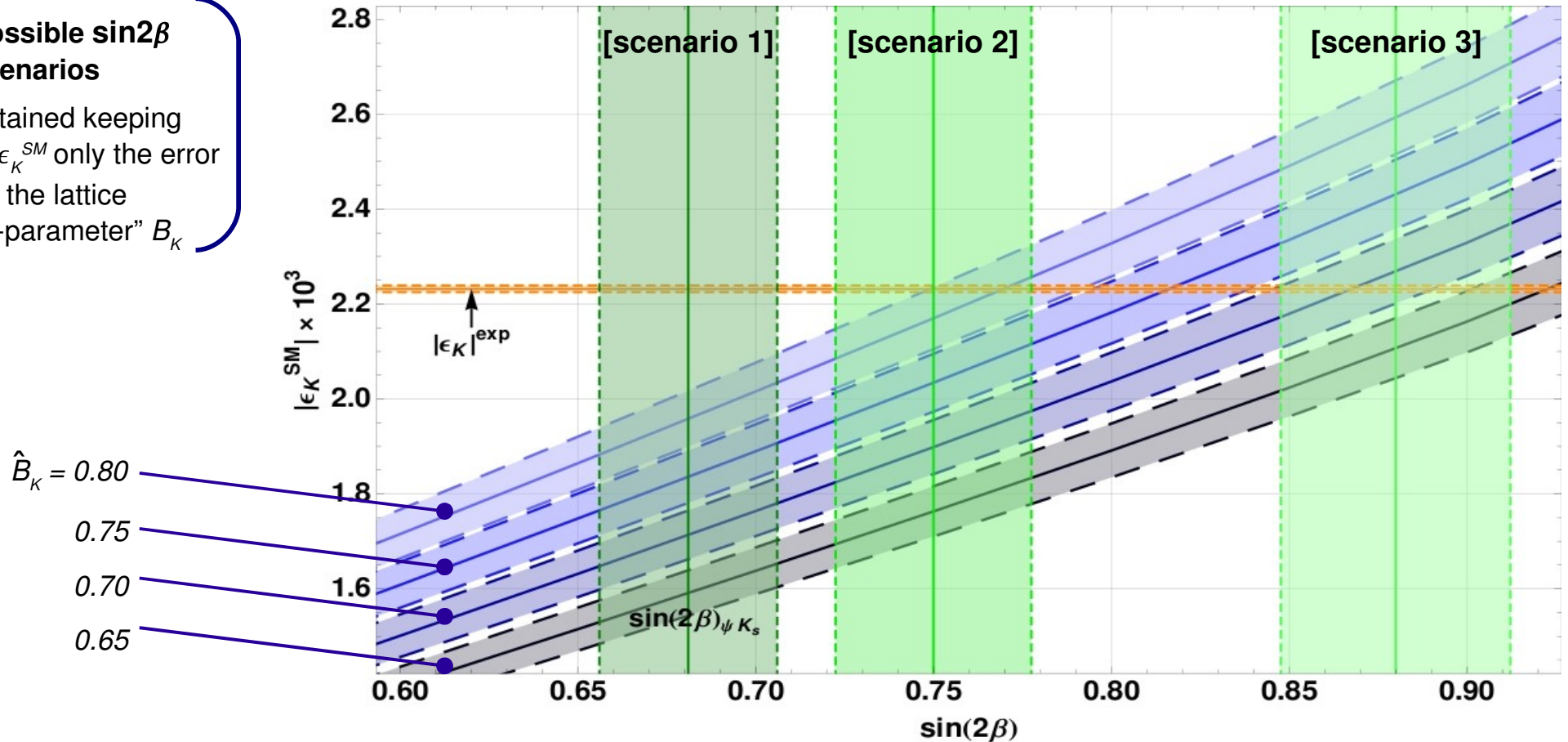
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Possible  $\sin 2\beta$  scenarios  
 obtained keeping in  $\epsilon_K^{SM}$  only the error on the lattice "B-parameter"  $B_K$



**More on scenario 1:**

$$\sin 2\beta_{J/\psi K_s} = \sin 2\beta$$

In this case one gets  $|\epsilon_K^{\text{SM}}| = (1.78 \pm 0.25) \times 10^{-3}$

to be compared with  $|\epsilon_K^{\text{exp}}| = (2.229 \pm 0.012) \times 10^{-3}$



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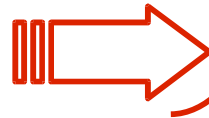
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*The simplest solution is a positive shift in the  $\epsilon_K$  loop function*

- *This solution is of MFV type. In fact, the CKM structure is preserved and non-SM physics only enters the short-distance S-function*

**[ Buras et al., 00 ]**

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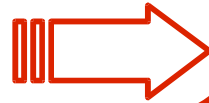
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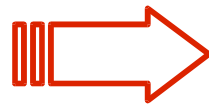


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Barring non-SM operators mediating mixing, the above shift would be universal, i.e. also affect  $B_d$  and  $B_s$  mass differences (and cancel in their ratio)



By demanding a 10% shift, for  $\epsilon_K$  to recover 1 sigma agreement with exp, one would get

$$\begin{array}{l} \Delta M_d^{\text{CMFV}} \approx (0.638 \pm 20\%) / \text{ps} \\ \Delta M_s^{\text{CMFV}} \approx (21.6 \pm 20\%) / \text{ps} \end{array} \quad \text{vs} \quad \begin{array}{l} \Delta M_d^{\text{HFAG}} = 0.507(5) / \text{ps} \\ \Delta M_s^{\text{HFAG}} = 17.77(12) / \text{ps} \end{array}$$

**More on scenario 3:**

$$\sin 2\beta_{J/\psi K_s} = \sin 2(\beta + \phi_d)$$

In this case the phase  $\beta$  cannot be accessed directly from the  $J/\psi K_s$  mode

One possible strategy to determine  $\beta$  is by using  $\epsilon_K$ ,  $\Delta m_d$  and  $\Delta m_s$  only.

See also:  
Lunghi, Soni, PLB08

We get:  $\sin 2\beta = 0.88^{+0.11}_{-0.12}$   
to be compared with  $\sin 2\beta_{J/\psi K_s} = 0.681 \pm 0.025$



At face value, this allows  
for  $\phi_d \approx -10^\circ$

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In spite of the presence of a new phase, the non-SM physics responsible for it does not need to be beyond MFV.

- In fact, MFV is the principle of minimal breaking of the SM  $[SU(3)]^5$  flavor group

The assumption that the breaking of the discrete CP symmetry be also minimal is in principle unrelated.

- However, the question on the predicted amount of CPV is of course highly model-dependent in this case

[ D'Ambrosio et al., 02 ]

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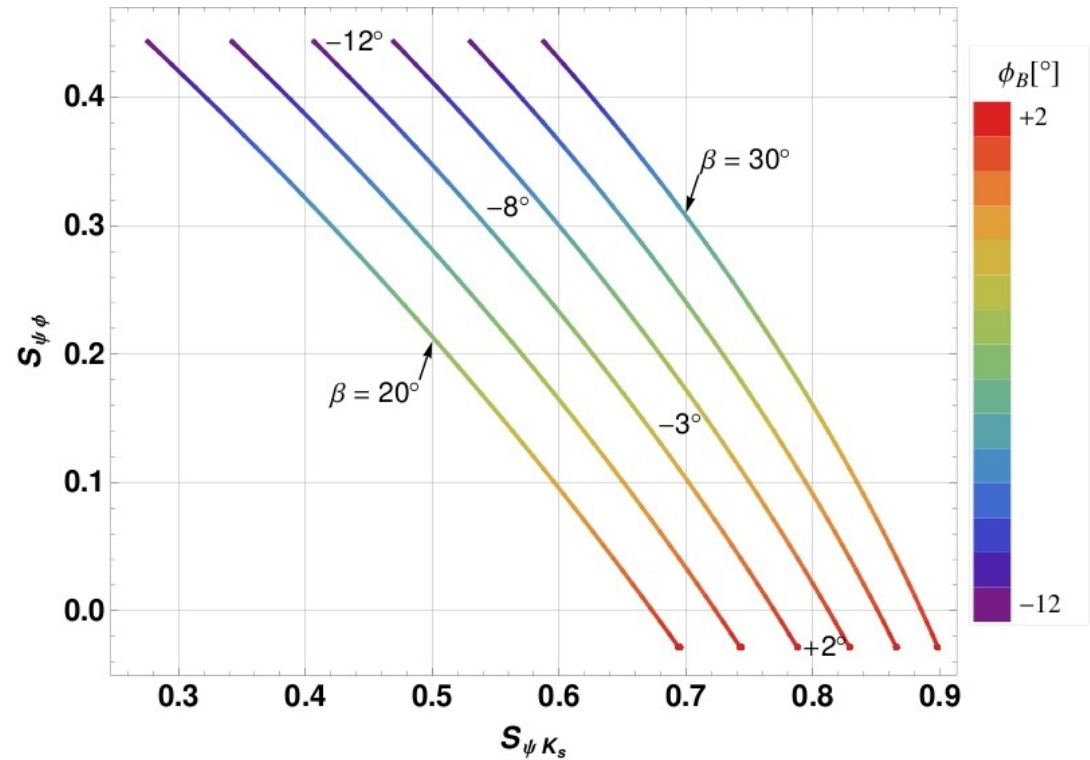
The case of  $\phi_d \approx \phi_s$

 What if

$$\phi_B = \phi_d \approx \phi_s \approx -9^\circ$$



$$\begin{cases} \beta_{\psi K_s} < \beta \approx 30^\circ \\ S_{\psi\phi} \approx 0.4 \end{cases}$$



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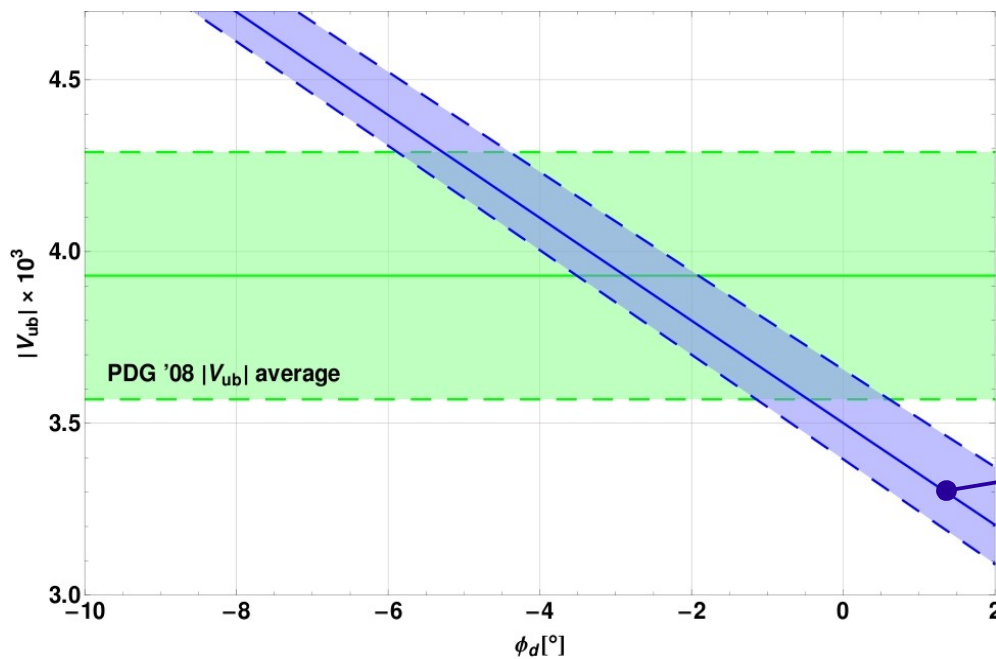
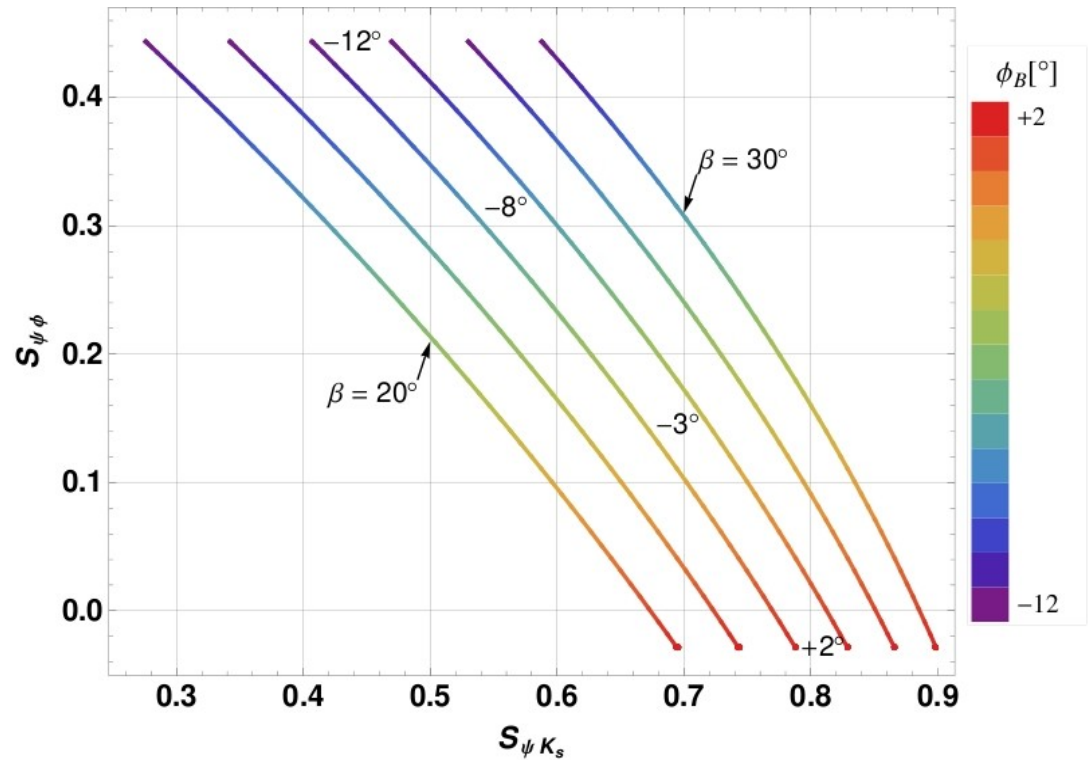
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A crucial test of this scenario "would come" from  $V_{ub}$

$|V_{ub}|$  range  
for the case  $\phi_d \approx \phi_s$

## How to (more generally) estimate $\kappa_\epsilon$

We saw that  $\kappa_\epsilon$  can be estimated from  $[\epsilon'/\epsilon]_{\text{exp}}$ , the main theoretical input being in  $\Omega$

$$\kappa_\epsilon = \frac{\sin \phi_\epsilon}{1/\sqrt{2}} \times (1 + \Delta_\epsilon) \quad \left( \Delta_\epsilon = -\frac{1}{\omega(1-\Omega)} \left[ \frac{\epsilon'}{\epsilon} \right]_{\text{exp}} \right)$$



Note that  $\Omega$  can be affected by non-SM contributions, because EW penguins are very sensitive to new physics.

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A convenient formula to evaluate  $\epsilon'/\epsilon$  and to understand its structure is provided by

the so-called Penguin-Box expansion  $\epsilon'/\epsilon = \text{Im}(V_{ts}^* V_{td}) \sum_i P_i \times X_i$

where

- $X_i = \text{Inami-Lim functions}$
- $P_i = r_i^{(0)} + r_i^{(6)} R_6 + r_i^{(8)} R_8$

#### Numerical coeffs:

encode the info on the Wilson coeff. functions of the  $\Delta S = 1$  Hamiltonian

Buras, Lautenbacher, 93  
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Encode the info on the matrix elements  $\langle Q_6 \rangle_0$  and  $\langle Q_8 \rangle_2$  with

$$Q_6 = (\bar{s}_\alpha d_\beta)_{V-A} \sum_q (\bar{q}_\beta q_\alpha)_{V+A}$$

$$Q_8 = \frac{3}{2} (\bar{s}_\alpha d_\beta)_{V-A} \sum_q e_q (\bar{q}_\beta q_\alpha)_{V+A}$$

leading QCD penguin

leading EW penguin

Buras, Lautenbacher, 93  
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## How to (more generally) estimate $\kappa_\epsilon$ : continued



Therefore,  $\epsilon'/\epsilon$  has the structure

$$\epsilon'/\epsilon = c_0 + c_6 R_6 + c_8 R_8 = [\epsilon'/\epsilon](\text{NP shifts}, R_6, R_8)$$

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### Procedure

- ✓ Evaluate  $c_0, c_6, c_8$ , including non-SM contributions, by appropriate shifts of the Inami-Lim functions
- ✓ Take for  $R_8$  the range  $R_8 = 1.0 \pm 0.2$  [ see analysis by Buras-Jamin, 03 ]
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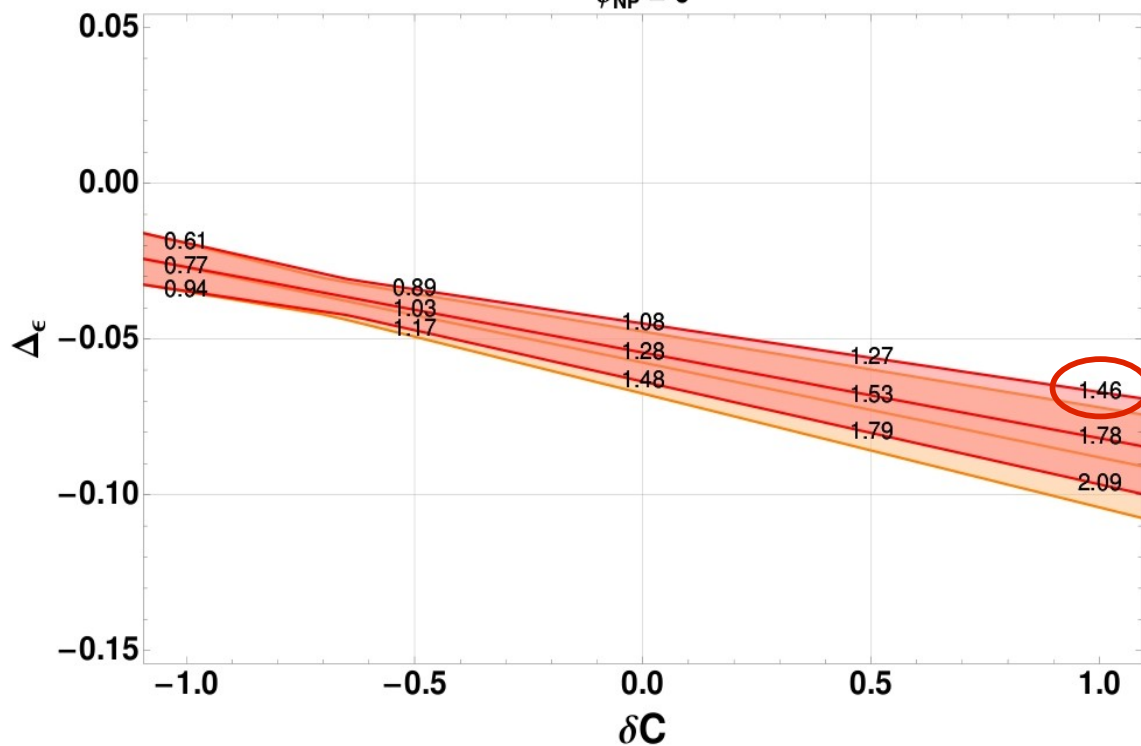
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- ✓ Estimate  $R_6$  by demanding equality of the above formula with  $[\epsilon'/\epsilon]_{\text{exp}}$
- ✓ Since
  - $c_6 R_6$  has only a  $\Delta I = 1/2$  component
  - $c_8 R_8$  has only a  $\Delta I = 3/2$  component
  - $-\omega \Delta_\epsilon$  is by definition the sum of the  $\Delta I = 1/2$  contributions

it follows that  $\Delta_\epsilon$  can be estimated through  $-1/\omega \cdot [\epsilon'/\epsilon](\text{NP shifts}, R_6, R_8 \rightarrow 0)$   
(with a small caveat on  $c_0$ )

More quantitatively:  
case of EW-penguin NP phase = 0

$\phi_{NP} = 0$



$\delta C$  = shift (normalized to the SM contribution) in the EW penguins.

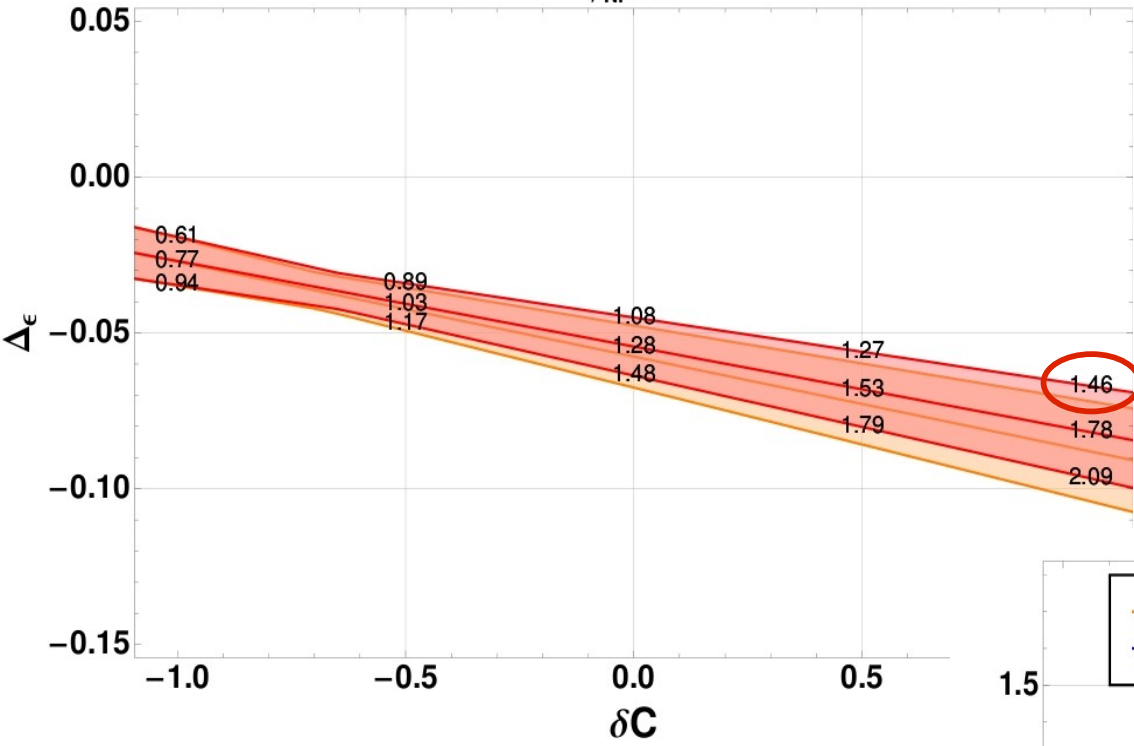
In this plot, the shift is assumed real

$R_6$ -value corresponding to the given  $\Delta_\epsilon$

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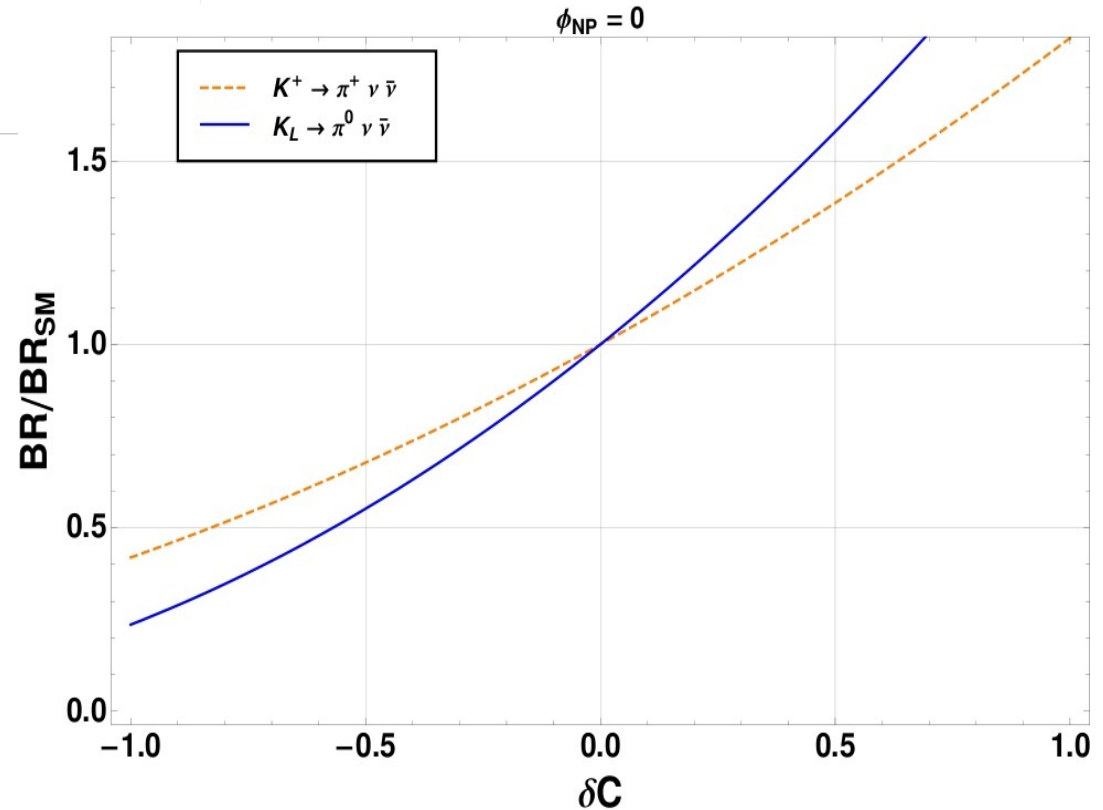
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In this plot, the shift is assumed real



$R_6$ -value corresponding to the given  $\Delta_\epsilon$

The shift predicts a very defined pattern of enhancement for rare  $K$  decays



## Conclusions


- ☑ *The correlation  $\epsilon_K - \sin 2\beta$  is a fundamental consistency check of SM CP violation. With regards to CP violation, it is the only one available at present.*
- ☑ *Our analysis shows that a (more) accurate SM formula for  $\epsilon_K$  implies a  $-8\%$  shift in the central value.*
- ☑ *Looking at the entailed prediction for  $\sin 2\beta$ , the above shift hints at a tension. While, with present errors, no statement above 2 sigma can be made, the issue warrants further investigation.*



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- ☑ Looking at the entailed prediction for  $\sin 2\beta$ , the above shift hints at a tension. While, with present errors, no statement above 2 sigma can be made, the issue warrants further investigation.
- ☑ Reaching firm(er) conclusions about the tension requires improvement in the theoretical input. To get an idea, the leading top-top contribution ( $\approx 75\%$ ) to  $\epsilon_K^{\text{SM}}$  goes as:

$$|\epsilon_K|^{\text{SM}} \propto \kappa_\epsilon \hat{B}_K |V_{cb}|^4 |V_{us}|^2 R_t^2 \sin 2\beta$$



$$\frac{\delta |\epsilon_K|^{\text{SM}}}{|\epsilon_K|^{\text{SM}}} \approx \sqrt{\underbrace{\left(\frac{\delta \hat{B}_K}{\hat{B}_K}\right)^2}_{5\%} + \underbrace{\left(4 \frac{\delta |V_{cb}|}{|V_{cb}|}\right)^2}_{11\%} + \underbrace{\left(2 \frac{\delta R_t}{R_t}\right)^2}_{8\%}} \approx 15\%$$

- ☑ If the tension becomes an inconsistency, scenarios providing a solution can be searched for already within MFV