

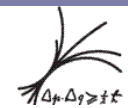


Deep Inelastic Scattering & Parton Densities at HERA

using data
from



- Historical prelude
- Deep inelastic scattering
Theory and Experiments
- HERA measurements:
Cross sections and
Structure functions
- Extraction of Parton Densities
- Open Issue(s) & Future PDF@HERA

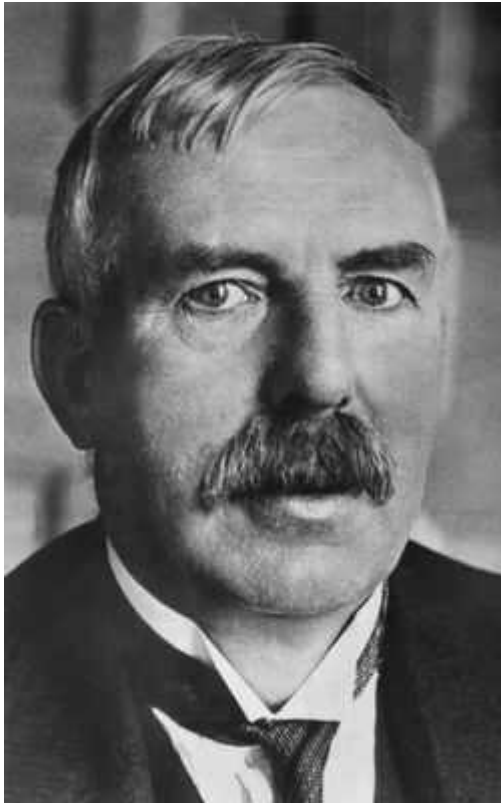


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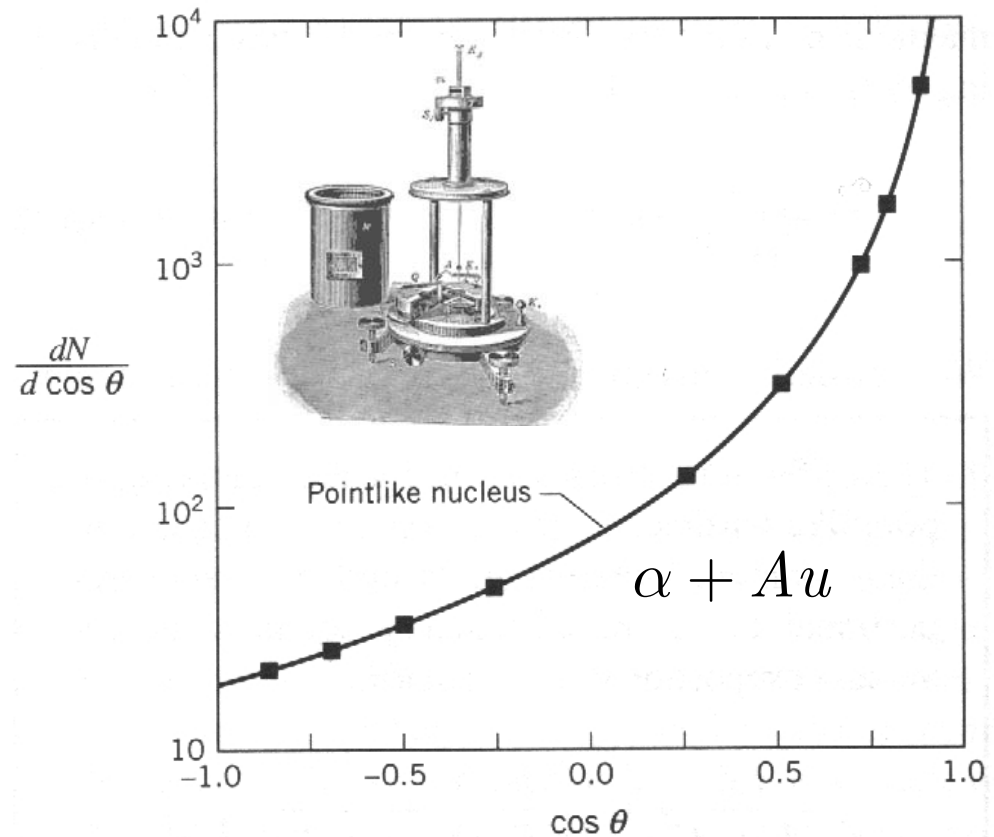
It all started with Ernest Rutherford ...

The mother of all scattering experiments



E. Rutherford

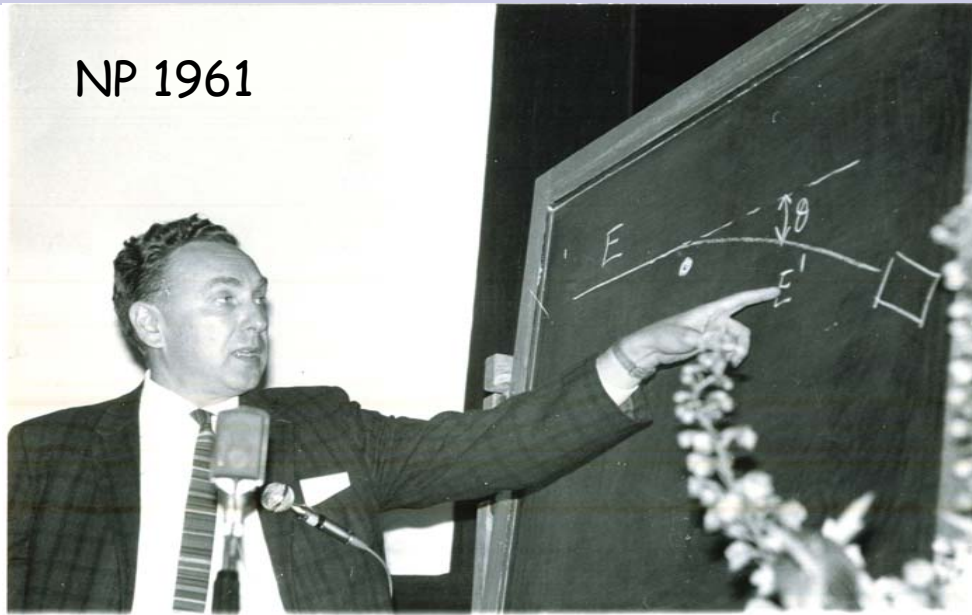
Nobel Price 1908 (for chemistry!)



$$\frac{d\sigma}{d\Omega} = (zZ\alpha)^2 \left(\frac{\hbar c}{4E_{\text{kin}}} \right)^2 \frac{1}{\sin^4(\theta/2)}$$

Pioneer of Electron Scattering: Robert Hofstadter

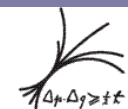
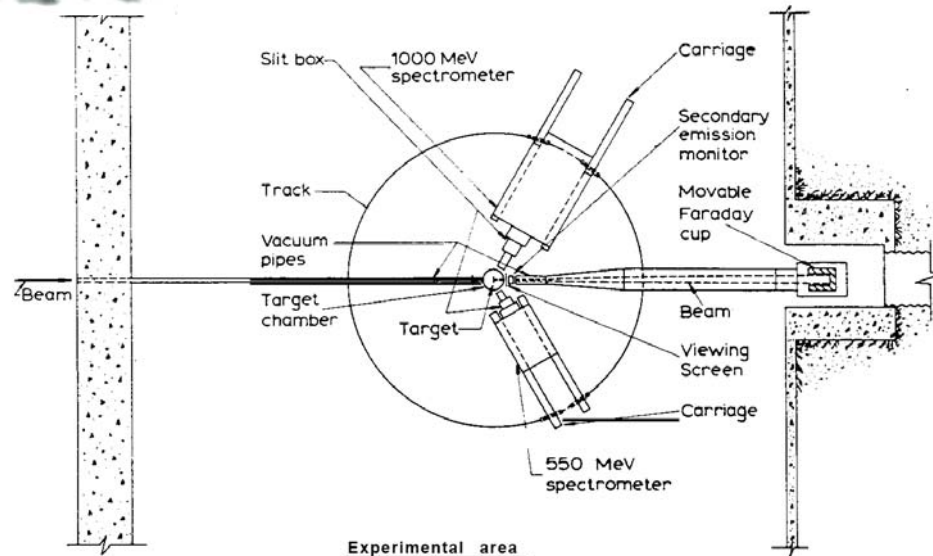
NP 1961



- electrons are elementary
- are sensitive only to charge
- can be easily produced and accelerated to high energy

$$\hbar c = 200 \text{ [MeV fm]}$$

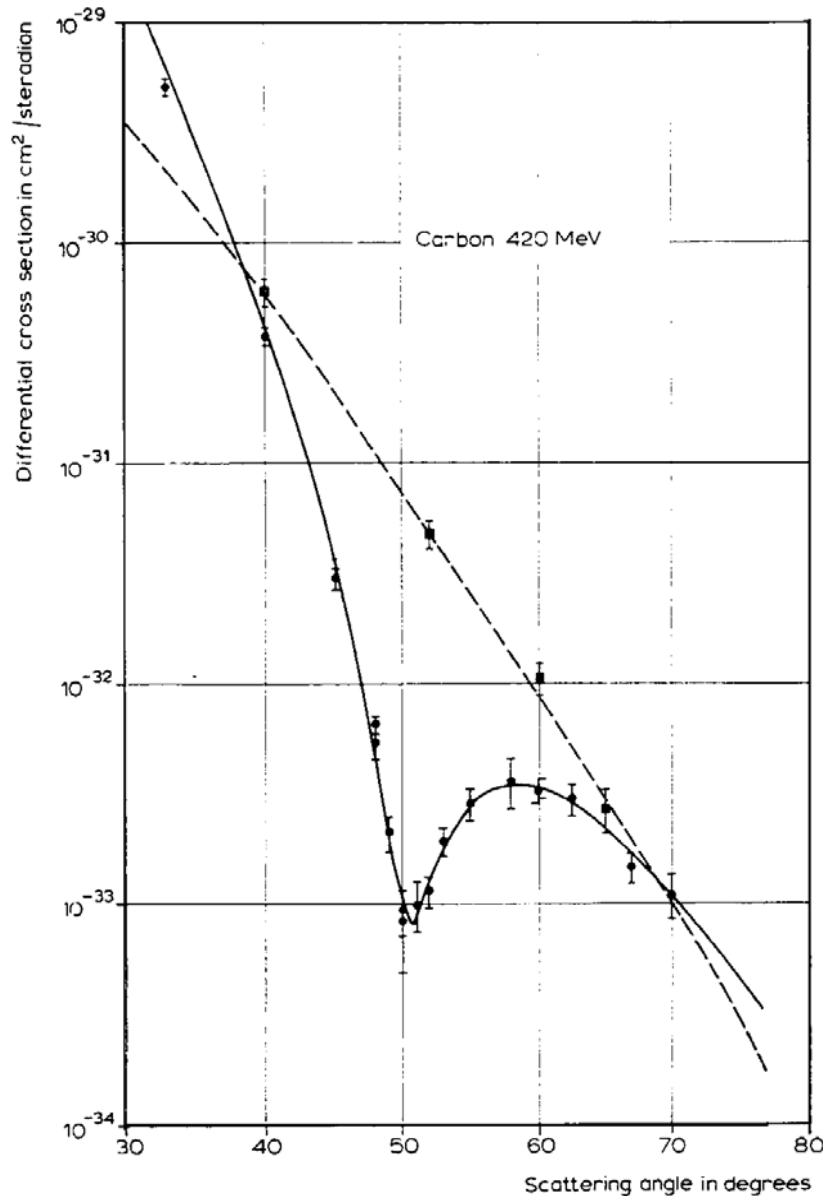
typical setup for an
electron scattering experiment



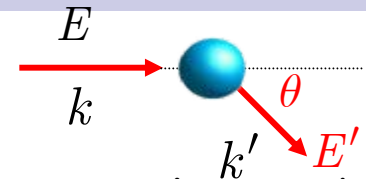
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Figures from R. Hofstadter's Nobel Lecture, Dec. 11, 1961



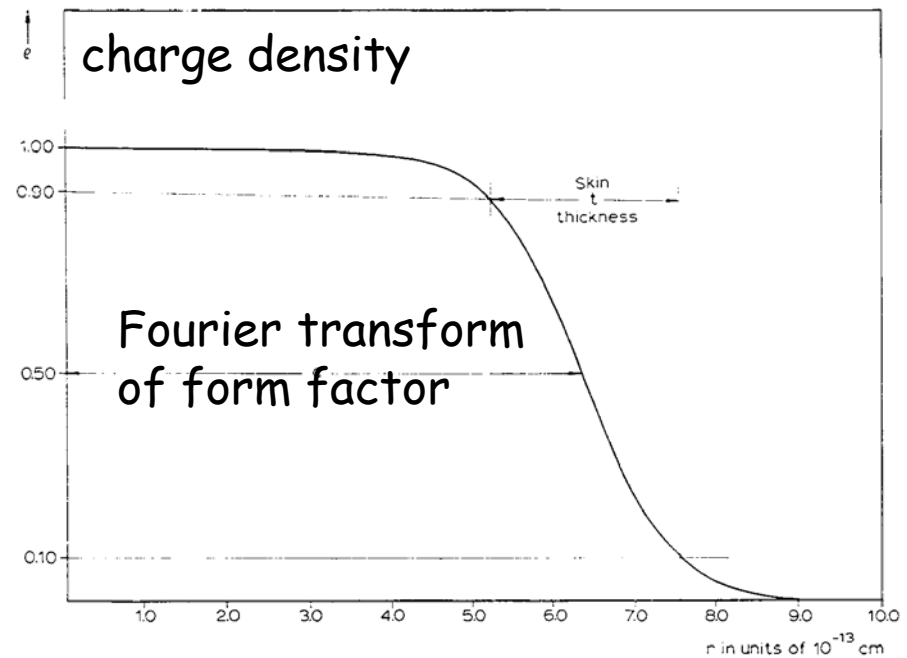
Elastic scattering:



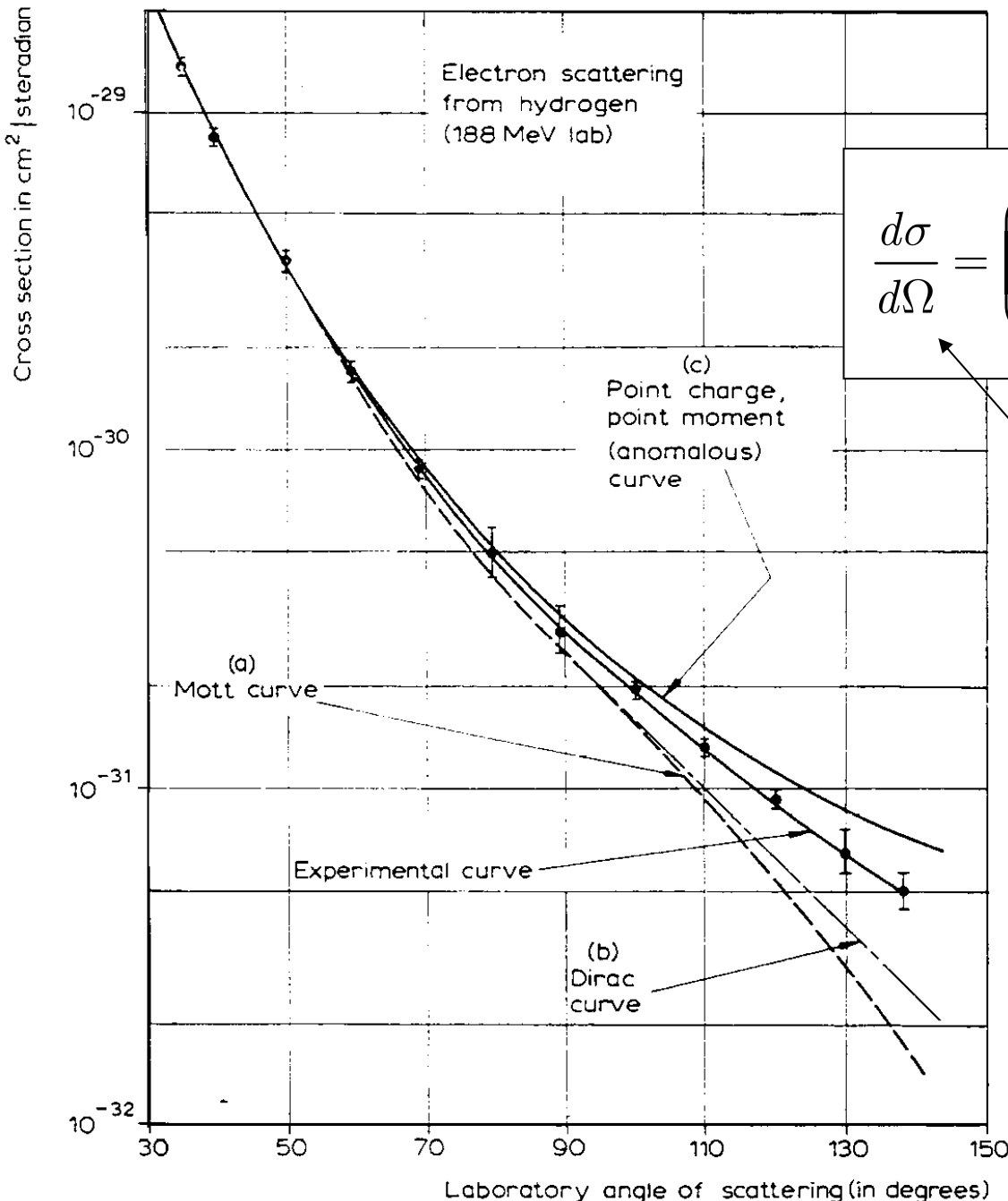
clear dip structure \rightarrow finite nuclear radius

$$\frac{d\sigma}{d\Omega} = \left(\frac{d\sigma}{d\Omega} \right)_{Ruth} \left(\frac{E'}{E} \right) \cos^2 \frac{\theta}{2} |F(q)|^2$$

$q = k - k'$ form factor



Something is „wrong“ with the proton ...

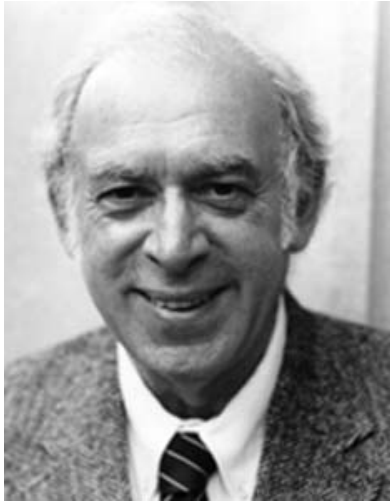


$$\frac{d\sigma}{d\Omega} = \left(\frac{d\sigma}{d\Omega} \right)_{Ruth} \frac{E'}{E} \left(\cos^2 \frac{\theta}{2} + \frac{q^2}{2M} \sin^2 \frac{\theta}{2} \right)$$

Cross section for electron scattering on a spin 1/2 Dirac proton

- the proton has an anomalous magnetic structure (NOT a Dirac particle)
- beam energy too small to resolve the proton structure

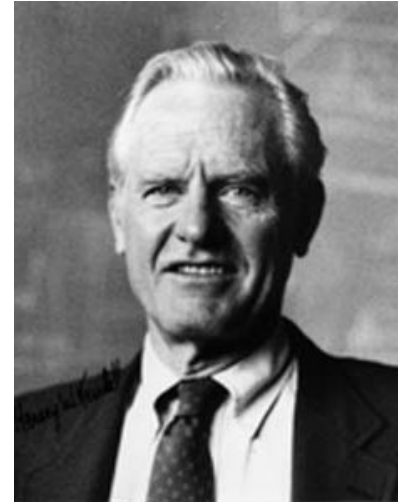
J. Friedman



R. Taylor



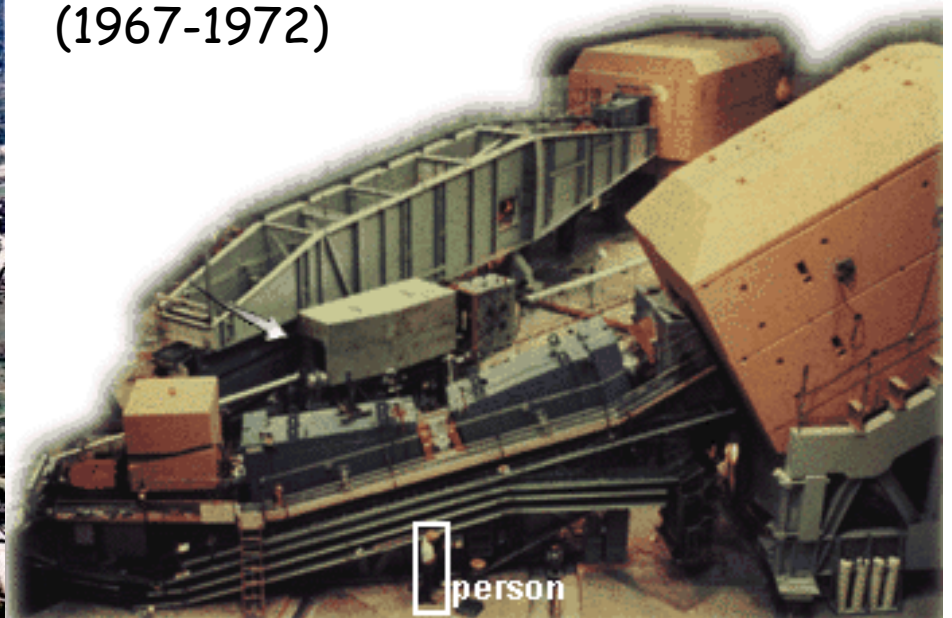
H. Kendall



NP 1990

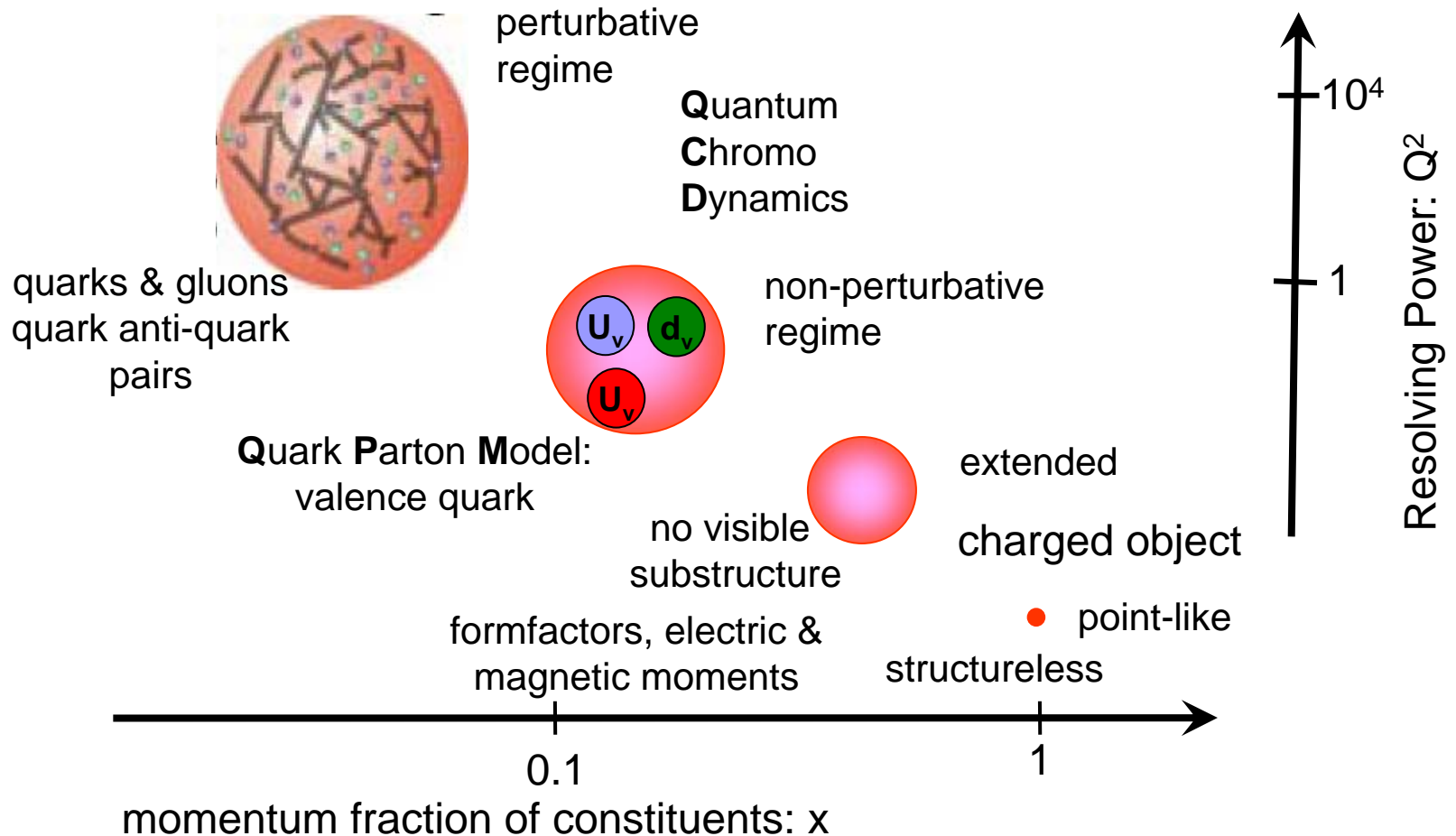


The proton is made out of quarks !
(1967-1972)



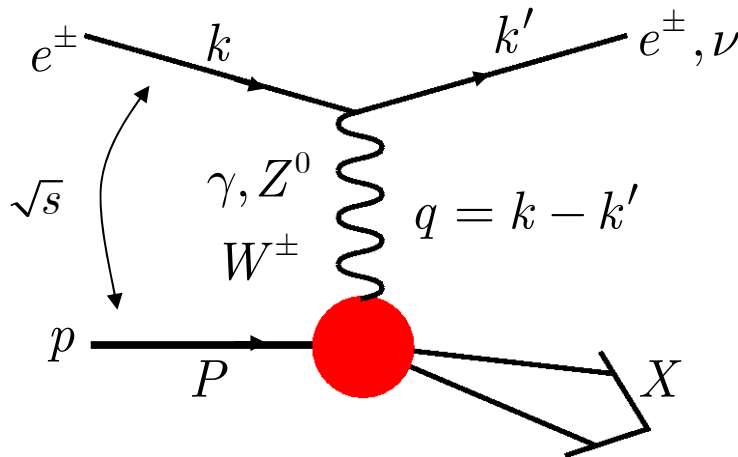
„Endstation A“ spectrometers

“Imaging” of the Proton

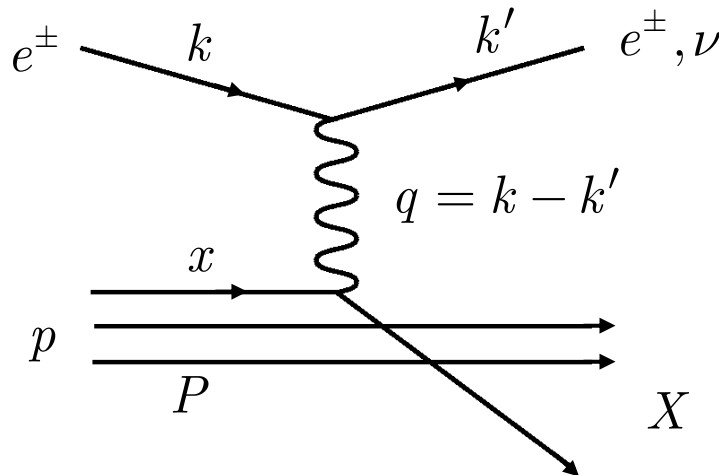


Low Q^2 : probing transition perturbative to non-perturbative QCD
 Low x : probing quarks and gluons at high density
 $\Delta E \Delta t \sim \hbar \rightarrow$ high Q^2 “high resolution but time averaged picture”

Deep Inelastic Scattering (DIS)



QPM



$$Q^2 = -(k - k')^2 \quad (\text{momentum transfer})^2$$

$$= -q^2$$

virtuality of γ^*, Z^0, W^\pm
 \rightarrow („size“ of the probe)⁻¹

$$x = \frac{Q^2}{2 P \cdot q}$$

fraction of the proton momentum carried by the charged parton

$$y = \frac{P \cdot q}{P \cdot k}$$

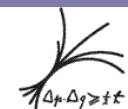
fraction of the electron energy carried by the virtual photon („inelasticity“)

$$s = (k + P)^2 \quad \text{center of mass energy of } ep \text{ system}$$

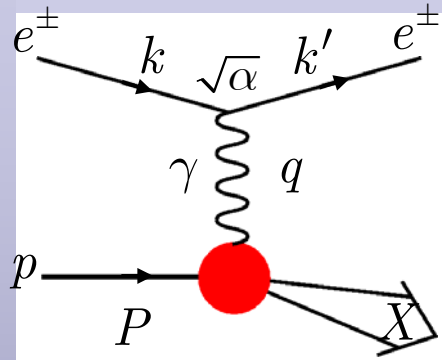
$$W^2 = M_X^2 \quad (\text{mass})^2 \text{ of } \gamma^* p \text{ system}$$

$$= (q + P)^2$$

$$Q^2 = sxy$$



Cross Section and Structure Functions



$$\frac{d^2\sigma}{d\Omega dE'} = \frac{\alpha^2}{2MQ^4} \frac{E'}{E} L_{\mu\nu} W^{\mu\nu}$$

$L_{\mu\nu}$ lepton tensor

$W_{\mu\nu}$ hadronic tensor

$$L_{\mu\nu} = 2 \left[k_\mu k'_\nu + k'_\mu k_\nu + \frac{q^2}{2} g_{\mu\nu} \right]$$

← minimal electromagnetic coupling

(unpolarized particles)

$$W_{\mu\nu} = \left(-g_{\mu\nu} + \frac{q_\mu q_\nu}{q^2} \right) 2F_1 + \left(P_\mu - \frac{P \cdot q}{q^2} q_\mu \right) \left(P_\nu - \frac{P \cdot q}{q^2} q_\nu \right) \frac{2}{P \cdot q} F_2$$

← most general tensor satisfying charge conservation

NC cross section :

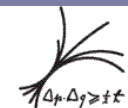
$$\frac{d^2\sigma(e^\pm p)}{dx dQ^2} = \frac{4\pi\alpha^2}{xQ^4} [xy^2 F_1 + (1-y)F_2]$$

$$F_L \equiv F_2 - 2xF_1$$

longitudinal structure function

$$= \frac{2\pi\alpha^2}{xQ^4} [Y_+ F_2 - y^2 F_L]$$

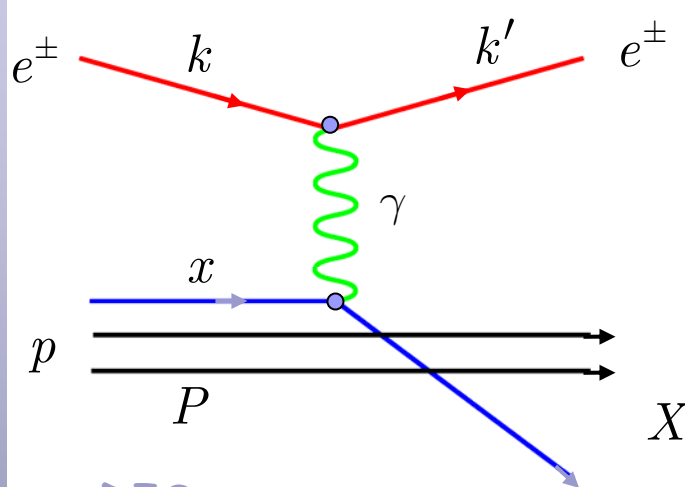
$$Y_\pm = 1 \pm (1-y)^2$$



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Structure Functions within the Quark-Parton-Model

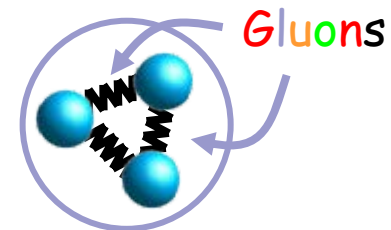
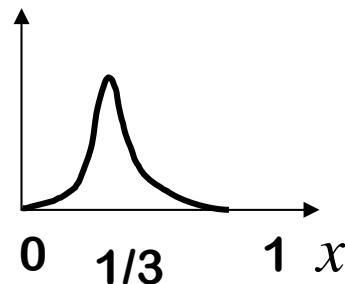
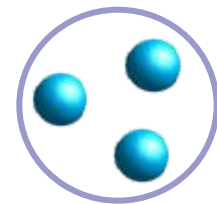
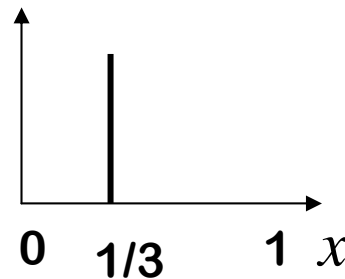
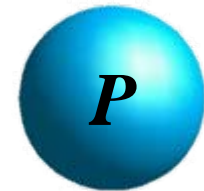


$$\frac{d^2\sigma(e^\pm p)}{dx dQ^2} = \frac{2\pi\alpha^2}{xQ^4} [Y_+ F_2 - y^2 F_L]$$

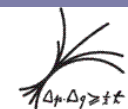
QPM: $F_2(x) = \sum_{i=u,d} e_i^2 x q_i(x)$ ← parton densities $xq_i(x)$ (pdf)

DIS =

- electron scatters off a charged constituent (parton) of the proton (= elastic scattering)
- identify the charged partons with **QUARKS** (= spin 1/2 fermions)

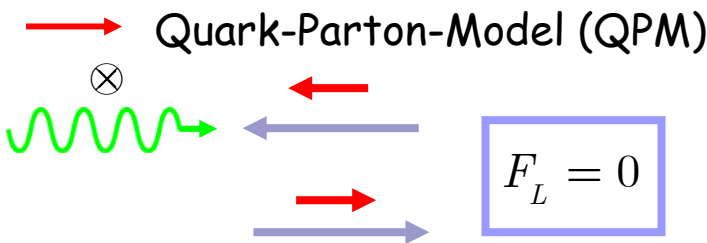


1/6 (Gluons carry half proton momentum)



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HERA - the world's largest electron microscope (Deutsches Elektronen-Synchrotron DESY, Hamburg, Germany)

Shutdown on June 30, 2007, 23:00

HERA start: 1992
upgraded in 2001: „HERA II“



N

$p : 920 \text{ GeV}$

$e^{\pm} : 27.5 \text{ GeV}$

polarized

$\sim 6.3 \text{ km}$
circumference

W

PETRA



e^{\pm}

HERA



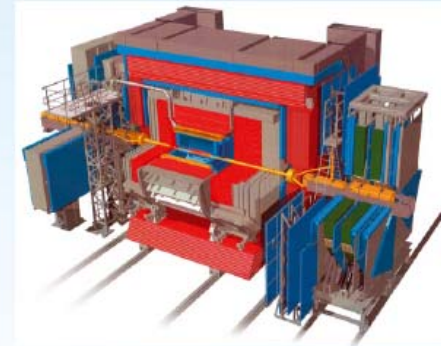
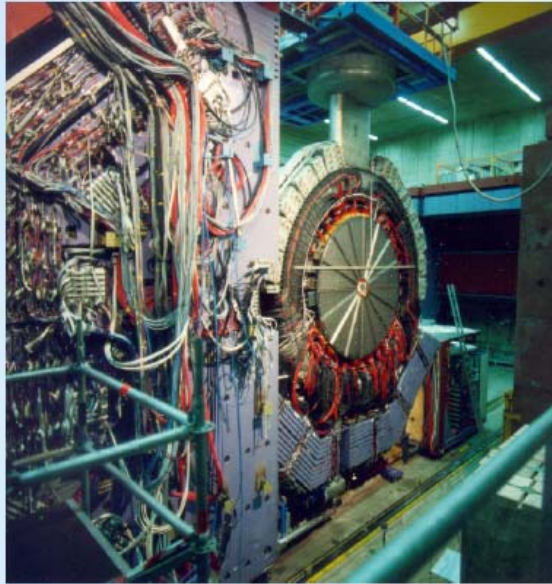
S

spatial resolution: $\sim 10^{-18} \text{ m}$

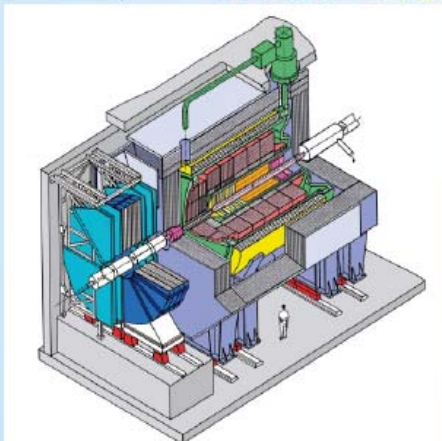
colliding beams equivalent to 50 TeV on fixed target

Collider Experiments at HERA

H1
went
for
LAr

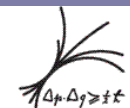
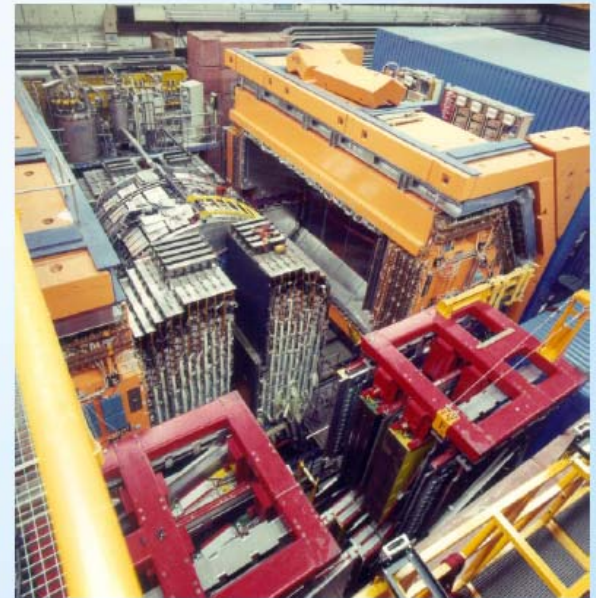


ZEUS
went
for
compen-
sation



A final salute to
our experiments

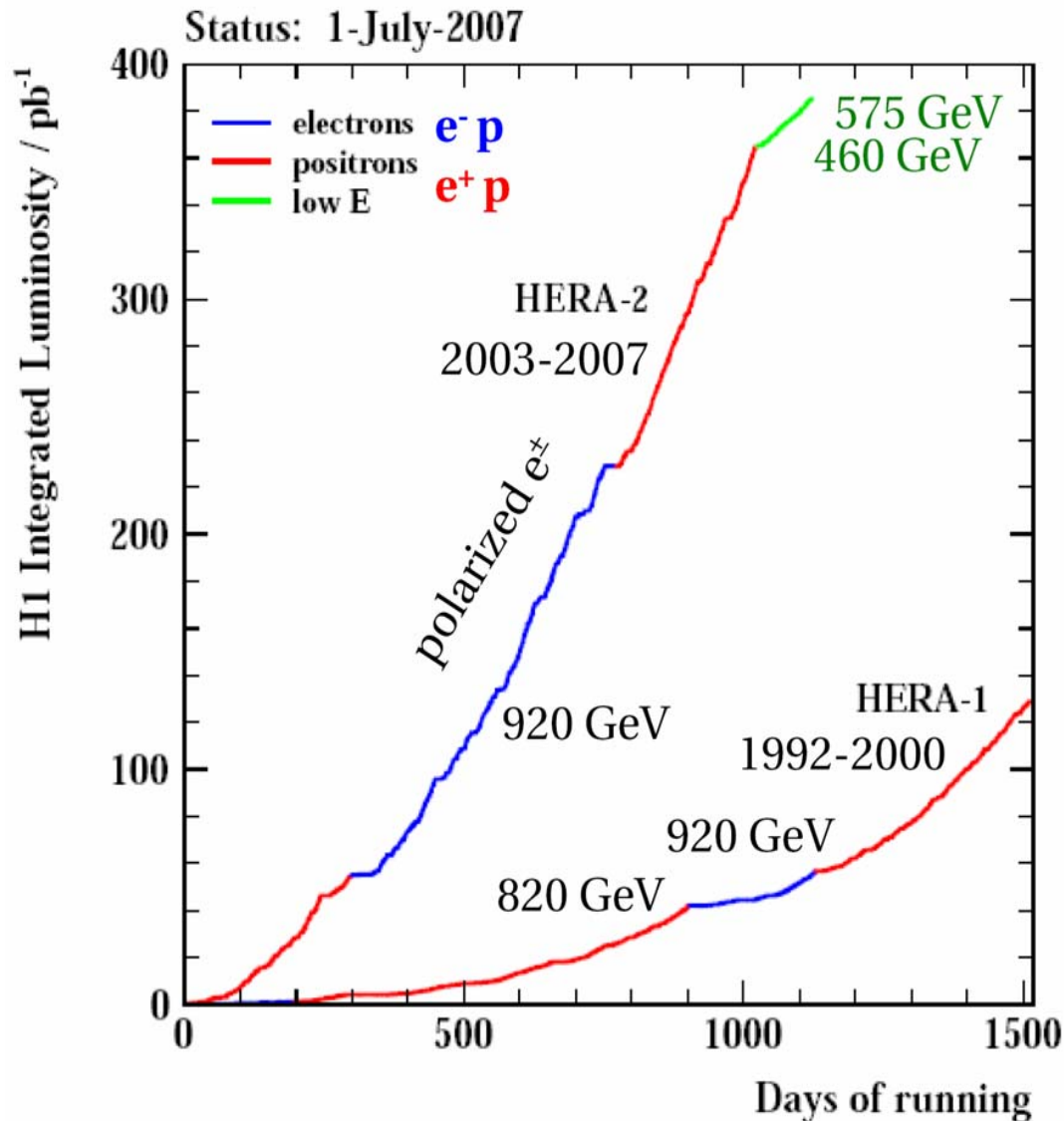
How we like to
remember
H1 and ZEUS



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Hera Luminosity



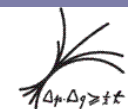
HERA I: 1992-2000

HERA II upgrade:

- luminosity
- longitudinal polarization of the lepton beams (spin rotator pairs around the interaction regions)
- massive upgrades also for the detectors



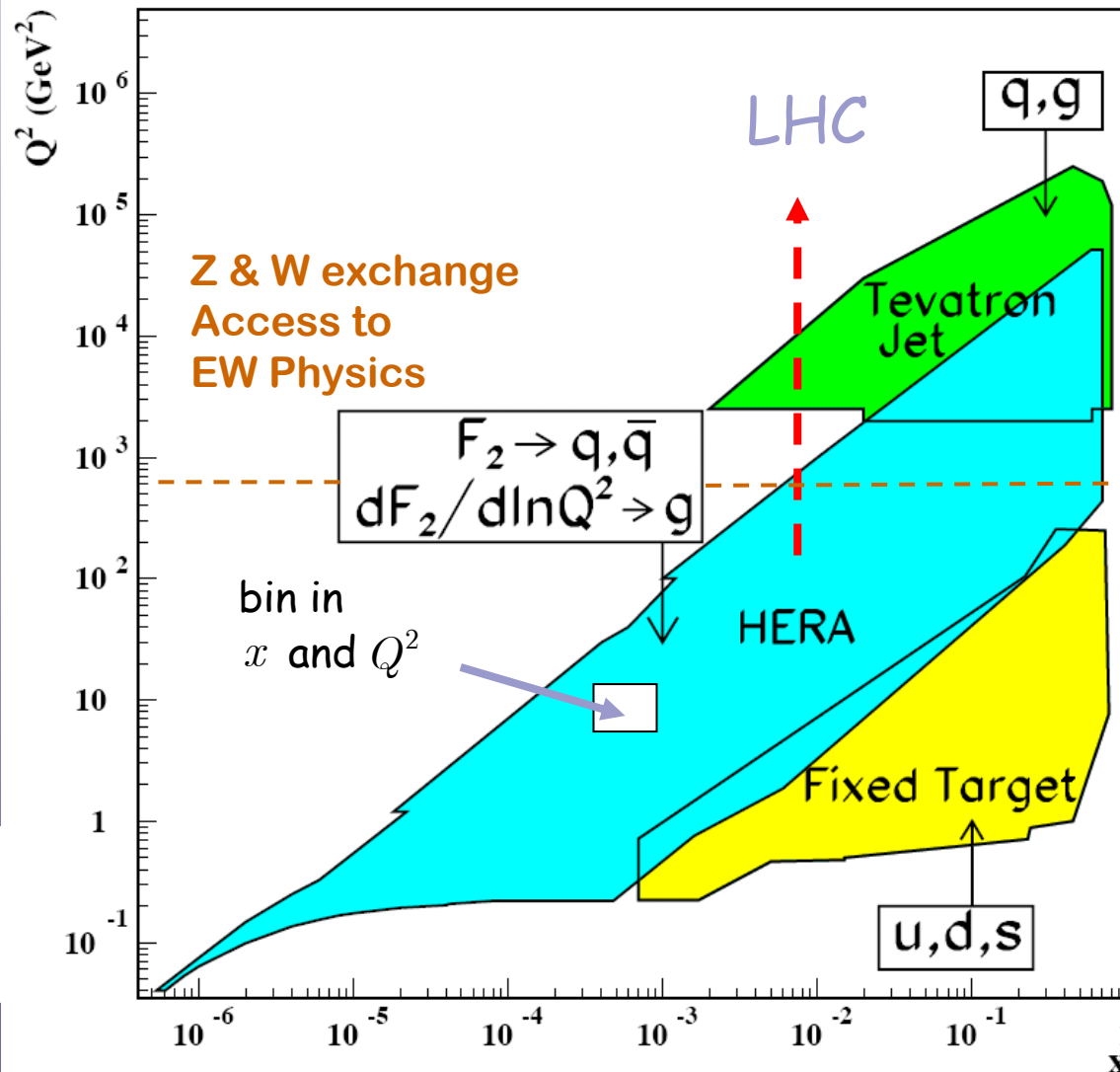
- running efficiently from 2003 onwards
- Luminosity
 $L = 500 \text{ pb}^{-1} \text{ per exp.}$



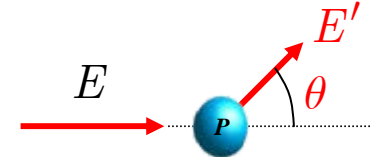
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The Kinematic Reach of HERA



Determination of kinematics („e“-method) :



$$Q^2 = 4EE' \cos^2\left(\frac{\theta}{2}\right)$$

$$y = 1 - \frac{E'}{E} \sin^2\left(\frac{\theta}{2}\right)$$

$$x = \frac{Q^2}{sy}$$

Determination of cross sections :

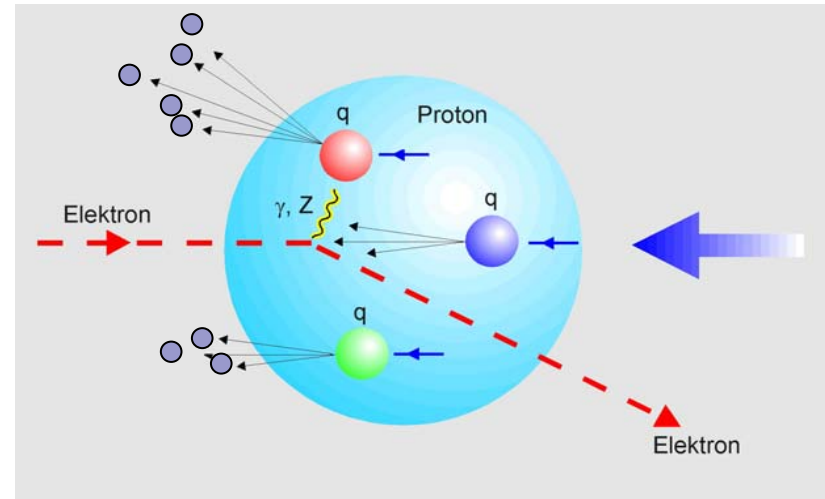
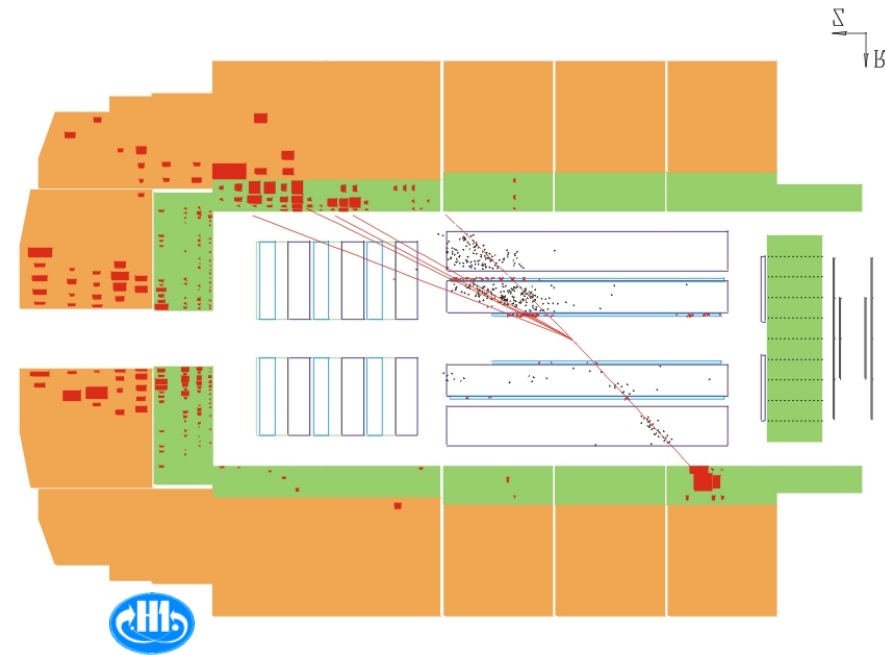
$$\frac{d^2\sigma}{dx dQ^2} \sim \frac{N - B}{\mathcal{L} \varepsilon}$$

backgr. \swarrow

efficiency \nwarrow

luminosity \nearrow

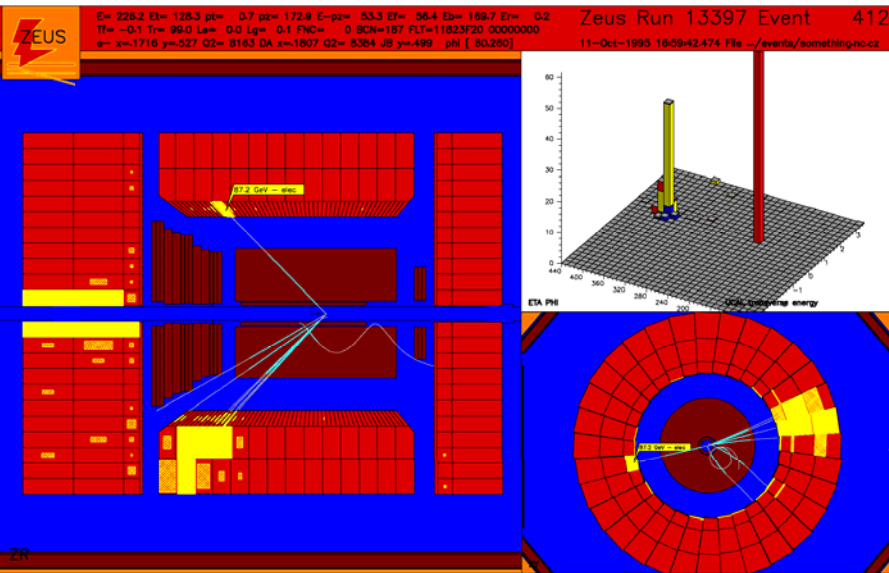
Electron Proton Scattering in Real Detectors



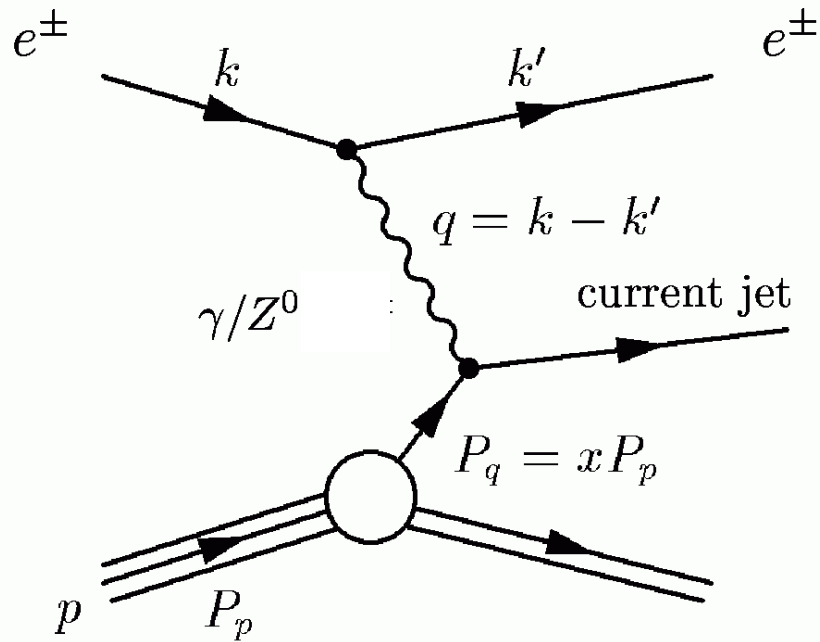
Neutral current events in

H1 (medium Q^2)

ZEUS (large Q^2)



Deep Inelastic Scattering



Kinematic Variables

- 4-momentum transfer resolving power

$$Q^2 = -q^2 = -(k - k')^2$$

- Björken scaling variable momentum fraction of struck parton

$$x = \frac{Q^2}{2p \cdot q}$$

- Inelasticity:

$$y = \frac{p \cdot q}{p \cdot k}$$

Center of mass energy \sqrt{s} : $s = (k + p)^2$ relation for fixed s: $Q^2 = sxy$

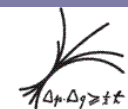
- Neutral current DIS cross section expressed by structure functions:

$$\frac{d^2 \sigma^{e^\pm p \rightarrow e^\pm X}}{dx dQ^2} = \frac{2\pi\alpha^2}{xQ^4} \underbrace{\left(1 + (1-y)^2\right)}_{Y_\pm = 1 \pm (1-y)^2} \cdot \left(F_2(x, Q^2) - \frac{y^2}{Y_+} F_L(x, Q^2) \mp \frac{Y_-}{Y_+} xF_3(x, Q^2) \right)$$

valence & sea quarks

gluons

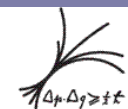
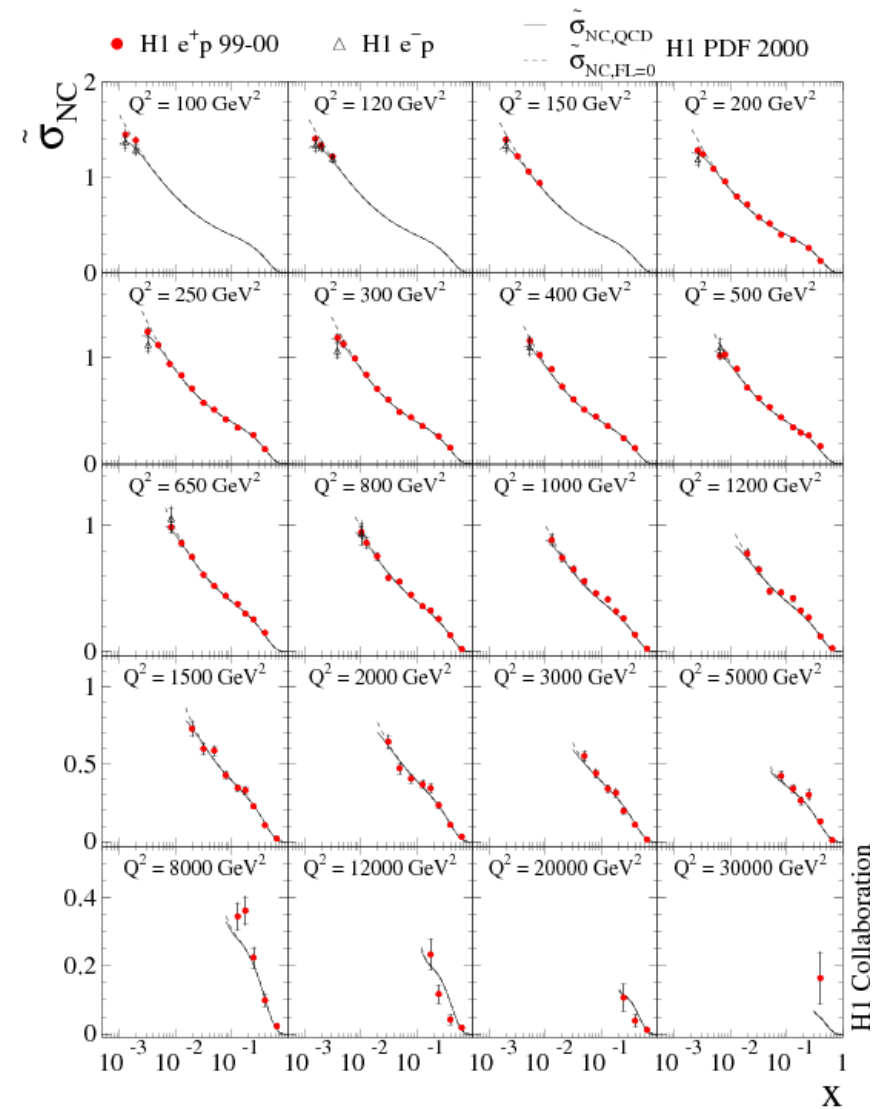
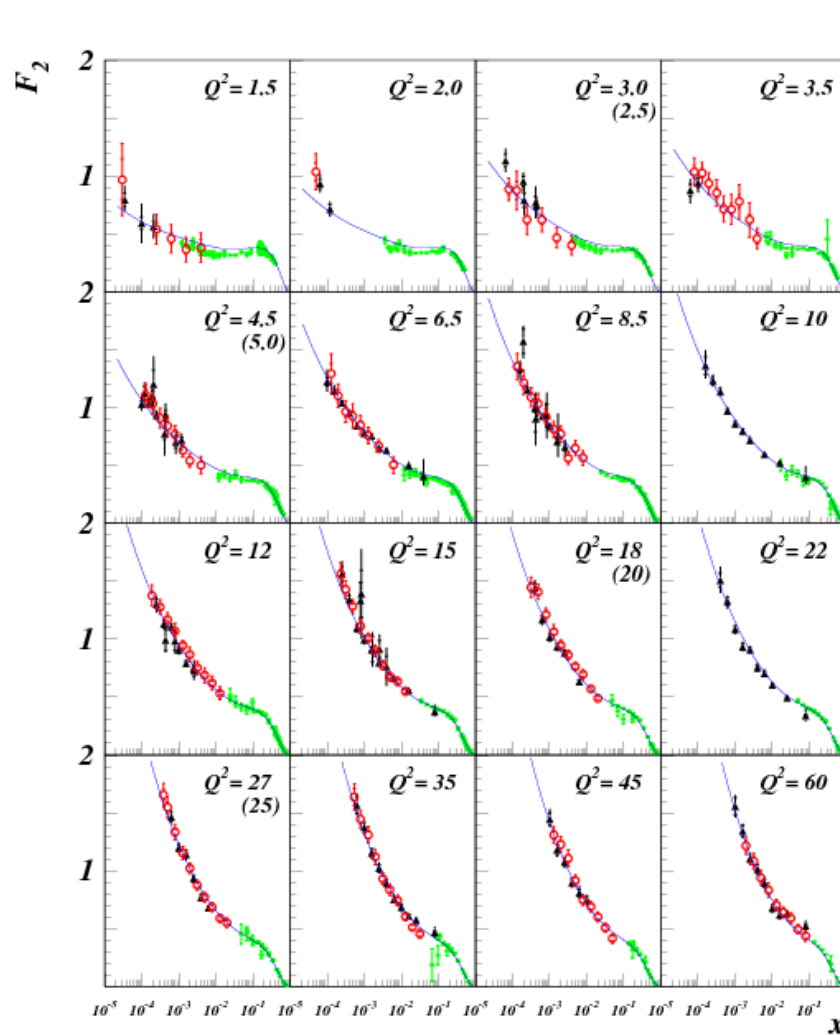
valence quarks



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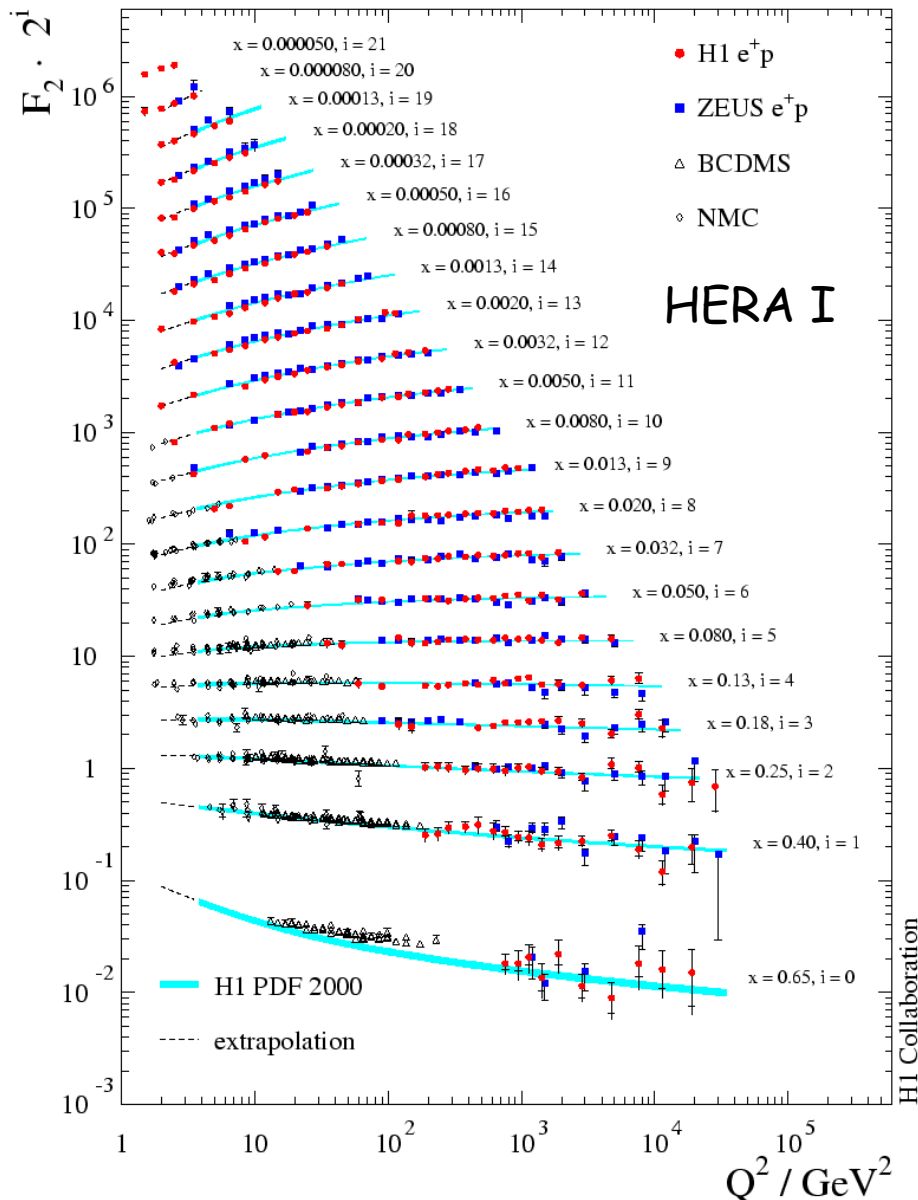
Reduced Cross Sections



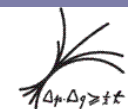
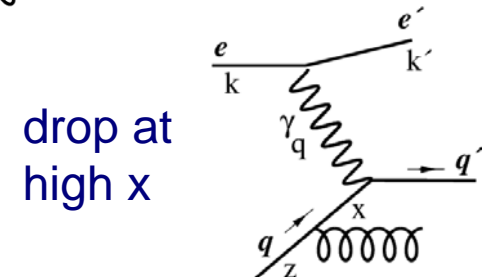
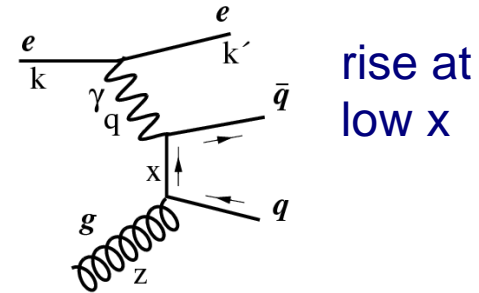
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Structure Function F_2

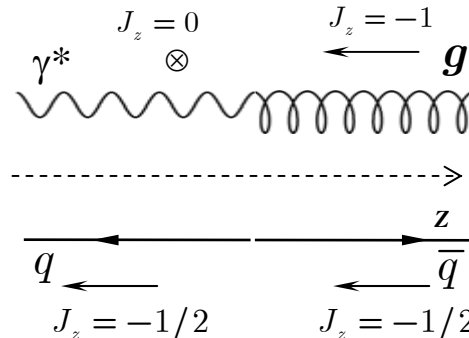
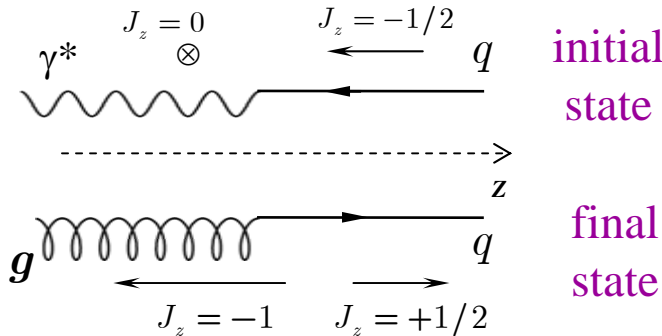
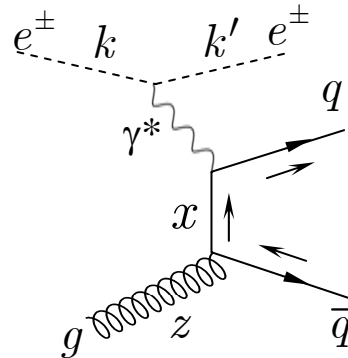
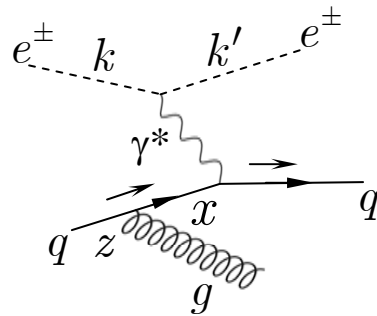


- H1 & ZEUS extended fixed target kinematic regime in x and Q^2 by 2 Orders
- Described by DGLAP
- Scaling violations \rightarrow QCD (QPM: $F_2 = \sum_i q_i(x)$)



The Longitudinal Structure Function F_L

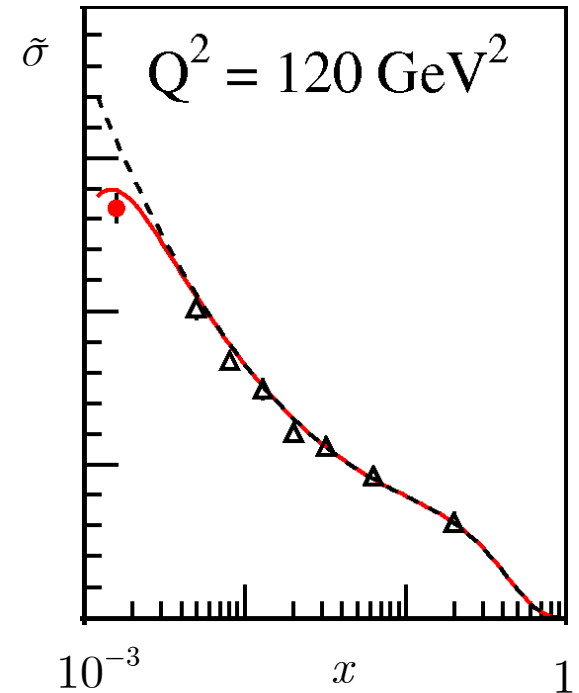
LO QCD :



$$\frac{d^2\sigma(e^\pm p)}{dx dQ^2} = \frac{2\pi\alpha^2}{xQ^4} [Y_+ F_2 - y^2 F_L]$$

in principle need
2 measurements at
different \sqrt{s}

F_L important at high y ($=$ low x)



$$F_L = \frac{Y_+}{y^2} (F_2^{\text{QCD}} - \tilde{\sigma})$$

extrapolated in Q^2
using DGLAP

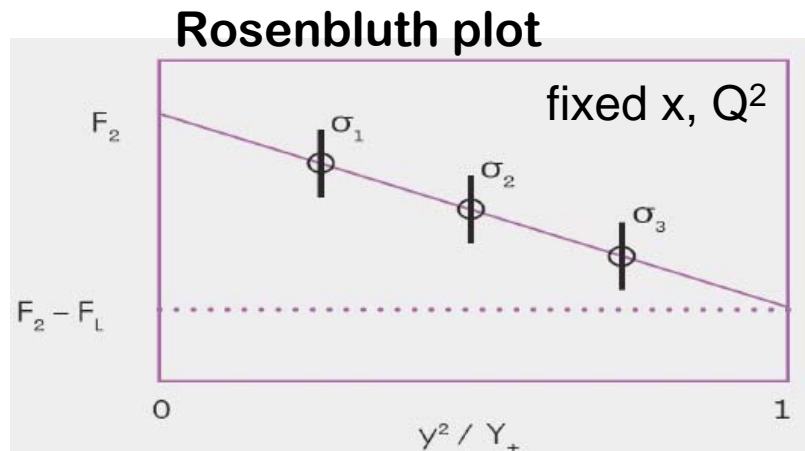
„Subtraction method“

Direct Measurement of F_L

- Neutral current DIS cross section expressed by structure functions:

$$\frac{d^2\sigma^{e^\pm p \rightarrow e^\pm X}}{dx dQ^2} = \frac{2\pi\alpha^2}{xQ^4} \underbrace{\left(1 + (1-y)^2\right)}_{Y_\pm = 1 \pm (1-y)^2} \cdot \left(F_2(x, Q^2) - \frac{y^2}{Y_+} F_L(x, Q^2) \right)$$

$\tilde{\sigma}$: Reduced cross section



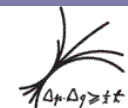
$$F_2(x, Q^2) = \sigma_r(x, Q^2, y=0)$$

$$F_L(x, Q^2) = -\frac{\partial \sigma_r(x, Q^2, y)}{\partial (y^2 / Y_+)}$$

Measure σ_r at fixed x, Q^2 but varying y

$$y = Q^2 / sx \quad \sqrt{s} = ep \text{ center-of-mass energy}$$

Varying $y \rightarrow$ varying $s \rightarrow$ dedicated low E_p runs at end of HERA



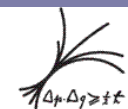
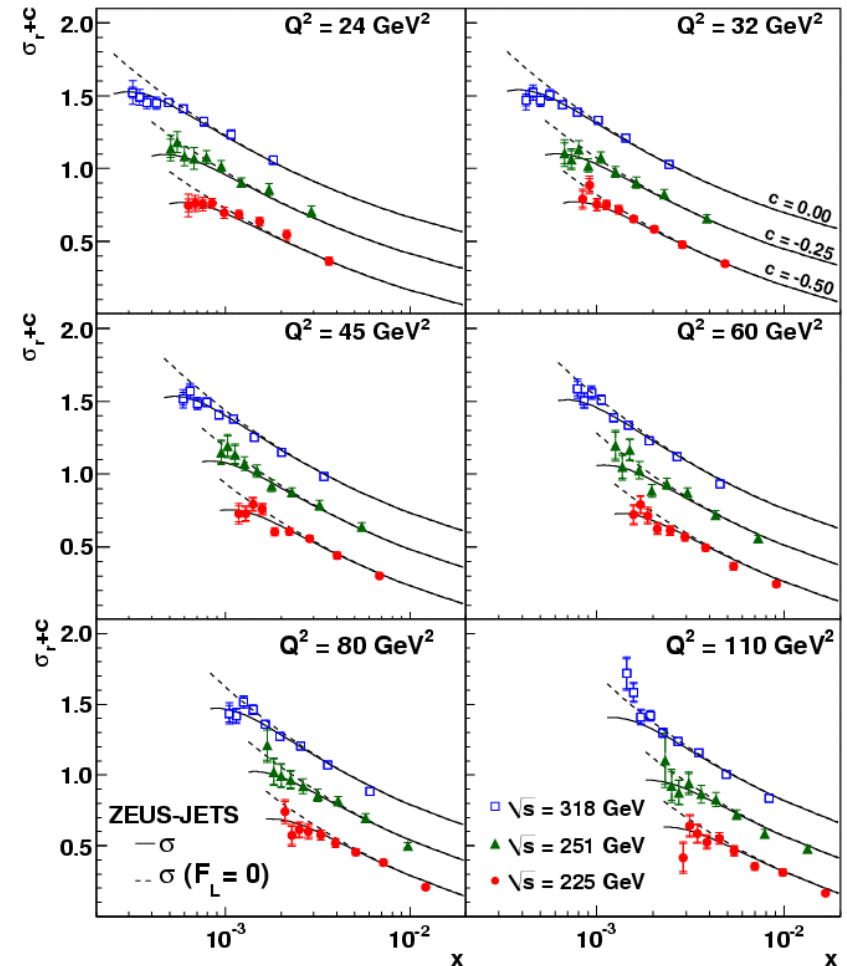
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Cross Sections at low x, high y

- σ_r measured at 3 different \sqrt{s}
- comparison to ZEUS-JETS PDF prediction with $F_L = F_L(\text{QCD})$ and $F_L=0$
- F_L causes suppression at low x
- Small effect but big enough to extract F_L and F_2

ZEUS



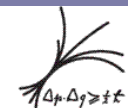
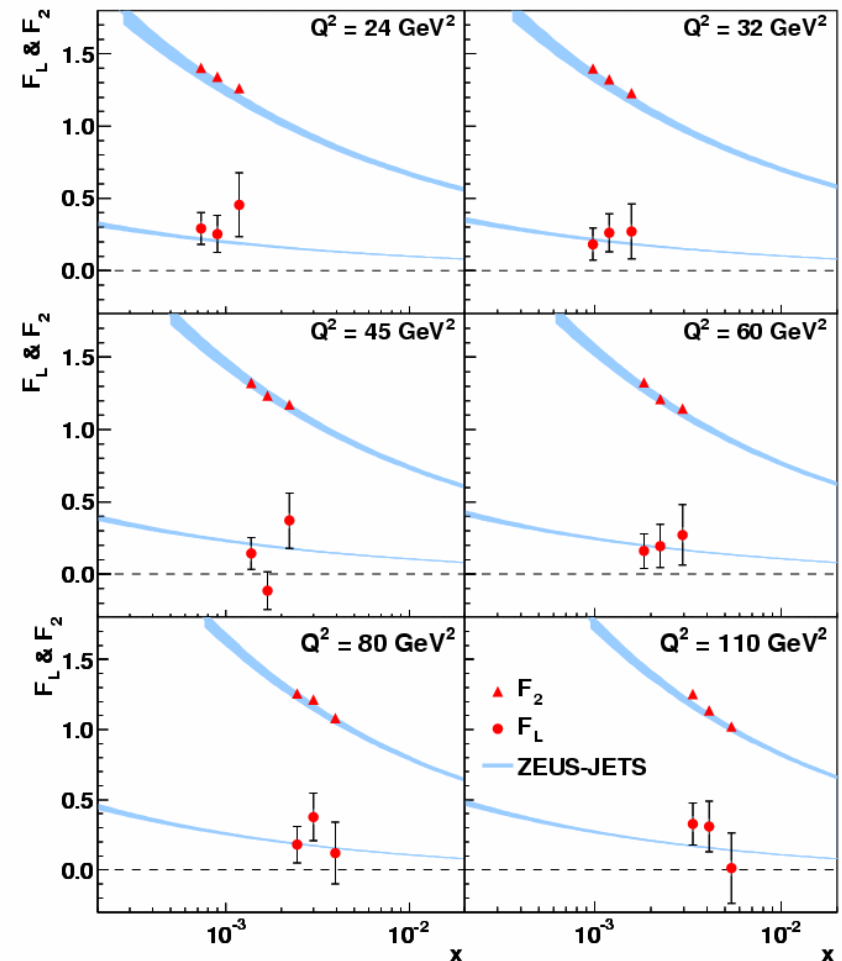
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Extracted Structure Functions

- Simultaneous extraction of $F_L(x, Q^2)$ and $F_2(x, Q^2)$
- Error bars: experimental errors (stat & syst)
- Data support non-zero F_L
- Consistent with expectation from QCD

ZEUS

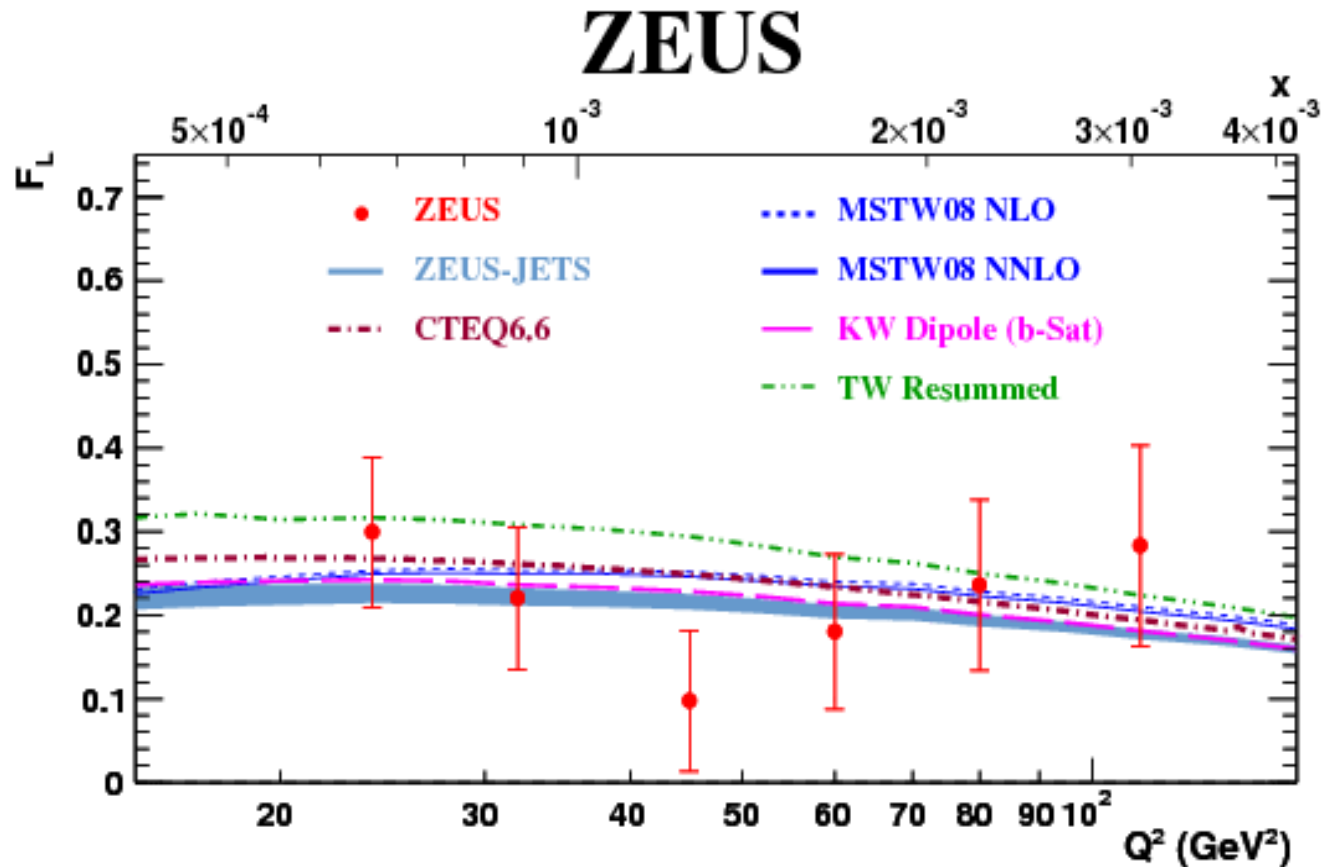


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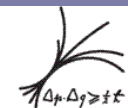
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Longitudinal Structure Function

- $F_L(Q^2)$ i.e. averaged x
- FL clearly non-zero



- Compared to various predictions, consistent with data
- Import Cross check of QCD



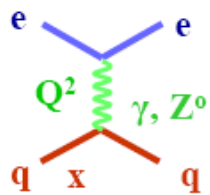
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Input to HERA PDF Fits

Inclusive DIS data: Neutral Current and Charged Current

■ NC
$$\frac{d^2 \sigma^{e^\pm p \rightarrow e^\pm X}}{dx dQ^2} = \frac{2\pi\alpha^2}{xQ^4} \underbrace{\left(1 + (1-y)^2\right)}_{Y_\pm = 1 \pm (1-y)^2} \cdot \left(\tilde{F}_2(x, Q^2) - \frac{y^2}{Y_+} \tilde{F}_L(x, Q^2) \mp \frac{Y_-}{Y_+} x\tilde{F}_3(x, Q^2) \right)$$

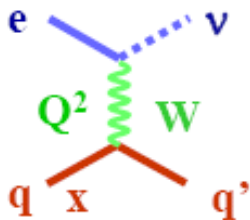


$$\tilde{F}_2 = \sum_i A_i(Q^2) [xq_i + x\bar{q}_i] \Rightarrow F_2^{em} = \frac{4}{9} x(u + \bar{u} + c + \bar{c}) + \frac{1}{9} x(d + \bar{d} + s + \bar{s})$$

$$x\tilde{F}_3 = \sum_i B_i(Q^2) [xq_i - x\bar{q}_i] \Rightarrow xF_3 = B_U x(u - \bar{u} + c - \bar{c}) + B_D x(d - \bar{d} + s - \bar{s})$$

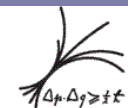
Electroweak Coefficient Functions $A_i(Q^2)$, $B_i(Q^2)$ (QED: $A_i = e_i^2$)

■ CC



$$\sigma_{CC}(e^+ p) \propto x \left[(1-y^2)(d+s) + (\bar{u} + \bar{c}) \right] \times (1 + P_e)$$

$$\sigma_{CC}(e^- p) \propto x \left[(u+c) + (1-y^2)(\bar{d} + \bar{s}) \right] \times (1 - P_e)$$

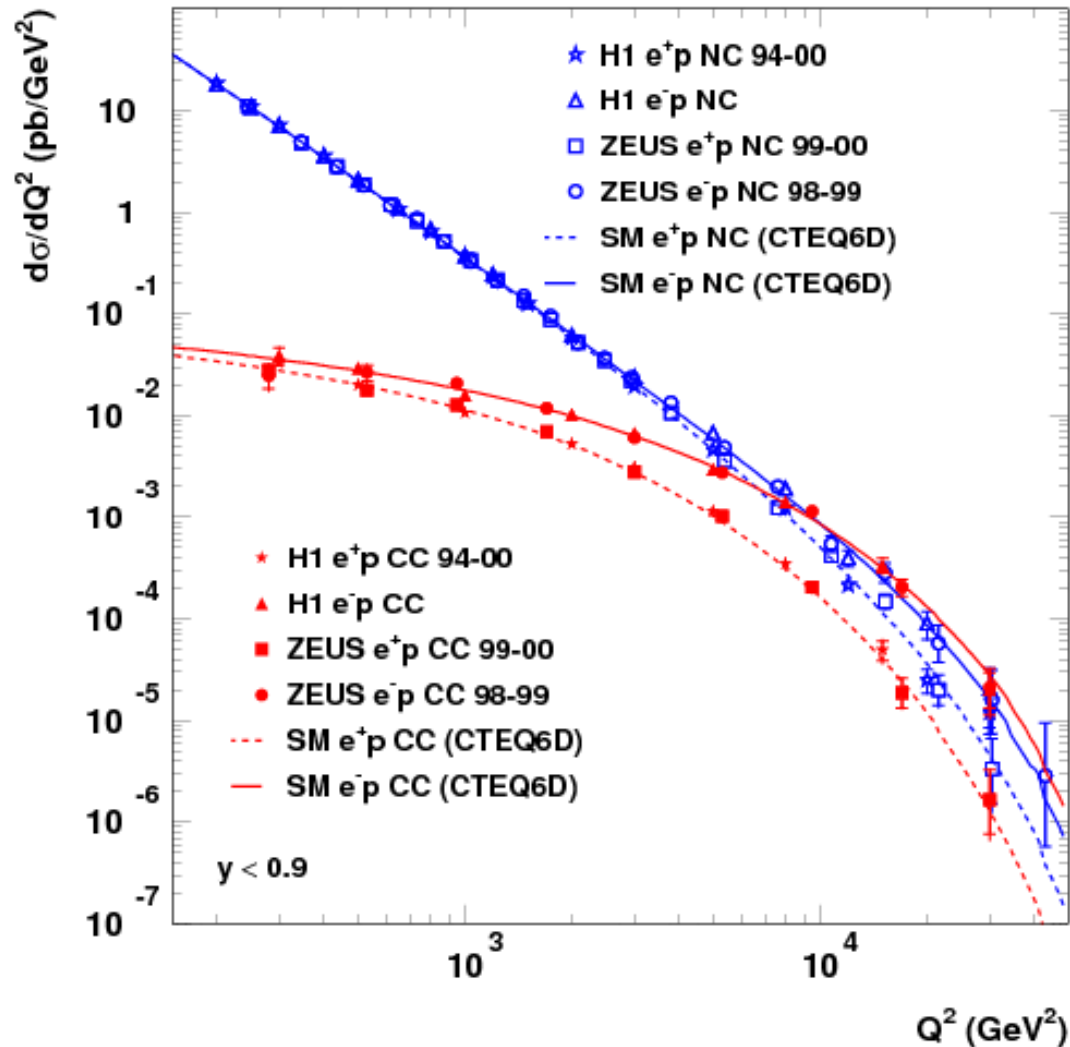


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Electroweak Unification at High Q^2 (NC & CC)

HERA

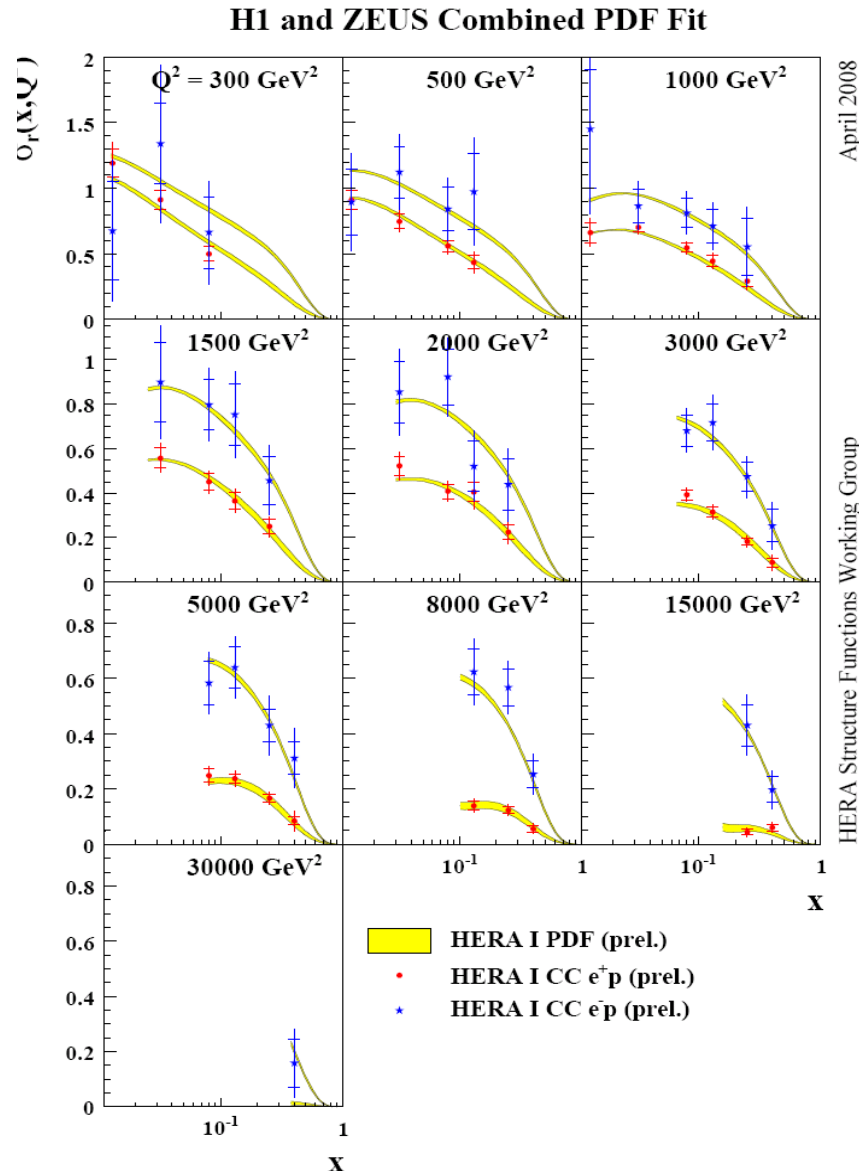


$\sigma_{NC} \gg \sigma_{CC}$
for $Q^2 \ll M_Z^2$
(photon exchange dominates)

$Q^2 \geq M_Z^2 : \sigma_{CC} \sim \sigma_{NC}$

manifest
electroweak unification

Charged Current Cross Section

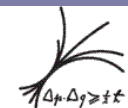


CC Cross section provide flavor sensible constraints at high x

$$\sigma_{CC}(e^+ p) \propto x \left[(1 - y^2) D + \bar{U} \right]$$

$$\sigma_{CC}(e^- p) \propto x \left[U + (1 - y^2) \bar{D} \right]$$

**Improved precision of σ_{CC}
By combining H1 and ZEUS**



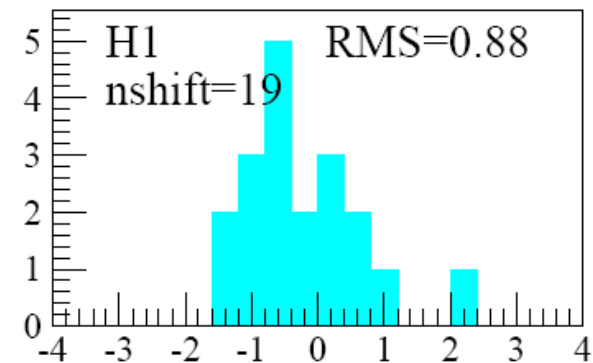
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PDF Fit: χ^2 -Definition

$$\chi^2 = \sum_k \sum_i^{n_{exp} N_k} \frac{\left(\sigma_{k,i}^{exp} - \sigma_{k,i}^{th} \cdot (1 - n_k \delta_k^{norm} - \sum_j^{n_s(k)} s_{k,j} \delta_{k,i,j}^{syst}) \right)^2}{\delta_{k,i}^{sta^2} + \delta_{k,i}^{unc^2}} + \sum_k^{n_{exp}} n_k^2 + \sum_k^{n_{exp}} \sum_j^{n_s(k)} s_{j,k}^2$$

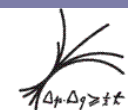
- We explicitly allow the data points to be shifted by the fit



$\sigma_{k,i}^{exp}$: measurement i of dataset k
 $\sigma_{k,i}^{th}$: prediction for data point i of dataset k
 n_{exp} : number of datasets
 N_k : number of data points in dataset k
 $\delta_{k,i}^{sta}$: statistical error

$\delta_{k,i}^{unc}$: uncorrelated systematic error
 $\delta_{k,i,j}^{sys}$: correlated systematic error of source j
 δ_k^{norm} : uncertainty of dataset normalization
 n_k : shift of normalization
 $s_{k,j}$: shift of systematic of source j

“Pascaud and Zomer” method, extended χ^2 (preprint LAL-95-05):



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Details on HERAPDF 0.1

Chosen form of the PDF parametrization at Q_0^2

$$xf(x) = Ax^B(1-x)^C(1+Dx+Ex^2+Fx^3\dots)$$

	A	B	C	D	E
gluon	sum rule				
u_v	sum rule				
d_v	sum rule	$= B(u_v)$			
U_{bar}	Lim $x \rightarrow 0 \ u/d \rightarrow 1$				
D_{bar}		$= B(U)$			

The number of parameters for each parton has been optimized

Optimization means starting with only BLUE parameters and adding D, E, F parameters until there is no further χ^2 advantage

PDFs fitted: gluon, u_v , d_v , $U_{bar}=u_{bar}+c_{bar}$, $D_{bar}=d_{bar}+s_{bar}+b_{bar}$

Sea flavour break-up at Q_0 : $s = fs \cdot D$, $c = fc \cdot U$, $AU_{bar} = (1-fs)/(1-fc)AD_{bar}$

Lim $x \rightarrow 0 \ u_{bar}/d_{bar} \rightarrow 1$

$fs = 0.33D$ ($s=0.5d$), $fc = 0.15U$ consistent with dynamical generation

$m_c=1.4$ GeV mass of charm quark $m_b=4.75$ GeV mass of beauty quark

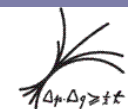
Zero-mass variable flavour number heavy quark scheme (for now)

$Q_0^2 = 4$ GeV² input scale

$Q_{min}^2 = 3.5$ GeV² minimum Q^2 of input data

$\alpha_s(M_Z) = 0.1176$ PDG2006 value

Renormalization and factorization scales = Q^2



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Comparison of Separate H1 and ZEUS PDF Fits

$$xf(x) = Ax^B(1-x)^C(1+Dx+Ex^2+Fx^3\dots)$$

Alternative form of PDF parametrization: H1 style

	A	B	C	D	E	F
gluon	sum rule					
U	$\lim_{x \rightarrow 0} \bar{u}/\bar{d} \rightarrow 1$			sum rule		
D		$= B(U)$		sum rule		
U_{bar}	$= A(U)$	$= B(U)$				
D_{bar}	$= A(D)$	$= B(U)$				

PDFs: gluon, $U=u+c$, $U_{\text{bar}}=u_{\text{bar}}+c_{\text{bar}}$, $D=d+s+b$, $D_{\text{bar}}=d_{\text{bar}}+s_{\text{bar}}+b_{\text{bar}}$

Sea flavour break-up at Q_0 : $s = f_s \cdot D$, $c = f_c \cdot U$ $AU = (1-f_s)/(1-f_c)AD$

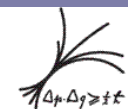
Alternative form of PDF parametrization: ZEUS style

	A	B	C	D	E
gluon	From Sum Rule				0.
u_v	From Sum Rule				
d_v	From Sum Rule	$= B_{uv}$			0.
$u_{\text{bar}} - d_{\text{bar}}$	from Z S 11 fit	from Z S 11 fit	from Z S 11 fit	0.	0.
Sea				0.	0.

PDFs: gluon, u_v , d_v , Sea = $u_{\text{sea}} + u_{\text{bar}} + d_{\text{sea}} + d_{\text{bar}} + s + s_{\text{bar}} + c + c_{\text{bar}}$

Sea flavour break-up at Q_0 : $s_{\text{bar}} = (d_{\text{bar}} + u_{\text{bar}})/4$, charm dynamically generated,

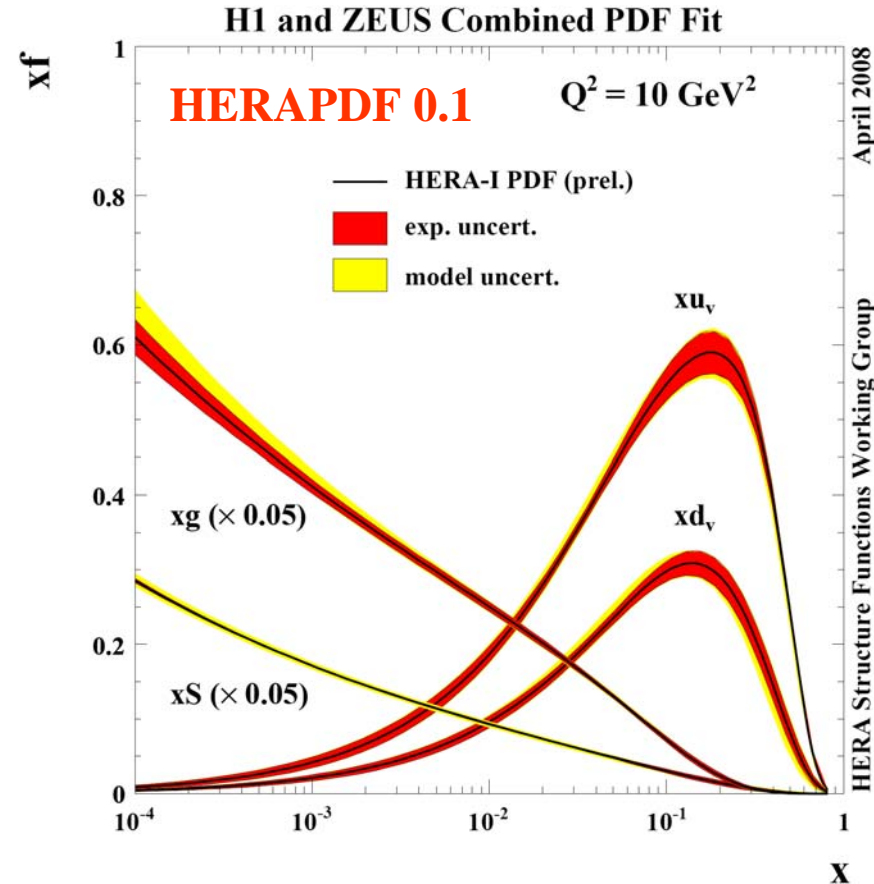
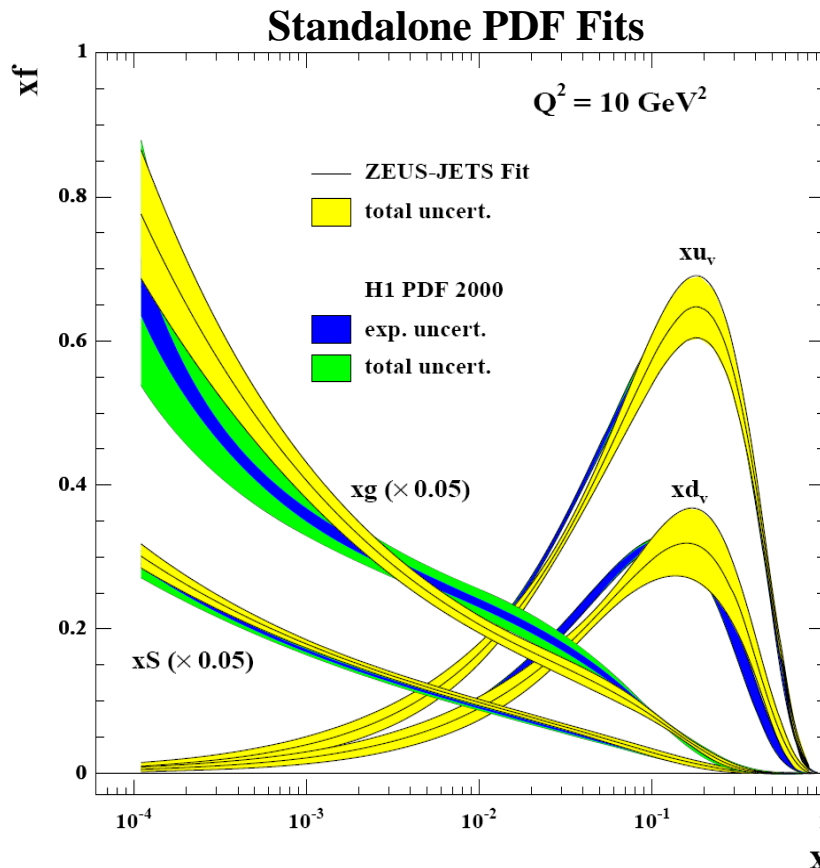
$d_{\text{bar}} - u_{\text{bar}}$ fixed to fit E866 data



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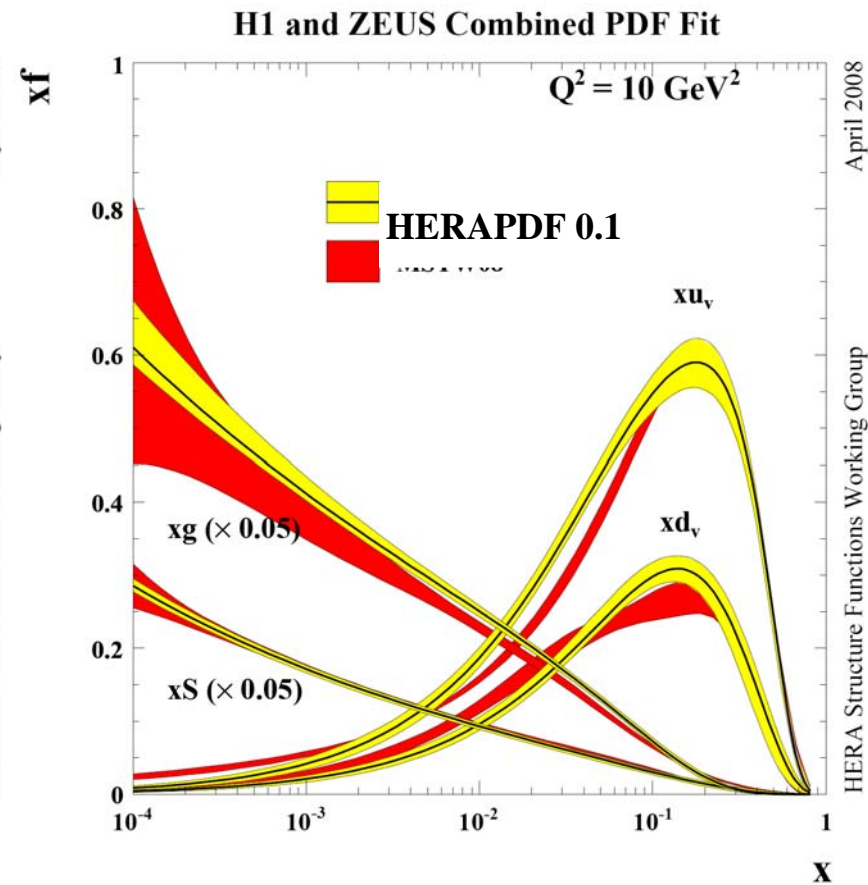
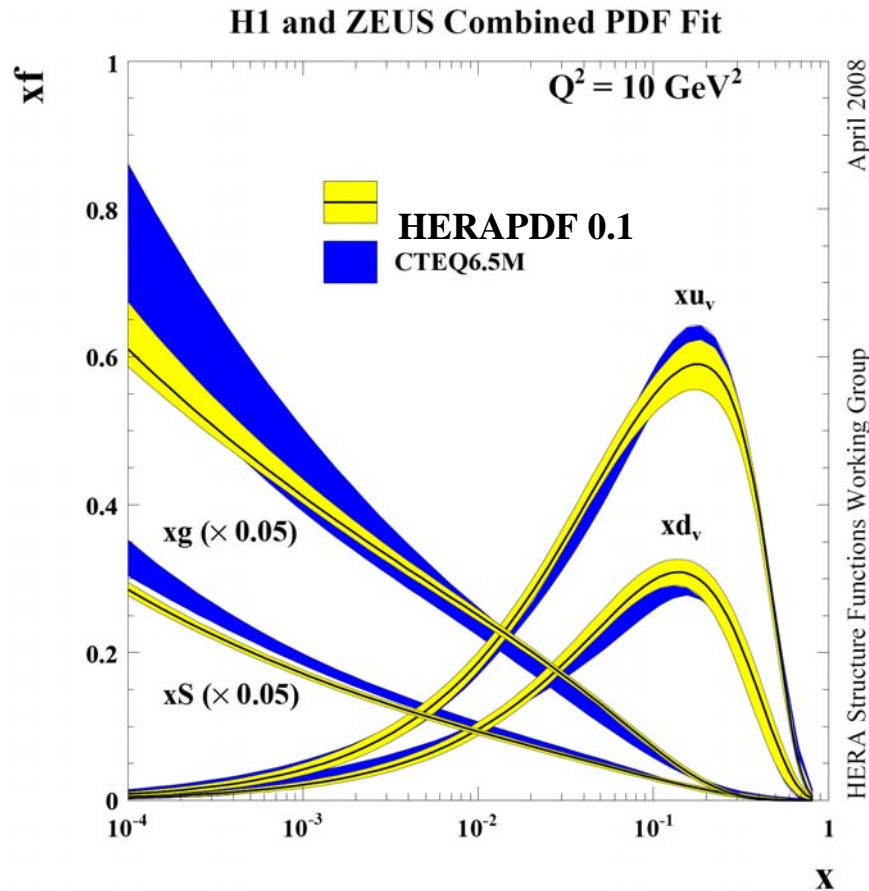
PDF Fits on HERA I Data



Impressive reduction of uncertainties of combined PDFs

Model uncertainty: variation of charm and bottom mass, starting scale Q_0^2 , Q_{\min}^2 of included data, strange and charm fraction at starting scale

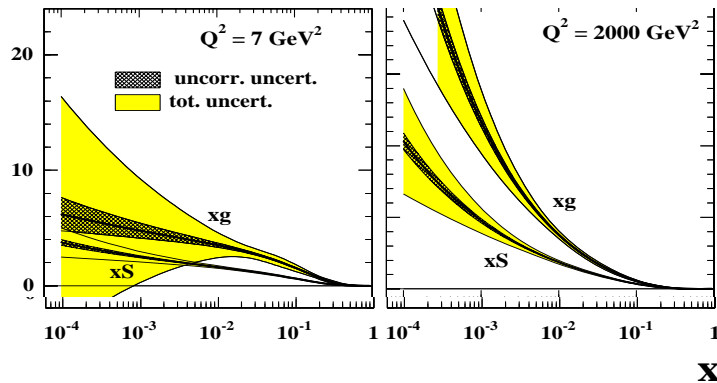
Comparison to Global Fits



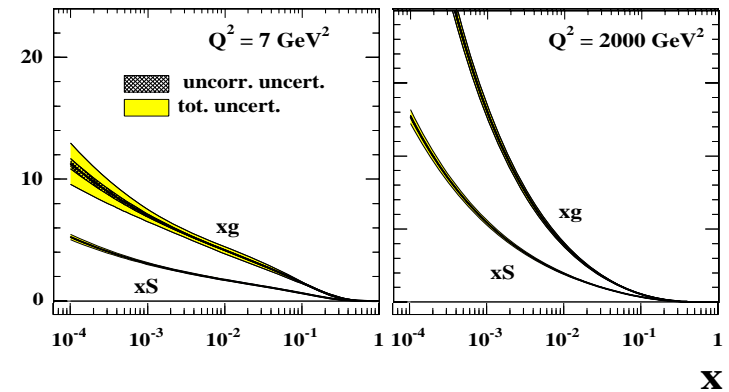
not a completely fair comparison:
HERA combined data were not available to global fitters

Impact of HERA data

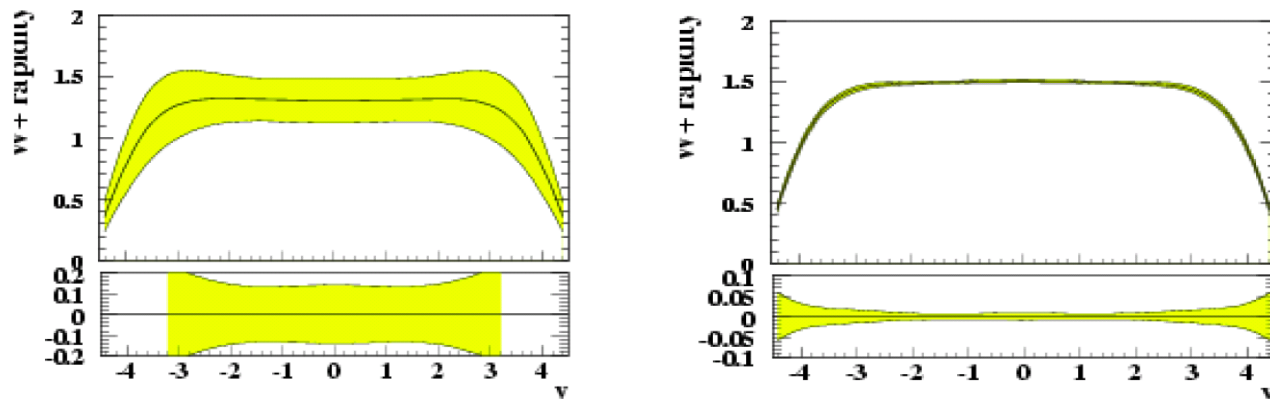
Pre-HERA uncert. in gluon/sea PDFs



HERA PDFs combining H1 and ZEUS



Example: W^+ production at the LHC (Study by A. Cooper-Sakar)



Note: Error bands are experimental uncertainties only
model uncertainty will become increasingly important

Open Issue

Light flavor decomposition of the $q\bar{q}$ sea

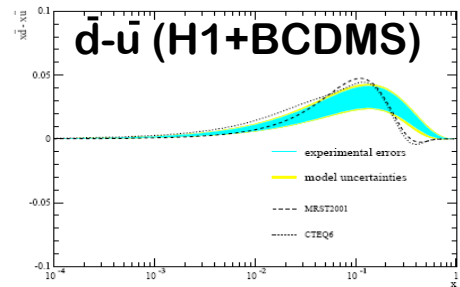
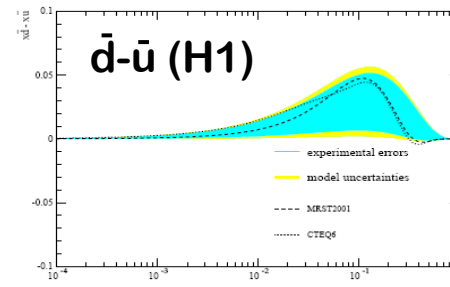
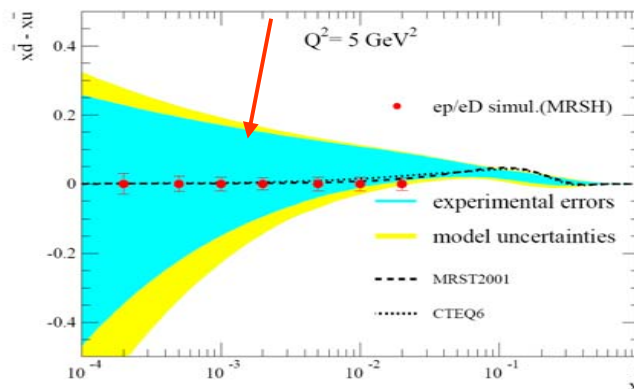
PDF fits conventionally assume

$$x\bar{d} - x\bar{u} \xrightarrow{x \rightarrow 0} 0$$

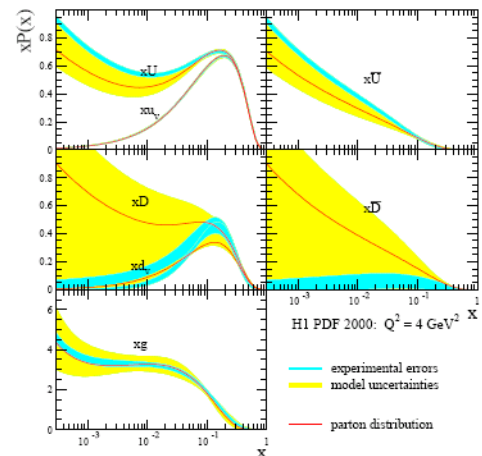
NMC found $d \neq u$ at medium x

Here is what happens when
the $x\bar{d} - x\bar{u}$ constraint at low x
is relaxed

A deuteron run at HERA
Could have disentangled
the light flavor sea



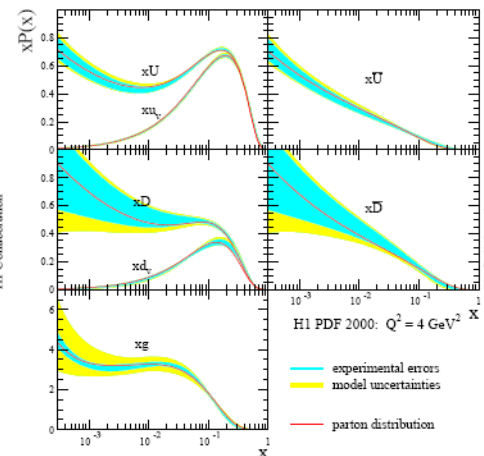
H1 only (HERA I)



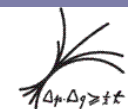
Poorly constrained

Attempt to fit U and D
Only one input: F_2 ep

H1 + BCDMS

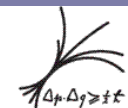


Fit stabilized by
fixed target data
(sum rules help)



Future HERA PDF Fits

- So far only part of the inclusive HERA I were used → HERAPDF0.1
- ➔ Incorporate all NC and CC from HERA I&II
- ➔ Include jet cross sections
 - ➔ constrain high x gluon
- ➔ Include charm and beauty
 - ➔ flavor decomposition of the sea
- ➔ Charged & Neutral current cross sections with polarized e^\pm beams
 - ➔ constrain valence quark region

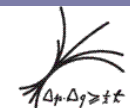


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Conclusion

- Deep Inelastic Scattering revealed the Structure of Nucleon
- Parton Densities of Proton obtained from inclusive cross section measurements



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Backup slides



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Charm & Beauty Structure

Charm and Beauty production in DIS
is driven by gluons in the proton

Charm tag: reconstruct D mesons

Beauty tag: displaced vertex, soft μ

