

# Characterization of a Segmented n-Type Broad Energy Germanium Detector

Martin Schuster

Max-Planck Institute for Physics



March 13, 2017



Physics Motivation

Experimental Setup

Detector Characteristics

Summary & Outlook



Physics Motivation

Experimental Setup

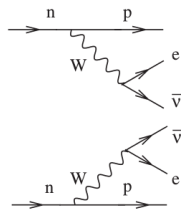
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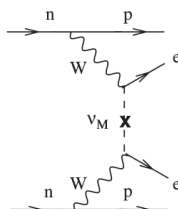
# Physics Motivation: $0\nu\beta\beta$ Decay

Could provide information on:

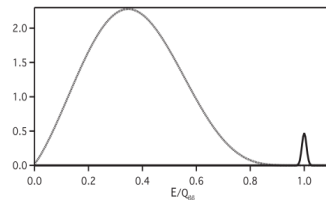
- ▶ Nature of the neutrino: Dirac or Majorana?
- ▶ Inverted or Normal Hierarchy
- ▶ Possibly hints on absolute mass scale from  $T_{1/2}$



(a)  $2\nu\beta\beta$



(b)  $0\nu\beta\beta$



(c) Combined spectrum



## Why Germanium?

- ▶  $^{76}\text{Ge}$  is a candidate for  $0\nu\beta\beta$  and Ge is a semiconductor
- ▶ Can act as source and detector simultaneously  
→ high detection efficiency
- ▶ Naturally good energy resolution
- ▶ Currently employed in experiments like GERDA and MAJORANA
- ▶ LEGEND collaboration is forming: Ton-scale germanium experiment



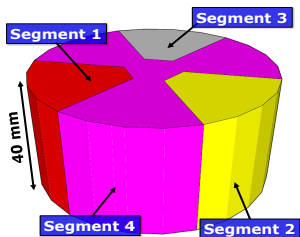
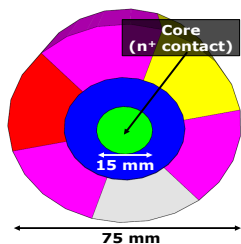
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# The Segmented Broad Energy Germanium Detector

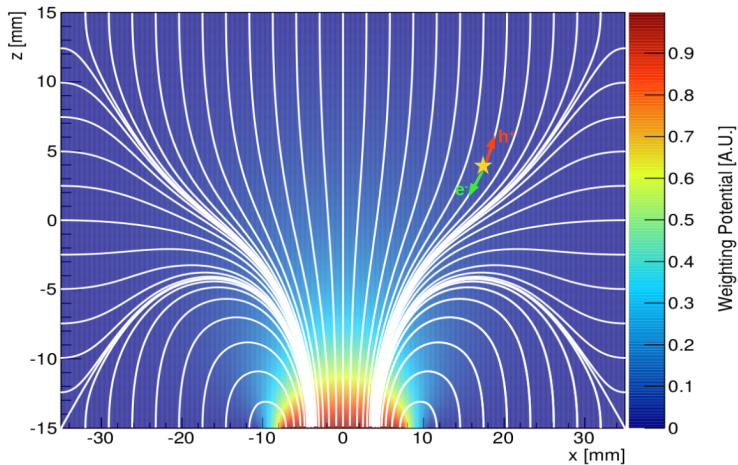


## Properties

- ▶ Special electric field due to point contact: Favorable for Pulse-Shape Analysis
- ▶ Segmentation provides
  - ▶ Information on  $\phi$  location of energy depositions
  - ▶ Complimentary background rejection

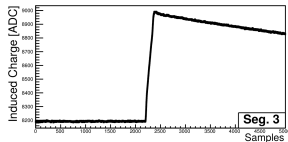
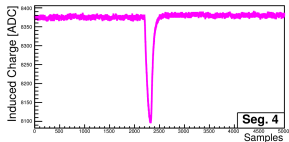
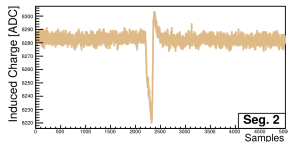
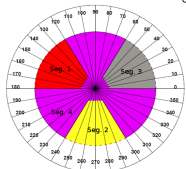
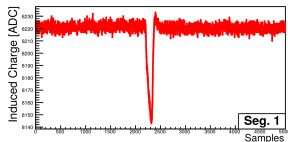
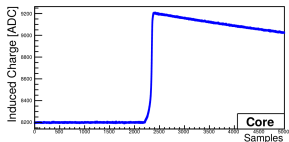


# Weighting Potential





# Example Pulse



Physics Motivation

Experimental Setup

**Detector Characteristics**

Summary & Outlook

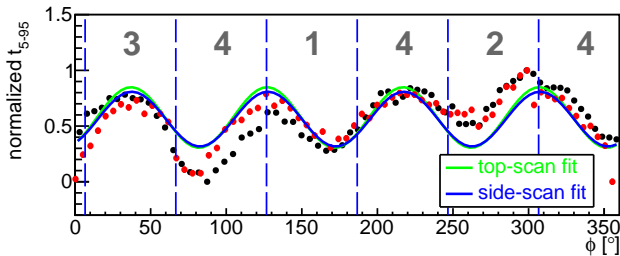
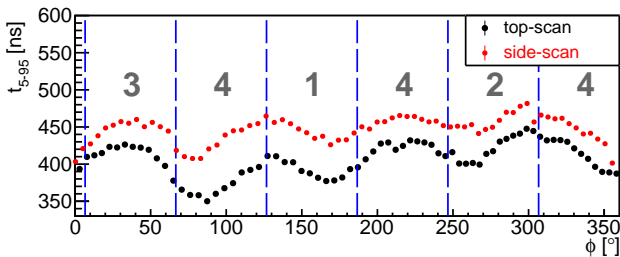
# Crystal Axes Determination

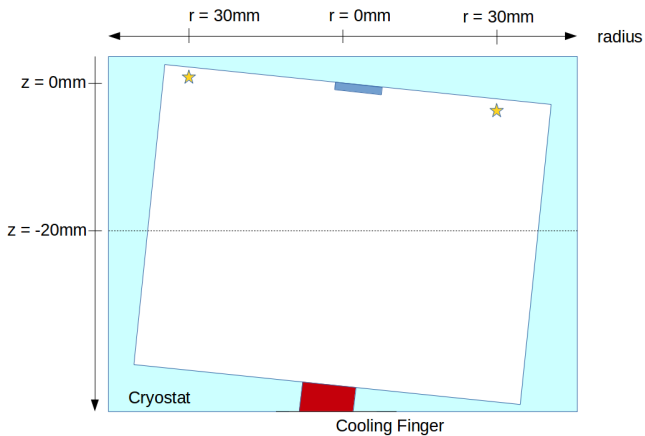
- ▶ Slow and fast crystal axes influence the drift-path and speed of the charge carriers
- ▶ The dependence of  $t_{5-95}$  on  $\phi$  is expected to be sinusoidal:

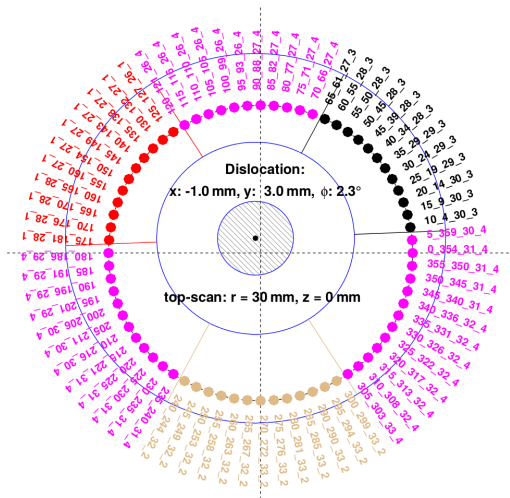
$$t_{5-95} = C + a \cdot \sin \left[ \frac{2\pi}{90} (\phi + \phi_{offset}) \right] .$$

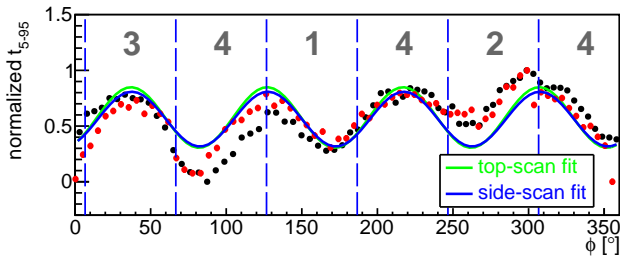
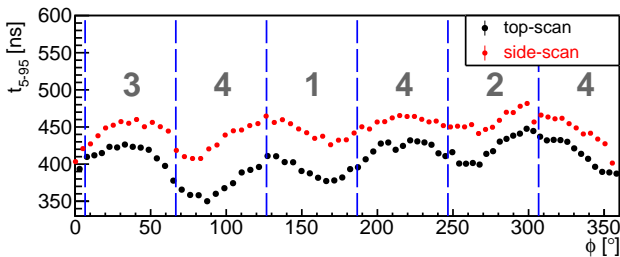
The parameters  $C$ ,  $a$  and  $\phi_{offset}$  were fitted to the data



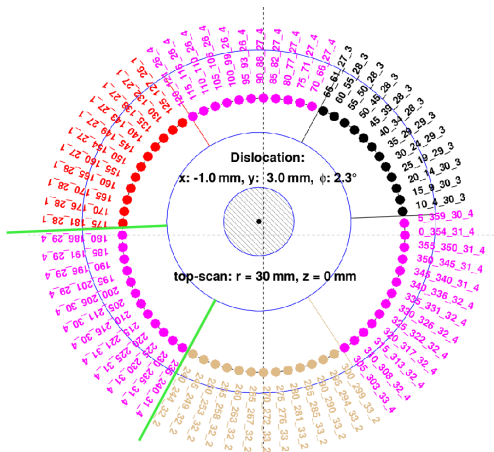




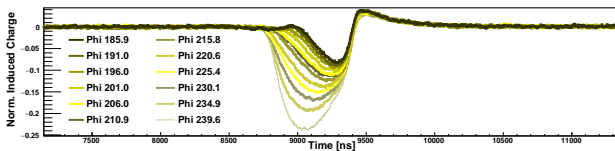
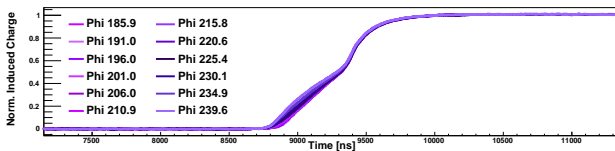
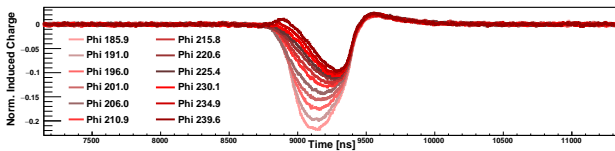


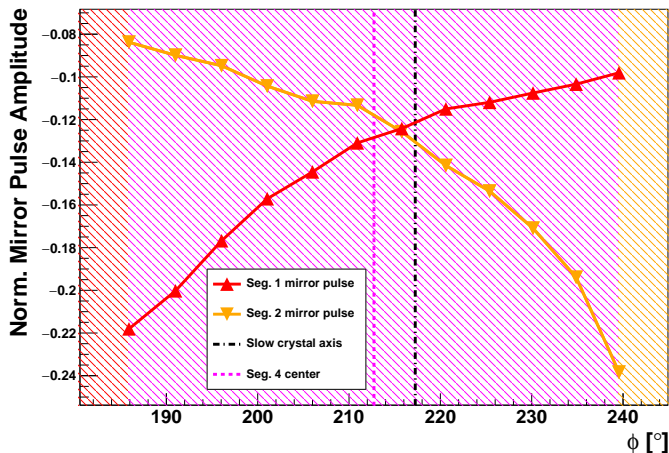


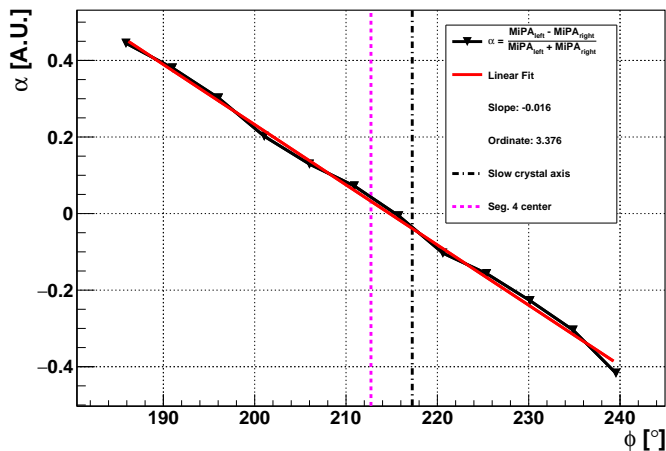
# Localizing $\phi$ Position using Mirror Pulses



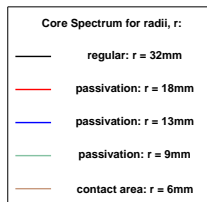
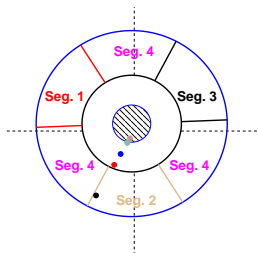
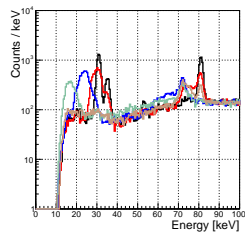
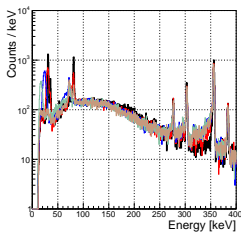


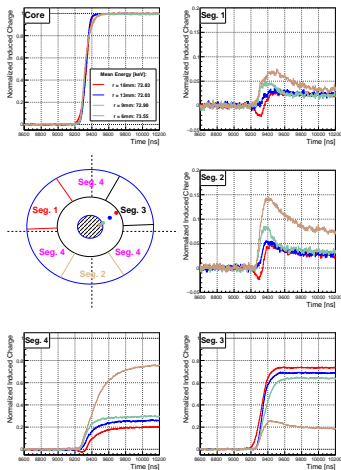






# Passivation Layer Studies





- ▶ Both segment 3 and 4 are collecting
- ▶ Truncated mirror pulses: Indication for charge trapping
- ▶ Increasing effect for smaller radii



## Near Surface Effects: Distorted Field Lines

- ▶ Drift paths are bent towards the surface, where the electric field is very weak
- ▶ At surface, the charge-drift slows down significantly
- ▶ This gives rise to the observed phenomena via two mechanics:
  - ▶ Incomplete charge collection: Charges cease to contribute to the pulse
  - ▶ Strongly delayed pulses: Pulses don't reach their amplitude during DAQ-window



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## Summary and Results

- ▶ Prototype detector, combining the point contact layout, which is favorable for PSA with a moderate segmentation.
- ▶ Localization of the crystal axes, despite tilted installation:  
Good agreement between top-scans and side-scans
- ▶ Reconstructing the  $\phi$ -position using mirror pulses of non-collecting segments: Proven to be feasible
- ▶ Indication of distorted drift paths for low energy events beneath the passivated area  
The effects show a significant  $r$  dependence, yet they are independent on  $\phi$





## Outlook

- ▶ The detector has been moved to new, electrically cooled cryostat
- ▶ Measurements of temperature dependence of passivation layer effects and drift times in general
- ▶ Expanding and refining the event position reconstruction, including Monte-Carlo based pulse-shape libraries
- ▶ Further studies on the effect of the crystal axes

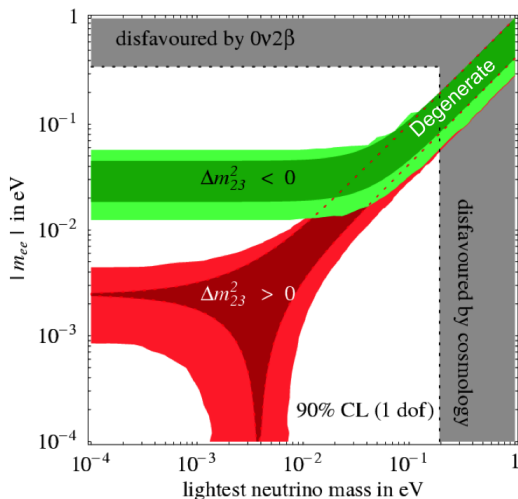


Thank you  
for your  
attention!



# BACKUP



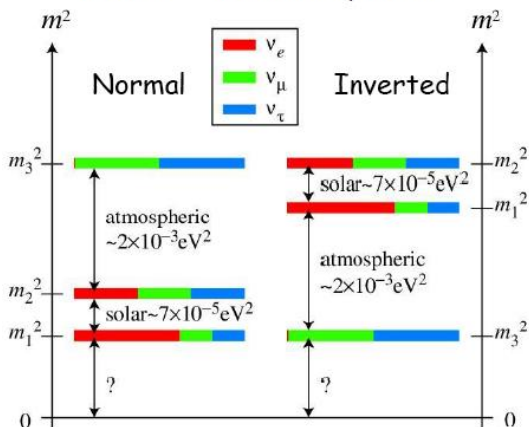


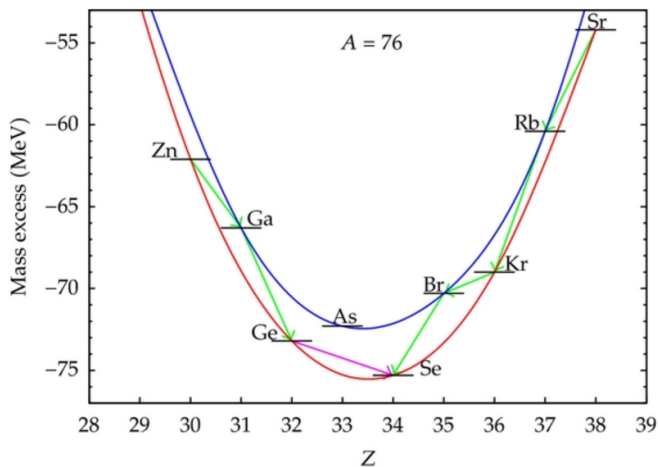
1

<sup>1</sup>I. Gil-Botella, Neutrino Physics (2015)



in the 3-neutrino picture

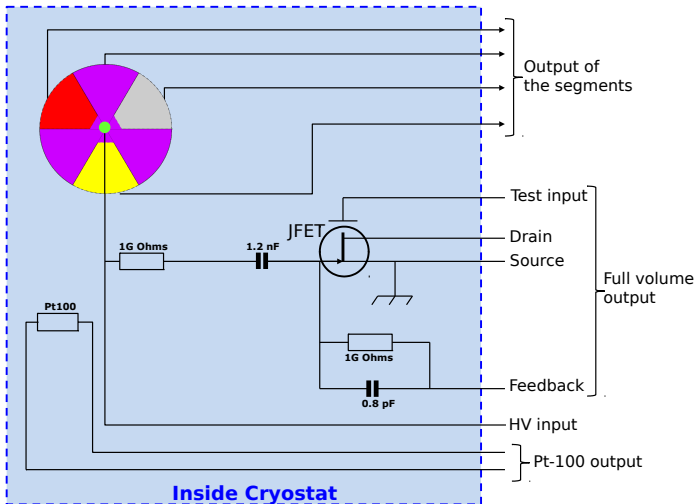




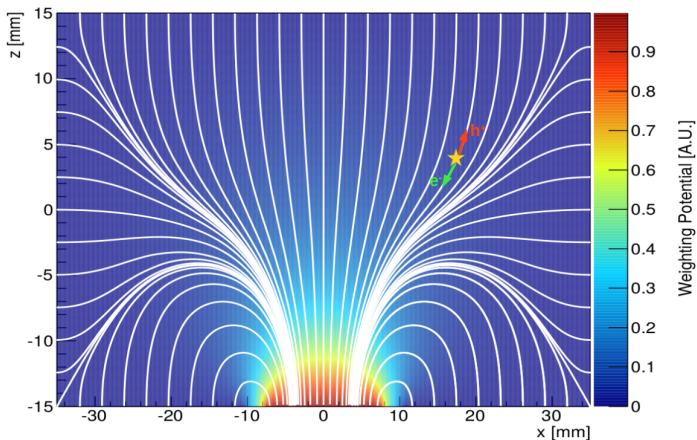
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<sup>2</sup>A. Giullani and A. Poves, Advances in High Energy Physics(2012)

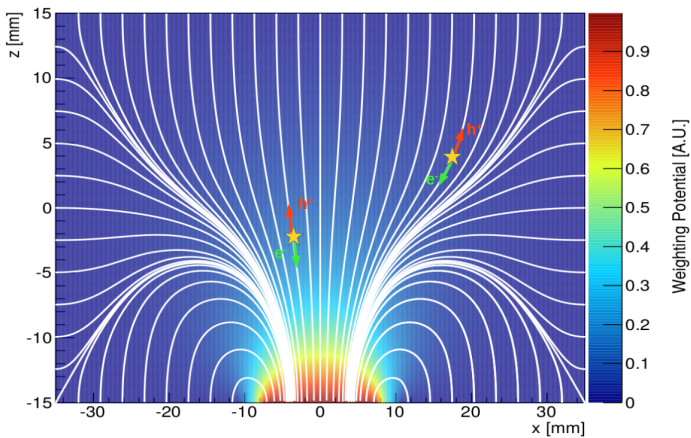




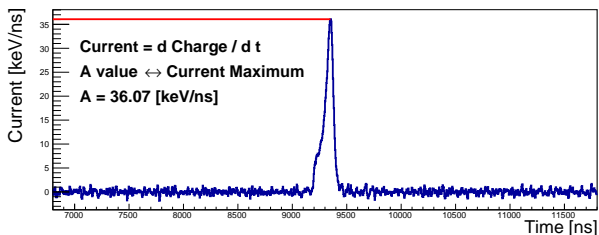
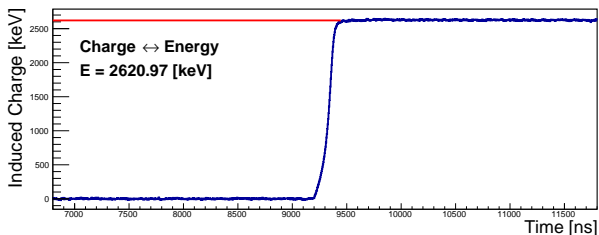
# Multi-Site- vs. Single-Site-Event Discrimination

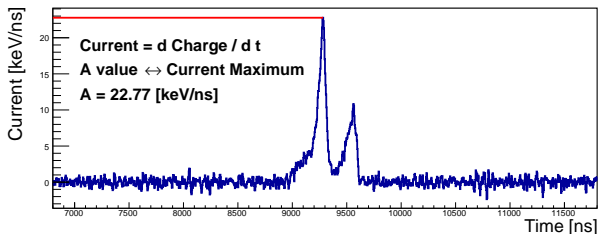
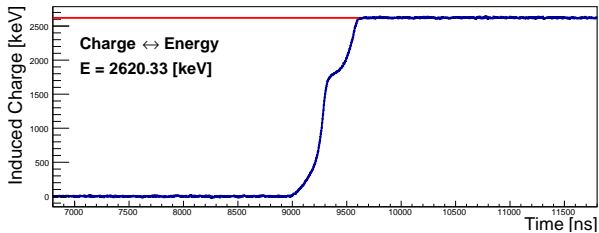


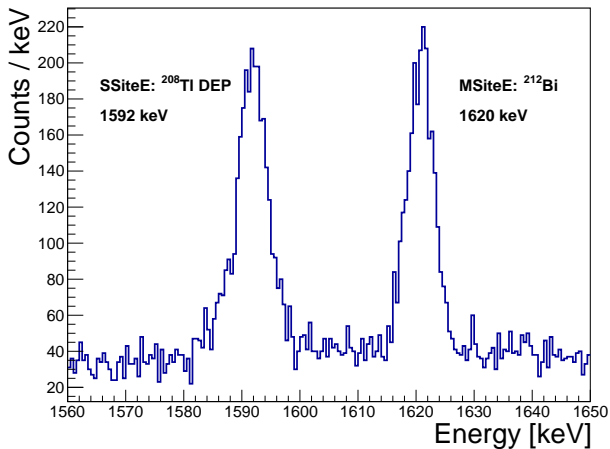


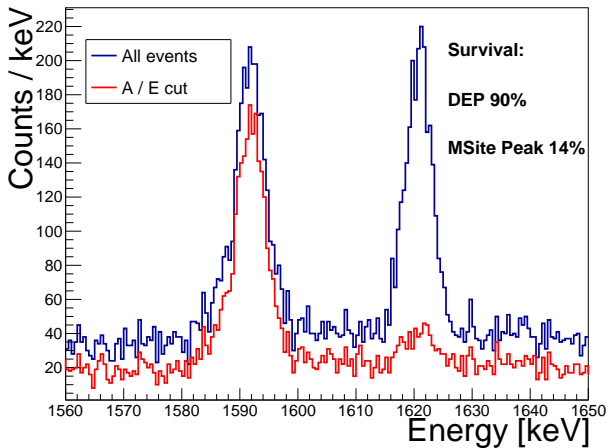


# A/E Method

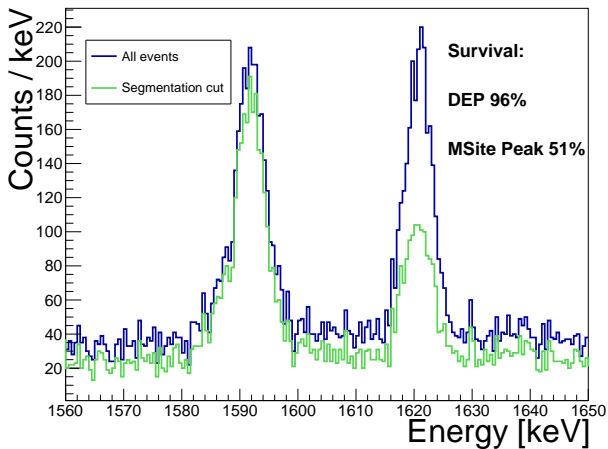


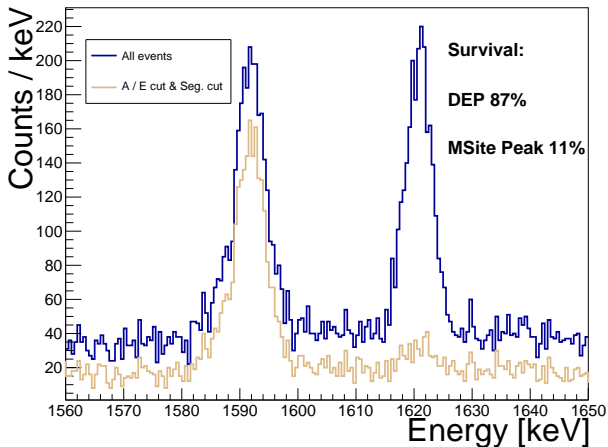




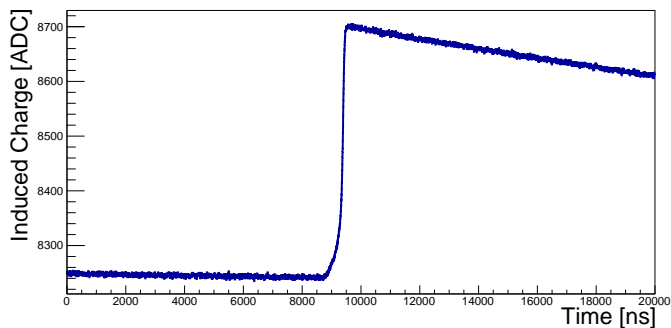


# Segmentation Cut



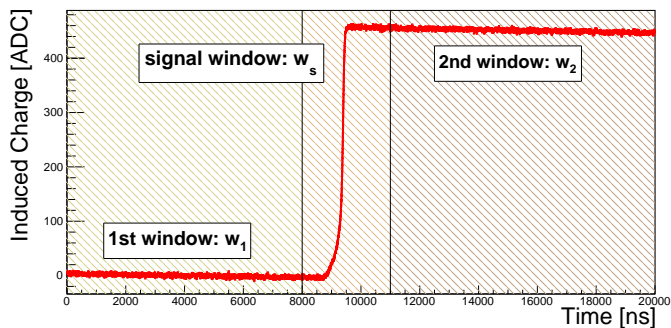


## Example Core Pulse as provided by the ADC





# Fixed Window Method for Pulse Amplitude Determination



The matrix  $c_{ij}^*$  links the true energy  $E_i^{true}$  and the measured pulse amplitudes  $M_i$ :

$$(E_0^{true} \quad \dots \quad E_4^{true}) \cdot \begin{pmatrix} c_{00}^* & \dots & c_{04}^* \\ \vdots & \ddots & \vdots \\ c_{40}^* & \dots & c_{44}^* \end{pmatrix} = (M_0 \quad \dots \quad M_4) \quad .$$

$E_i^{true}$  can be determined using the inverse matrix  $c_{ij}^{*-1}$ :

$$(E_0^{true} \quad \dots \quad E_4^{true}) = (M_0 \quad \dots \quad M_4) \cdot \begin{pmatrix} c_{00}^* & \dots & c_{04}^* \\ \vdots & \ddots & \vdots \\ c_{40}^* & \dots & c_{44}^* \end{pmatrix}^{-1} \quad .$$



$$M_0 = E_0^{true} \cdot c_{00}^* + E_1^{true} \cdot c_{10}^* + \dots + E_4^{true} \cdot c_{40}^* \quad (1)$$

By introducing the assumption that the cross-talk from the segments to the core is the same for all the segments,  $c_{10}^* = c_{i0}^*$ ,  $i \in \{2, 3, 4\}$  and using the identity  $\sum_{i=1}^4 E_i^{true} = E_0^{true}$ ,  $M_0$  becomes:

$$M_0 = E_0^{true} \cdot (c_{00}^* + c_{10}^*) \quad (2)$$

$$= E_0^{true} \cdot c^{core} \quad (3)$$

The value of the thus defined  $c^{core}$  can be determined by the standard procedure of fitting the core spectrum with the known photon lines using all events.



The cross-talk from the core to the segments can mathematically be unified with the segment to segment cross-talk. For demonstration, consider the measured pulse amplitude for segment 1,  $M_1$ . Again making use of the identity  $\sum_{i=1}^4 E_i^{true} = E_0^{true}$ :

$$M_1 = \sum_{i=0}^4 E_i^{true} \cdot c_{i1}^* \quad (4)$$

$$= E_0^{true} \cdot c_{01}^* + \sum_{i=1}^4 E_i^{true} \cdot c_{i1}^* \quad (5)$$

$$= \sum_{i=1}^4 E_i^{true} \cdot c_{01}^* + \sum_{i=1}^4 E_i^{true} \cdot c_{i1}^* \quad (6)$$

$$= \sum_{i=1}^4 E_i^{true} (c_{01}^* + c_{i1}^*) = \sum_{i=1}^4 E_i^{true} \cdot c_{i1} \quad (7)$$



Reduced matrix  $c_{ij}$ ,  $i, j \in \{1, 2, 3, 4\}$  and the equations:

$$M_0 = E_0^{true} \cdot c^{core} \quad (8)$$

$$(M_1 \quad \dots \quad M_4) = (E_1^{true} \quad \dots \quad E_4^{true}) \cdot \begin{pmatrix} c_{11} & \dots & c_{14} \\ \vdots & \ddots & \vdots \\ c_{41} & \dots & c_{44} \end{pmatrix} \cdot \quad (9)$$

The matrix elements  $c_{ij}$  can be determined by investigating single-segment events.



For single-segment  $i$  events,  $E_i^{true} = E_0^{true}$ ,  $E_{j \neq i}^{true} = 0$ . Using segment 1 as an example, the measured pulse amplitude for single-segment 1 events,  $M_1$ , can then be written as:

$$M_1 = E_1^{true} \cdot c_{11} + E_2^{true} \cdot c_{21} + E_3^{true} \cdot c_{31} + E_4^{true} \cdot c_{41} \quad (10)$$

$$= E_1^{true} \cdot c_{11} + 0 \cdot c_{21} \dots \quad (11)$$

In general, for single-segment  $i$  events,  $M_i = E_i \cdot c_{ii}$ ,  $i \in \{1, 2, 3, 4\}$ . Then the matrix elements  $c_{ii}$  can be determined:

$$c_{ii} = \frac{M_i}{E_i^{true}} = \frac{M_i}{E_0^{true}} = \frac{M_i \cdot c^{core}}{M_0} \quad (12)$$

$$= R^{ssi} \cdot c^{core} \quad (13)$$

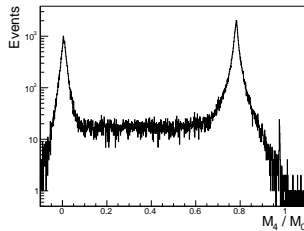
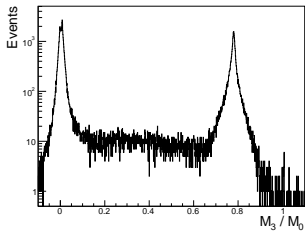
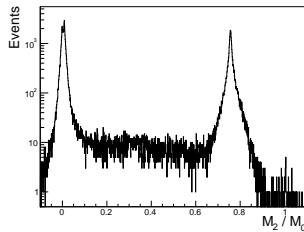
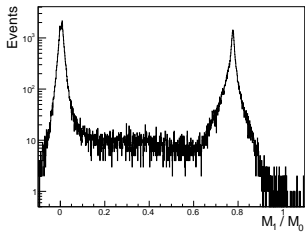


After the cross-talk correction and energy calibration procedure, Tab 1 gives an overview of the energy resolutions for selected  $\gamma$ -lines and all the segments.

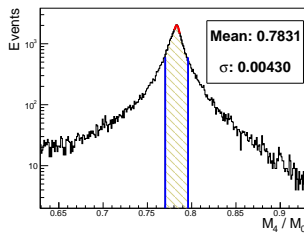
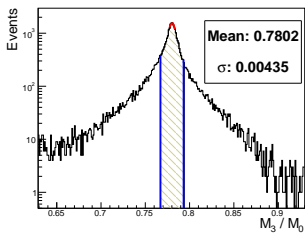
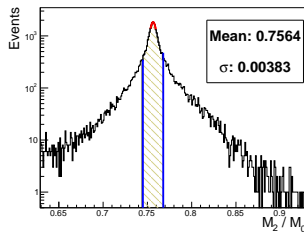
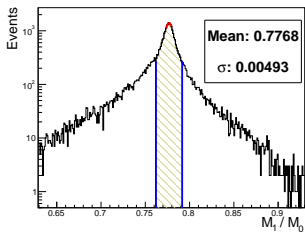
Source	$\gamma$ -line [keV]	FWHM [keV]				
		Core	Seg. 1	Seg. 2	Seg. 3	Seg. 4
$^{133}\text{Ba}$	81	3.29	5.15	4.05	4.07	7.26
	356	2.64	4.58	3.47	3.52	6.04
$^{60}\text{Co}$	1173	4.57	4.88	4.16	4.34	8.05
	1332	4.93	5.00	4.39	4.14	7.84
$^{228}\text{Th}$	2614	7.65				

**Table:** Energy resolutions as absolute FWHM in keV. The resolutions for especially characteristic  $\gamma$ -lines for the respective sources are given.









The remaining matrix elements can also be determined studying single-segment events. For single-segment  $i$  events, the measured pulse amplitudes in the other segments  $j, j \neq i$ , can be written as:

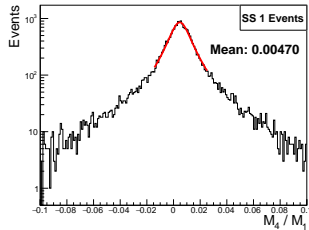
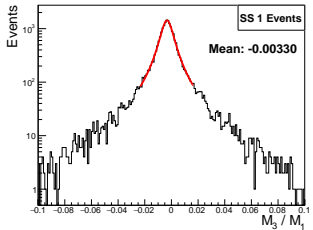
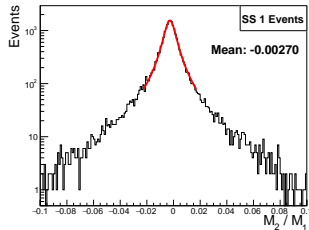
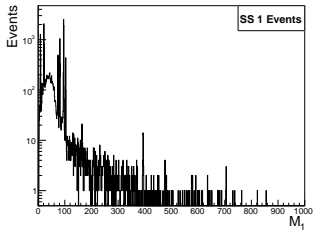
$$M_j = E_i^{true} \cdot c_{ij} \quad , \quad (14)$$

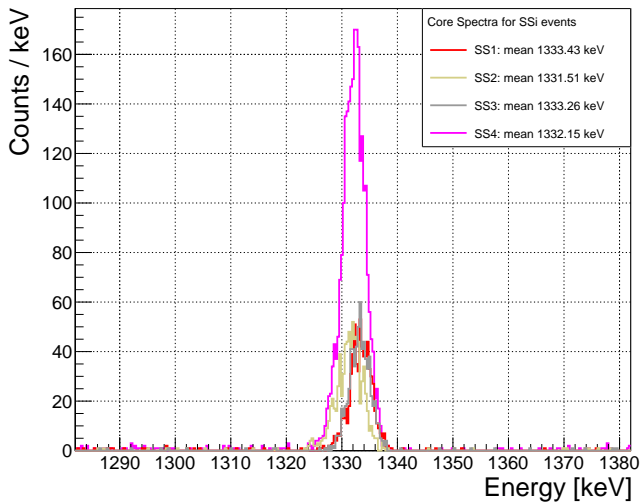
and the ratio between  $M_j$  and  $M_i$ ,  $m_{ji} = M_j/M_i$ , leads to:

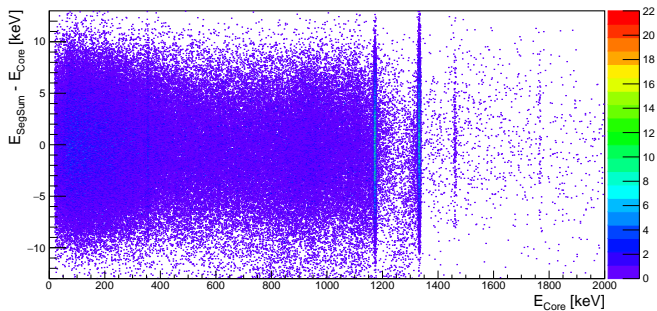
$$m_{ji} = \frac{E_i^{true} \cdot c_{ij}}{E_i^{true} \cdot c_{ii}} = \frac{c_{ij}}{c_{ii}} \quad (15)$$

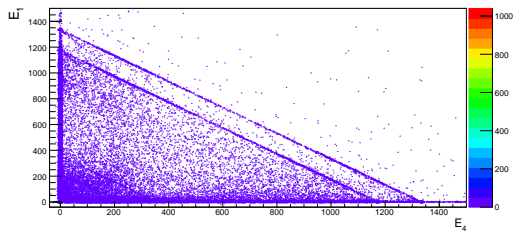
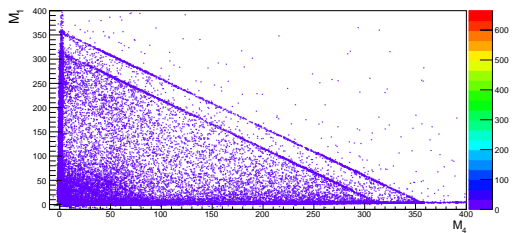
$$\Rightarrow c_{ij} = m_{ji} \cdot c_{ii} \quad . \quad (16)$$











Segment	Boundary(i, j)	Top-scan	Side-scan
1	$\phi_{4,1}$	$124.7 \pm 0.6$	$120.0 \pm 0.8$
	$\phi_{1,4}$	$181.1 \pm 0.6$	$181.1 \pm 0.8$
2	$\phi_{4,2}$	$244.0 \pm 0.6$	$242.5 \pm 0.7$
	$\phi_{2,4}$	$302.5 \pm 0.6$	$302.1 \pm 0.8$
3	$\phi_{4,3}$	$7.1 \pm 0.5$	$3.8 \pm 0.7$
	$\phi_{3,4}$	$64.8 \pm 0.6$	$63.5 \pm 0.7$
4	$\phi_{3,4}$	$66.6 \pm 0.6$	$62.8 \pm 0.7$
	$\phi_{4,1}$	$123.0 \pm 0.6$	$120.2 \pm 0.8$
	$\phi_{1,4}$	$182.5 \pm 0.5$	$180.8 \pm 0.7$
	$\phi_{4,2}$	$242.4 \pm 0.6$	$242.3 \pm 0.6$
	$\phi_{2,4}$	$303.2 \pm 0.8$	$302.0 \pm 0.7$
	$\phi_{4,3}$	$6.2 \pm 0.5$	$3.4 \pm 0.6$



Scan	$\phi_{offset}$	$\phi_{\langle 110 \rangle}$
Top	$-14.52 \pm 0.24$	$37.02 \pm 0.24$
Side	$-15.05 \pm 0.30$	$37.55 \pm 0.30$

**Table:** The values of  $\phi_{offset}$  and  $\phi_{\langle 110 \rangle}$ , see Eq. 11, as determined by fits to the data.





