Flavored Axions

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Recent theoretical motivation

Flaxion: a minimal extension to solve puzzles in the standard model Yohei Ema (Tokyo U.), Koichi Hamaguchi, Takeo Moroi, Kazunori Nakayama (Tokyo U. & Tokyo U., IPMU). Dec 16, 2016. 23 pp. Published in JHEP 1701 (2017) 096 UT-16-36, IPMU16-0189 DOI: 10.1007/JHEP01(2017)096 e-Print: arXiv:1612.05492 [hep-ph] I PDF

Minimal axion model from flavor Lorenzo Calibbi (Beijing, Inst. Theor. Phys.), Florian Goertz (CERN), Diego Redigolo (Tel Aviv U. & Weizmann Inst.), Robert Ziegler (Karlsruhe U., TTP), Jure Zupan (CERN & Cincinnati U.). Dec 23, 2016. 6 pp. Published in Phys.Rev. D95 (2017) no.9, 095009 TTP16-058, CERN-TH-2016-261 DOI: 10.1103/PhysRevD.95.095009 e-Print: arXiv:1612.08040 [hep-ph] I PDF

The Minimal Flavour Violating Axion F. Arias-Aragon, L. Merlo (Madrid, Autonoma U. & Madrid, IFT). Sep 20, 2017. 11 pp. Published in JHEP 1710 (2017) 168 DOI: <u>10.1007/JHEP10(2017)168</u> e-Print: <u>arXiv:1709.07039</u> [hep-ph] I PDF

Supersymmetric Flaxion Yohei Ema, Daisuke Hagihara (Tokyo U.), Koichi Hamaguchi, Takeo Moroi (Tokyo U., IPMU & Tokyo U.), Kazunori Nakayama (Tokyo U., IPMU & Department of Physics, The University of Tokyo, Tokyo 113-0033, Japan & Tokyo U.). Feb 21, 2018. 31 pp. Published in JHEP 1804 (2018) 094 UT-18-1, IPMU18-0036 DOI: 10.1007/JHEP04(2018)094 e-Print: arXiv:1802.07739 [hep-ph] I PDF

<u>A Common Source for Scalars: Axiflavon-Higgs Unification</u> <u>Tommi Alanne, Simone Blasi, Florian Goertz (Heidelberg, Max Planck Inst.</u>). Jul 26, 2018. 8 pp. e-Print: <u>arXiv:1807.10156</u> [hep-ph] I <u>PDF</u>

Phenomenological motivation

Astrophobic Axions

Luca Di Luzio (Durham U., IPPP), Federico Mescia (ICC, Barcelona U.), Enrico Nardi (Frascati), Paolo Panci, Robert Ziegler (CERN). Dec 13, 2017. 6 pp. Published in Phys.Rev.Lett. 120 (2018) no.26, 261803 IPPP-17-102, CERN-TH-2017-256 DOI: <u>10.1103/PhysRevLett.120.261803</u> e-Print: <u>arXiv:1712.04940</u> [hep-ph] I PDF

Note that these are the invisible version of the VARIANT axion models discussed to avoid some experimental constraints on the original PQ axion... is a fun story here to tell ;-)

The photo-philic QCD axion Marco Farina (Rutgers U., Piscataway), Duccio Pappadopulo (New York U., CCPP & New York U.), Fabrizio Rompineve (New York U., CCPP & U. Heidelberg, ITP & New York U.), Andrea Tesi (Chicago U., EFI & Chicago U.). Nov 29, 2016. 21 pp. Published in JHEP 1701 (2017) 095 DOI: 10.1007/JHEP01(2017)095 e-Print: arXiv:1611.09855 [hep-ph] | PDF

Experimental Targets for Photon Couplings of the QCD Axion Prateek Agrawal (Harvard U., Phys. Dept.), JiJi Fan (Brown U.), Matthew Reece (Harvard U., Phys. Dept.), Lian-Tao Wang (Chicago U., EFI & Chicago U., KICP). Sep 18, 2017. 23 pp. Published in JHEP 1802 (2018) 006 DOI: 10.1007/JHEP02(2018)006 e-Print: arXiv:1709.06085 [hep-ph] I PDF

B. Doebrich from NA62 gave a recent talk where she shows some slides and axion-flavon sensitivity

https://indico.in2p3.fr/event/17826/attachments/49465/62902/marseille_2018.pdf

The Axion

- Motivation: Strong CP Problem
- Peccei-Quinn (PQ) mechanism [Peccei, Quinn, '77]
- pseudo-Nambu Goldstone boson [Weinberg '78, Wilczek '78]
- axion mass: $m_a = 5.70(6)(4) \,\mu \text{eV} \left(\frac{10^{12} \,\text{GeV}}{f_a}\right)$
- Lagrangia $a \rightarrow a + \text{const.}$

$$\mathcal{L} \supset \frac{1}{2} \partial_{\mu} a \,\partial^{\mu} a - \frac{\alpha_s}{8\pi} \left(\bar{\theta} + \frac{a}{f_a} \right) G^b_{\mu\nu} \tilde{G}^{b,\mu\nu} - \frac{\alpha}{8\pi} C_{a\gamma} \frac{a}{f_a} F_{\mu\nu} \tilde{F}^{\mu\nu} + \dots$$

solves strong CPaperoblem(x) + $\overline{\theta} f_a$ axion searches $\langle \overline{a} \rangle = 0$ allows for axion dark matter



Axion Models



Hadronic (KSVZ) Axion Model



KSVZ [Kim '79; Shifman, Vainshtein, Zakharov '80]

Axion Models



Axion-Photon Interaction: Primakoff



\rightarrow "see" the "invisible" axion

Axion-Photon Interaction: Effective Lagrangian

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - J^{\mu}A_{\mu} + \frac{1}{2}\partial_{\mu}a\partial^{\mu}a - \frac{1}{2}m_{a}^{2}a^{2} - \frac{g_{a\gamma}}{4}F_{\mu\nu}\widetilde{F}^{\mu\nu}a$$

- EM field-strength tensor $\, F_{\mu
 u} = \partial_\mu A_
 u \partial_
 u A_\mu$
- Dual EM field-strength tensor $\widetilde{F}^{\mu\nu} = \frac{1}{2} \varepsilon^{\mu\nu\alpha\beta} F_{\alpha\beta}$
- Vector Potential $A^{\mu} = (A_0, \mathbf{A})$

• 4-Current
$$J^{\mu}=(
ho,{f J})$$

• Electric Field

$$\mathbf{E} = -\boldsymbol{\nabla}A_0 - \dot{\mathbf{A}}$$

- Magnetic Field
- $\mathbf{B} = \mathbf{\nabla} \times \mathbf{A}$

$$g_{a\gamma} = -\frac{\alpha}{2\pi f_a} C_{a\gamma} = -2.04(3) \times 10^{-16} \text{ GeV}^{-1} \left(\frac{m_a}{1\,\mu\text{eV}}\right) C_{a\gamma}$$

$$C_{a\gamma} = \frac{\mathcal{E}}{\mathcal{N}} - 1.92(4)$$

MADMAX Goal - probe QCD axion DM scenarios

 $\nu_a \, [\text{GHz}]$



accidental inital misalignment $m_{A} < 100 \ \mu eV$

Scenario 2

 $m_A = 26.2 \pm 3.4 \,\mu \text{eV}$

[Klaer, Moore,'17]

 $m_a = 5.70(6)(4) \,\mu \text{eV} \left(\frac{10^{12} \,\text{GeV}}{f_a}\right)$

10⁻¹⁰ CAST helioscope PDG 10⁻¹¹ ORGAN Axion Coupling IG $_{A_{\gamma\gamma}}$ I (GeV $^{-1})$ Ringwald, ct. 10⁻¹² MADMAX HAYSTAC dielectric haloscope 10⁻¹³ **RBF** Andreas UF 10⁻¹⁴ KSVZ QCD Axion Coupling **ADMX** 10⁻¹⁵ DFSZ cavity haloscopes 10⁻¹⁶ 10⁻⁶ 10⁻⁴ 10⁻⁵ Axion Mass m_A (eV) N = 1Scenario 2 DESY $m_A = 26.2 \pm 3.4 \,\mu \text{eV}$ [Klaer, Moore,'17]

Existing exclusion limits today

Constraints on the Peccei-Quinn (PQ) scale f_{PQ}



Bounds from Axion Searches Cosmological Axion Bounds Astrophysical Axion Bounds

 $\begin{cases} \text{Peccei-Quinn Scale} \\ f_a \gtrsim 6 \times 10^8 \,\text{GeV} \end{cases}$

 $\begin{array}{l} \left(\begin{array}{c} \text{Axion Mass} \\ m_a \end{array} \simeq 0.6 \end{array} \right. \text{meV} \left(10^{10} \, \text{GeV} / f_{\rm PQ} \right) \end{array} \end{array}$

Astrophysical Bounds - Coupling Dependent

• axion-photon coupling - evolution of Horizontal Branch stars



• axion-electron coupling - white dwarf cooling



• axion-nucleon coupling - supernova neutrino burst duration $C_N < \frac{0.004 \,\mathrm{eV}}{0.004 \,\mathrm{eV}}$ $C_N < \frac{0.004 \,\mathrm{eV}}{m_a}$



♦ Precision Flavor Experiments: An intere Astrophobic Axionse of the actrophotic akions that they n Sarily have flavor-violating couplings to quarks leptons. In the lepton sector the strongest its on the axion mass arise from $\mu - e$ transit $0.2\,\mathrm{eV}$ suppression of axion-related onger b r model energy transfer away from lab. I). S astrophysical objects LH rota $(\epsilon_{11}^{e_L} = 1/(1+z))$, in order to have this const Combined SN/WD body with erefore $\epsilon_{22}^{22} \gtrsim 2.5 \cdot 10^{-4}$. This implies that the corresp ing (large) PMNS mixing angle must come the neutrino sector. Moreover, since $\sum_{i} \epsilon_{ii}^{e_L}$ $\epsilon_{33}^{e_L} = z/(1+z)$ is fixed and we product

$$\beta_{\tau \to ea} = 6.7 \cdot 10^{-6} \left(\frac{m_a}{0.2 \,\mathrm{eV}}\right)^2 \,,$$

Note: outside of the natural e abov axion dark matter region

$$10^{-7} \text{ eV} \lesssim m_a \lesssim 10^{-4} \text{ eV} \mathcal{B}_{K^+ \to \pi^+ a} = 1.0 \cdot 10^{-9} \left(\frac{m_a}{0.2 \text{ eV}}\right)^2 \left(\frac{\epsilon_{11}^d \epsilon_{22}^d}{\lambda^{10}}\right),$$

where $\epsilon^d = \epsilon^{d_R}$ for M4 and $\epsilon^d = \epsilon^{d_L}$ for M1 while M3 is automatically safe (cf. Tab. I). Coming to the present bound $\mathcal{B}_{K^+ \to \pi^+ a} < 7.3 \cdot 10^{-11}$ this implies the constraint $\epsilon^{d_R}_{22} \lesssim 10^{-9}$ in rs

Non-Universal DFSZ Axion Models

is an efficient suppression by UKINI angles, so the bounds from flavor physics become comp



e sense that Yukawa hierarchies are $\alpha_{ud} =$ of the Yukaw by • UM(nlm) al DESTA REPORTS a BRO in the generations given by the 1 freedom from $\mathcal{O}(1)$ uncertainties, we simlistributered-604(inhal OESZA finitions pronchanges for the generations therefore the same pre idom sign, resulting in a range allows for nucleophobic axion models relaxed SN bound flavon contair allows for electrophobic axion models - relaxed WD bound $\frac{1}{N} \in \begin{bmatrix} 2.4, 3.0 \end{bmatrix}$, (15) • PQ = Froggatt-Nielsenr(FN)_{Qi} H ϕ [0.5, 1.1]. to be compare q_{avon} th the usual $\operatorname{ndow} \left| \frac{\overline{j}}{l} \, \widehat{y}_{ij}^u \, \overline{Q_i} \, H u_{Rj} \right| \longrightarrow y_{ij}^u \left(\frac{\phi}{M} \right)^{q_{Q_i} - q_{u_j}} \overline{Q_i} H u_{Rj}$ $\det nange a_{1}s$ on the suppression of the suppression of the explains hierarchical flavor structure of quarks and leptons Being a QC Eq. (14) over store dependent in ator is domi-[2.4, 3.0] with a cosmological $m_u \operatorname{det} m_{ijd}^{q} \neq w_{ij}^{q} \approx^{q_{Q_i}} 4^{q_{4j}} \langle \mathcal{W} \rangle ile$ Example: $q_Q = \{3, 2, 0\}, q_u = \{-5, -1, 0\}$ or is $\log \det \underline{m}_{e} / \det m_{e} \approx -\frac{1}{\epsilon^{\frac{3}{3}}} / t_{e} \approx -\frac{1}{\epsilon^{\frac{3}{3}}} / t_{$ corresponding $\rightarrow m^u \sim \begin{pmatrix} \epsilon^8 & \epsilon^4 & \epsilon^3 \\ \epsilon^7 & \epsilon^3 & \epsilon^2 \\ \epsilon^5 & \epsilon^4 & 1 \end{pmatrix}$ flation the or

tions, one find

arXiv:1612.05492 [pdf, ps, other]

Flaxion: a minimal extension to solve puzzles in the standard model

Yohei Ema, Koichi Hamaguchi, Takeo Moroi, Kazunori Nakayama Comments: 23 pages, 1 figure; v2: version published in JHEP Subjects: High Energy Physics - Phenomenology (hep-ph)

arXiv:1612.08040 [pdf, other]

The <u>Axiflavon</u>

Lorenzo Calibbi, Florian Goertz, Diego Redigolo, Robert Ziegler, Jure Zupan Comments: 6 pages, typos corrected, references added Subjects: High Energy Physics - Phenomenology (hep-ph)



[from Koichi Hamaguchi's talk @ PACIFIC 2018]

Puzzles of the SM explained by a simple extension

Flaxion Dark Matter

Case 1: U(1) is broken after inflation.

-> Domain Wall ! In the flaxion scenario, typically $N_{DW} \neq 1$, and this possibility is **excluded**.



Case 2: U(1) was already broken during inflation.

Quantum fluctuation during inflation leads to **DM isocurvature** perturbation, which is severely constrained [Planck,'15].

-> Strong bound on inflation scale.

$$H_{\rm inf} \lesssim 3 \times 10^7 \,{
m GeV} \, \theta_i^{-1} \left(\frac{10^{12} \,{
m GeV}}{f_a}\right)^{0.19}$$



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[from Koichi Hamaguchi's talk @ PACIFIC 2018]



[from Koichi Hamaguchi's talk @ PACIFIC 2018]

Lagrangian

$$\mathcal{L} = -\frac{|\partial\phi|^2}{(1-|\phi|^2/\Lambda^2)^2} - \lambda_{\phi} \left(|\phi|^2 - v_{\phi}^2\right)^2 + y_{ij}^d \left(\frac{\phi}{M}\right)^{n_{ij}^d} \overline{Q}_i H d_{Rj} + y_{ij}^u \left(\frac{\phi}{M}\right)^{n_{ij}^u} \overline{Q}_i \widetilde{H} u_{Rj} + y_{ij}^l \left(\frac{\phi}{M}\right)^{n_{ij}^l} \overline{L}_i H l_{Rj} + y_{i\alpha}^{\nu} \left(\frac{\phi}{M}\right)^{n_{i\alpha}^{\nu}} \overline{L}_i \widetilde{H} N_{R\alpha} + \frac{1}{2} y_{\alpha\beta}^N \left(\frac{\phi}{M}\right)^{n_{\alpha\beta}^N} M \overline{N_{R\alpha}^c} N_{R\beta} + \text{h.c.}$$



Conclusions

- flavored axion models non-universal axion-fermion couplings
- astrophysical constraints can be relaxed new (mini)IAXO opportunities
- but: there a non-axion dark matter explanation will be required
- flavor precision experiments provide stringent tests of such models
- flaxion/axiflavon models allow for solutions of many SM puzzles
- but: hierarchy problem may require SUSY in a somewhat intrigued way
- CERN NA62: K -> pi a will test such models soon
- MADMAX may test such models independently from NA62

Fundamental Constituents of Matter



Axions → Extremely Weakly Interacting Particles (EWIPs)