

Inclusive semileptonic B decays

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Our world is full of tensions

Outline

- Inclusive semileptonic decays, moments, fits, V_{cb}
- Do not use the word *precision* in vain
how reliable are the mass determinations?
- Better not to forget unitarity
a remark on heavy quark sum rules
- The inclusive V_{ub} determination
- Love your fellow theorists like yourself
comparison of various approaches
- Conclusions

Inclusive semileptonic B decays: basic features

- **Simple idea:** inclusive decay do not depend on final state, factorize long distance dynamics of the meson. OPE allows to express it in terms of matrix elements of local operators
- The Wilson coefficients are perturbative, matrix elements of local ops parameterize non-pert physics: **double series in $\alpha_s, \Lambda/m_b$**
- Lowest order: decay of a free b , linear Λ/m_b absent. Depends on $m_{b,c}$, 2 parameters at $O(1/m_b^2)$, 2 more at $O(1/m_b^3)$...

$$\mu_\pi^2(\mu) = \frac{1}{2M_B} \left\langle B \left| \bar{b} (i\vec{D})^2 b \right| B \right\rangle_\mu \quad \mu_G^2(\mu) = \frac{1}{2M_B} \left\langle B \left| \bar{b} \frac{i}{2} \sigma_{\mu\nu} G^{\mu\nu} b \right| B \right\rangle_\mu$$

The total s.l. width in the OPE

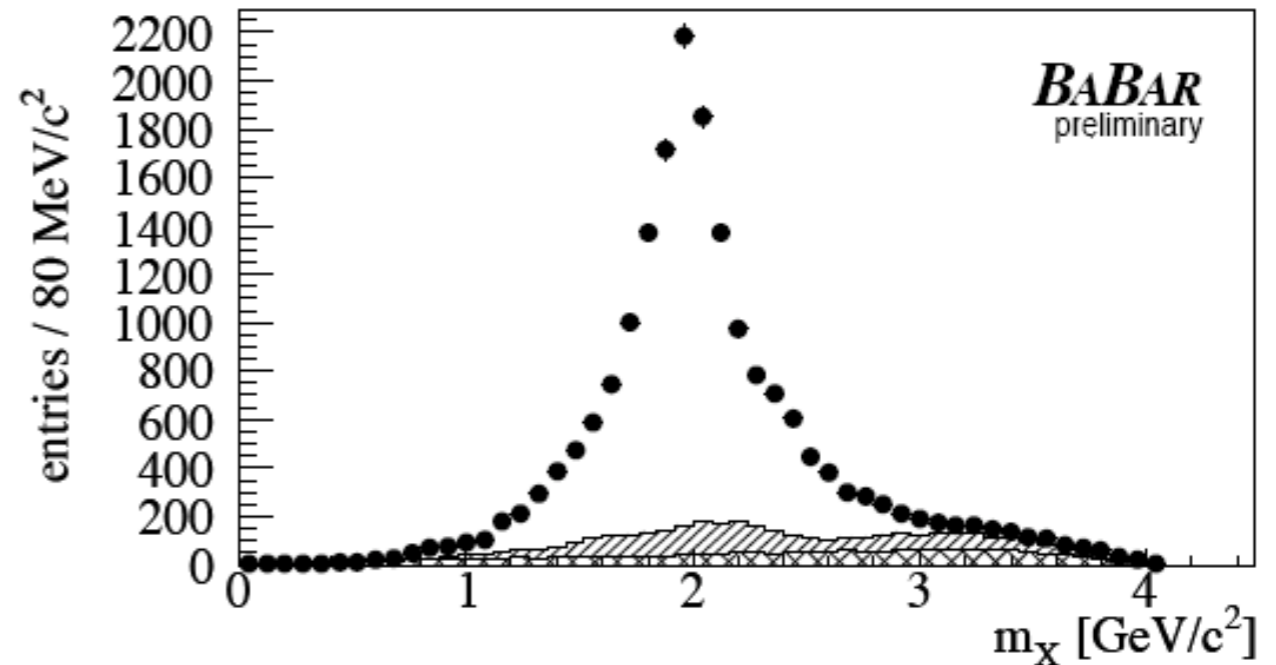
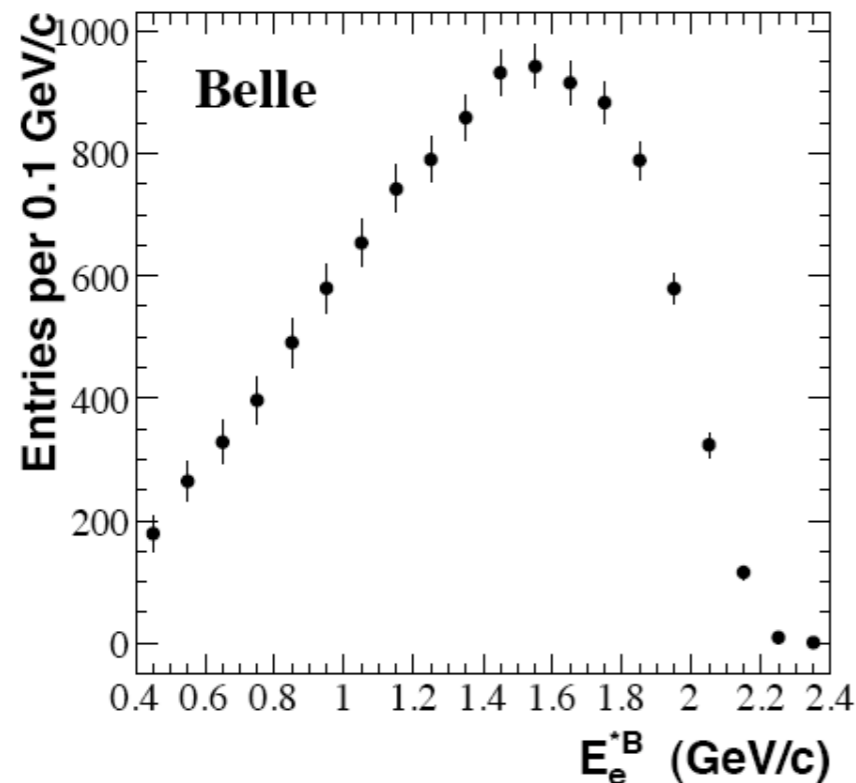
$$\Gamma[\bar{B} \rightarrow X_c e \bar{\nu}] = \frac{G_F^2 m_b^5 |V_{cb}|^2 g(r)}{192\pi^3} \left[1 + \frac{\alpha_s}{\pi} p_c^{(1)}(r, \mu) + \frac{\alpha_s^2}{\pi^2} p_c^{(2)}(r, \mu) \right. \\ \left. - \frac{\mu_\pi^2}{2m_b^2} + \left(\frac{1}{2} - \frac{2(1-r)^4}{g(r)} \right) \frac{\mu_G^2 - \frac{\rho_{LS}^3 + \rho_D^3}{m_b}}{m_b^2} \right. \\ \left. + \left(8 \ln r - \frac{10r^4}{3} + \frac{32r^3}{3} - 8r^2 - \frac{32r}{3} + \frac{34}{3} \right) \frac{\rho_D^3}{g(r) m_b^3} \right] \\ + O\left(\alpha_s \frac{\mu_{\pi,G}^2}{m_b^2}\right) + O\left(\frac{1}{m_b^4}\right)$$

$$r = \frac{m_c^2}{m_b^2}$$

OPE valid for inclusive enough measurements, away from perturbative singularities \implies moments

At present the implementations for moments include $O(\alpha_s^2 \beta_0, 1/m_b^3)$ terms

Fitting OPE parameters to the moments



Total **rate** gives $|V_{cb}|$, global **shape** parameters (moments of the distributions) tell us about B structure, m_b and m_c

OPE parameters describe universal properties of the B meson and of the quarks

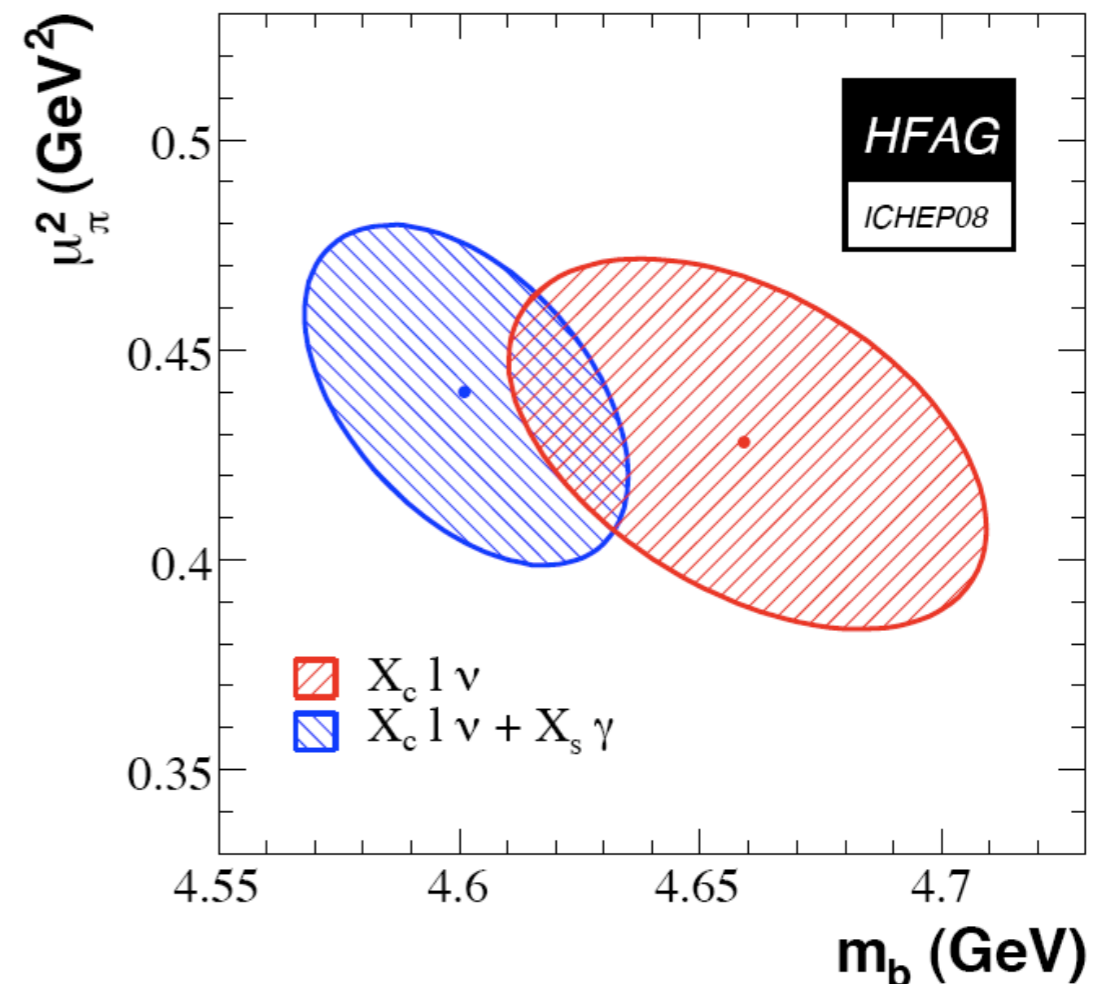
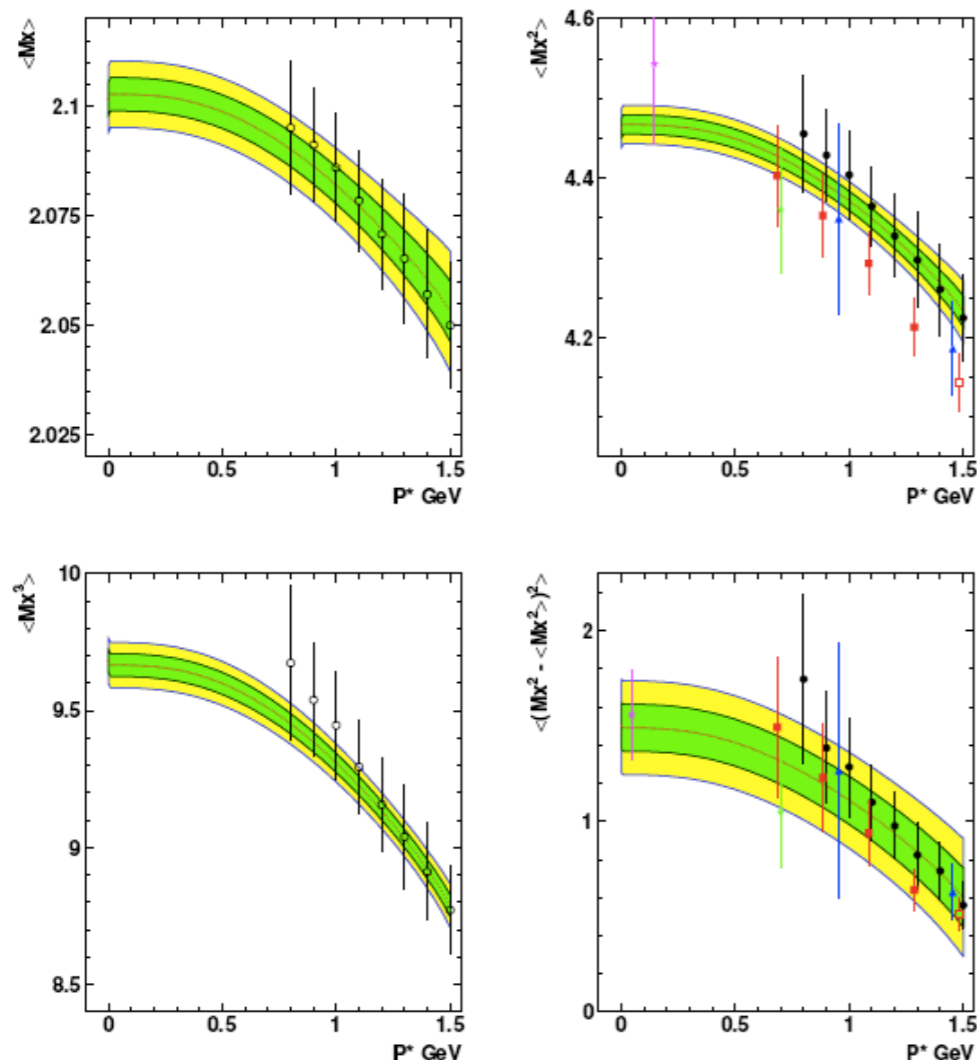
Global fit (kinetic scheme)

Inputs	$ V_{cb} \cdot 10^3$	m_b^{kin}	χ^2/ndf
$b \rightarrow c$ & $b \rightarrow s\gamma$	41.67(44)(58)	4.601(34)	29.7/57
$b \rightarrow c$ only	41.48(48)(58)	4.659(49)	24.1/46

Based on PG, Uraltsev, Benson et al

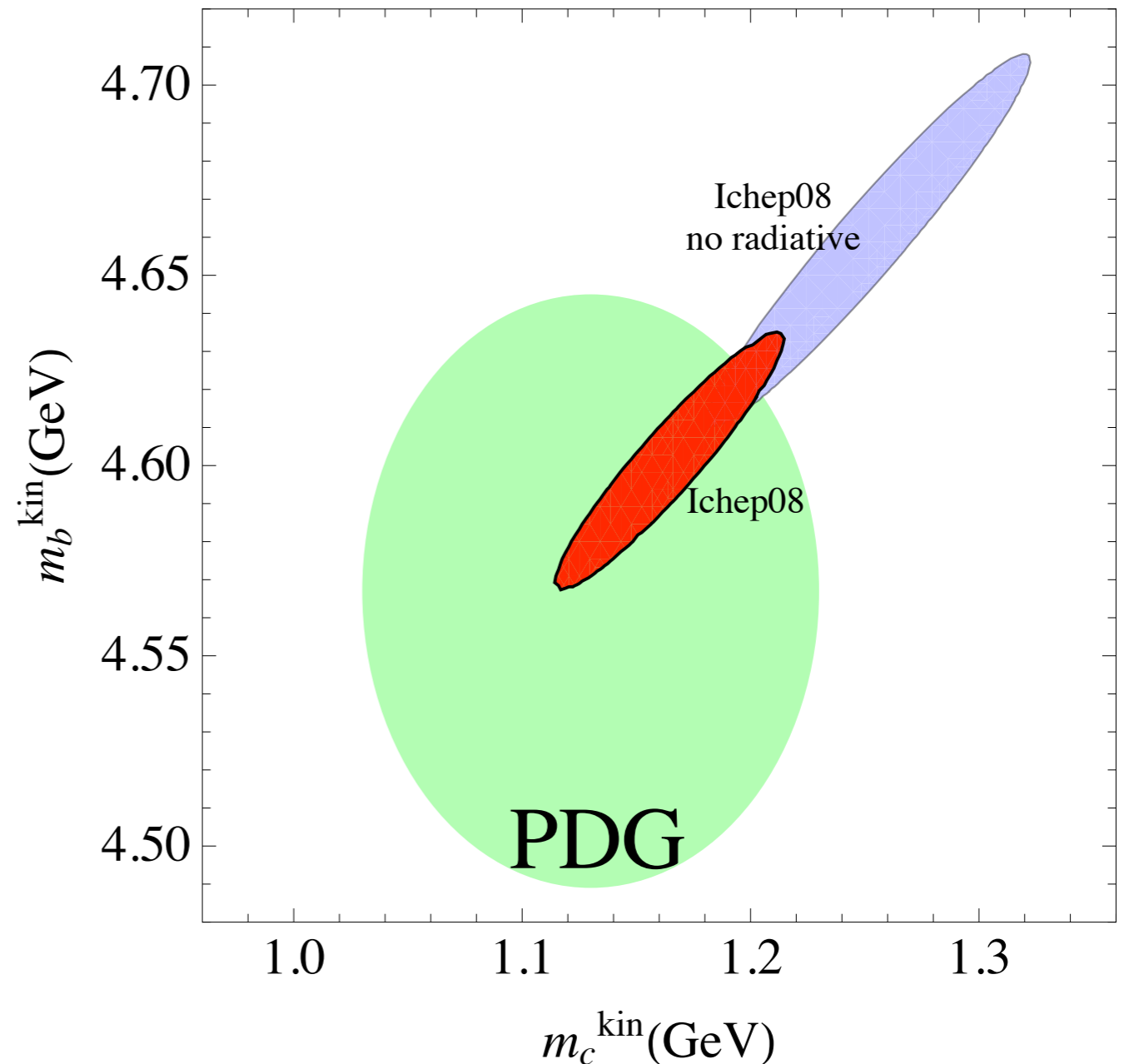
In the kinetic scheme the contributions of gluons with energy below $\mu \approx 1 \text{ GeV}$ are absorbed in the OPE parameters

Here scheme means also a number of different assumptions inclusion of different data, and a recipe for theory errors



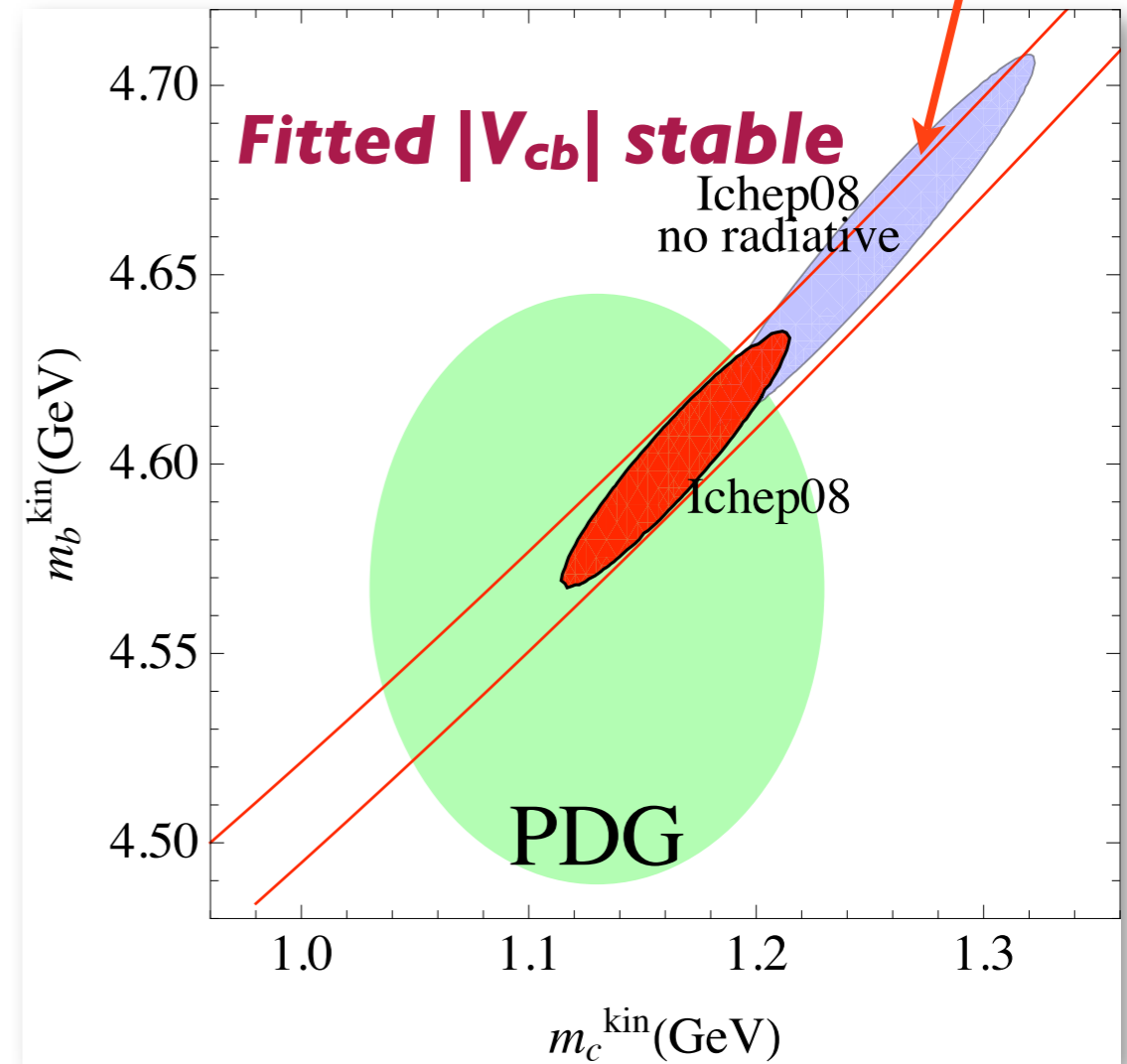
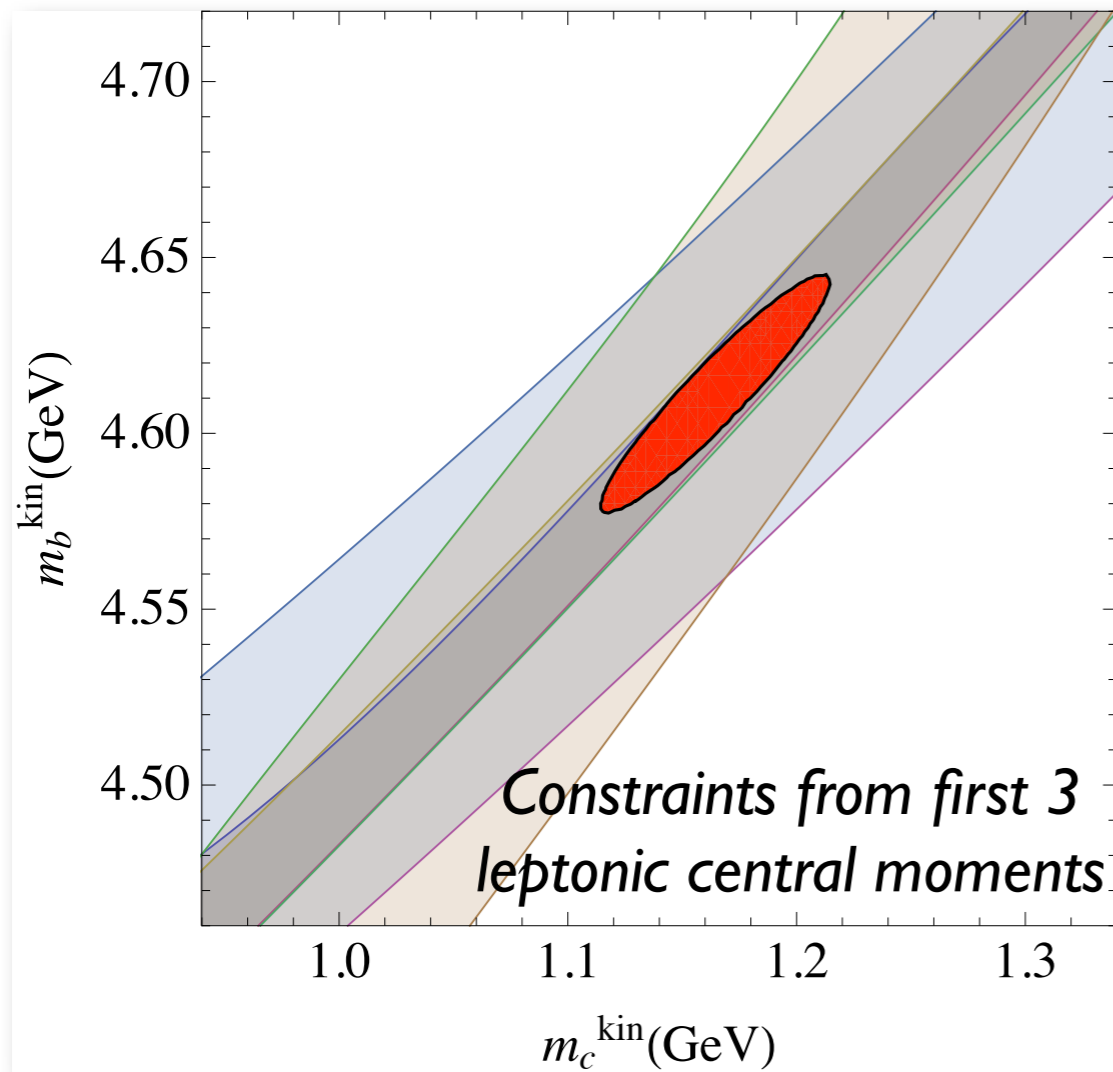
Fits & Quark Masses

- ▶ Assumes duality but it self-consistently checks it
- ▶ Very close result for $|V_{cb}|$ in $1S$ scheme (Bauer et al)
- Higher order power corr. under control Mannel et al
- New pert $O(\alpha_s^2) \Rightarrow -0.5\%$ in $|V_{cb}|$
Melnikov, Czarnecki, Pak
- Part of $O(\alpha_s/m_b^2)$ Becher et al
- ▶ New calculations give generally small contributions, will be included
- ▶ In the global HFAG fit the $B \rightarrow X_s \gamma$ moments **change significantly** $m_{b,c}$ determination. Without radiative moments the masses are too high!



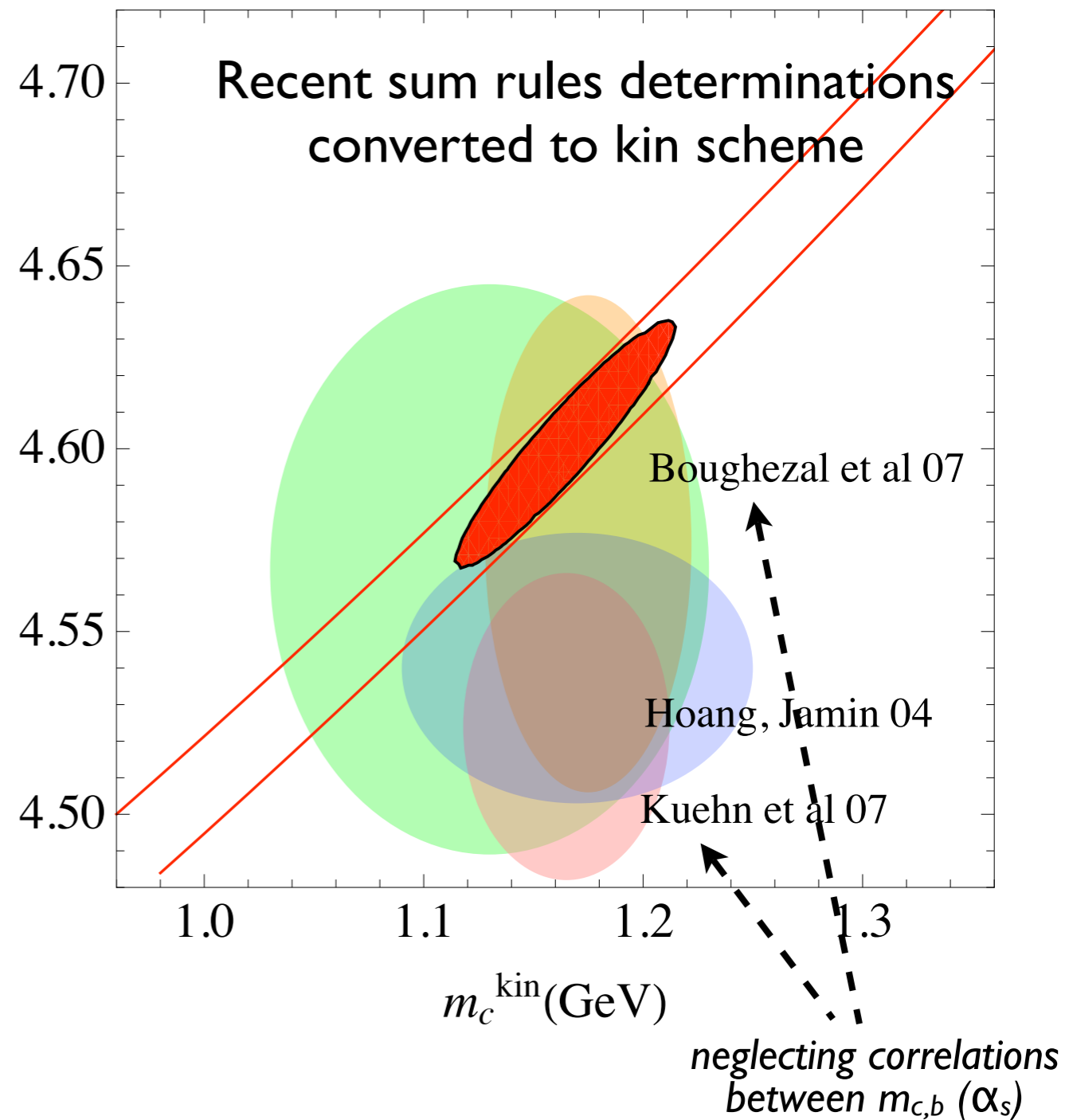
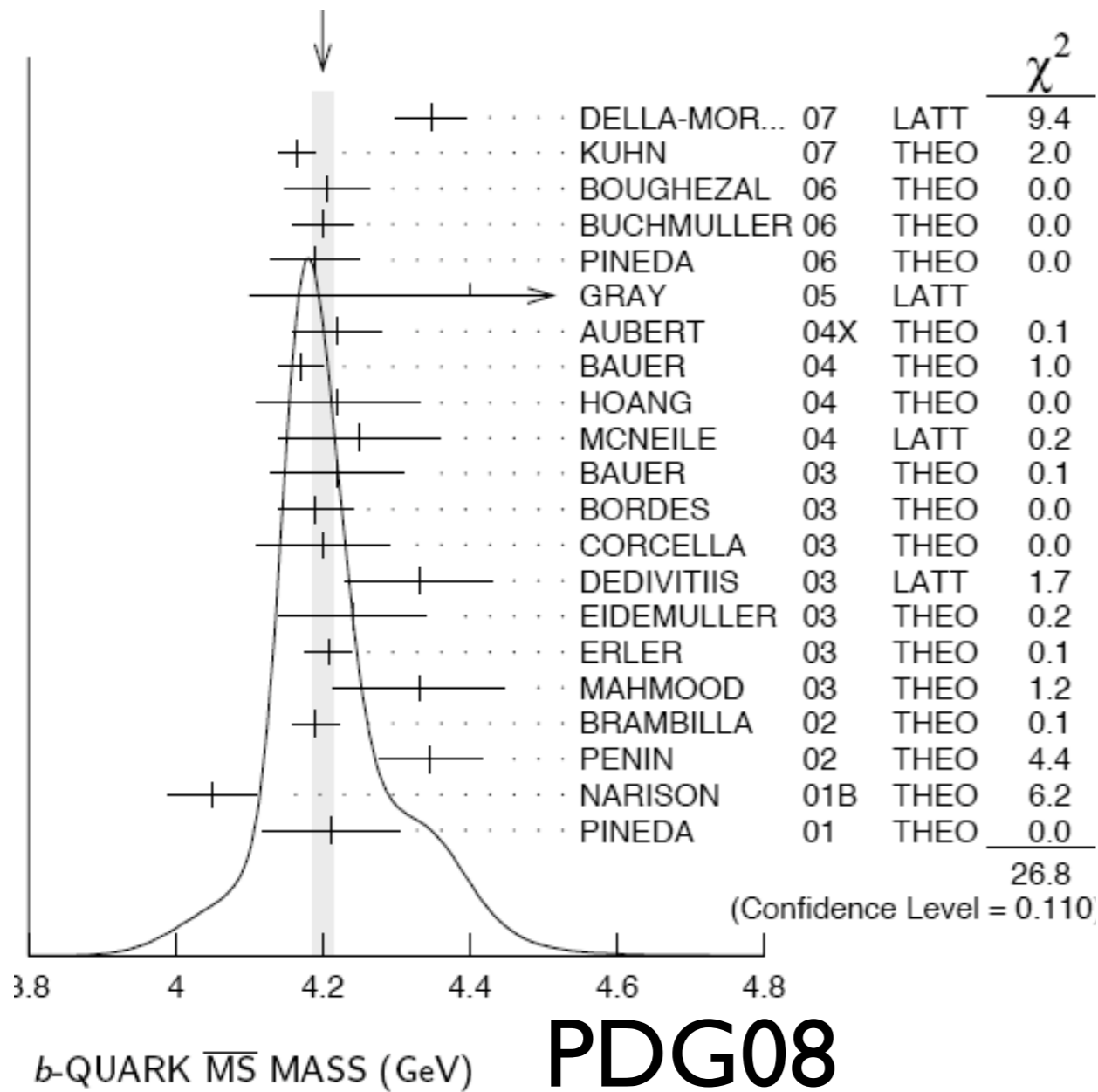
A strip in the m_b - m_c plane

Constant values
of s.l. width
at fixed V_{cb}



- ▶ Semileptonic moments identify a strip in (m_b, m_c) plane along which the minimum is **shallow**.
- ▶ V_{ub} inclusive studies require m_b , μ_π , etc with correlations
- ▶ m_b, m_c and OPE parameters necessary for $BR(B \rightarrow X_s \gamma)$ etc

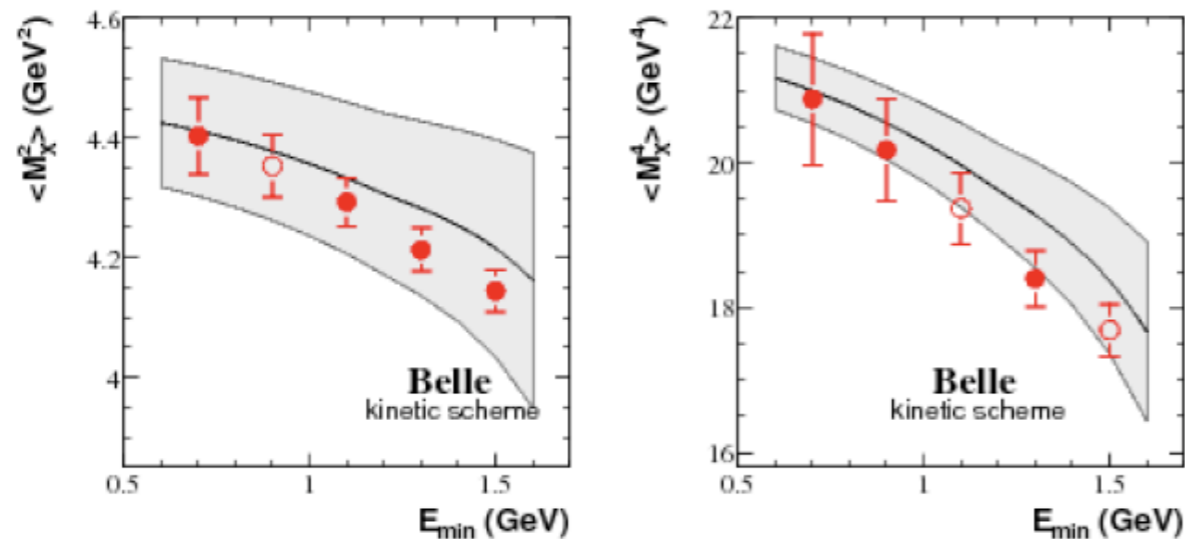
Mass determinations



How reliable are mass determinations?

In collaboration with C. Schwanda

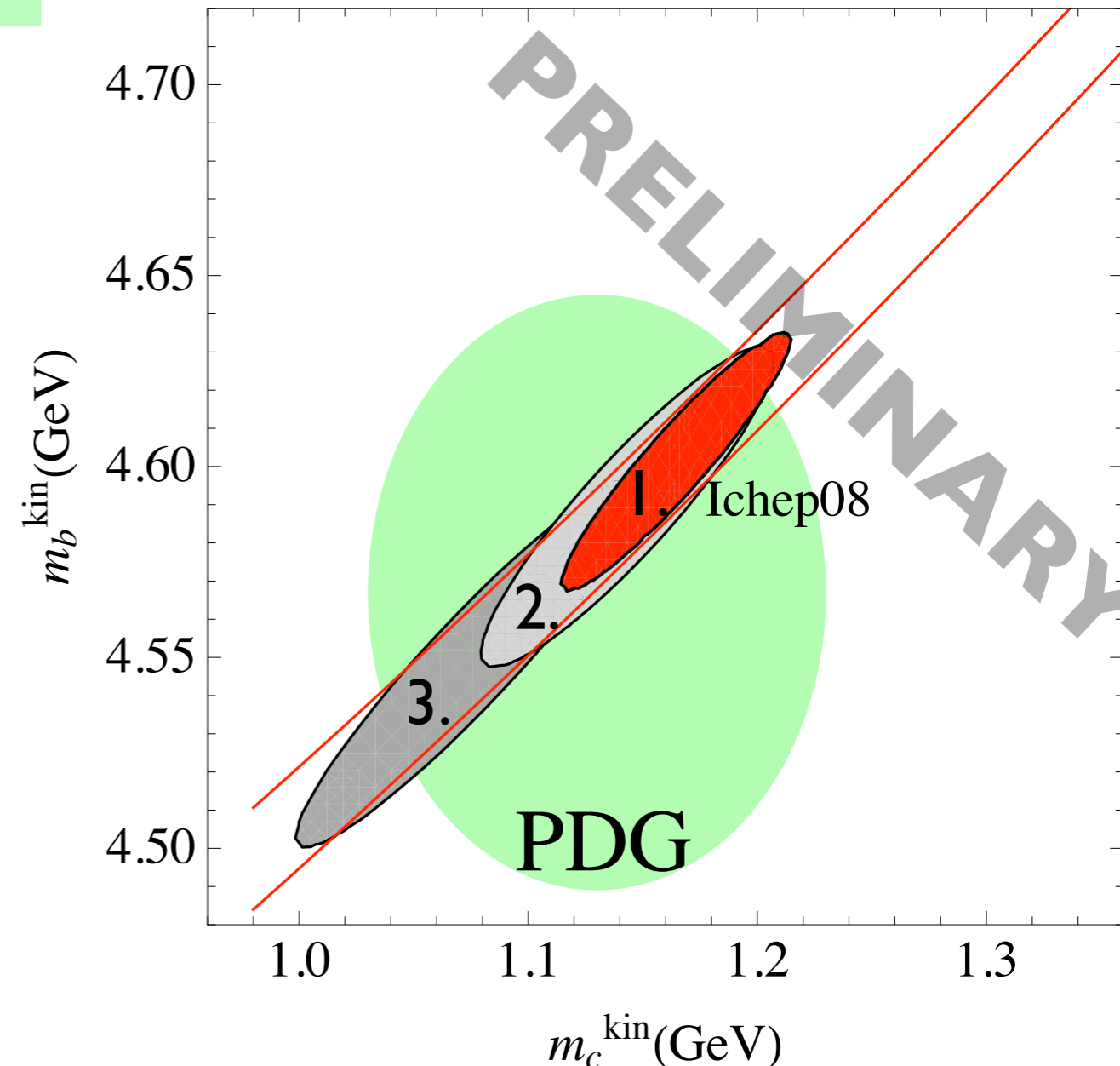
I. Theoretical correlations



Correlations between theory errors of moments with different cuts difficult to estimate. Examples:

1. 100% correlations
2. corr. computed from low-order
3. experimental correlations

always assume different central moments uncorrelated



2. How important are radiative moments?

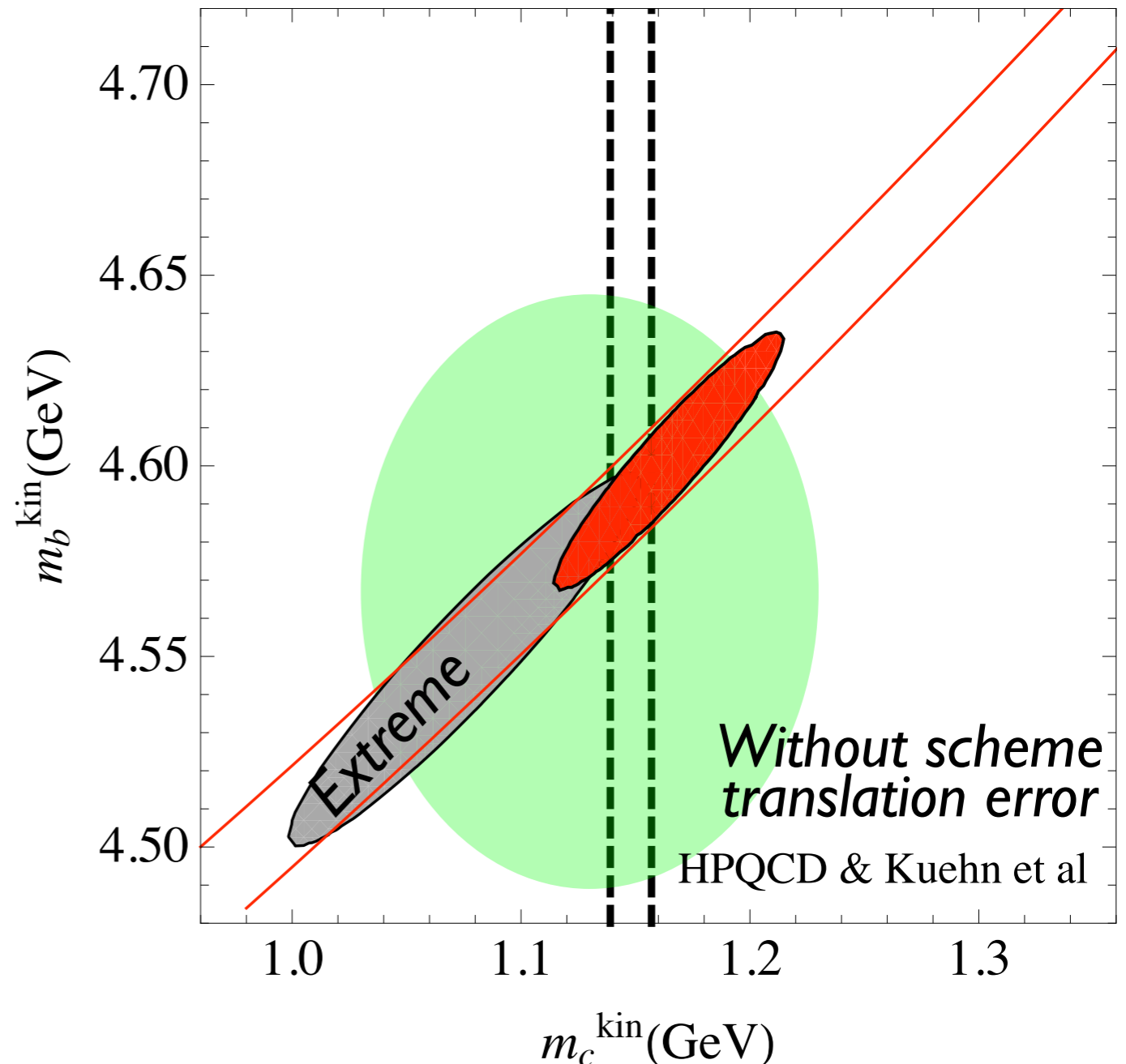
3. Can we include other constraints?

OPE fails for $bs\gamma$, but only at $O(\alpha_s)$ with operators $\neq O_7$. Unlikely to be relevant for normalized moments, but it must be studied

At the moment the role of radiative moments in the fits is almost **identical** to using PDG07 bound $m_b(m_b)=4.20(7)\text{GeV}$

Though $bs\gamma$ are important as independent checks, the inclusion of additional constraints is in principle very useful. But which ones?

Fits to $m_c(2-3\text{GeV})$ coming soon: no scheme translation error

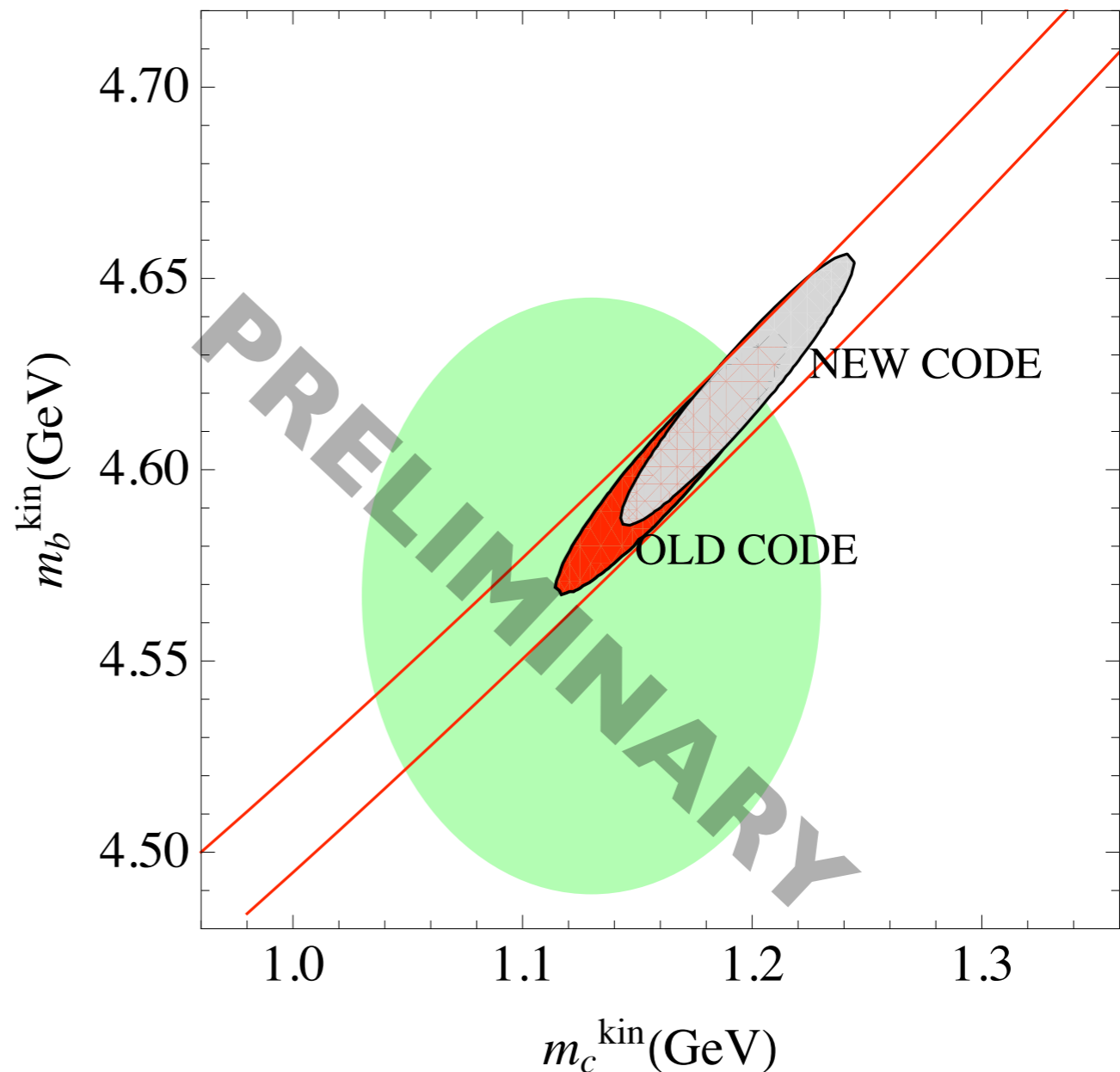


4. Fitting properly radiative moments

Since the $bs\gamma$ moments are measured with a relatively high cut on E_γ , a purely local OPE is insufficient.

SF can be implemented at NLO + BLM (Benson et al) but depends on m_b, μ^2_π, \dots

For the first time the fit is performed with the full parameter dependence.



Exclusive decays: $B \rightarrow D^* l \nu$

At zero recoil, where rate vanishes, the ff is

$$\mathcal{F}(1) = \eta_A(1 + \delta_{1/m^2})$$

Recent progress in the measurement of slopes and shape parameters *Despite extrapolation, exp error ~2%*

Main problem is normalization $F(1)$: **requires non-perturbative methods**

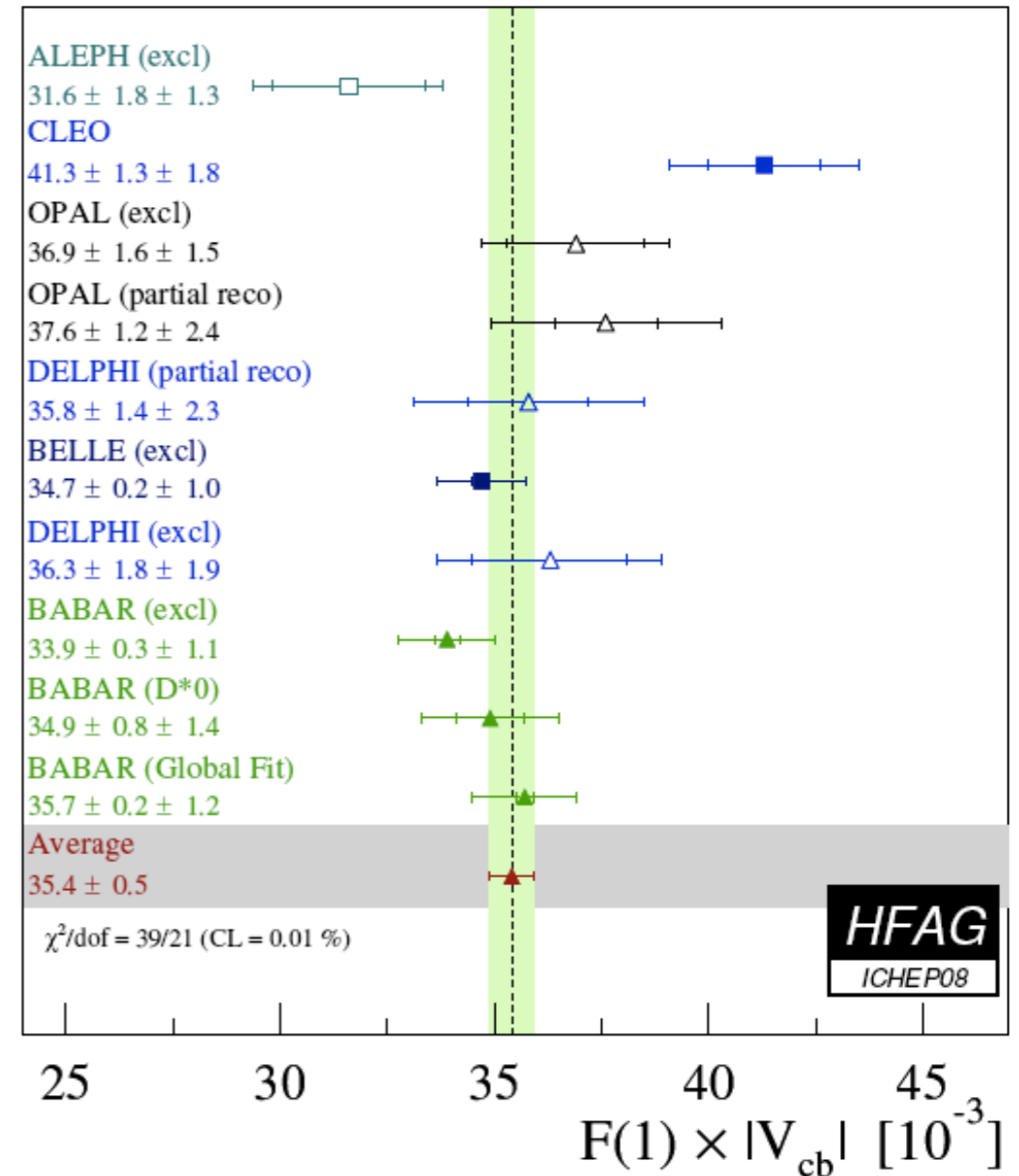
New and **only** unquenched Lattice QCD:

$F(1) = 0.921(24)$ Laiho et al 2008, HQET, double ratio

$$|V_{cb}| = 38.2(0.5)(1.1) \times 10^{-3}$$

$\sim 2.4\sigma$ from inclusive determination which would imply $F(1) = 0.857(13)$

$B \rightarrow D l \nu$ gives consistent but much less precise results



Lattice promising alternative: step scaling, w dependence, only quenched de Divitiis et al

Heavy Quark Sum Rules for $B \rightarrow D^* l \nu$

Heavy Quark Sum rules provide a (unitarity) bound on $F(1)$:

$$F_{D^*}^2 + \sum_{f \neq D^*} |F_{B \rightarrow f}|^2 = \xi_A^{\text{pert}} - \frac{\mu_G^2}{3m_c^2} - \underbrace{\frac{\mu_\pi^2 - \mu_G^2}{4}}_{>0} \left(\frac{1}{m_c^2} + \frac{1}{m_b^2} + \frac{2}{3m_c m_b} \right) - \Delta_{\frac{1}{m_Q^3}} + \Delta_{\frac{1}{m_Q^4}} + \dots$$

> 0

Numerically (preliminary)

$$\sqrt{\xi_A^{\text{pert}}} \simeq 0.96 \quad -\Delta_{\frac{1}{m_Q^2}} - \Delta_{\frac{1}{m_Q^3}} = -0.14(2) \quad \boxed{F_{D^*} \lesssim 0.90}$$

The quark masses and B meson expectation values measured in inclusive decays strongly disfavor $F(1) > 0.9$

$$\sum_{f \neq D^*} |F_{B \rightarrow f}|^2 = \chi \cdot \left(\Delta_{\frac{1}{m_Q^2}} + \Delta_{\frac{1}{m_Q^3}} + \dots \right) \quad F_{D^*} \simeq \sqrt{\xi_A^{\text{pert}}} - \frac{1}{2}(1 + \chi)\Delta$$

No reason to expect $\chi=0$, typically $\chi > 0.5$ or $F(1) < 0.87$

Uraltsev, Mannel, PG...

The total $B \rightarrow X_u l \bar{\nu}$ width in the OPE

$$\Gamma[\bar{B} \rightarrow X_u e \bar{\nu}] = \frac{G_F^2 m_b^5 |V_{ub}|^2}{192\pi^3} \left[1 + \frac{\alpha_s}{\pi} p_u^{(1)}(\mu) + \frac{\alpha_s^2}{\pi^2} p_u^{(2)}(r, \mu) - \frac{\mu_\pi^2}{2m_b^2} - \frac{3\mu_G^2}{2m_b^2} \right. \\ \left. + \left(\frac{77}{6} + 8 \ln \frac{\mu_{\text{WA}}^2}{m_b^2} \right) \frac{\rho_D^3}{m_b^3} + \frac{3\rho_{LS}^3}{2m_b^3} + \frac{32\pi^2}{m_b^3} B_{\text{WA}}(\mu_{\text{WA}}) \right] \\ + O\left(\alpha_s \frac{\mu_{\pi, G}^2}{m_b^2}\right) + O\left(\frac{1}{m_b^4}\right)$$

Weak
Annihilation

Yes, life would be MUCH easier with the total width...

The problems with cuts

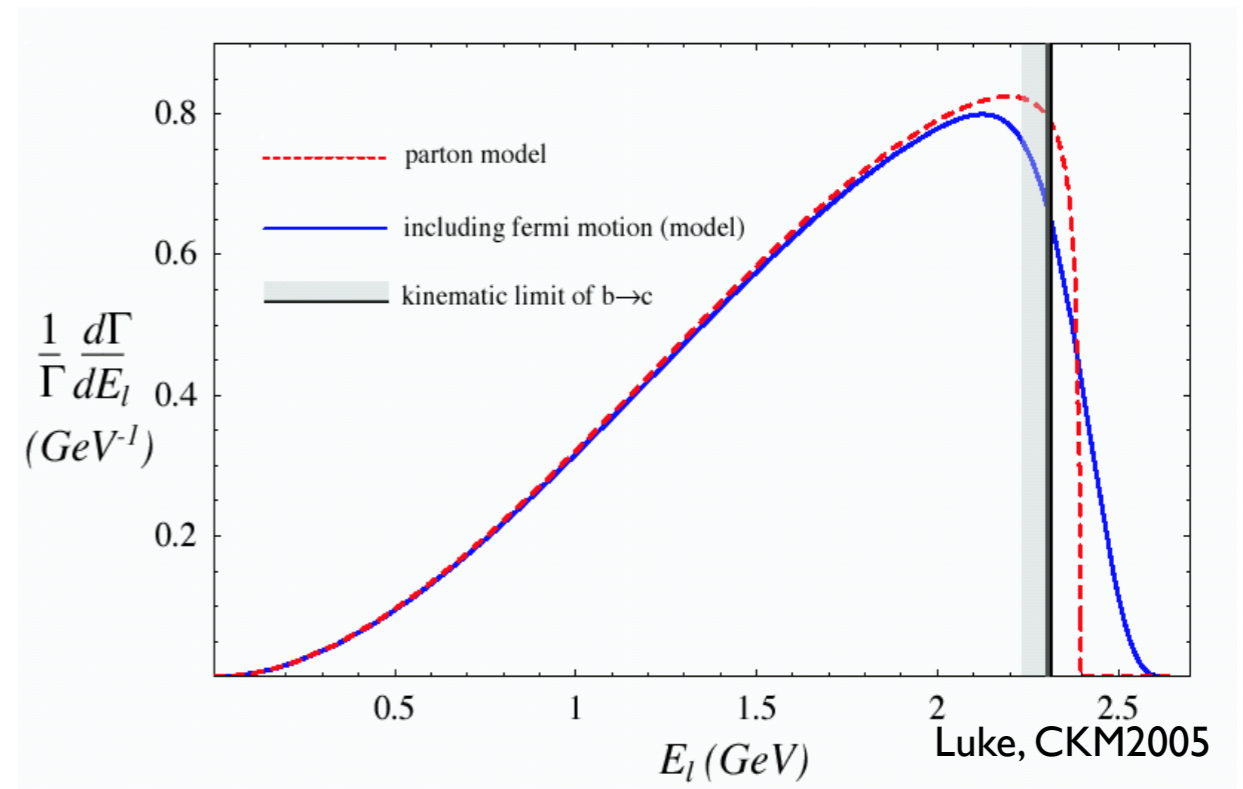
$|V_{ub}|$ from total BR($b \rightarrow ul\nu$) like incl $|V_{cb}|$ but we need kinematic cuts to avoid the $\sim 100x$ larger $b \rightarrow cl\nu$ background:

$$m_X < M_D \quad E_l > (M_B^2 - M_D^2) / 2M_B \quad q^2 > (M_B - M_D)^2 \dots$$

or combined (m_X, q^2) cuts

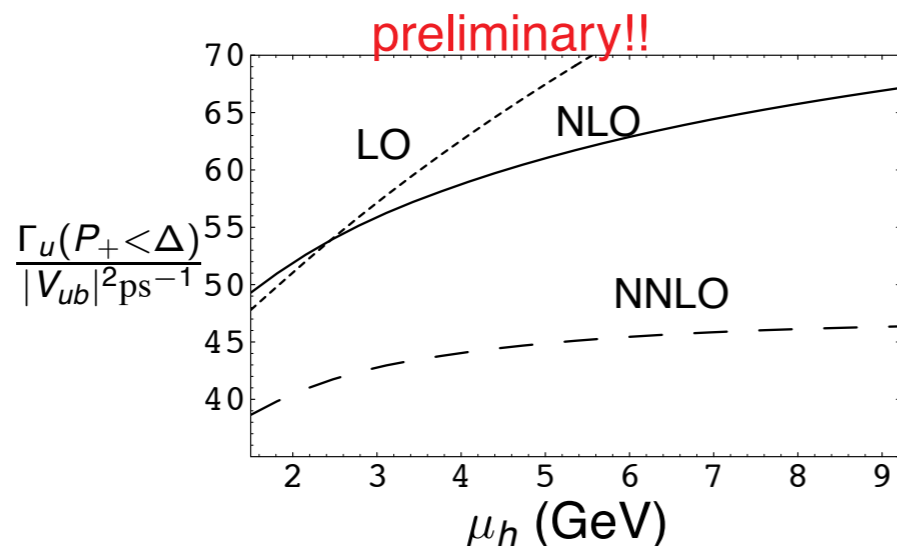
*The cuts destroy convergence of the OPE that works so well in $b \rightarrow c$.
OPE expected to work only away from
pert singularities*

Rate becomes sensitive to “local”
b-quark wave function properties
like Fermi motion Dominant non-
pert contributions can be resummed
into a **SHAPE FUNCTION** $f(k_+)$



Perturbative calculations

Partial rate for $P_+ < \Delta = M_D^2/M_B$



- ▶ NNLO result is smaller and less dependent on μ_h than NLO
- ▶ would lead to higher $|V_{ub}|$ compared to NLO (preliminary)

Some of the shift is due to different S at LO, NLO, NNLO

Ben Pecjak, ICHEP08

Complete $O(\alpha_s)$ implemented by all groups De Fazio-Neubert

Complete running coupling NNLO $O(\alpha_s^2\beta_0)$ Gardi,Ridolfi, PG

in GGOU & DGE lead to -5% & +2%, resp. in $|V_{ub}|$

2008

Asatrian, Greub, Pecjak

Bonciani, Ferroglia,

Beneke, Huber, Li

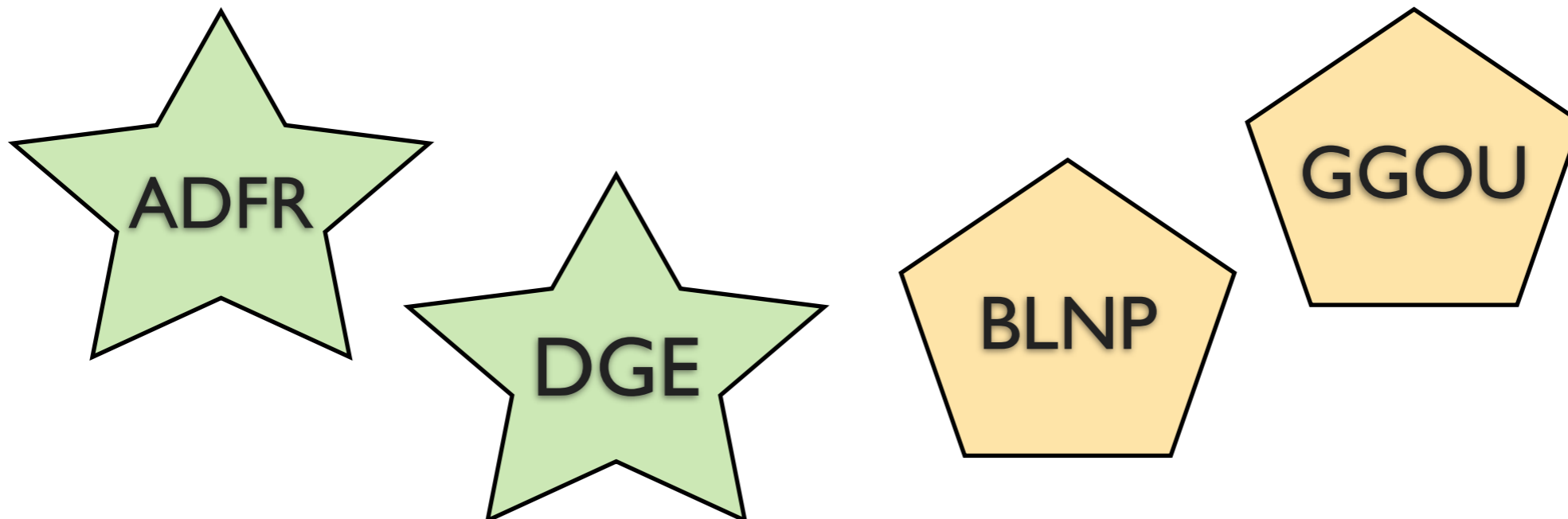
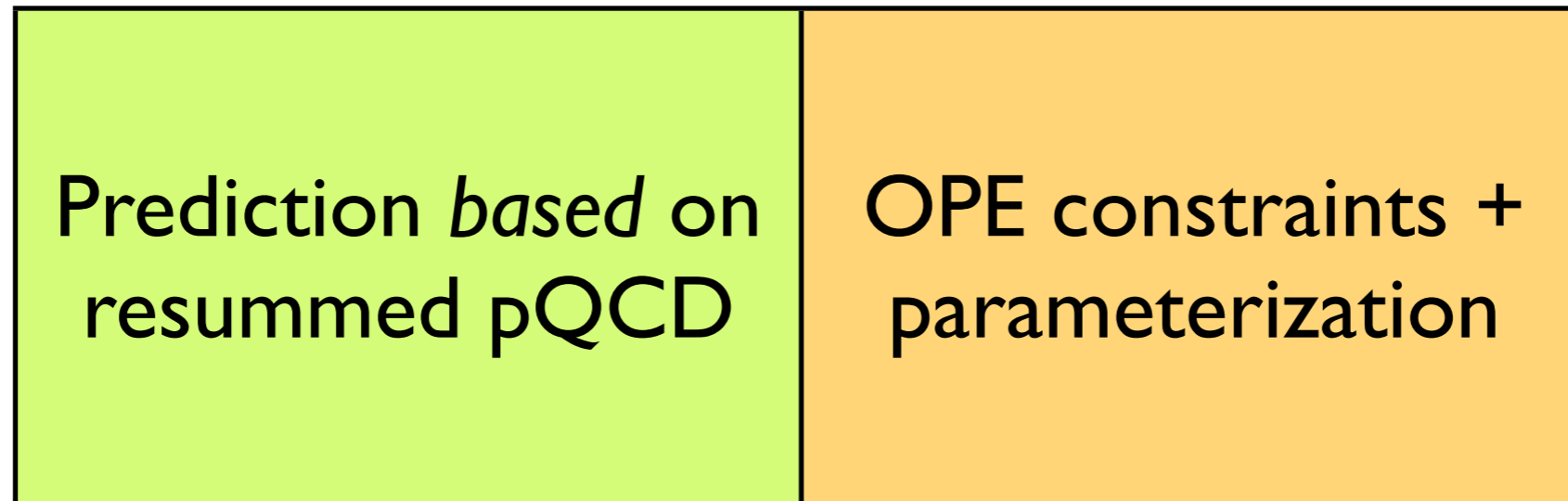
G. Bell

in SCET-HQET

corresponds to fixed order

$O(\alpha_s^2)$ in the SF region

How to access the SF?



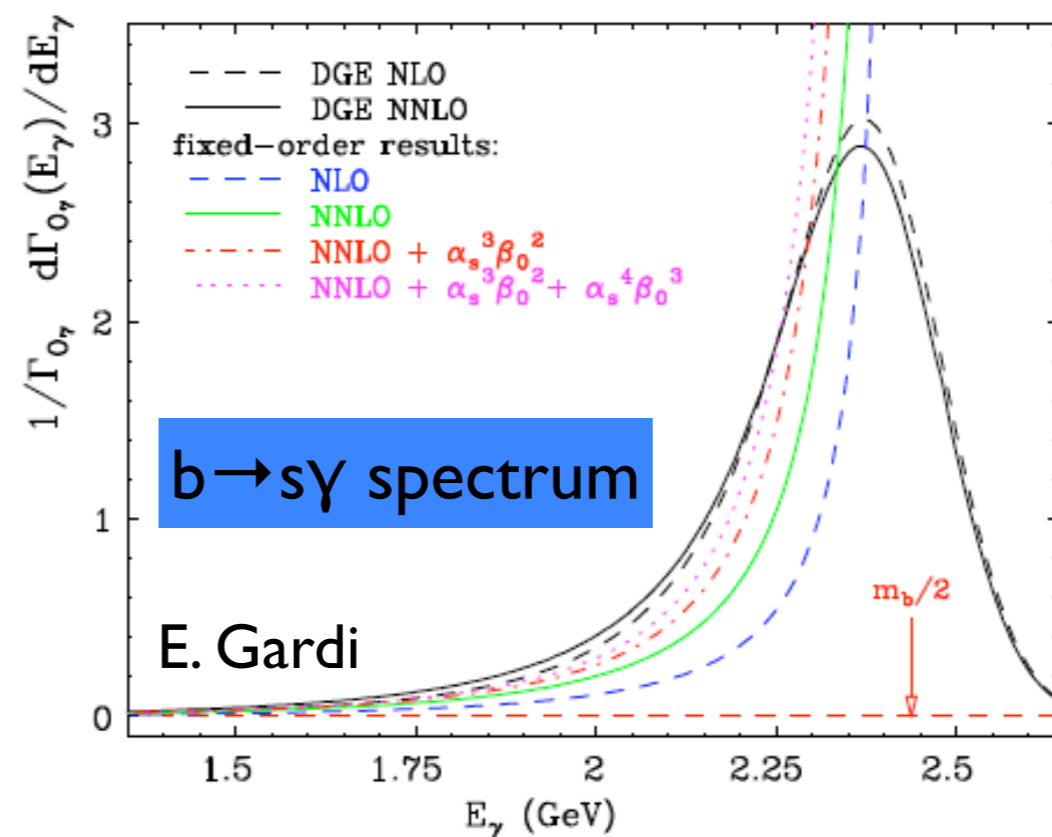
SF from perturbation theory

Resummed perturbation theory is qualitatively different: **Support properties; stability!** (E. Gardi)

b quark SF emerges from resummed pQCD but needs an IR prescription and power corrections for $b \rightarrow B$

Dress Gluon Exponentiation (DGE) by Gardi et al employs renormalon resummation to define Fermi motion. Power corrections can be partly accommodated.

Aglietti et al (ADFR) use Analytic Coupling in the IR, a model



The SF in the OPE

Local OPE has also threshold singularities and SF can be equivalently introduced resumming dominant singularities *Bigi et al, Neubert*

Fermi motion can be parameterized within the OPE like PDFs in DIS. At leading order in m_b only a single universal function of one parameter enters (SF).

Unlike resummed pQCD, **the OPE does not predict the SF**, only its first few moments. One then **needs an ansatz for its functional form**.

$$\int dk_+ k_+^n F_i(k_+, q^2) = \text{local OPE prediction} \Leftarrow \text{moments fits}$$

Two very different implementations:
PG, Giordano, Ossola, Uraltsev (GGOU)
Bosch, Lampe, Neubert, Paz (BLNP)

The SF in GGOU

Leading SF resums leading twist effects, $m_b \rightarrow \infty$
universal, q^2 indep



Finite m_b distribution functions include all $1/m_b$ effects, *non-universal*
no need for subleading SFs

$$F(k_+) \longrightarrow F_i(k_+, q^2, \mu)$$

Structure function ($i = 1, 2, 3$) q^2 dependence cutoff dependence (gluons with $E_g < \mu$)

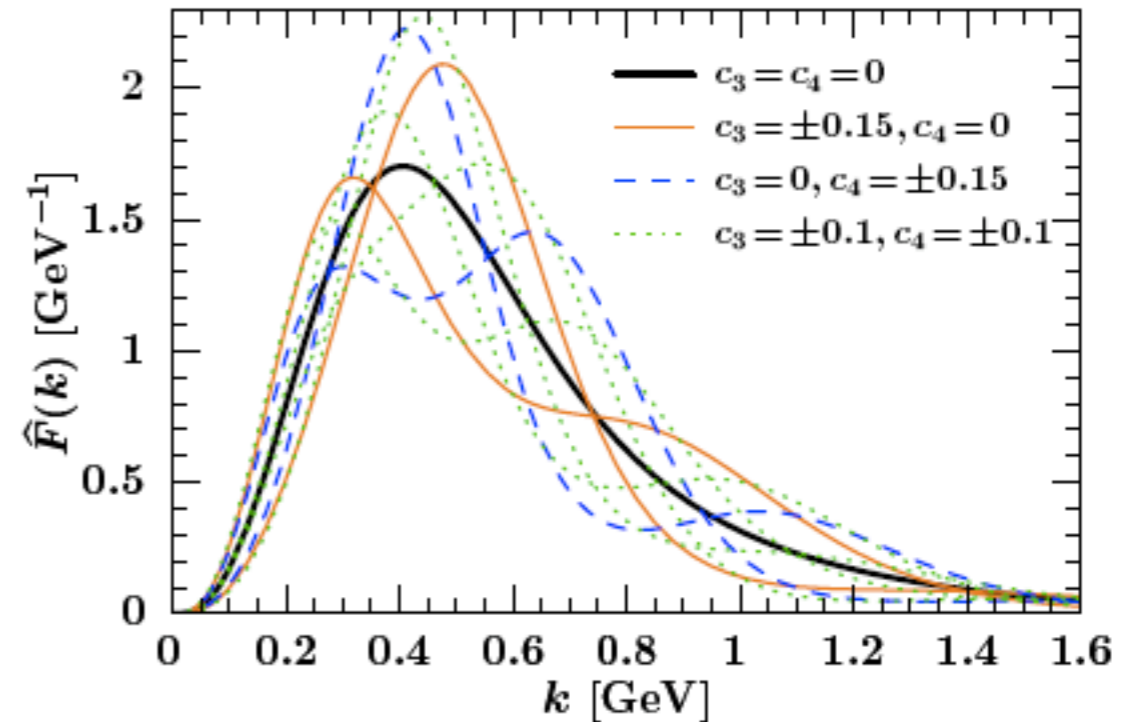
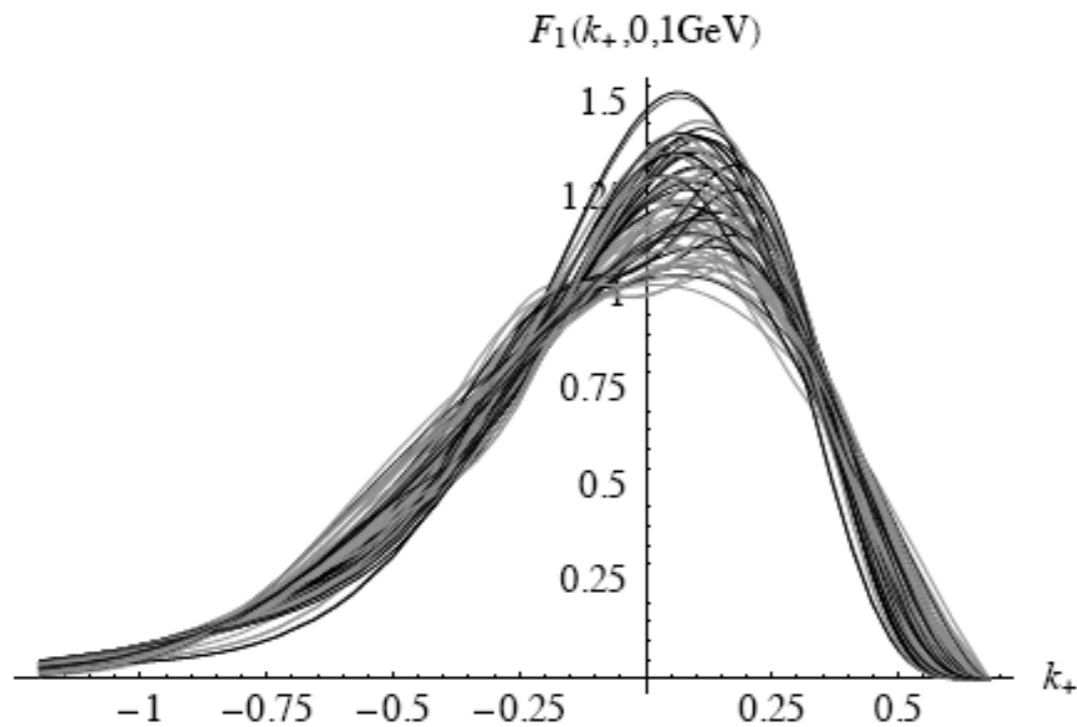
$$\frac{d^3\Gamma}{dq^2 dq_0 dE_\ell} = \frac{G_F^2 |V_{ub}|^2}{8\pi^3} \left\{ q^2 W_1 - \left[2E_\ell^2 - 2q_0 E_\ell + \frac{q^2}{2} \right] W_2 + q^2 (2E_\ell - q_0) W_3 \right\}$$

$$W_i(q_0, q^2) = m_b^{n_i}(\mu) \int dk_+ F_i(k_+, q^2, \mu) W_i^{\text{pert}} \left[q_0 - \frac{k_+}{2} \left(1 - \frac{q^2}{m_b M_B} \right), q^2, \mu \right]$$

This factorization formula perturbatively defines the distribution functions
see also Benson, Bigi, Uraltsev for bsy

$$\int dk_+ k_+^n F_i(k_+, q^2) = \text{local OPE} \quad \text{Importance of subleading effects}$$

Functional forms



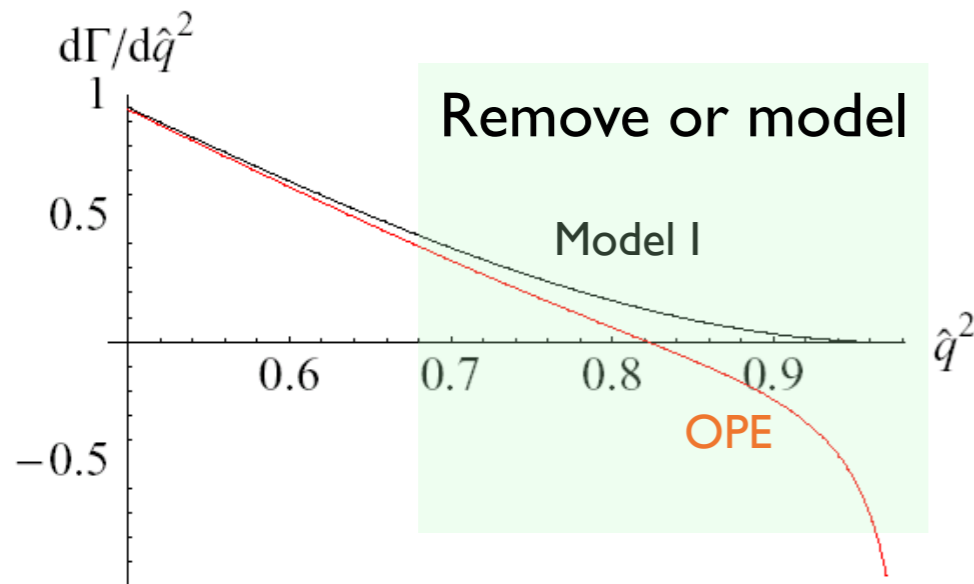
About 100 forms considered in GGOU, large variety, double max discarded. Small uncertainty (1-2%) on V_{ub}

Recent more systematic method by Ligeti et al. arXiv:0807.1926
Plot shows 9 SFs that satisfy all the first three moments

The high q^2 tail

At high q^2 higher dimensional operators are not suppressed leading to pathological features. Origin in the non-analytic square root

$$\frac{d\Gamma}{dq_0 dq^2} \propto \sqrt{q_0^2 - q^2} \quad \Rightarrow \quad \frac{d\Gamma}{dq^2} \sim - \sum_{n=1}^{\infty} \frac{(-1)^n b_n(\hat{q}^2)}{(1 - \hat{q}^2)^{n-2}} \left(\frac{\bar{\Lambda}}{m_b}\right)^n$$



In the integrated rate the $1/m_b^3$ singularity is removed by the WA operator: needs modelling for q^2 spectrum

$$\delta\Gamma \sim \left[C_{\text{WA}} B_{\text{WA}}(\mu_{\text{WA}}) - \left(8 \ln \frac{m_b^2}{\mu_{\text{WA}}^2} - \frac{77}{6} \right) \frac{\rho_D^3}{m_b^3} + \mathcal{O}(\alpha_s) \right]$$

WA matrix element B_{WA} parameterizes global properties of the tail, affects V_{ub} depending on cuts, tends to decrease V_{ub} , may pollute all present determinations

Comparing the existing approaches at common m_b

(HFAG ichep08, CKM08)

$$m_b^{\text{kin}} = 4.601(35)\text{GeV} \implies m_b(m_b) = 4.23\text{GeV}$$

$$\mu_\pi^2 = 0.440(40)\text{GeV}^2$$

very strong dependence on m_b ,
twice larger than in total rate

$|V_{ub}|$ from DGE

Gardi & Andersen

Main features of the spectra are reproduced $\implies |V_{ub}|$ stable, small errors and good χ^2

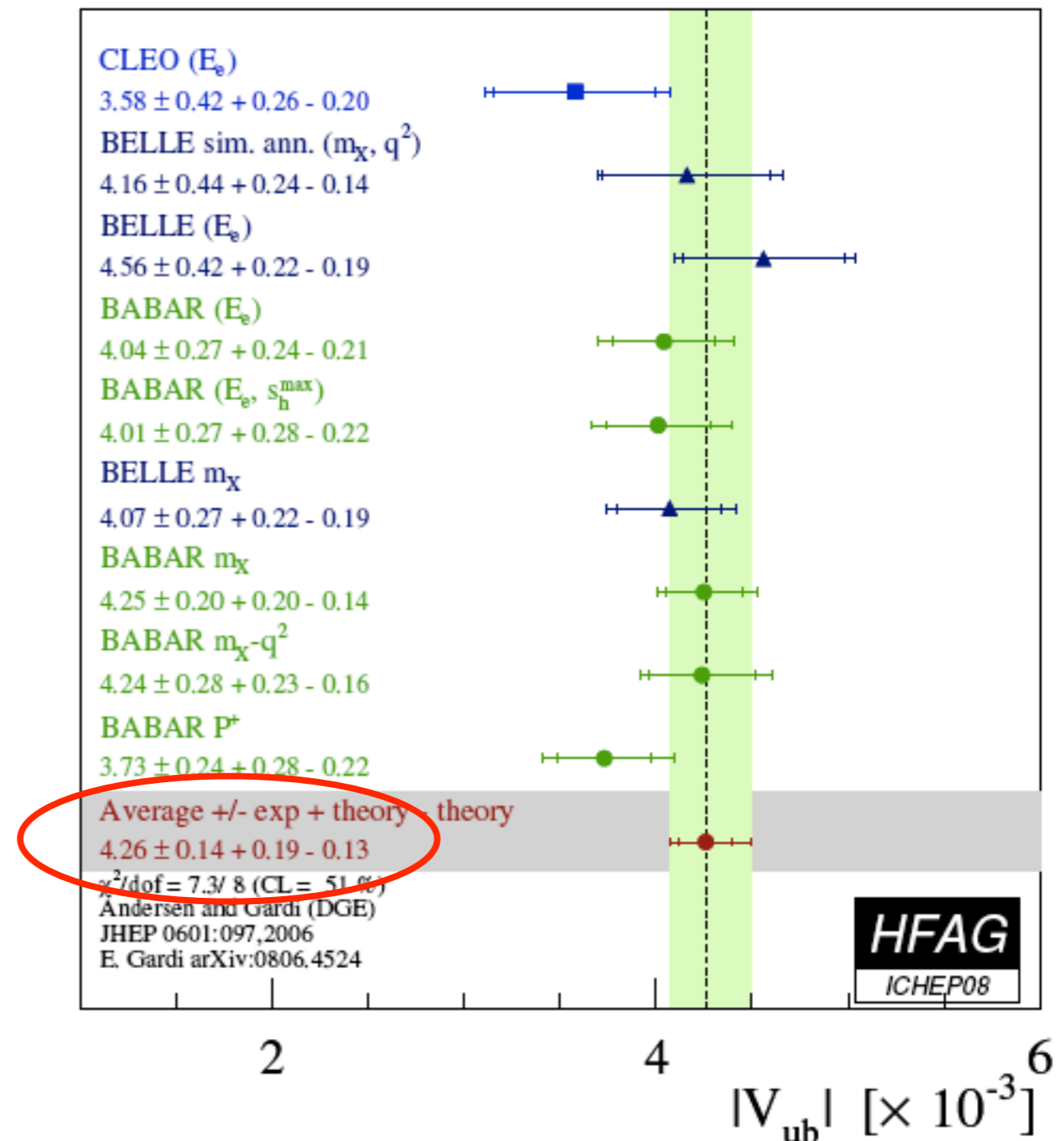
NNLL and $O(\alpha_s^2\beta_0)$ implemented

Power corrections in the SF region are included here only in theor. err. No subleading SF.

Matches to local OPE.

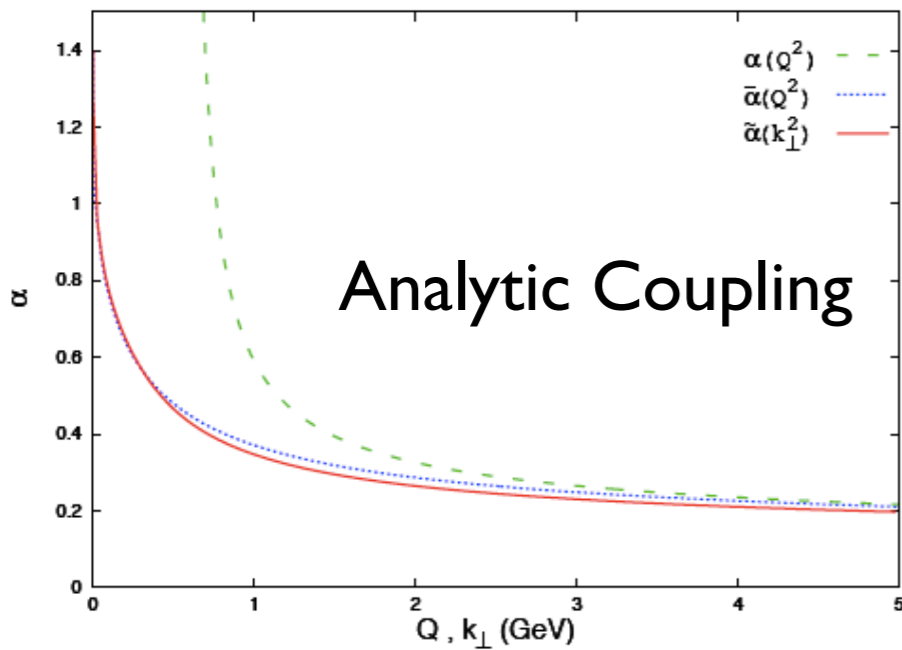
Only input other than α_s
 $m_b(m_b)=4.24(4)$ from global fit

5-6% total error, mostly m_b



$|V_{ub}|$ from ADFR

Aglietti, Di Lodovico, Ferrera, Ricciardi



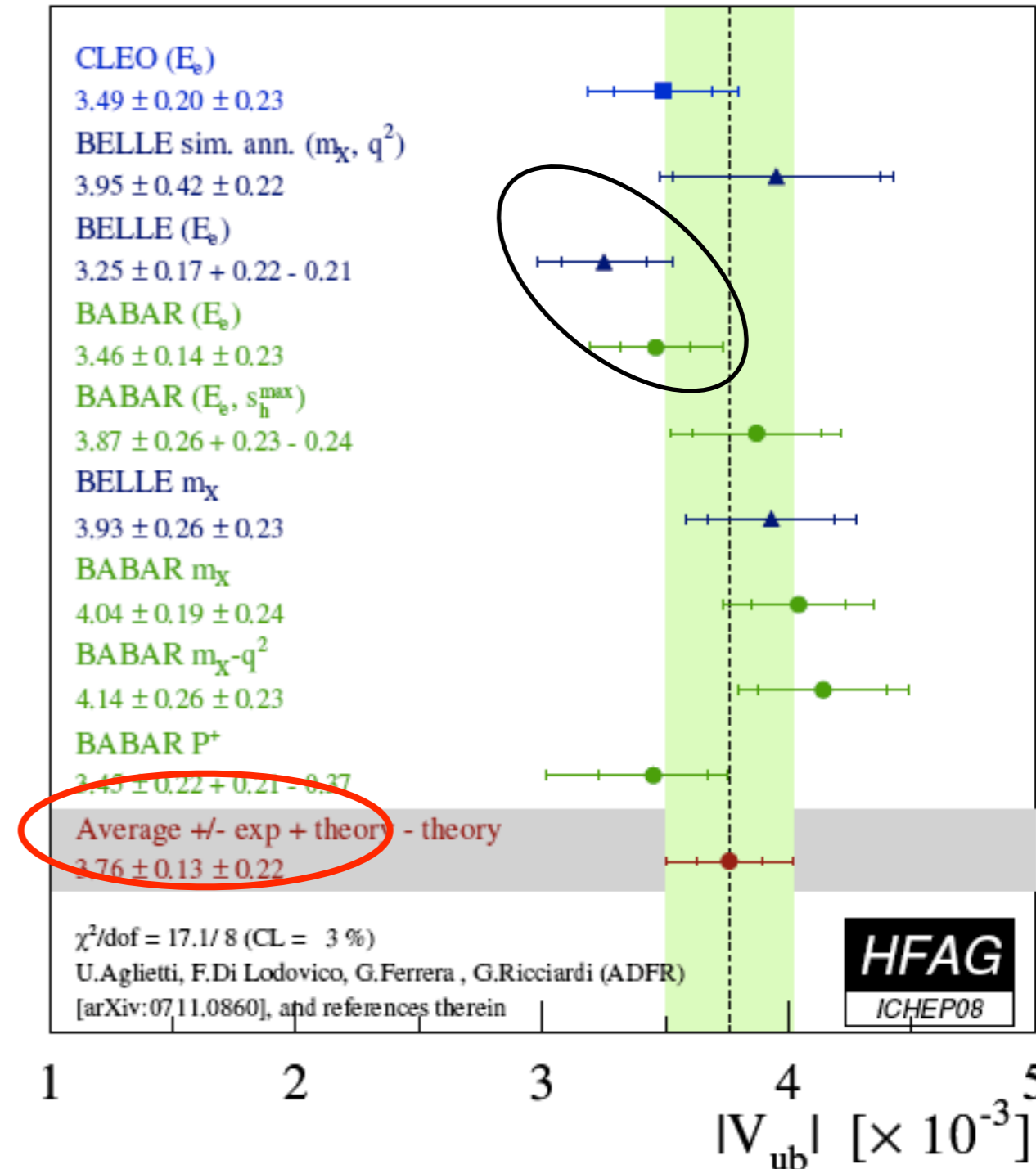
Worse consistency here.

NNLO resummation, NLO constants

Consider E_l cuts higher than 2.3 GeV because their E_l apparently does not reproduce data (see later)

employs M_B in on-shell calculation of spectra: no renormalon cancellation, no convergence to OPE. no model error

~7% total error, mostly m_c



$|V_{ub}|$ in BLNP

Bosch, Lange, Neubert, Paz

$$\tilde{W}_1^{(0)}(P_+, y) = U_y(\mu_h, \mu_i) H(y, \mu_h) \int_0^{P_+} d\hat{\omega} m_b J(y, m_b(P_+ - \hat{\omega}), \mu_i) \hat{S}(\hat{\omega}, \mu_i)$$

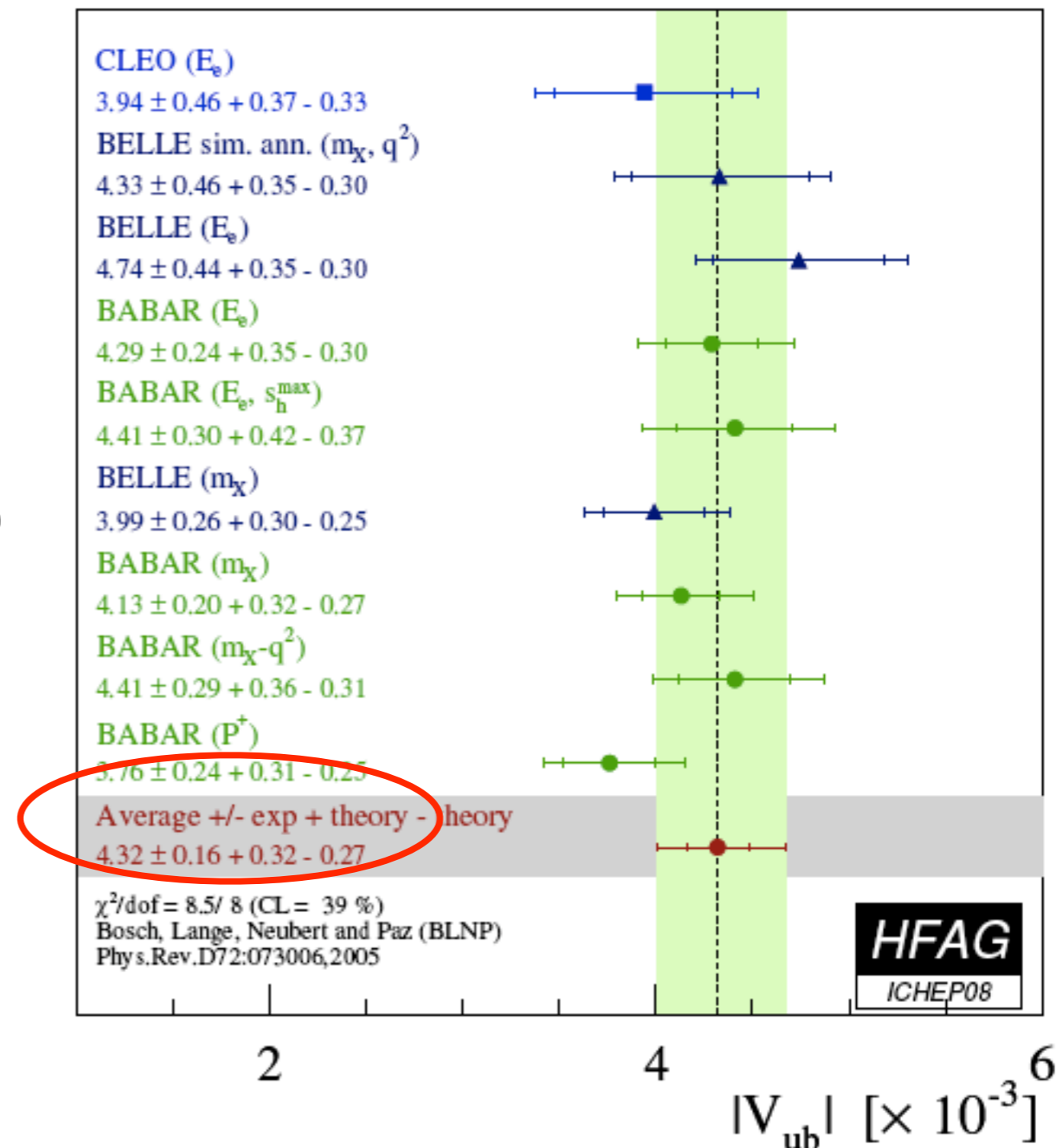
$$d\Gamma = HJ \otimes \hat{S} + \frac{1}{m_b} H'_i J'_i \otimes \hat{S}'_i + \dots$$

Good consistency. Uses elegant multiscale OPE that resums soft-collinear logs, but many largely unconstrained subleading SFs

NNLL resummation, only $O(\alpha_s, \Lambda^2/m_b^2)$ matching to OPE, 3 ffs for leading SF, extensive modelling of SSF.

m_b and μ_π^2 in SF scheme obtained from global fit in the kin scheme

~7-8% total error, main error HQE parameters



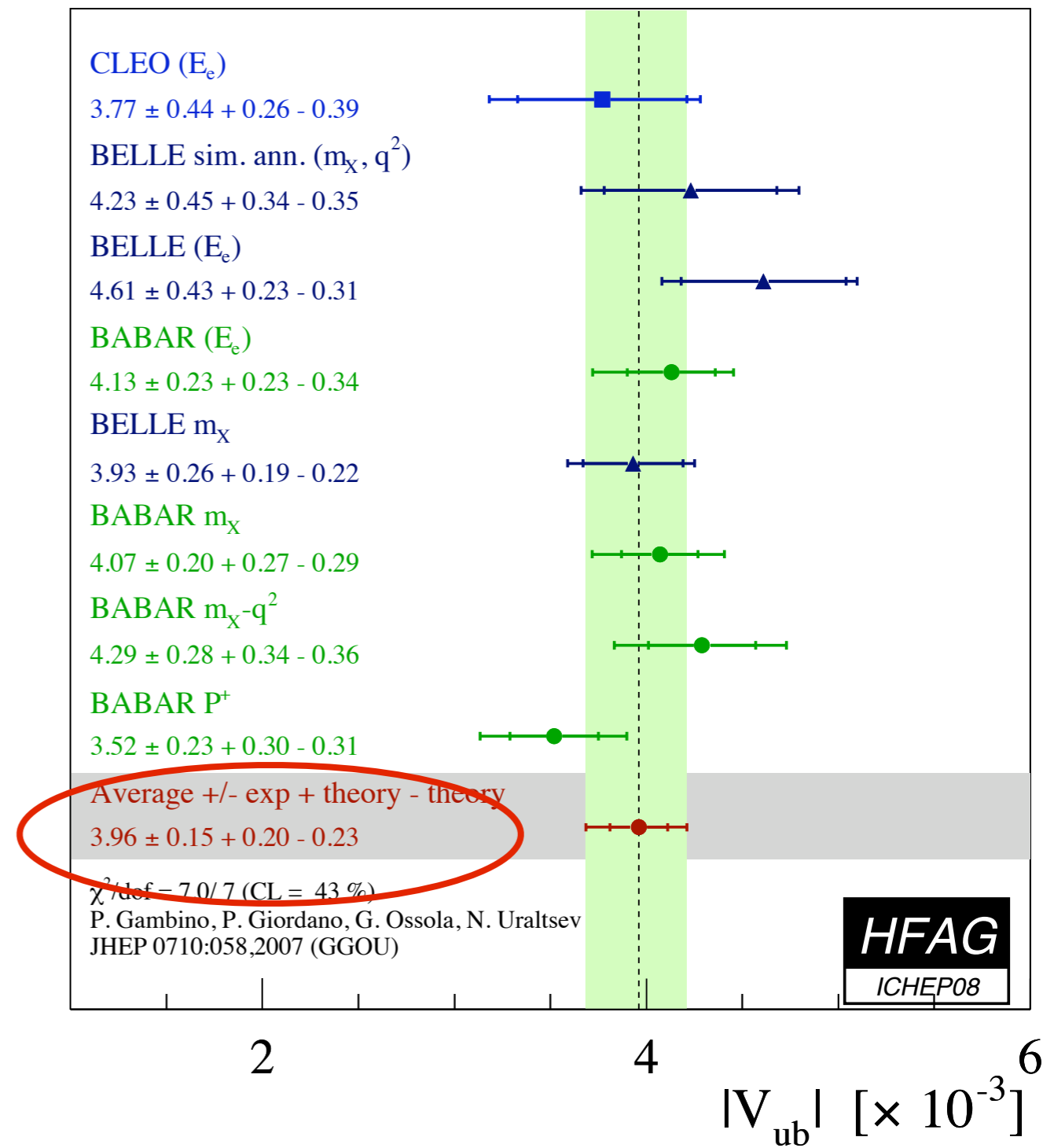
$|V_{ub}|$ in the kinetic scheme -GGOU

PG, Giordano, Ossola, Uraltsev

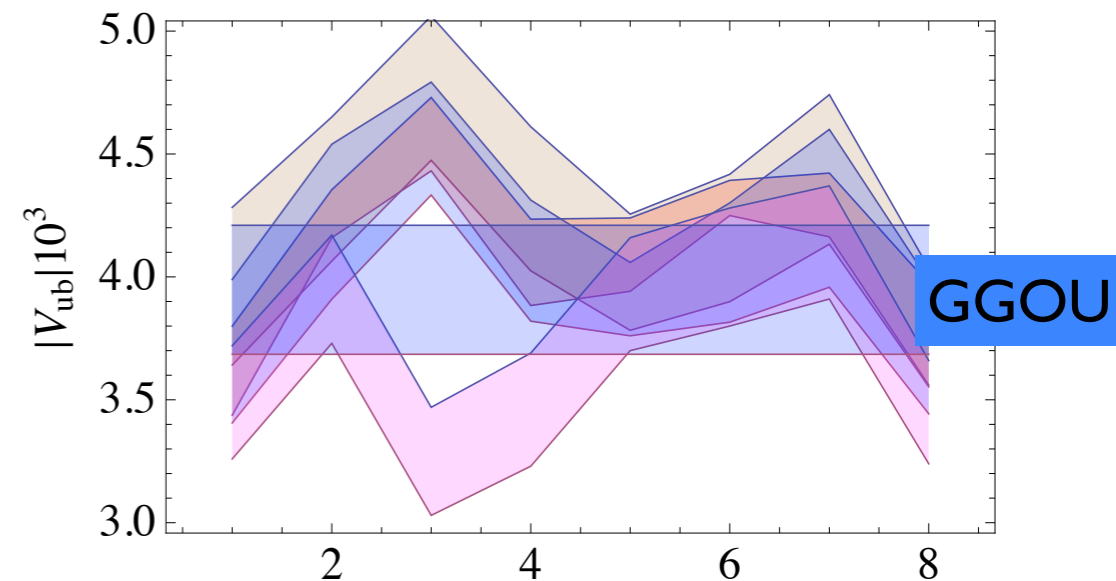
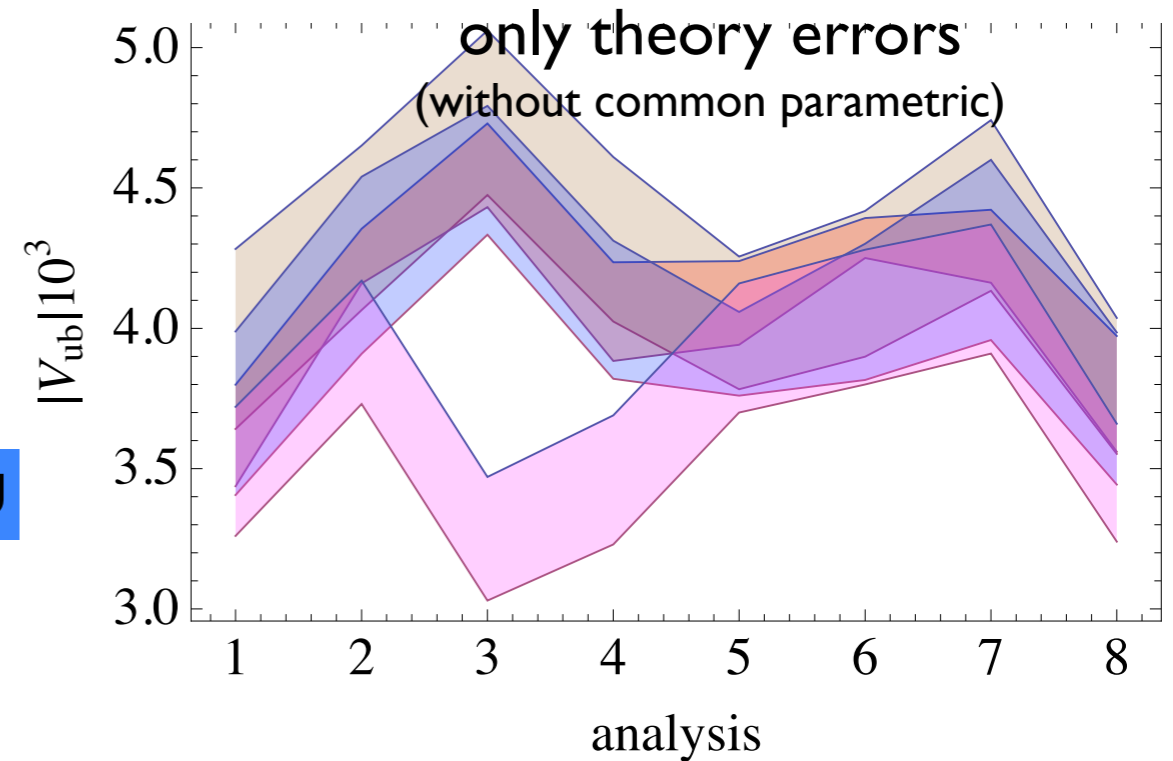
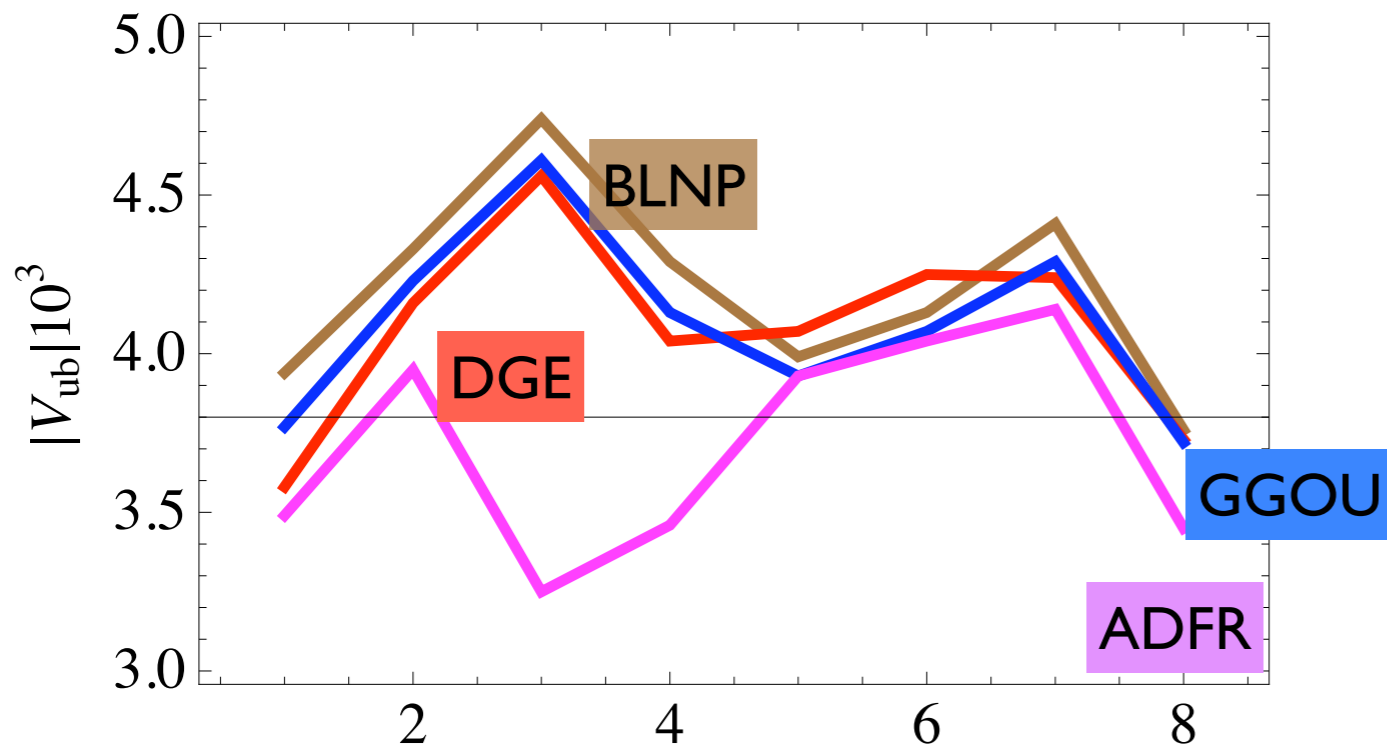
Good consistency & small th error.
 OPE in a scheme with Wilsonian IR cutoff
 $\sim 1 \text{ GeV}$, all subleading $1/m_b$ and $O(\alpha_s^2 \beta_0)$
 terms consistently included,
 careful treatment of high q^2 tail.

Inputs from global fit to the moments

+6.3-7.0% total error



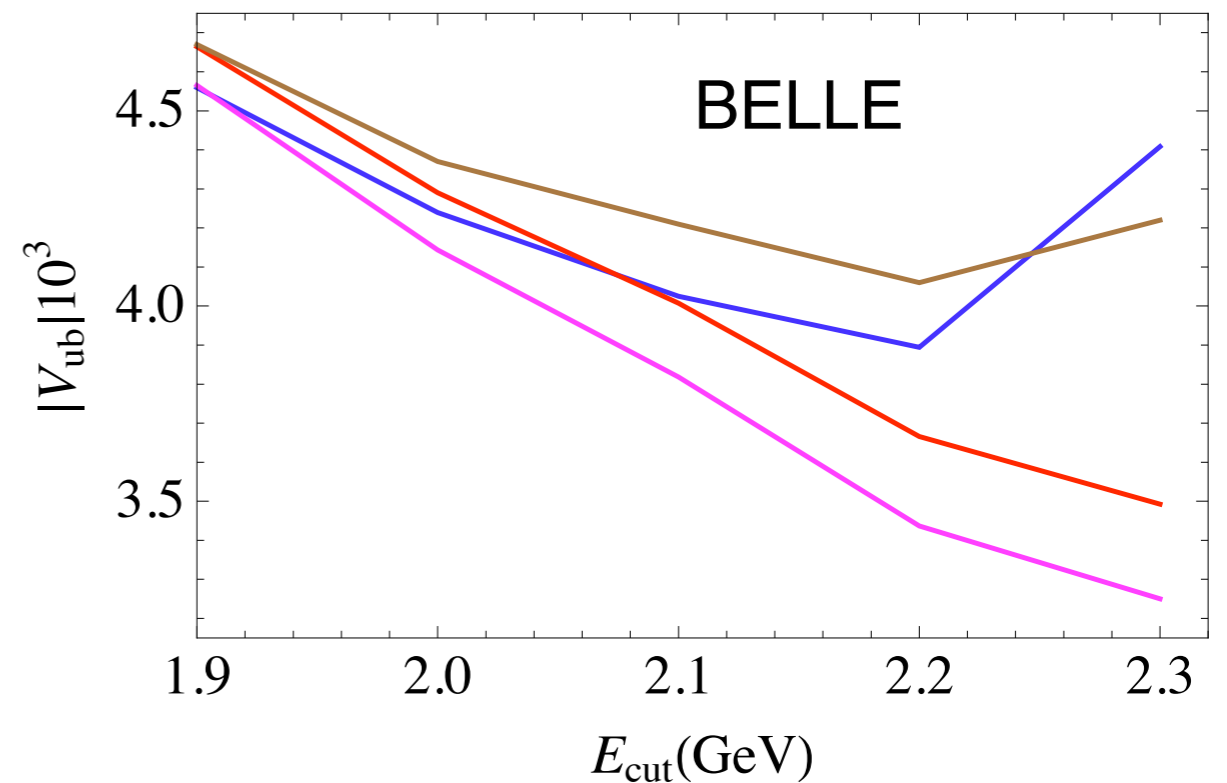
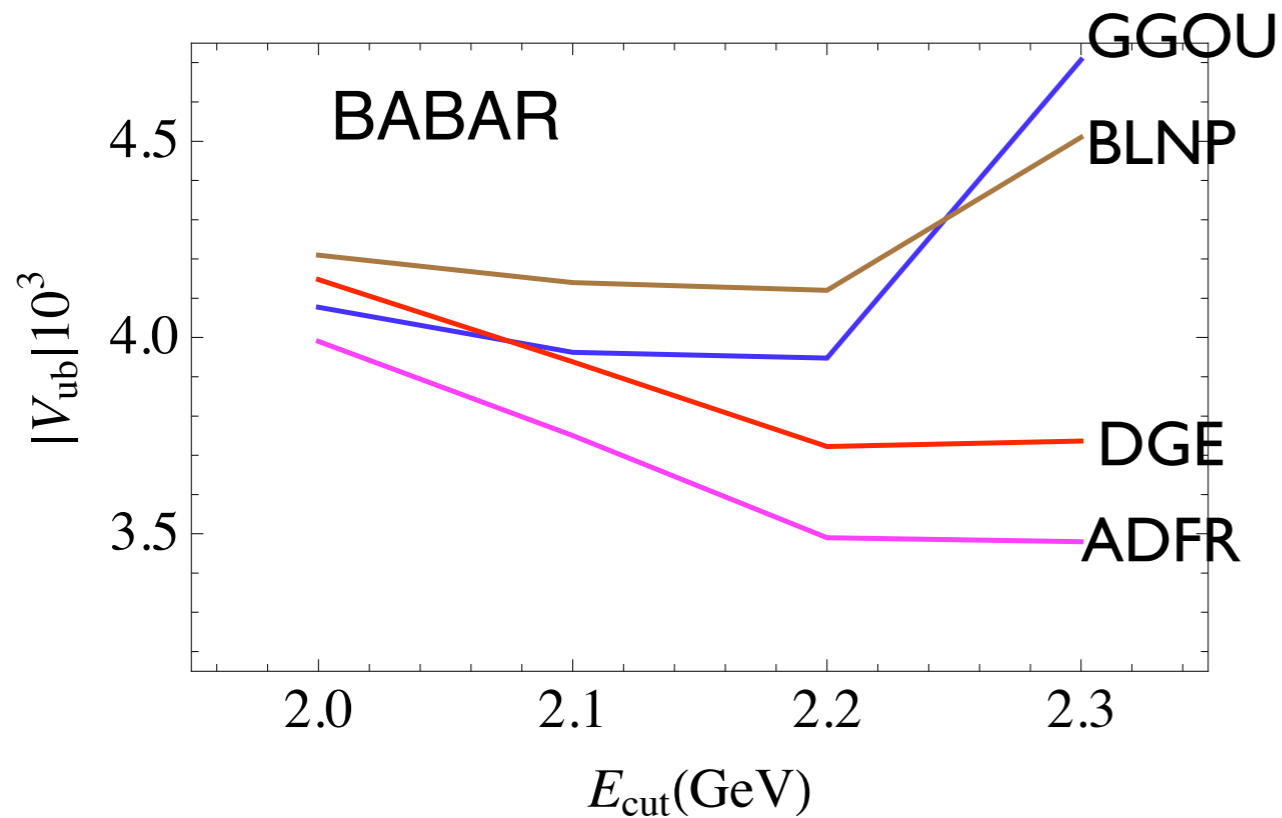
A global comparison



- * common inputs (except ADFR)
- * Overall good agreement with one exception
SPREAD WITHIN THE ERRORS!
- * Systematic offset of central values:
normalization? to be investigated
- * Very different methods, common systematics?
WA, inputs, pert corrections

Why do central values differ up to 9-10%?

The lepton spectrum

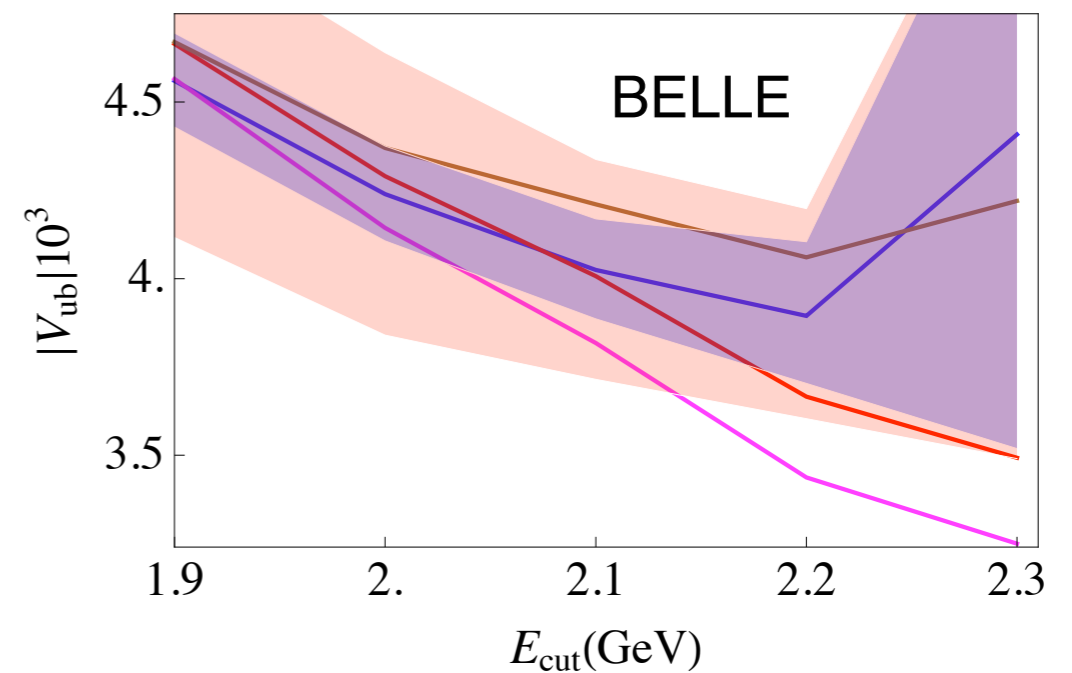
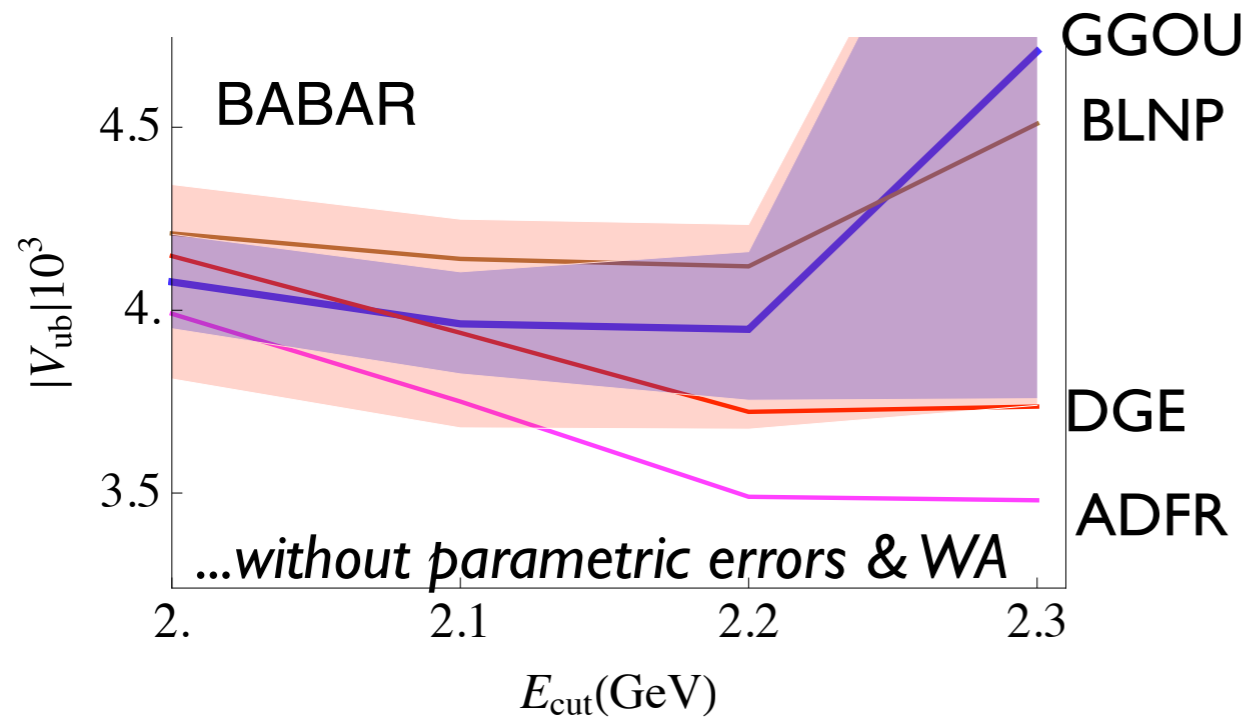


Babar E_l determination

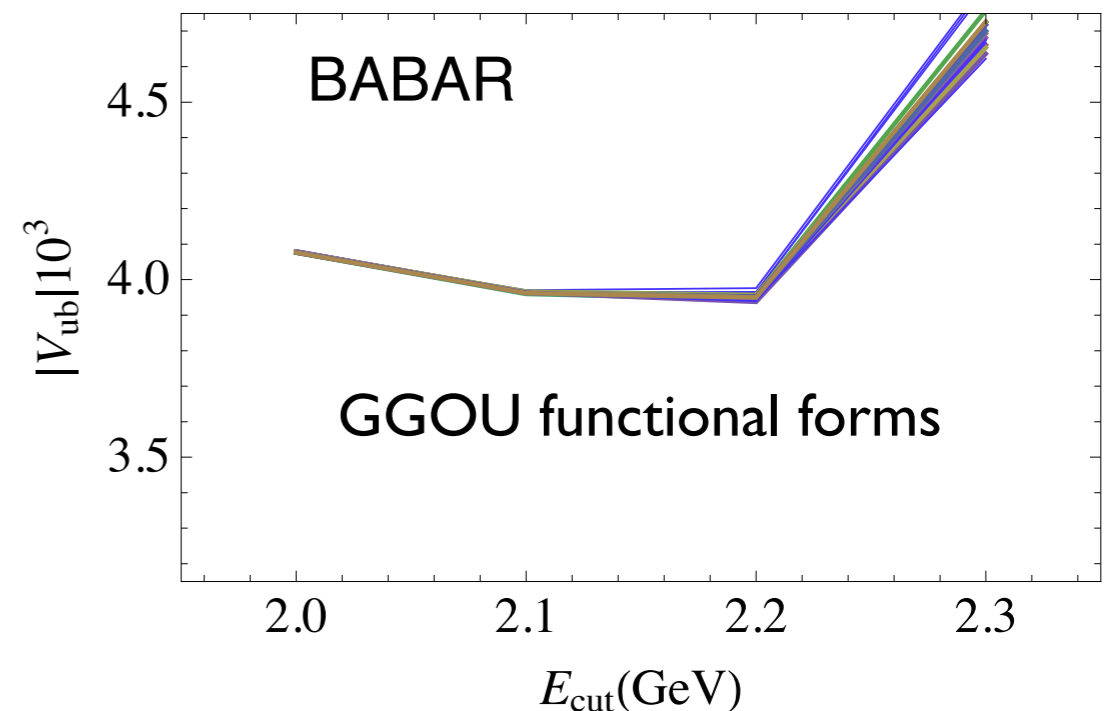
Belle E_l determination

Common inputs, $m_b^{\text{kin}}=4.60\text{GeV}$ or $m_b(m_b)=4.24\text{GeV}$.
Exp analyses depend strongly on generator (inconsistent!!!!)

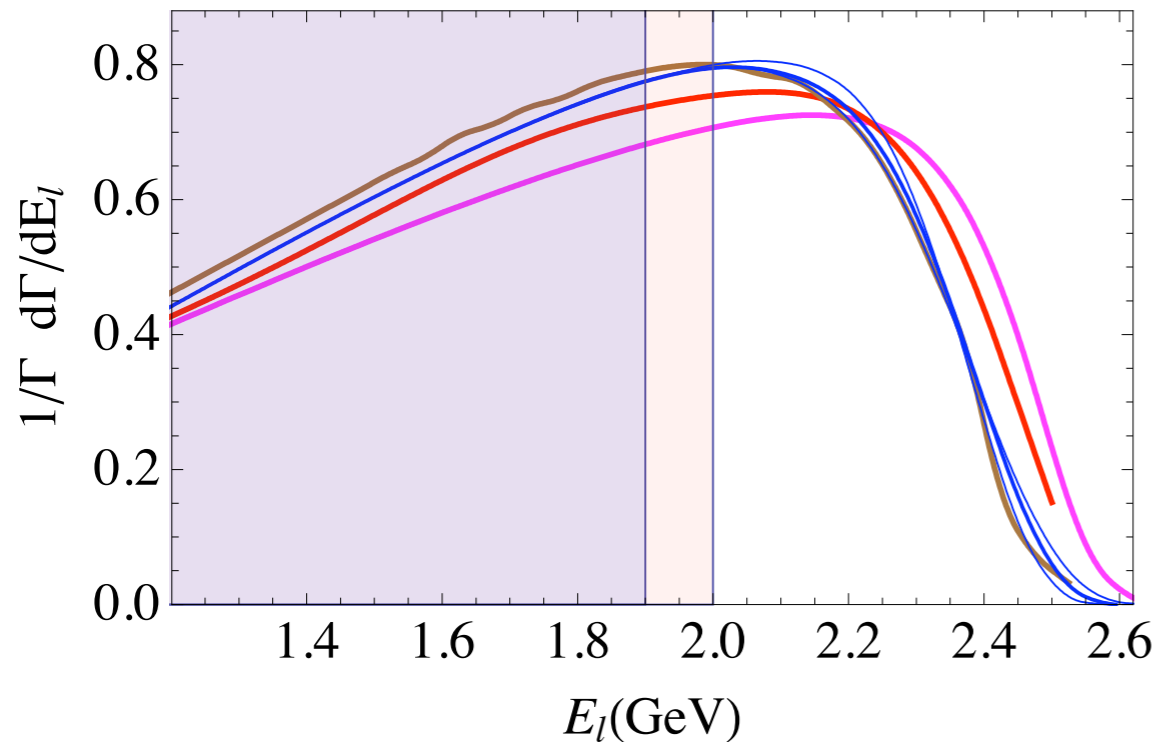
The lepton spectrum



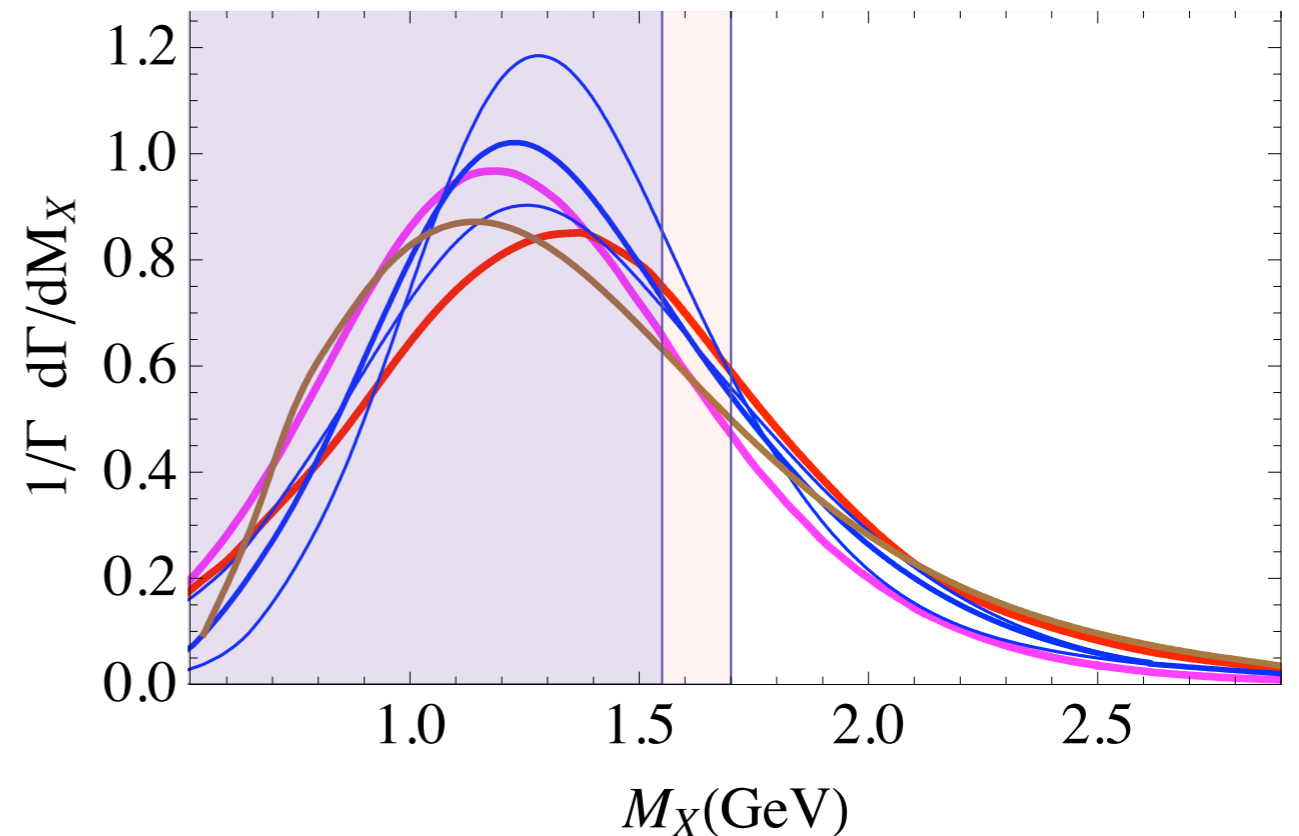
The spectrum does provide information:
 OPE based methods close to each other up to 2.2 GeV, resummed methods show larger slope, only seem to behave in same way



The leptonic and M_X spectra



DGE GGOU BLNP ADFR



The leptonic spectrum is not sensitive to the SF except quite close to the endpoint. At 1.5 GeV all methods should agree (it's pQCD after all)

- *Not all observables are equally clean. eg high q^2 tail is sensitive to WA*
- Need spectra: only way to test frameworks (see E_l spectrum).
- *More inclusive measurements, less dependence on m_b*
- Theory errors are partly parametric: m_b dependence is crucial

	Average $ V_{ub} \times 10^3$
DGE	$4.26(14)_{\text{ex}}^{+19}_{-13}$
BLNP	$4.31(16)_{\text{ex}}^{+32}_{-27}$
GGOU	$3.96(15)_{\text{ex}}^{+20}_{-23}$

2.1, 1.9, 1.3 σ from $B \rightarrow \pi l \nu$
(MILC-FNAL)

3.1, 2.4, 1.5 σ from UFit
(because of $\sin 2\beta$)

NEW preliminary Belle Multivariate analysis only $E_l > 1 \text{ GeV}$

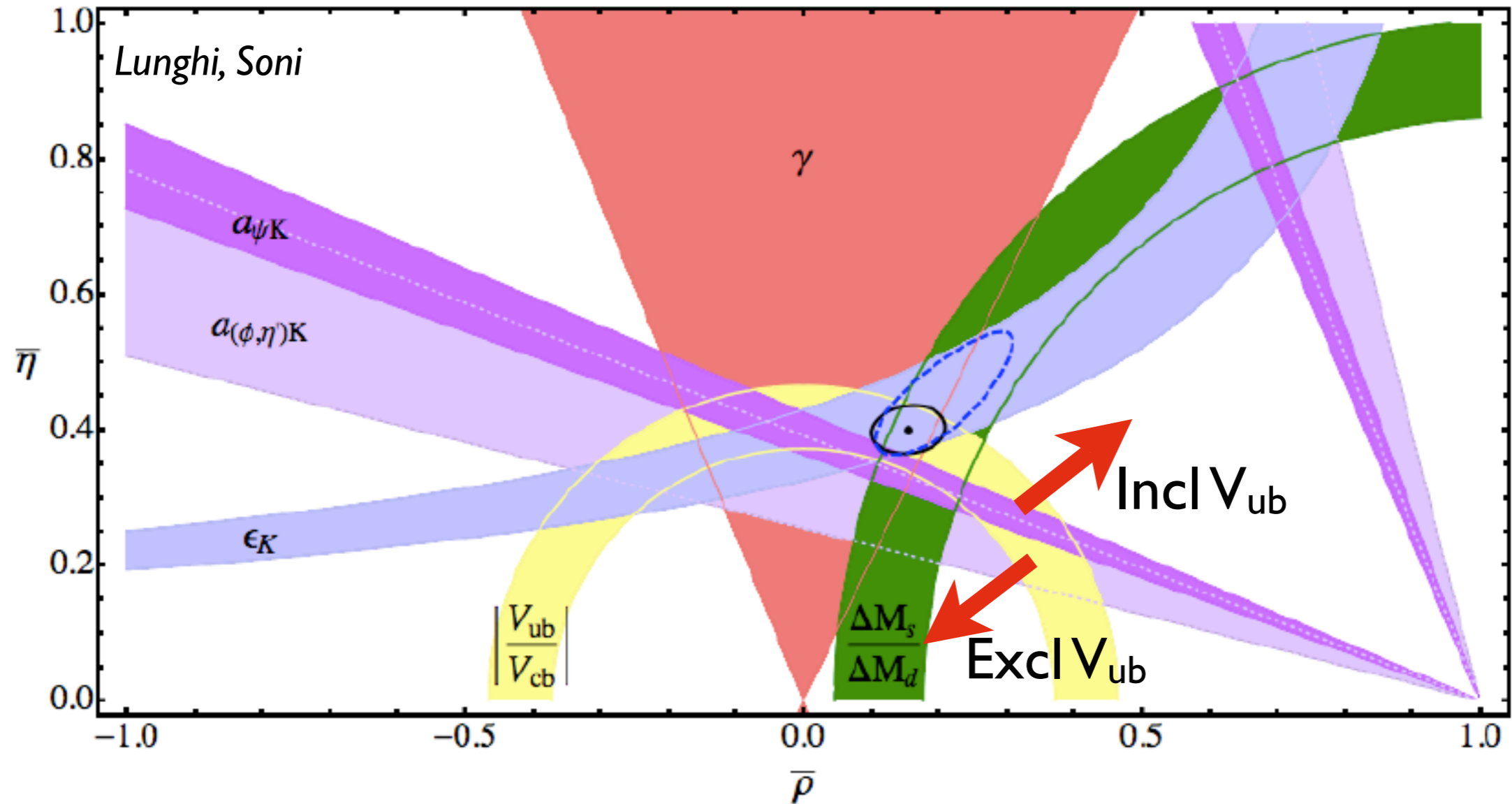
$$|V_{ub}| = (4.45 \pm 0.26^{+0.13}_{-0.22}) \times 10^{-3}_{\text{GGOU}}$$

2.1 σ from excl, 2.5 σ from UFit
probably a bit less after fit upgrade

This includes about 90% of the rate really inclusive measurement, no need for SF. Only crucial input m_b needs to be confirmed!

NEW PHYSICS?
eg LR models Chen, Nam





Recent lattice results for B_K and previously neglected contributions lead to 15% smaller ϵ_K , in $\sim 1.8-2\sigma$ conflict with $\exp \sin 2\beta$. Buras, Guadagnoli

Perhaps $\sin 2\beta$ is simply too low...

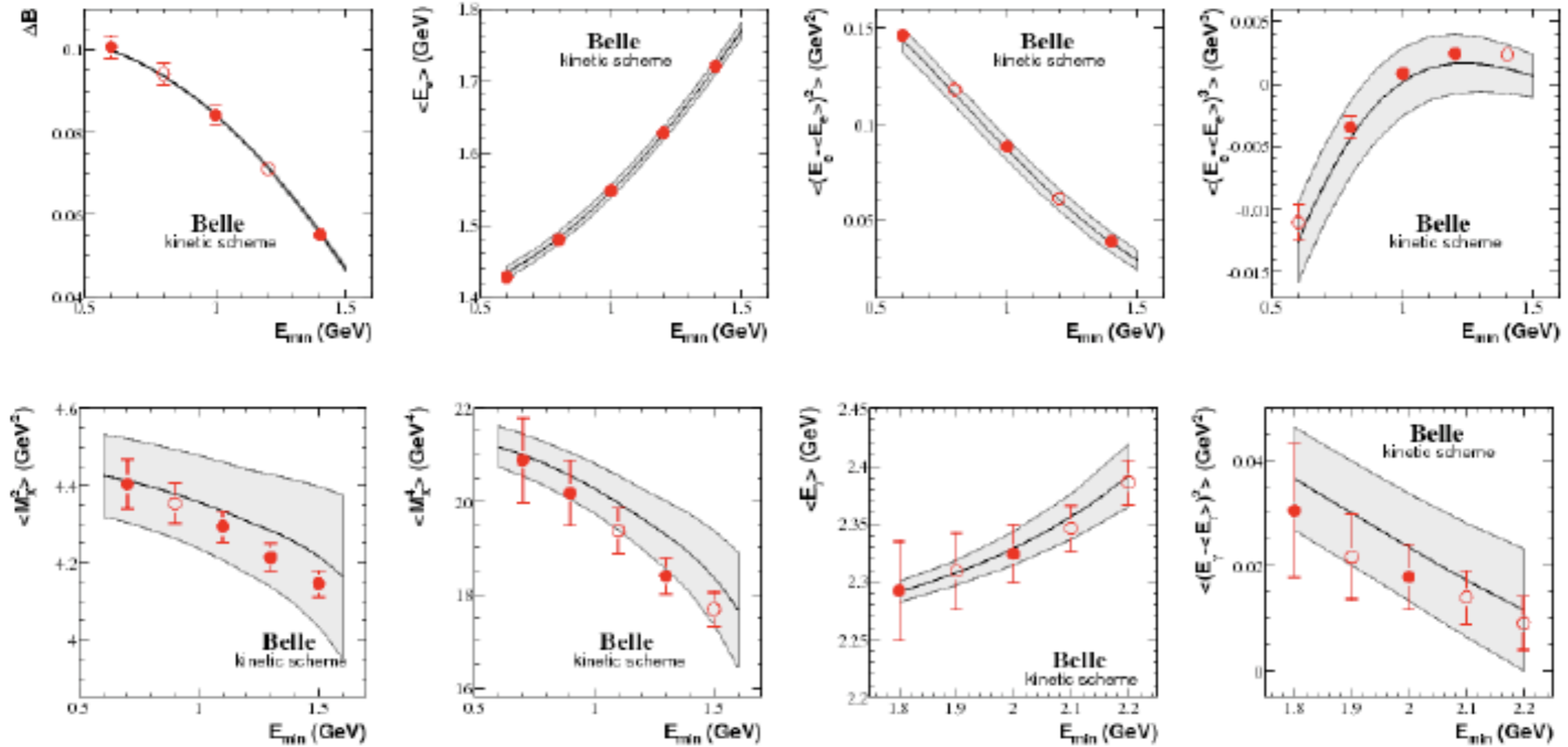
too early to say

Conclusions

- Inclusive $|V_{cb}|$ seems OK, good prospects for th error reduction
- m_b - m_c well determined by semileptonic fits, individual errors somewhat underestimated. Expect improvements by summer
- Latest FNAL result for $B \rightarrow D^*$ f.f. clashes with heavy quark sum rules
- Approaches to V_{ub} agree and seem consistent with data, need spectra and varying cuts
- No real problem between excl-incl V_{ub} , 1.5σ with UTfit, but new Belle multivariate result implies even larger V_{ub} with small th uncertainty

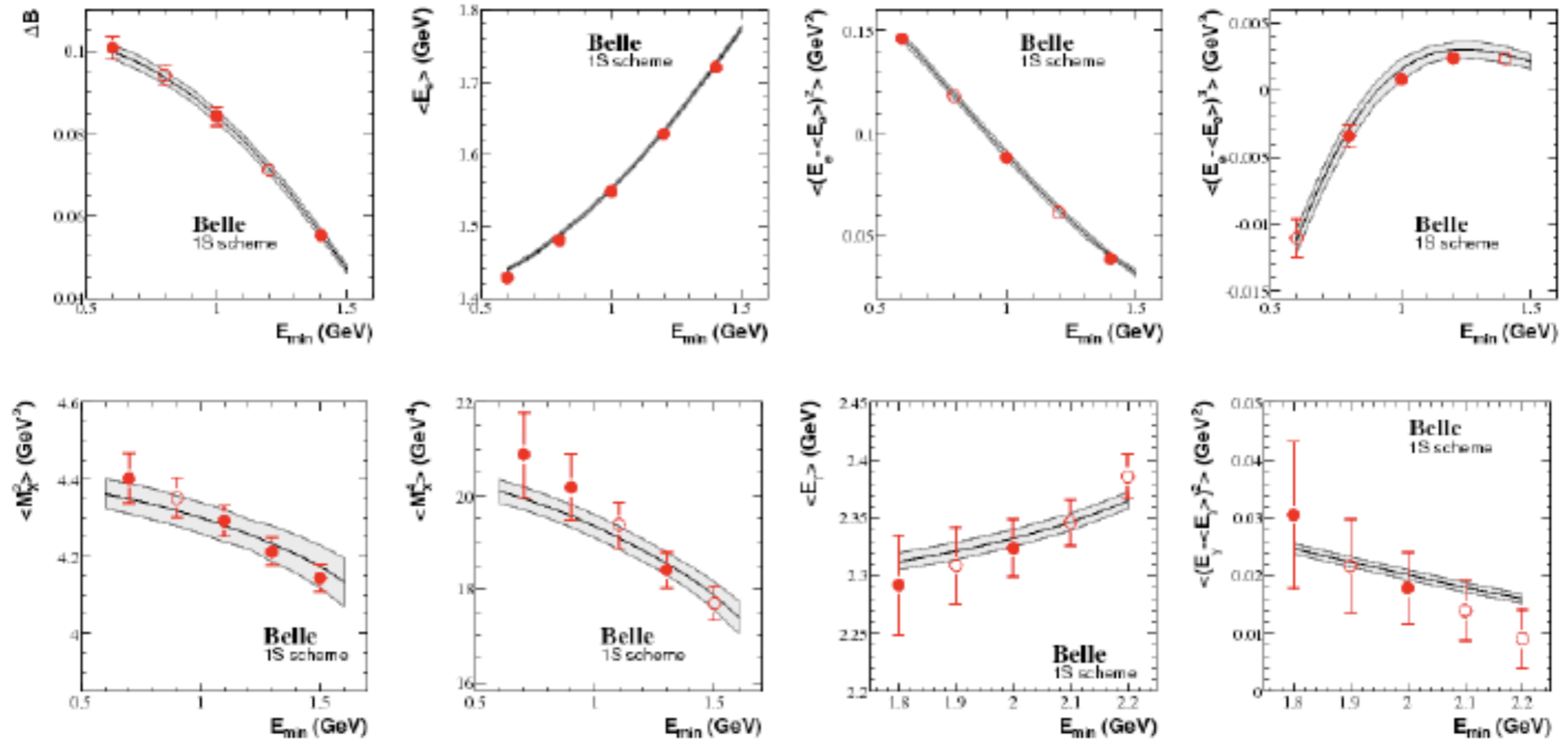
Back-up slides

Fit result in the kinetic scheme



$$\chi^2/\text{ndf.} = 4.7 / (25-7)$$

Fit result in the 1S scheme



$$\chi^2/\text{ndf.} = 7.3 / (25-7)$$

Constraining Weak Annihilations

WA may pollute all present estimates, and tend to **decrease** the extracted V_{ub} . Need **upper cut on q^2** to remove this uncertainty and/or constrain WA from $B^{0/+}$ and q^2 spectrum

