

## Principles

(1)

$$L_{A\gamma\gamma} = -g_A \frac{\alpha}{4\pi} \frac{A(x)}{f_A} \vec{E} \cdot \vec{B}$$

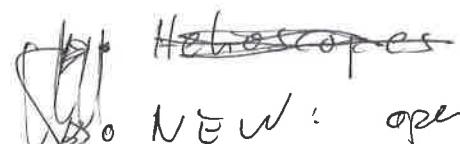
axion  $\rightarrow$  photon mixing in B-field  
 $\Rightarrow$  conversion of axion to photon  
 in magnetic field (or vice versa)

## Poincaré effect:



- Microwave cavity experiment

DM



- Heteroscopes
- New: open ~~resonator~~ quasi-resonator approach
- Photon regeneration: light shining through the wall
- Heteroscopes
- Polarization effect

• Dark Matter axions  
Resonant conversion (cavities) (2)  
Axions at rest?

The axion acts like a classical field  
leading to oscillations of  $E$ .  
⇒ modifies Maxwell equations!  
For the interesting mass/frequency range (DM):

$$16\text{ Hz} \approx 100\text{ GHz}$$

with  $\omega \approx 10^{-3}\text{ c}$ : quasi non-propagating.

extension of experiment  $l \approx 1\text{ m}$ ,  $1\text{ GHz}$

$$\Rightarrow \text{time} \approx \frac{sl}{\omega} \approx \frac{1\text{ m}}{10^{-3}\text{ c}} = 10^4 \frac{\text{m}}{\text{c}} \approx \frac{10^3}{3 \cdot 10^8} \approx 3 \cdot 10^{-6}$$

for axion to pass experiment

Period of axion induced photon  $\frac{1}{100\text{ GHz}} - \frac{1}{6\text{ GHz}} = 10^{11} - 10^{13}\text{ s}$

→ While axion is passing through experiment  
 $\approx 10^3$  to  $10^5$  phases, respectively

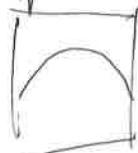
→ Photon part of wave function does resonantly  
coherently couple to experimental infrastructure

→ Synchronization in resonance: constructive

Superposition of  $\approx 10^3$  periods of

photon part of wave function with  $E$  field part

→ axion field "pumps" cavity if in resonance



(3)

Resonance increases expected output power  
(pumping of cavity):

$$P_{\text{sig}} = g_{A\bar{\nu}\nu}^2 \left( \frac{s_a}{m_a} \right) B^2 \cdot V \cdot Q_{L,A} C$$

$s_a$  = DM density

$m_a$  = Axion mass

$V$  = Volume of cavity

$Q_{L,A}$  = loaded quality factor of cavity  
or of axion signal, whichever is lower

$C$  = overlap integral of external  $B$  field  
with oscillating  $E$ -field of the mode

$$= \frac{\int_V d^3x \vec{E}_w \cdot \vec{B}_0 |^2}{\vec{B}_0^2 V \int_V d^3x \epsilon_0 |E_w|^2}$$

dielectric const.

$$\Rightarrow P_{\text{sig}} \approx 0.5 \cdot 10^{-26} W \left( \frac{V}{500 \mu} \right) \left( \frac{B_0}{7 T} \right)^2 C \left( \frac{g_{A\bar{\nu}\nu}}{0.36} \right)^2 \left( \frac{s_a}{0.5 \cdot 10^{-24} \text{ g/cm}^3} \right)$$

$$\left( \frac{m_a}{2\pi/\text{GHz}} \right) \cdot Q_{L,A}$$

Can calculate sensitivity of experiment  
using Dickey's radiometer formula:

(how long does it take to detect signal  $P_{\text{sig}}$  with  $\frac{s}{N}$ )

$$\frac{S}{N} = \frac{P_{sig}}{k_B T_{sys}} \sqrt{\frac{t_{scan}}{2V}}$$

with  $\Delta\sigma \approx 10^{-6} \cdot V_a$

with  $\frac{S}{N} = 5$  (standard assumptions)

$T_{sys} = 100\text{K}$  (LiHe cooled cryostat + noise temp of device)

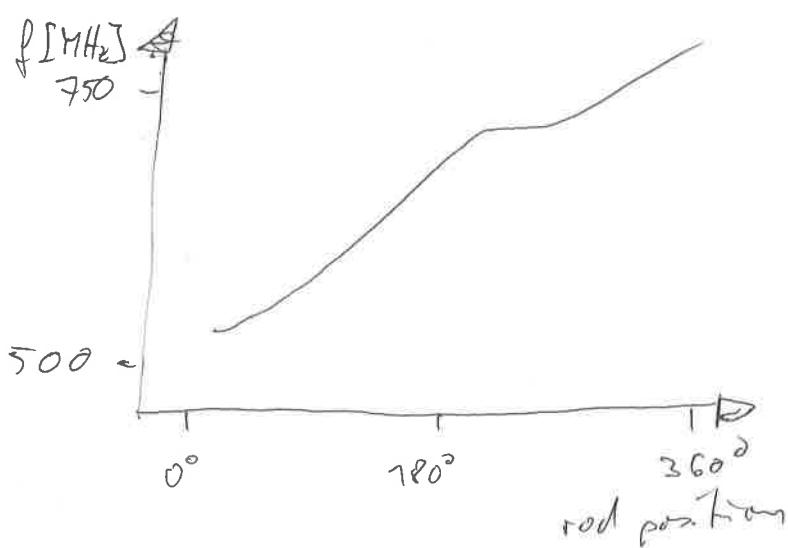
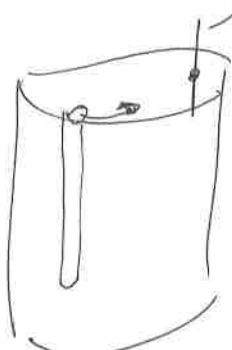
$t_{scan} = 100\text{s}$  (High  $Q_L \Rightarrow$  once; scan small range  
 $\Rightarrow$  many measurements, ie  $t_{scan} \leq 100\text{s}$ )

$$\Rightarrow P_{sens} = 5 \cdot k_B T_{sys} \sqrt{\frac{2V}{t_{scan}}} \approx 10^{-23} \text{W}$$

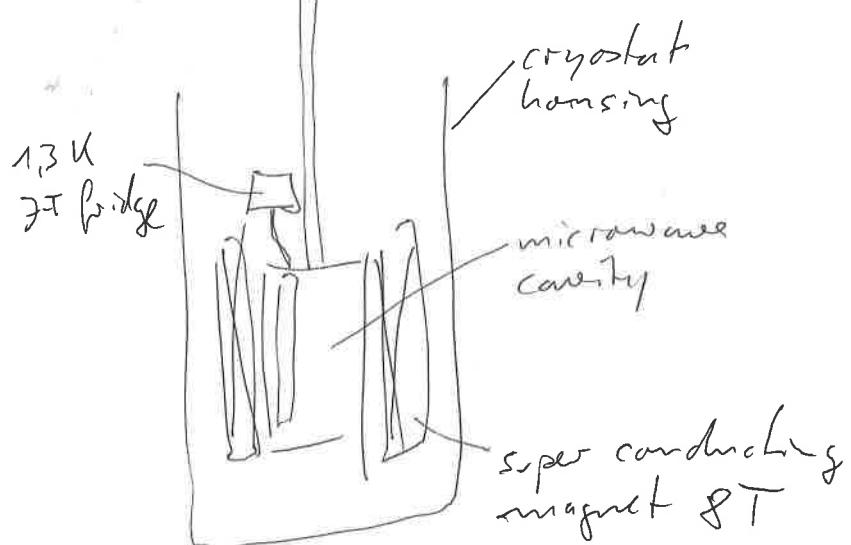
## ADMX

dielectric

Use resonant cavity with tunable rods  
 antenna to couple to resonant frequency  
 by turning rod within cavity;  
 change resonance frequency of  
 cavity:



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$$T_{\text{phys}} \approx 1.5 \text{ K}$$

$$\Rightarrow T_{\text{sys}} \approx 3 \text{ K}$$

$$Q_c \approx 10^5$$

$$V \approx 100 \text{ l}$$

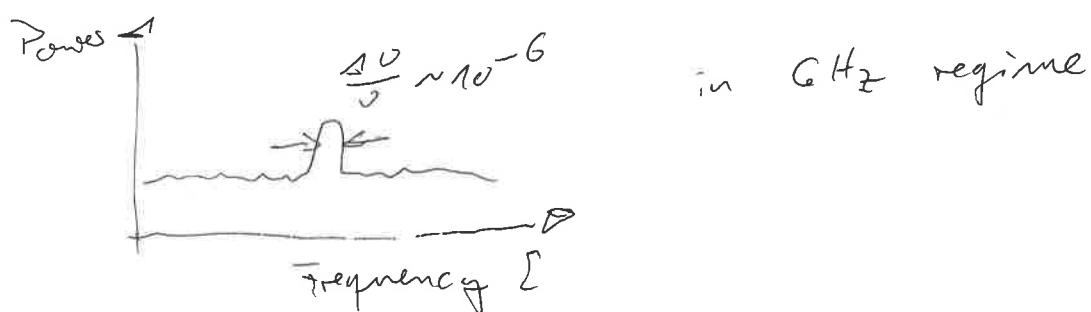
$$B_0 = 8 \text{ T}$$

$$C \approx 0.5$$

$\Rightarrow$  expected power for DM axion signal:

$$P_{\text{sig, DM}} \approx 10^{-23} \text{ Watt}$$

Look for (gaussian) peak over background



$\Rightarrow$  Sensitivity for DM axions in ~6 GHz to 10 GHz

( $\sim 1 \mu\text{eV}$  to  $20 \mu\text{eV}$  Axion mass)

using SQUID amplification and heterodyne mixing

Higher frequencies: sensitivity getting worse due to

- smaller volumes (higher frequency, large wavelength)
- $Q_c$  decreasing for smaller volumes

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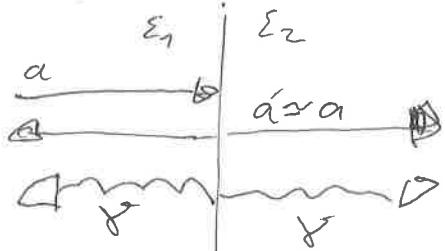
## Madmax

During passage of axions through B. field:  
axion-photon mixing (Primakoff effect)

⇒ photon part of wave func. "feels" transitions b/w different media, i.e. difference in refractive index, dielectric constant

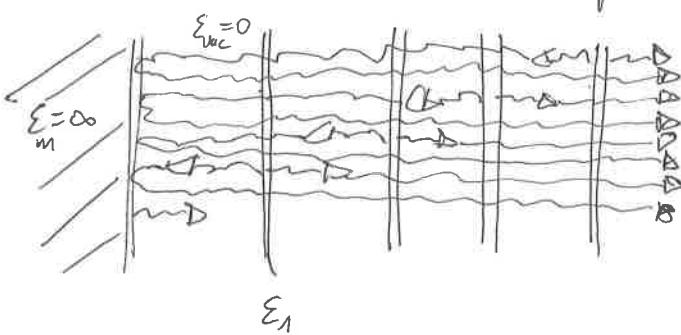
⇒ Conservation of momentum:

Emission of two photons perpendicular to surfaces in both directions



$$\left(\frac{P}{A}\right)_{\text{mirror}} = 2 \cdot 10^{-27} \frac{V}{m^2} \left(\frac{B_{11}}{10T}\right)^2 (g_{ggg} m_a)^2$$

Due to coherence (same as for microwave cavities):  
constructive interference possible of γs emitted at different surfaces



⇒ 2 effects:

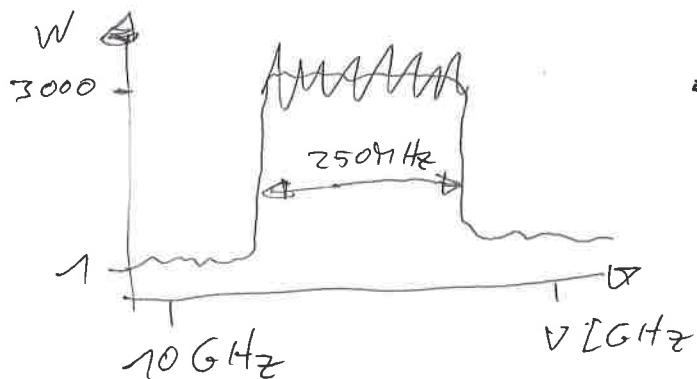
1. Decay and constructive interference
2. Additional radiation by extra surfaces

⇒ Power boost  $\propto W = \frac{P_{\text{cavity}}}{P_{\text{mirror}}} \times \text{Number of discs} \cdot \epsilon_1$

⇒ To get reasonable power for detection ( $\approx 10^{-23} W$ ):  
Need  $\epsilon \approx 25$   $N \approx 80$   $T = 10T$   $A = 1m^2$

Simulations (EM - 1D):

$$T_{05} \quad \epsilon \approx 25 \quad N=20 \quad \Delta V = 250 \text{ MHz}$$

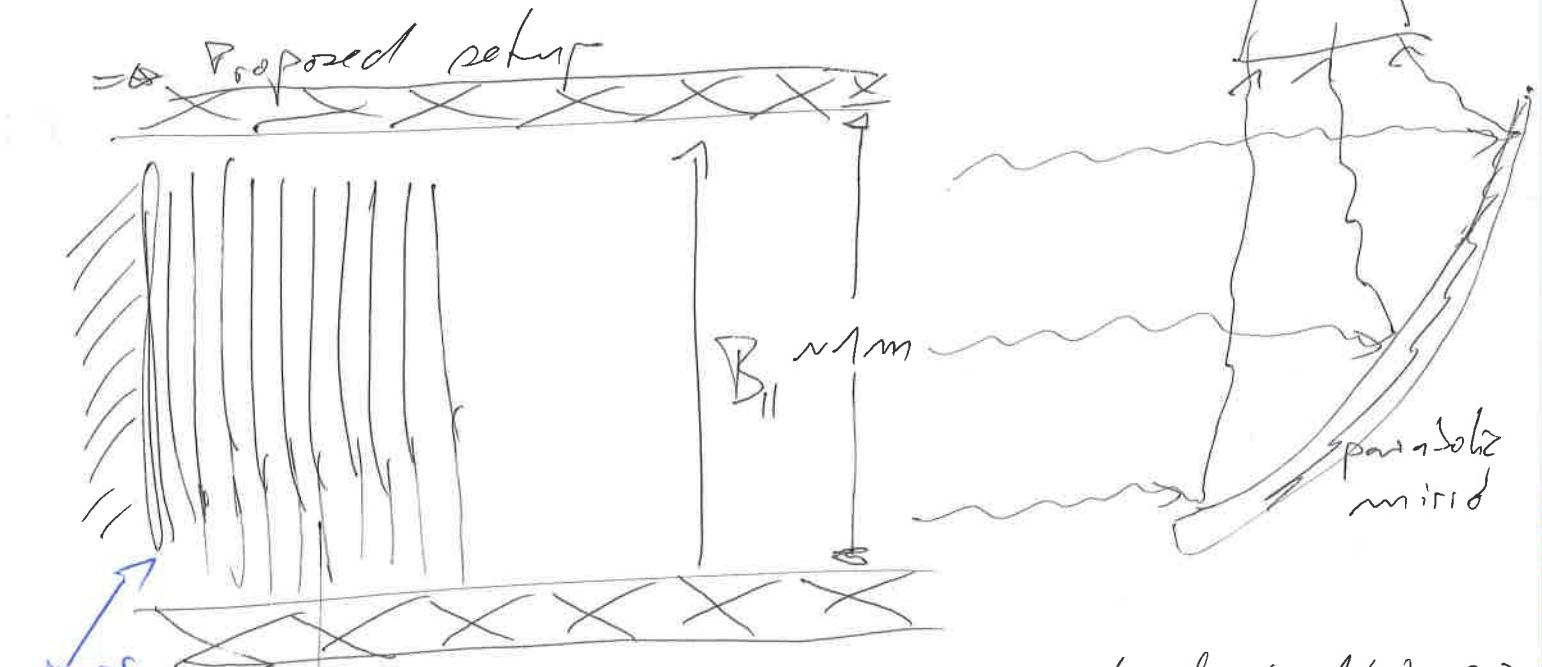


$\Rightarrow$  For  $P_{sig} \approx 10^{-23} \text{ W}$   
need  $\omega \approx 10^4$

take simulation above

Heterodyne mixing  
Receiver

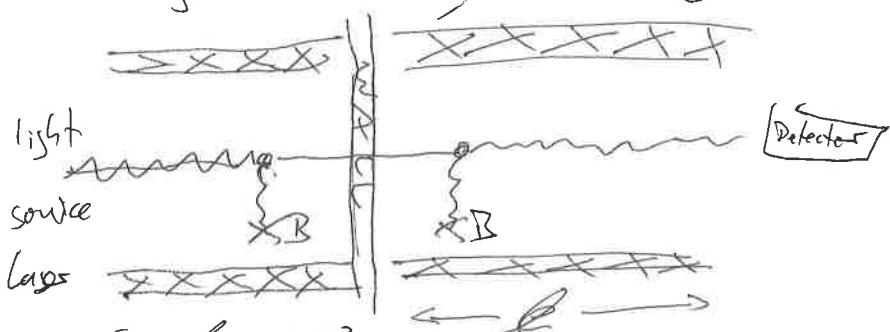
$$\Rightarrow N=80 \quad B_{||}=10 \text{ T} \quad \epsilon=25 \quad A=1 \text{ m}^2$$



80 dielectric plates  $1 \text{ m}^2$  made of  $\text{LaAlO}_3 \epsilon=25$   
positioning tunable  $\Rightarrow$  scan different  
frequency ranges  $\Rightarrow P_{sig} \approx 10^{-23} \text{ W}$  ✓

~~DA axes~~  
 $\Rightarrow$  Scan range  $10 \text{ GHz} - 100 \text{ GHz}$  ( $\sim 40 \mu\text{eV} - 400 \mu\text{eV}$ )

Photon regeneration in lab:  
light shining through the wall:



axions with momentum

$$P_{\gamma \rightarrow a} = \frac{1}{4} (g_{\text{gyr}} \beta l)^2 + (ql)$$

$$\rightarrow \frac{1}{4} (g_{\text{gyr}} \beta l)^2 \frac{\sin(\frac{1}{2} ql)}{\frac{1}{2} ql}$$

$(P_\gamma - P_a)$

on

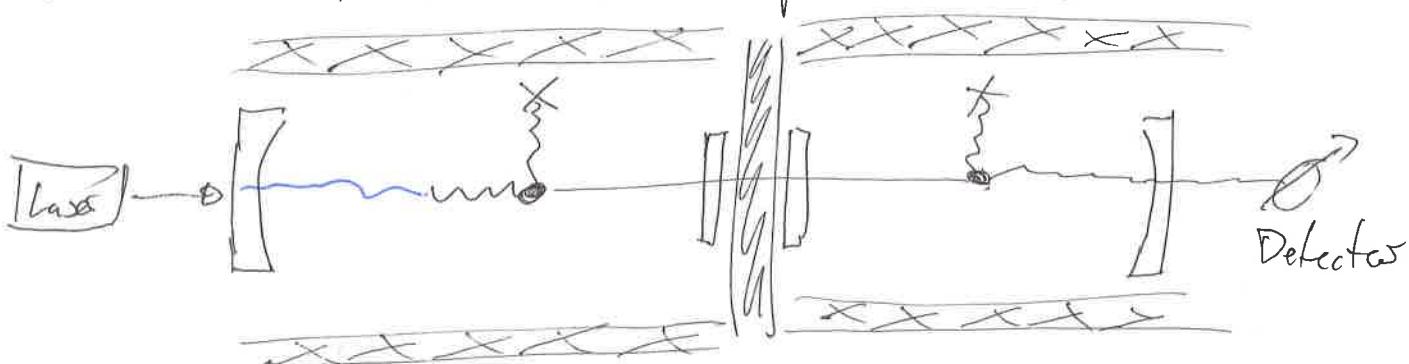
$\downarrow$  4<sup>th</sup> power

$$\Rightarrow P_{a \rightarrow \gamma \rightarrow a} = \frac{1}{16} (g_{\text{gyr}} \beta l)^4 + (ql)^2$$

$$= \underline{6 \cdot 10^{-38}} \left( \frac{g_{\text{gyr}}}{10^{-10} \text{GeV}^{-1}} \right)^4 \left( \frac{\beta}{\pi} \right)^4 \left( \frac{l}{10 \text{m}} \right)^4$$

$\Rightarrow$  long strings of dipole magnets

$\Rightarrow$  implement Optical resonators to recycle light  
 $\circlearrowleft$  to boost re-conversion to photons behind wall



$$\Rightarrow \text{powers } P_{a \rightarrow \gamma \rightarrow a} = 6 \cdot 10^{-38} F_{\text{PD}} F_{\text{RC}} \left( \frac{g_{\text{gyr}}}{10^{-10} \text{GeV}^{-1}} \right)^4 \left( \frac{\beta}{\pi} \right)^4 \left( \frac{l}{10 \text{m}} \right)^4$$

Finesse of interferometers

Magnet

# Any Light Particle Search (ALPS)

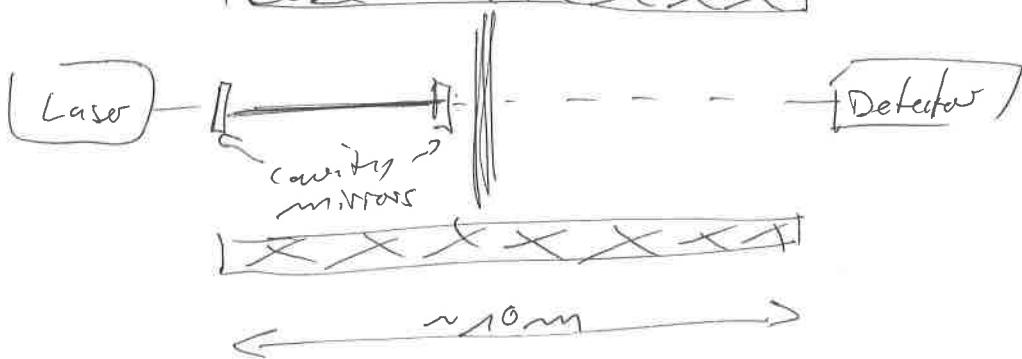
(9)

@ DESY use old HERA magnets  $\sim 5.3\text{ T}$

$\sim 10\text{ m length}$

$\sim 50\text{ mm aperture}$

Phase I: already concluded 2010



Limit:  $\text{g}_{\text{grav}} \rightarrow \text{few} \cdot 10^{-8} \text{ GeV}^{-1}$  for  $M_a \leq 1 \text{ meV}$

Phase II:

- use 10+10 magnets  $\Rightarrow$  total length  $\approx 100\text{ m}$
- Regeneration cavity to increase back-conversion probability
- Single photon counter using superconducting transition edge sensor (TES)

$\Rightarrow$  Increase of sensitivity by

$$F_{PC} \cdot F_{RC} \cdot N^4 \simeq 5000 \cdot 40.000 \cdot 10^4 \\ = 2 \cdot 10^8$$

Challenges:

- magnets from HERA are bent  $\rightarrow$  straighten (possible!)
- Optics: operate two  $\sim 100\text{ m}$  long optical resonators with high finesse that are mode matched ( $\rightarrow$  alignment precision!)  $\Rightarrow$  Experience from GW -
- Detectors: Single photon count @ 1064 nm (laser frequency)

# Helioscopes:

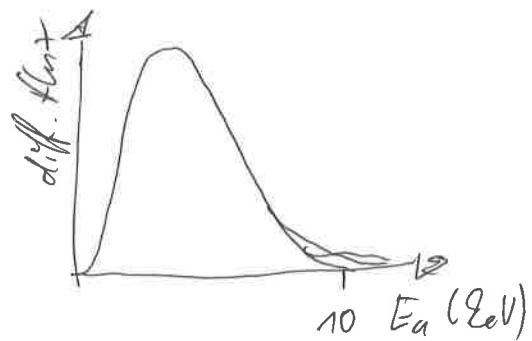
- Look for axions produced in the sun:

Robust prediction: if axions exist, they are emitted by the sun

$\Rightarrow$  No model dependence (like DM)!

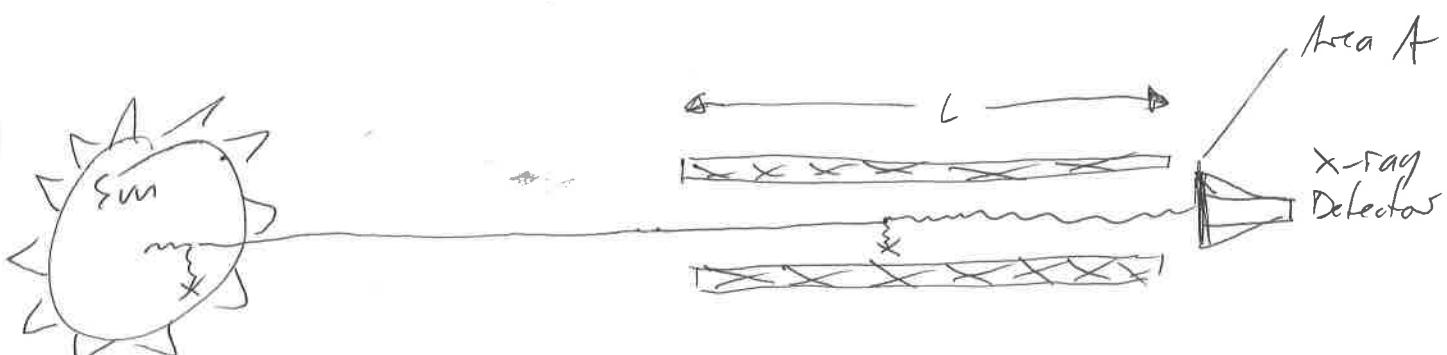
Production mechanism in sun:

Again: Primakoff effect



Detection mechanism:

Again: Primakoff effect



$$P_{\text{ax}} = 2.6 \cdot 10^{-17} \left( \frac{B}{10 \text{T}} \right)^2 \left( \frac{L}{10 \text{m}} \right)^2 \left( g_{g\gamma} \cdot 10^{10} \text{GeV} \right)^2 F$$

CAST @ CERN:

- Use decommissioned LTC test magnet  $L=10 \text{m}$   $B=9 \text{T}$
- Moving platform: point towards sun
- X-ray detectors to detect ~keV x-rays from  $a \rightarrow \gamma$  conversion

Future plan:

1A×0

East Crust

gap  $\leq 10^{-10} \text{ GeV}^{-1}$  for  $m_N \leq 1 \text{ eV}$

(RP)

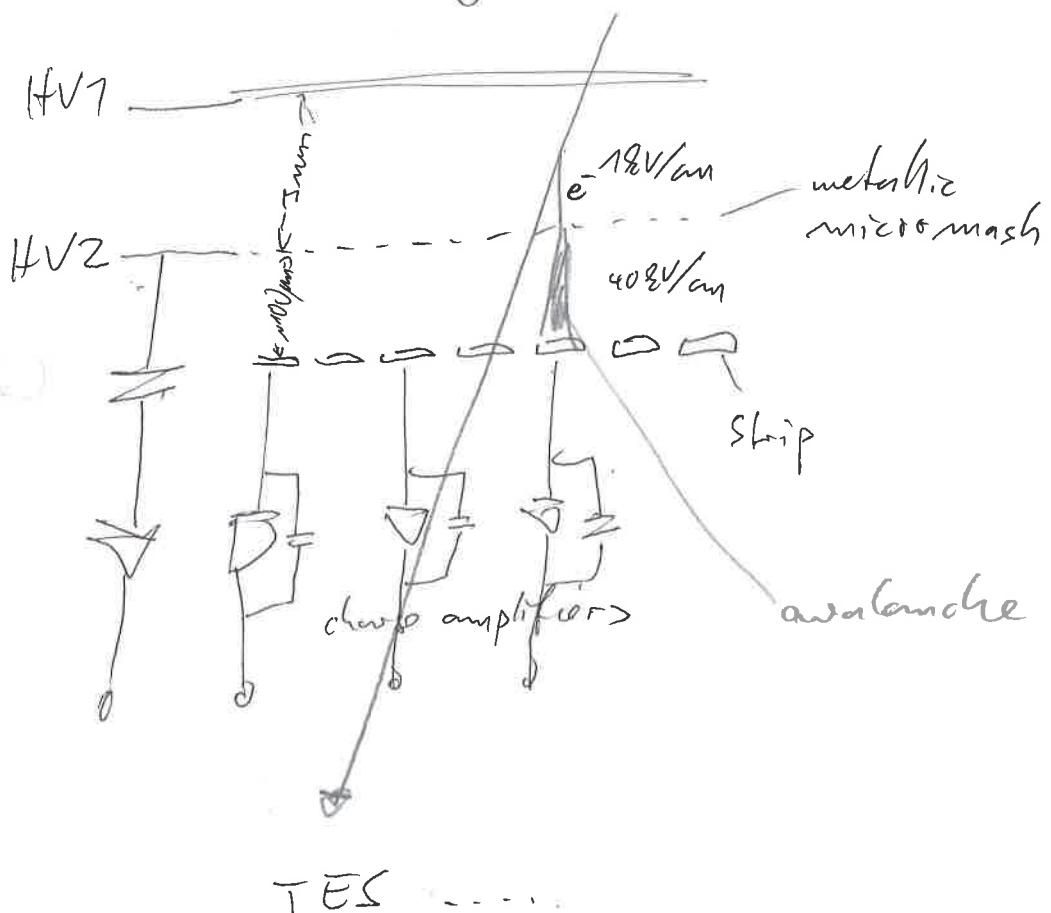
conceptual design:

- Use large toroidal 8-coil magnet  $L \approx 20 \text{ m}$   $\varnothing 60 \text{ cm}$

Figure of merit:  $B^2 \cdot L^2 \cdot A = 21 T_B^2 \text{ m}^4$

\$\$\$

- Use x-ray optics to maximize  $A$  while reducing overall detector size, i.e. background
- Use "Ultra-Low Background" x-ray detectors  
Micromegas from clean material



→ projected sensitivity gap  $\sim$  few  $10^{-11} \text{ GeV}^{-1}$

# Experimental situation projections

(12)

