

L_∞ algebras and conformal field theory

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What?

L_∞ algebras
and conformal
field theory

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Symmetries in
physics

Symmetries in
the Standard
Model

String theory

Symmetries
underlying string
theory

2d conformal
field theory

L_∞ in 2d
CFT

Two dimensional conformal field theory with consistent
extended \mathcal{W} -symmetry \Leftrightarrow Theory has L_∞ symmetry

Outline

L_∞ algebras
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- Why should we care about symmetries at all?

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- Why should we care about symmetries at all?
- ⇒ Symmetries constrain physical theories!

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Unification of Special Relativity and Quantum Mechanics

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Global Poincaré invariance

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Global Poincaré invariance

Fundamental forces
**strong, weak and
electromagnetic**
are **gauge fields**

Gauge invariance

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Unification of Special Relativity and Quantum Mechanics

Global Poincaré invariance

Fundamental forces
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Gauge invariance

→ Action for Standard Model is Poincaré and gauge invariant.

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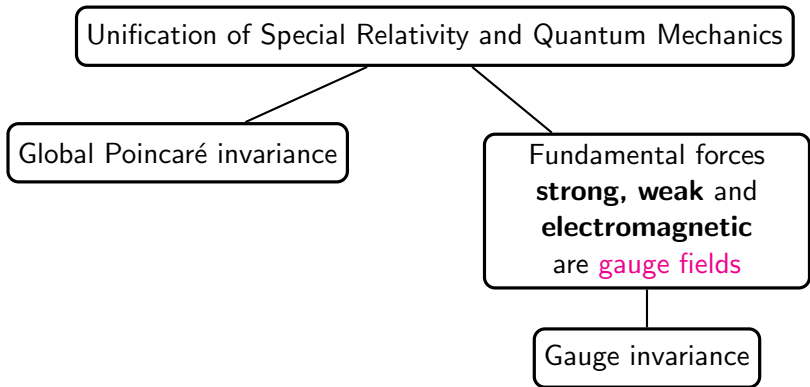
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Symmetries in the Standard Model



→ Action for Standard Model is Poincaré and gauge invariant.

- Standard Model is not enough, as it doesn't include General Relativity.

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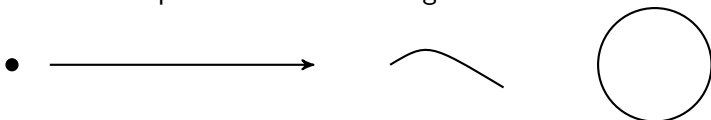
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Idea: Point particle becomes string



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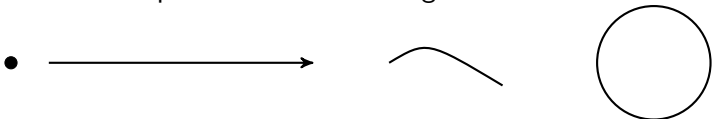
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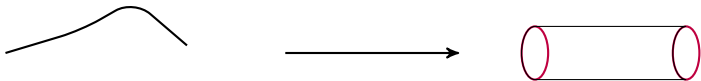
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Idea: Point particle becomes string



Worldline becomes worldsheet



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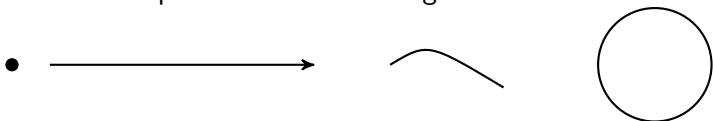
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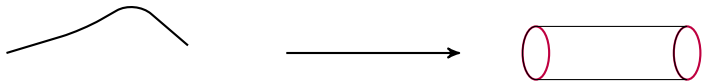
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Idea: Point particle becomes string



Worldline becomes worldsheet



→ Oscillation modes of the string describe different particles

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- What are the symmetries underlying string theory?

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- What are the symmetries underlying string theory?
 - Usually world sheet approach with no spacetime action
- Hard to unravel symmetries of the theory

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- What are the symmetries underlying string theory?
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Idea: Since there is no spacetime action, use scattering diagrams to learn something.

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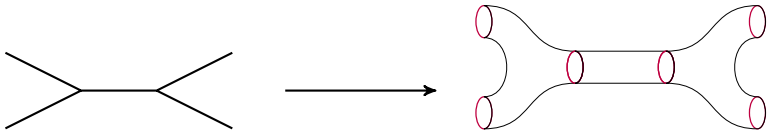
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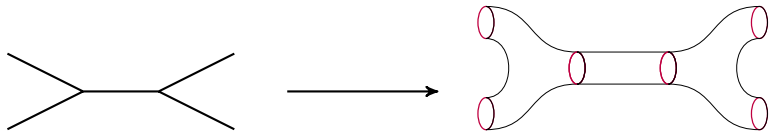
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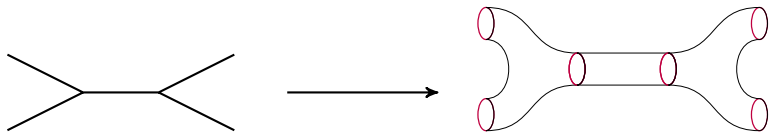
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- Get all possible diagrams from gluing "pair of pants" and cylinder

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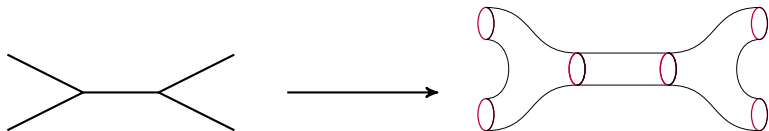
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- Get all possible diagrams from gluing "pair of pants" and cylinder

→ Symmetry in gluing \Rightarrow **Loop L_∞ algebra**

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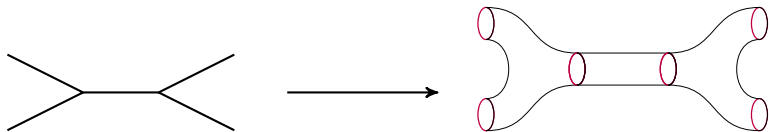
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- Get all possible diagrams from gluing "pair of pants" and cylinder
- Symmetry in gluing \Rightarrow **Loop L_∞ algebra**
- Gluing together only tree diagrams \Rightarrow **L_∞ algebra**

Symmetries underlying string theory

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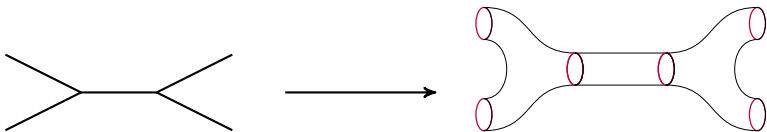
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- Get all possible diagrams from gluing "pair of pants" and cylinder
- Symmetry in gluing \Rightarrow **Loop L_∞ algebra**
- Gluing together only tree diagrams \Rightarrow **L_∞ algebra**
- Use the L_∞ algebra to write down a spacetime action for string theory
- Classical Closed Bosonic String Field Theory action, which has **L_∞ gauge symmetry**

Symmetries underlying string theory

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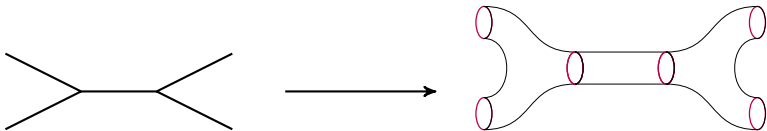
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- Get all possible diagrams from gluing "pair of pants" and cylinder
- Symmetry in gluing \Rightarrow **Loop L_∞ algebra**
- Gluing together only tree diagrams \Rightarrow **L_∞ algebra**
- Use the L_∞ algebra to write down a spacetime action for string theory
- Classical Closed Bosonic String Field Theory action, which has **L_∞ gauge symmetry**
- L_∞ algebra is generalization of Lie algebra

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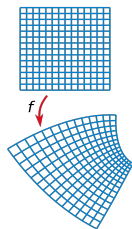
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- Two dimensional field theory which is invariant under all angle preserving maps.

2d conformal field theory

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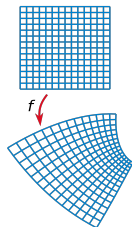
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- Two dimensional field theory which is invariant under all angle preserving maps.



- Conformal transformations are generated by energy-momentum tensor (spin 2)

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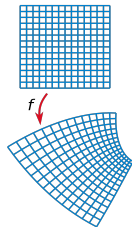
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- Two dimensional field theory which is invariant under all angle preserving maps.



- Conformal transformations are generated by energy-momentum tensor (spin 2)
- Add symmetries generated by a spin 3 field W_3 , a spin 4 field W_4, \dots

2d conformal field theory

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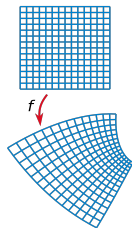
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- Two dimensional field theory which is invariant under all angle preserving maps.



- Conformal transformations are generated by energy-momentum tensor (spin 2)
- Add symmetries generated by a spin 3 field W_3 , a spin 4 field W_4, \dots
- Upon **crossing symmetry** for the transformations one gets 2d CFT with additional W_3, W_4, \dots symmetry

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- L_∞ symmetry in classical Yang-Mills gauge theory, Einstein gravity

(O.Hohm, B.Zwiebach arXiv:1701.08824)

- L_∞ symmetry in classical Yang-Mills gauge theory,
Einstein gravity

(O.Hohm, B.Zwiebach arXiv:1701.08824)

- 1) Constraints for classical field theory with W_N -symmetry
are satisfied \Leftrightarrow W_N -theory has L_∞ symmetry

(R.Blumenhagen, M.Fuchs, MT JHEP07(2017)060)

- L_∞ symmetry in classical Yang-Mills gauge theory, Einstein gravity

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- 1) Constraints for classical field theory with W_N -symmetry are satisfied \Leftrightarrow W_N -theory has L_∞ symmetry

(R.Blumenhagen, M.Fuchs, MT JHEP07(2017)060)

- 2) Constraints for quantum field theory with W_N -symmetry are satisfied \Leftrightarrow W_N -theory has Quantum L_∞ symmetry

(R.Blumenhagen, M.Fuchs, MT JHEP10(2017)163)

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- Relation of Quantum L_∞ to String theory (loop L_∞)? → global vs. local

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- Relation of Quantum L_∞ to String theory (loop L_∞)? \rightarrow global vs. local
- The role of L_∞ in String theory? Describes e.g. R-flux/octonion algebra

(O.Hohm, V.Kupriyanov, D.Lüst, MT arXiv:1709.10004)

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Thank you!

Yang-Mills gauge theory

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- Gauge field A in the adjoint representation

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- Gauge field A in the adjoint representation
- infinitesimal transformation: $\delta_\lambda A = d\lambda + [A, \lambda]$

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- Gauge field A in the adjoint representation
- infinitesimal transformation: $\delta_\lambda A = d\lambda + [A, \lambda]$
- Gauge algebra closes:

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- Gauge field A in the adjoint representation
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$$[\delta_{\lambda_1}, \delta_{\lambda_2}] = \delta_{[\lambda_1, \lambda_2]}$$

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- infinitesimal transformation: $\delta_\lambda A = d\lambda + [A, \lambda]$
- Gauge algebra closes:

$$[\delta_{\lambda_1}, \delta_{\lambda_2}] = \delta_{[\lambda_1, \lambda_2]} \quad \xRightarrow{\text{Jacobi}} \quad \sum_{\text{cycl}} [\delta_{\lambda_1}, [\delta_{\lambda_2}, \delta_{\lambda_3}]] = 0$$

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- Rewrite equations:

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- infinitesimal transformation: $\delta_\lambda A = d\lambda + [A, \lambda]$
- Gauge algebra closes:

$$[\delta_{\lambda_1}, \delta_{\lambda_2}] = \delta_{[\lambda_1, \lambda_2]} \quad \xRightarrow{\text{Jacobi}} \quad \sum_{\text{cycl}} [\delta_{\lambda_1}, [\delta_{\lambda_2}, \delta_{\lambda_3}]] = 0$$

- Rewrite equations:

$$\delta_\lambda A = \ell_1(\lambda) + \ell_2(\lambda, A),$$

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- Gauge field A in the adjoint representation
- infinitesimal transformation: $\delta_\lambda A = d\lambda + [A, \lambda]$
- Gauge algebra closes:

$$[\delta_{\lambda_1}, \delta_{\lambda_2}] = \delta_{[\lambda_1, \lambda_2]} \quad \xRightarrow{\text{Jacobi}} \quad \sum_{\text{cycl}} [\delta_{\lambda_1}, [\delta_{\lambda_2}, \delta_{\lambda_3}]] = 0$$

- Rewrite equations:

$$\delta_\lambda A = \ell_1(\lambda) + \ell_2(\lambda, A), \quad [\delta_{\lambda_1}, \delta_{\lambda_2}] = \delta_{-\ell_2(\lambda_1, \lambda_2)}$$

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- Self-interacting massless spin 3 field with consistent gauge algebra?

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- Self-interacting massless spin 3 field with consistent gauge algebra?

→ Gauge algebra does not close in Lie algebra!

(BBvD '85)

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- Bosonic closed string field theory

(B.Zwiebach '92)

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$$\delta_\Lambda \Phi \sim l_1(\Lambda) + l_2(\Lambda, \Phi) + l_3(\Lambda, \Phi, \Phi) + l_4(\Lambda, \Phi^3) + \dots$$

$$[\delta_{\Lambda_1}, \delta_{\Lambda_2}] \sim \delta_{-l_2(\Lambda_1, \Lambda_2) - l_3(\Lambda_1, \Lambda_2, \Phi) - l_4(\Lambda_1, \Lambda_2, \Phi, \Phi) + \dots}$$

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$$\delta_\Lambda \Phi \sim l_1(\Lambda) + l_2(\Lambda, \Phi) + l_3(\Lambda, \Phi, \Phi) + l_4(\Lambda, \Phi^3) + \dots$$

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⇒ L_∞ algebra

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Definition

A L_∞ algebra consists of a graded vector space X and multilinear maps $\{\ell_n\}_{n \geq 1}$, $\ell_n : \underbrace{X \otimes \cdots \otimes X}_n \rightarrow X$, satisfying generalized Jacobi identities.

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- Classical: $\hbar \rightarrow 0$ in quantum CFT

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- Classical: $\hbar \rightarrow 0$ in quantum CFT
- Virasoro algebra: Energy momentum tensor T , spin 2

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- Classical: $\hbar \rightarrow 0$ in quantum CFT
- Virasoro algebra: Energy momentum tensor T , spin 2
- \mathcal{W}_N algebra: Generators $\{T, W_3, \dots, W_N\}$, W_i spin i field

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- Classical: $\hbar \rightarrow 0$ in quantum CFT
- Virasoro algebra: Energy momentum tensor T , spin 2
- \mathcal{W}_N algebra: Generators $\{T, W_3, \dots, W_N\}$, W_i spin i field

Theorem (classical)

$\{T, W_3, \dots, W_N\}$ form a classical \mathcal{W}_N algebra iff their symmetry transformations form a L_∞ algebra.

(R.Blumenhagen, M.Fuchs, MT JHEP07(2017)060)

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Classical \mathcal{W}_3 symmetry transformations:

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Classical \mathcal{W}_3 symmetry transformations:

$$\delta_\eta W = \frac{c}{360} \partial^5 \eta + \alpha \left(\frac{1}{6} \partial^3 \eta T + \frac{1}{4} \partial^2 \eta \partial T + \frac{3}{20} \partial \eta \partial^2 T + \frac{1}{30} \eta \partial^3 T \right) + \beta \left(\partial \eta (TT) + \frac{1}{2} \eta \partial (TT) \right)$$

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Classical \mathcal{W}_3 symmetry transformations:

$$\begin{aligned}\delta_\eta W &= \frac{c}{360} \partial^5 \eta + \alpha \left(\frac{1}{6} \partial^3 \eta T + \frac{1}{4} \partial^2 \eta \partial T + \frac{3}{20} \partial \eta \partial^2 T \right. \\ &\quad \left. + \frac{1}{30} \eta \partial^3 T \right) + \beta \left(\partial \eta (TT) + \frac{1}{2} \eta \partial (TT) \right) \\ &= \ell_1^W(\eta) + \ell_2^W(\eta, T) - \frac{1}{2} \ell_3^W(\eta, T, T)\end{aligned}$$

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Classical \mathcal{W}_3 symmetry transformations:

$$\begin{aligned}\delta_\eta W &= \frac{c}{360} \partial^5 \eta + \alpha \left(\frac{1}{6} \partial^3 \eta T + \frac{1}{4} \partial^2 \eta \partial T + \frac{3}{20} \partial \eta \partial^2 T \right. \\ &\quad \left. + \frac{1}{30} \eta \partial^3 T \right) + \beta \left(\partial \eta (TT) + \frac{1}{2} \eta \partial (TT) \right) \\ &= \ell_1^W(\eta) + \ell_2^W(\eta, T) - \frac{1}{2} \ell_3^W(\eta, T, T)\end{aligned}$$

$$\delta_\varepsilon T = \ell_1^T(\varepsilon) + \ell_2^T(\varepsilon, T)$$

$$\delta_\eta T = \ell_2^T(\eta, W)$$

$$\delta_\varepsilon W = \ell_2^W(\varepsilon, W)$$

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→ Demand closure

$$[\delta_{\epsilon_1}, \delta_{\epsilon_2}] = \delta_{-\ell_2(\epsilon_1, \epsilon_2) - \ell_3(\epsilon_1, \epsilon_2, \mathbf{W}) + \frac{1}{2}\ell_4(\epsilon_1, \epsilon_2, \mathbf{W}, \mathbf{W}) + \dots}$$

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$$[\delta_{\epsilon_1}, \delta_{\epsilon_2}] = \delta_{-\ell_2(\epsilon_1, \epsilon_2) - \ell_3(\epsilon_1, \epsilon_2, \mathbf{W}) + \frac{1}{2}\ell_4(\epsilon_1, \epsilon_2, \mathbf{W}, \mathbf{W}) + \dots}$$

■ E.g. $[\delta_{\eta_1}, \delta_{\eta_2}] = \delta_{-\ell_2(\eta_1, \eta_2) - \ell_3(\eta_1, \eta_2, T)}$

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→ Demand closure

$$[\delta_{\epsilon_1}, \delta_{\epsilon_2}] = \delta_{-\ell_2(\epsilon_1, \epsilon_2) - \ell_3(\epsilon_1, \epsilon_2, \mathbf{W}) + \frac{1}{2}\ell_4(\epsilon_1, \epsilon_2, \mathbf{W}, \mathbf{W}) + \dots}$$

- E.g. $[\delta_{\eta_1}, \delta_{\eta_2}] = \delta_{-\ell_2(\eta_1, \eta_2) - \ell_3(\eta_1, \eta_2, T)}$

- L_∞ relations fix $\alpha = 2$, $\beta = \frac{32}{5c}$

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■ Generators $\{T, W\}$

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- Generators $\{T, W\}$
- Symmetry transformations get correction terms proportional to \hbar

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- Generators $\{T, W\}$
- Symmetry transformations get correction terms proportional to \hbar
- Specify a product between operators \rightarrow Normal ordered product in 2D CFT

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CFT

- Generators $\{T, W\}$
 - Symmetry transformations get correction terms proportional to \hbar
 - Specify a product between operators \rightarrow Normal ordered product in 2D CFT
- \Rightarrow Cross relations among the fundamental identities of the L_∞ algebra

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- Generators $\{T, W\}$
 - Symmetry transformations get correction terms proportional to \hbar
 - Specify a product between operators \rightarrow Normal ordered product in 2D CFT
- \Rightarrow Cross relations among the fundamental identities of the L_∞ algebra
- \Rightarrow Quantum L_∞ algebra (R.Blumenhagen, M.Fuchs, MT JHEP10(2017)163)