

Germanium Detectors

–

Working Principle and Pulse Shape Analysis

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IMPRS Young Scientist Workshop
Ringberg Castle - 07/2019



Max-Planck-Institut für Physik
(Werner-Heisenberg-Institut)

Outline

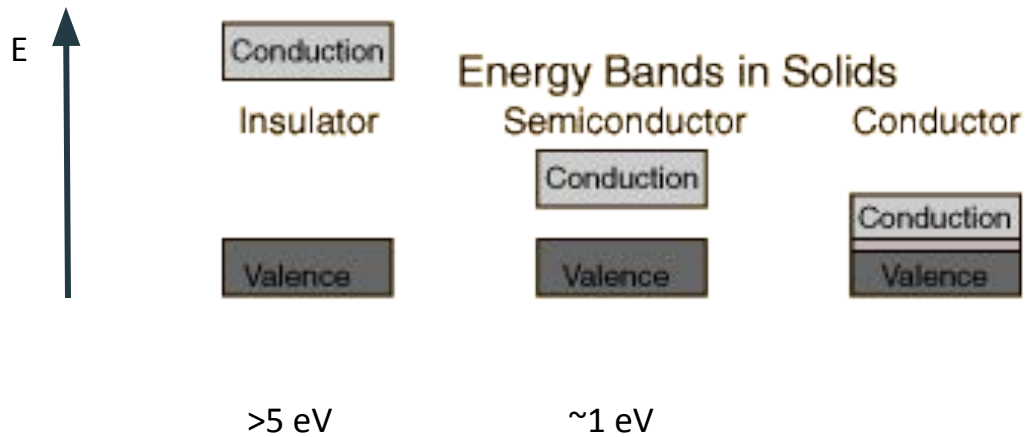
- Motivation: Germanium Detectors in $0\nu\beta\beta$ searches (LEGEND)
- Working Principle of Solid State Detectors
 - Semiconductors
 - p-n junction
- Charge Drift & Simulation
- Signal Formation
 - Weighting Potential
 - Pulses
- Using Pulse Shapes to reduce background: Multi-site vs. Single-site Event Discrimination
- Summary

Germanium in $0\nu\beta\beta$ searches

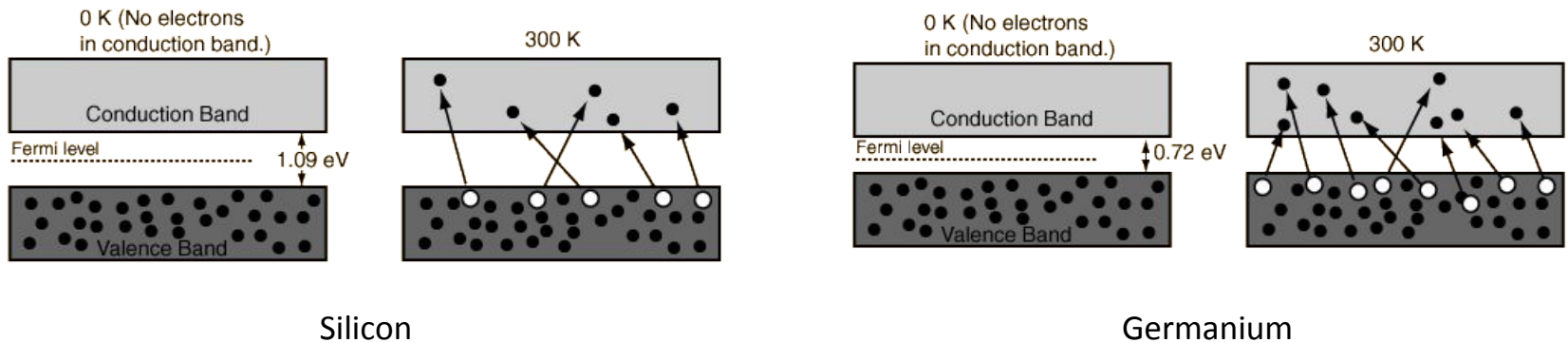
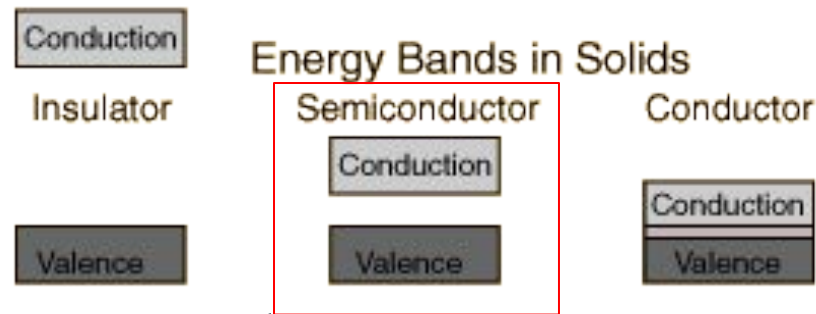
- Germanium has an isotope, ^{76}Ge , that undergoes $2\nu\beta\beta$ decay and is therefore also a candidate for $0\nu\beta\beta$
- As a semiconductor (enriched) Germanium can be used to produce solid state detectors
 - Source = Detector
- Solid state detectors come with many favorable properties for rare event searches
 - Excellent Energy Resolution
 - Low intrinsic radioactivity
 - Fast Signal collection
 - Compact geometries

Working Principle of solid state detectors

Band structure of solids

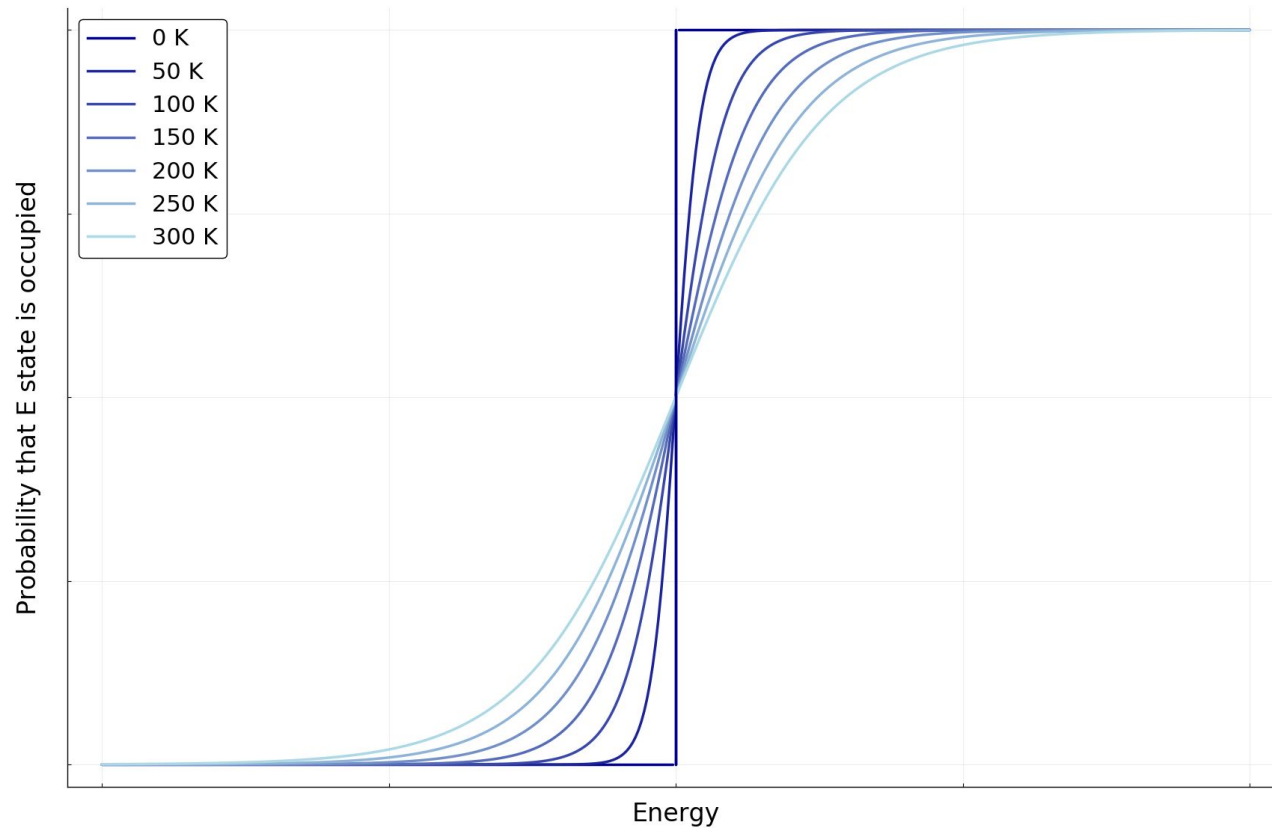


Band structure of solids

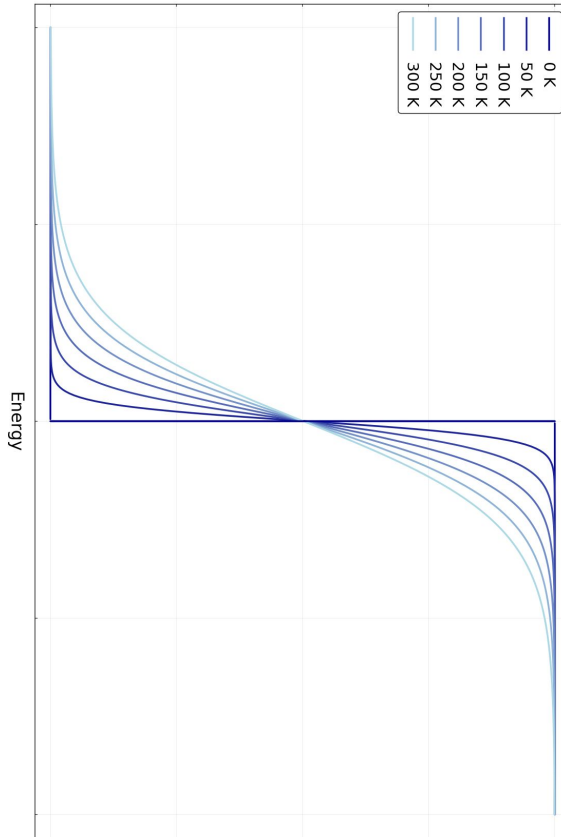


Fermi Level and Fermi Function

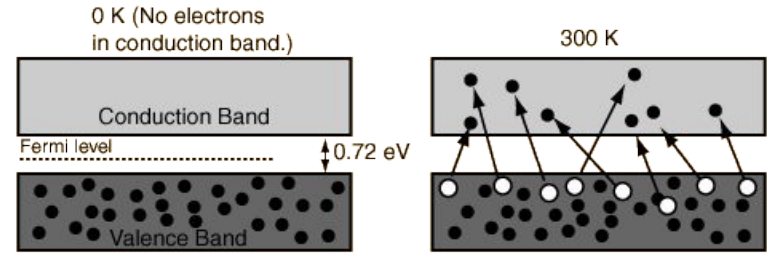
$$f(E) = \frac{1}{e^{(E - E_F)/kT} + 1}$$



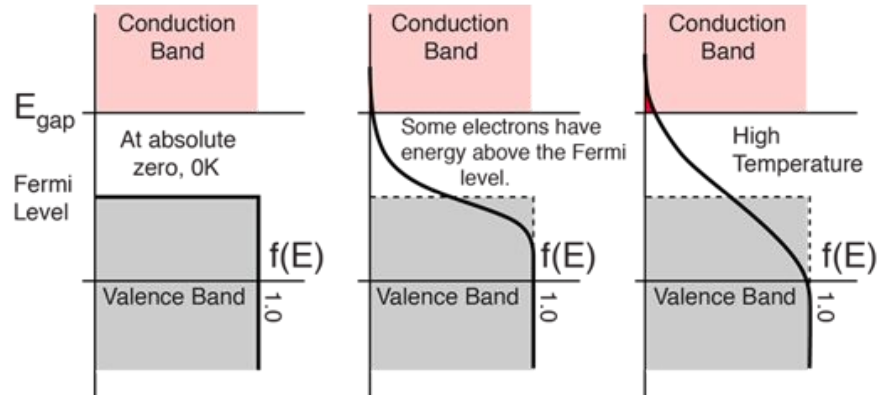
Probability that E state is occupied



+



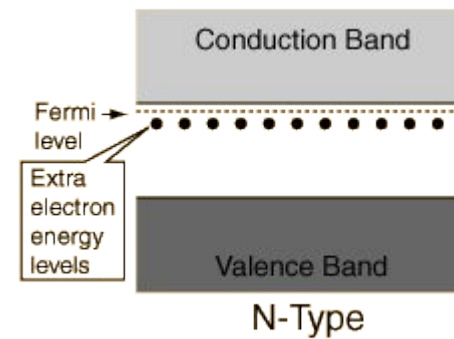
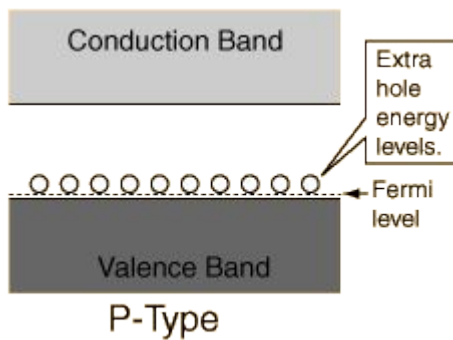
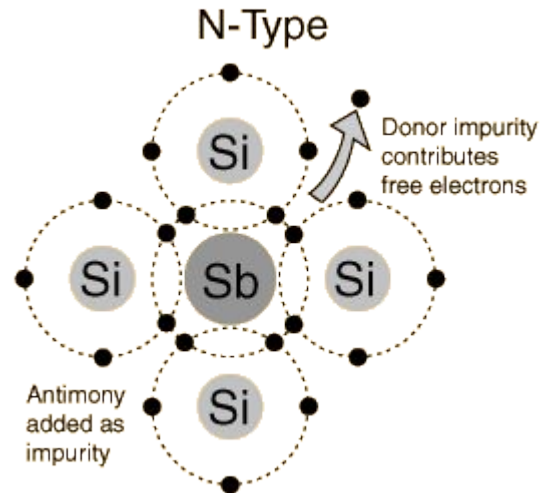
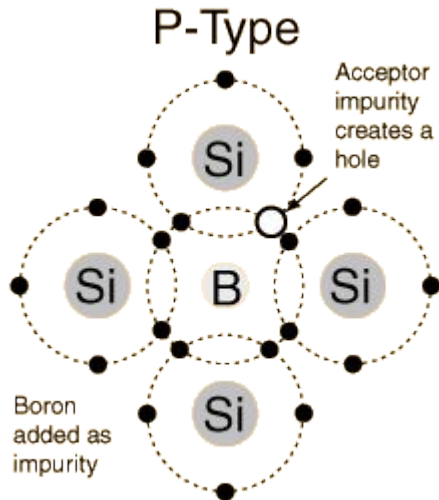
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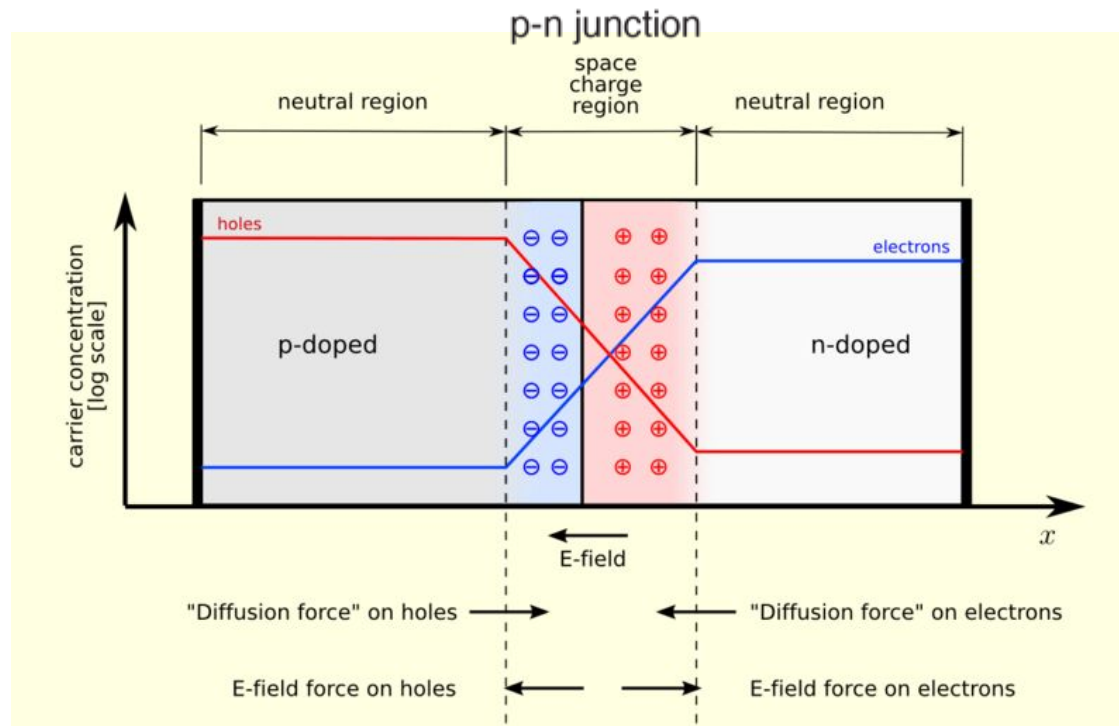
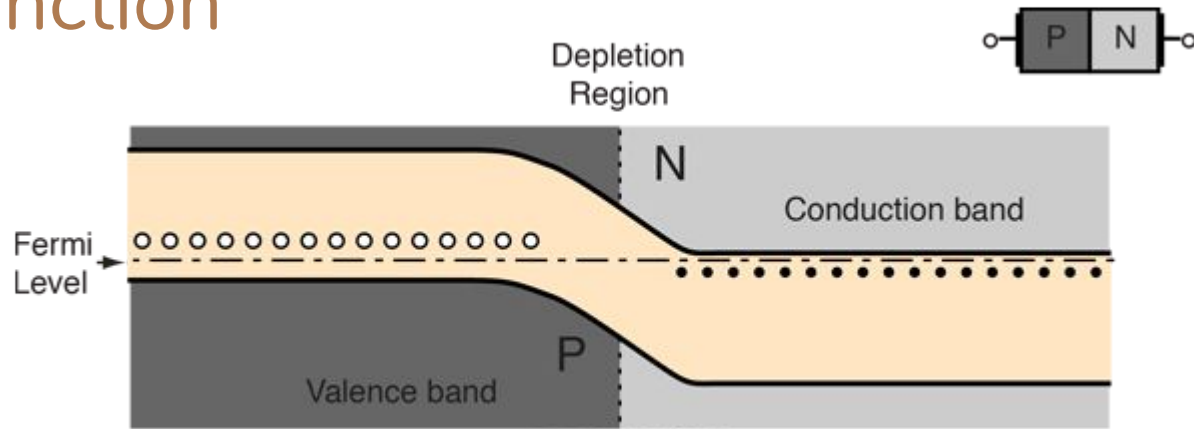
No electrons can be above the valence band at 0K, since none have energy above the Fermi level and there are no available energy states in the band gap.

At high temperatures, some electrons can reach the conduction band and contribute to electric current.

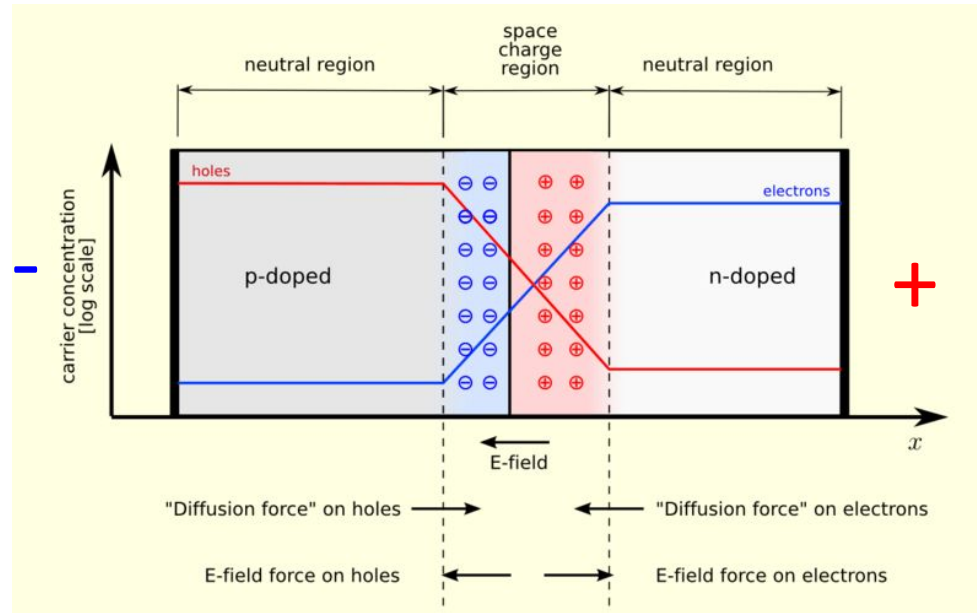
Doping of Semiconductors



p-n junction



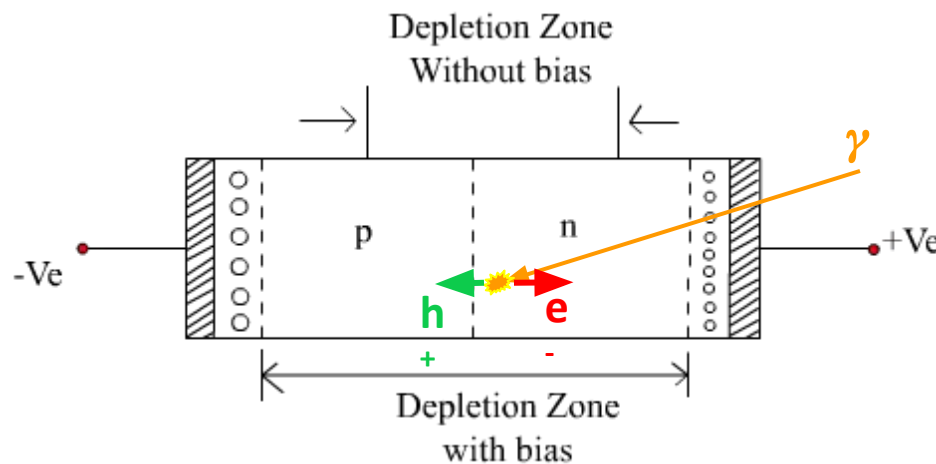
Reverse Bias



A p-n junction is a **diode**.

Applying a **reverse bias** (+ on n-side, - on p-side) increases the **depleted region**
 → fixed space charges, E-Field, no free charge carriers

The **depleted region** is the “active region” of the detector
 → Full depletion!

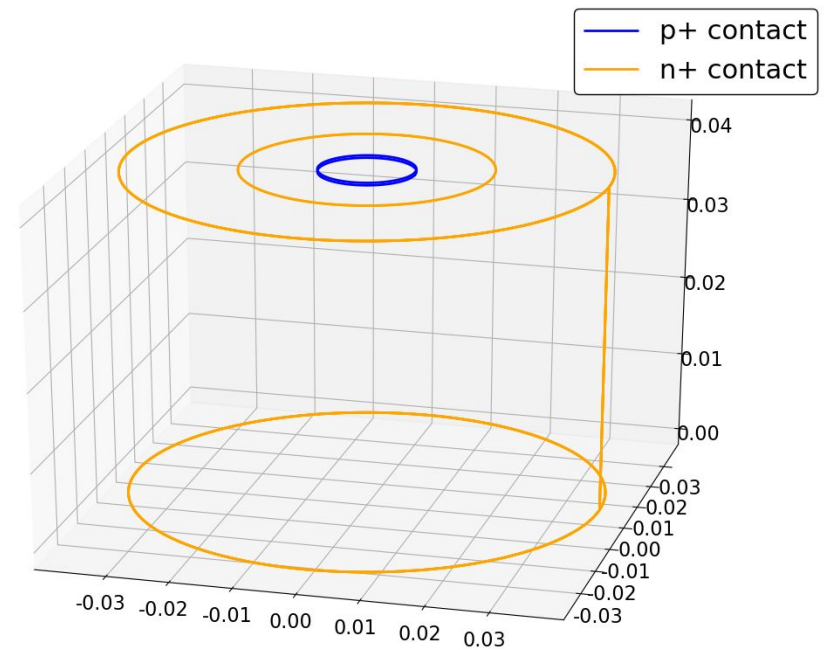


Incident ionizing radiation creates electron - hole pairs in an amount proportional to the deposited energy

The electrons and holes created in the depletion region move to the respective electrodes
 → Detectable Current!

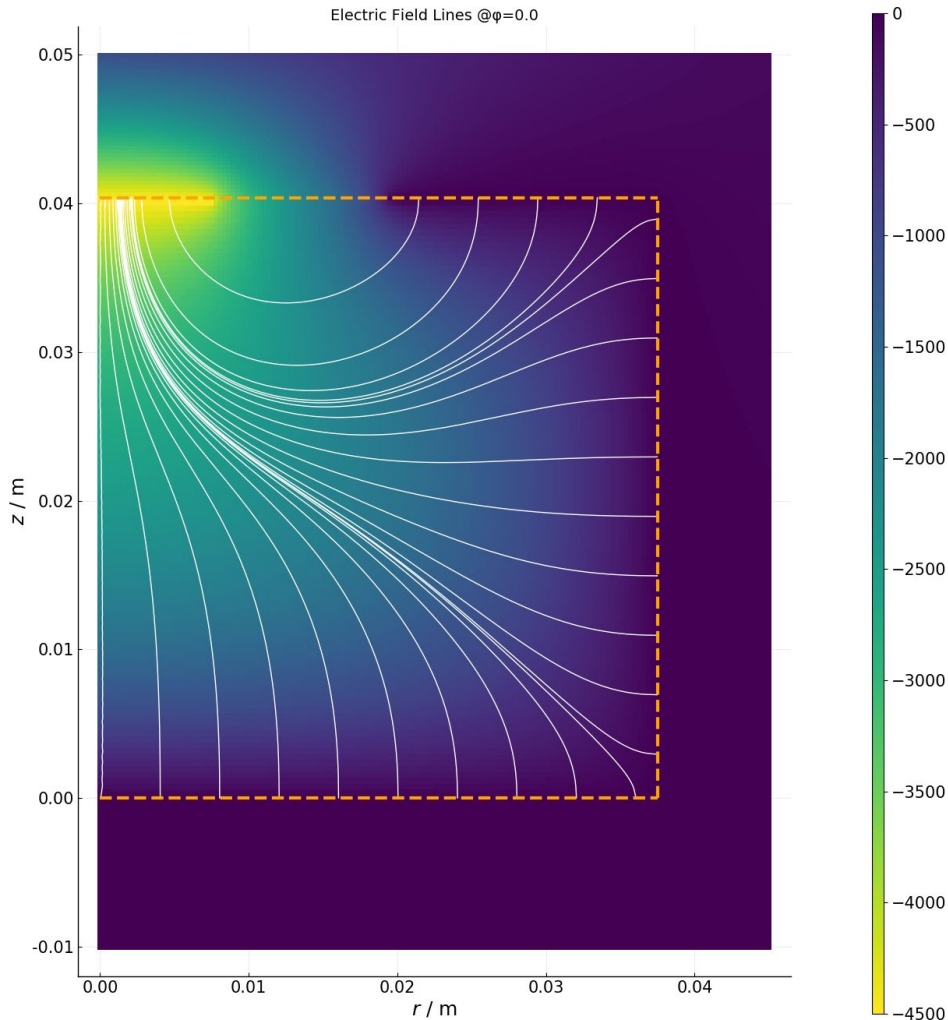
Charge Drift

Real World Detectors - p-type Point Contact



<https://github.com/JuliaHEP/SolidStateDetectors.jl>

Electric Potential & Charge Drift



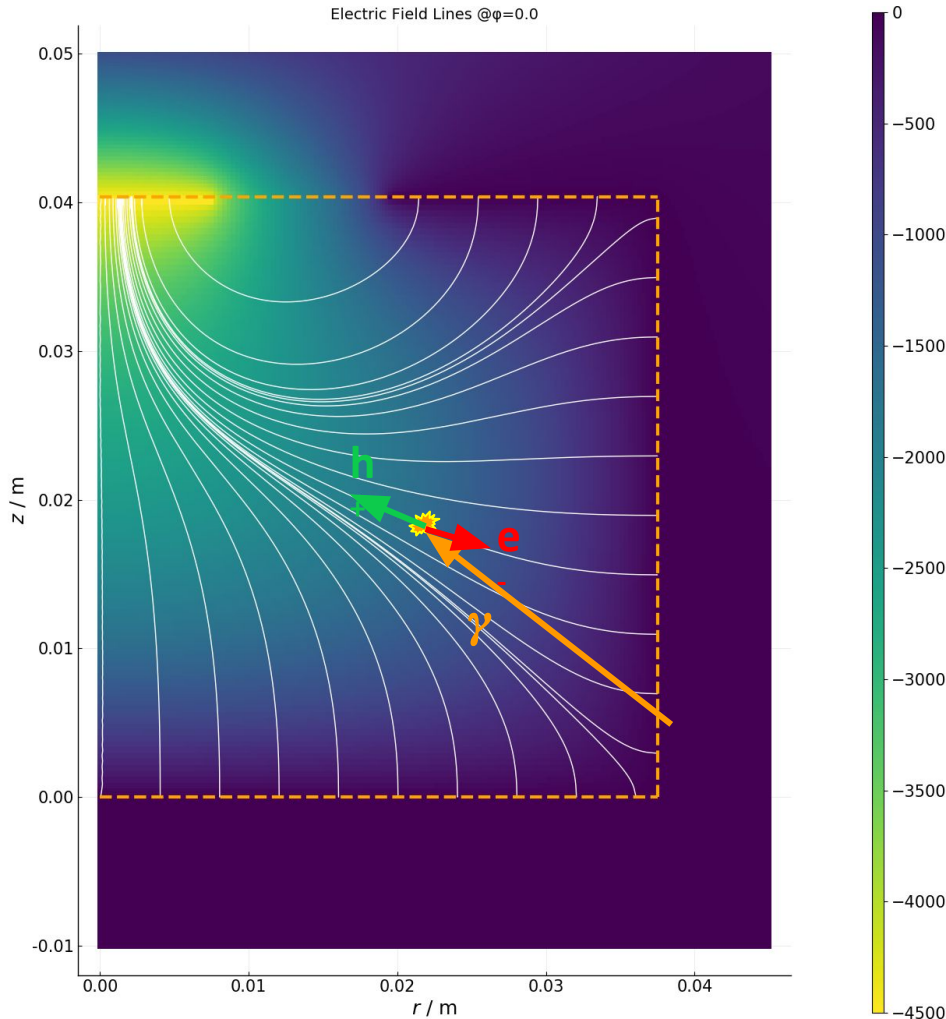
Potential: Solve Gauss' Law

$$\nabla(\epsilon_r(\vec{\mathbf{r}})\nabla\varphi(\vec{\mathbf{r}})) = \frac{\rho(\vec{\mathbf{r}})}{\epsilon_0}, \quad \vec{\mathbf{r}} = (r, \phi, z)$$

Electric Field:

$$\vec{\mathcal{E}}(\vec{\mathbf{r}}) = -\nabla\varphi(\vec{\mathbf{r}})$$

Electric Potential & Charge Drift



Potential: Solve Gauss' Law

$$\nabla(\epsilon_r(\vec{\mathbf{r}})\nabla\varphi(\vec{\mathbf{r}})) = \frac{\rho(\vec{\mathbf{r}})}{\epsilon_0}, \quad \vec{\mathbf{r}} = (r, \phi, z)$$

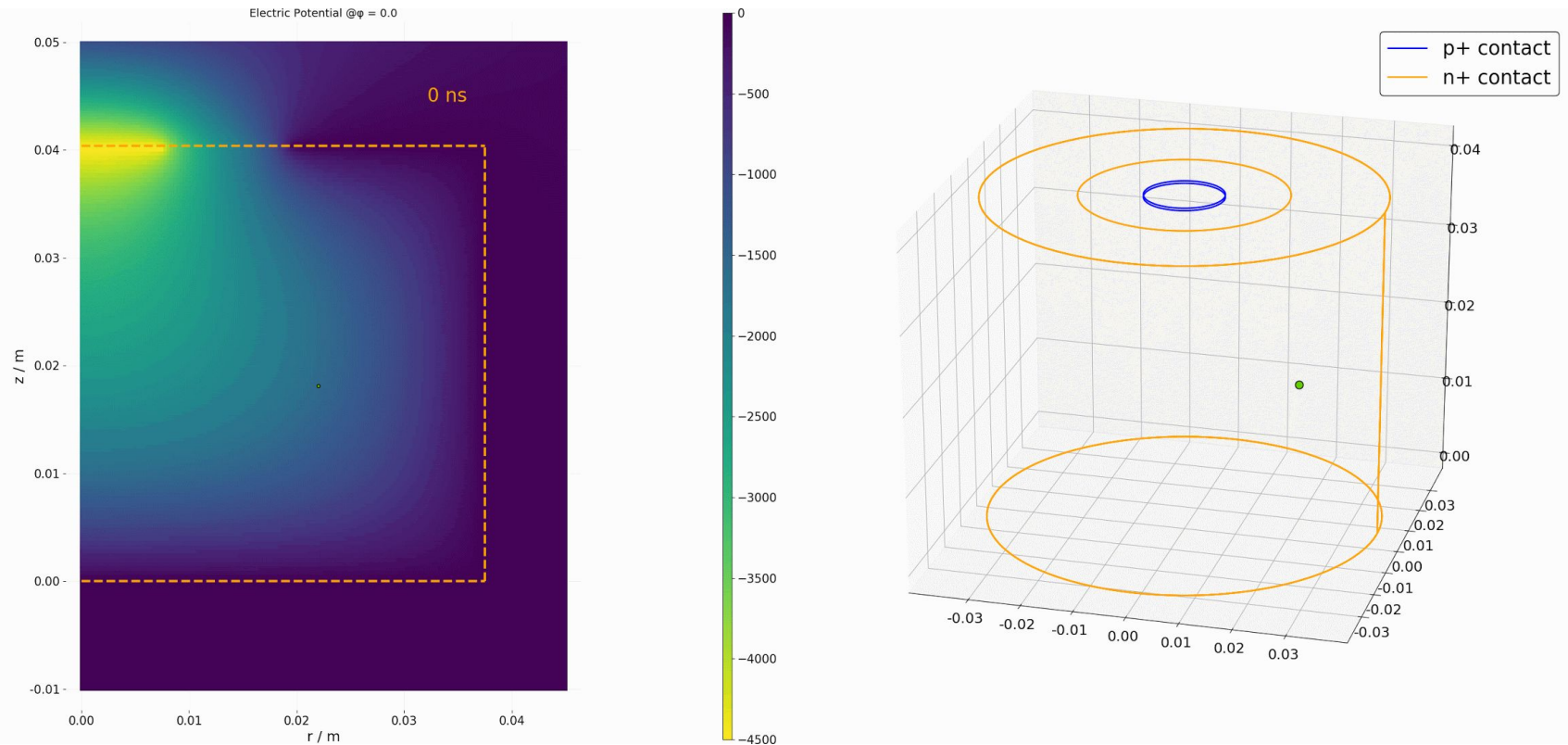
Electric Field:

$$\vec{\mathcal{E}}(\vec{\mathbf{r}}) = -\nabla\varphi(\vec{\mathbf{r}})$$

Charge Drift Model:

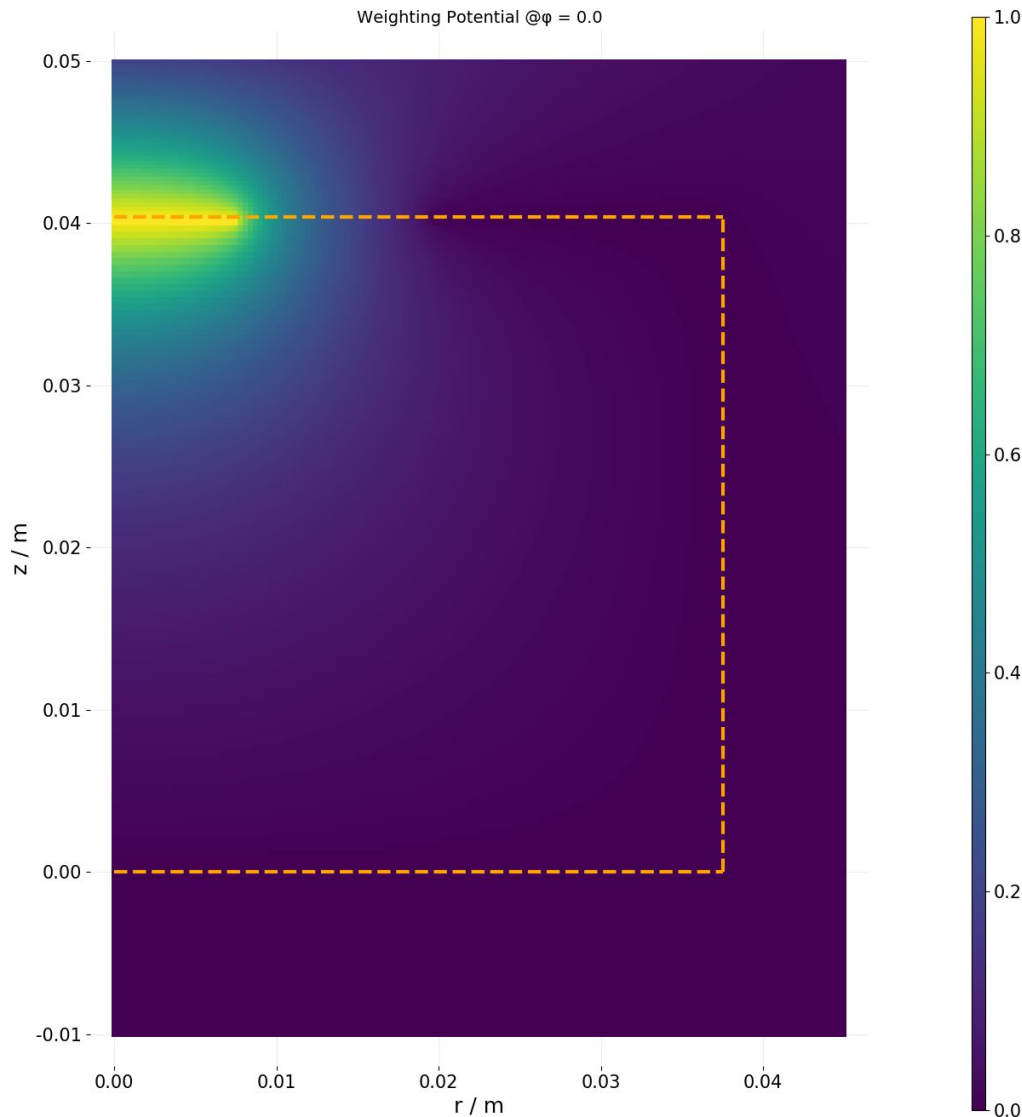
- Effects of Crystal Axes
- Different Mobilities for \mathbf{e}^- and \mathbf{h}^+
- ...

Electric Potential & Charge Drift



Signal Formation

Weighting Potential & Signal Formation



Shockley-Ramo Theorem:
Charge induced on given electrode by motion of charge carriers according to:

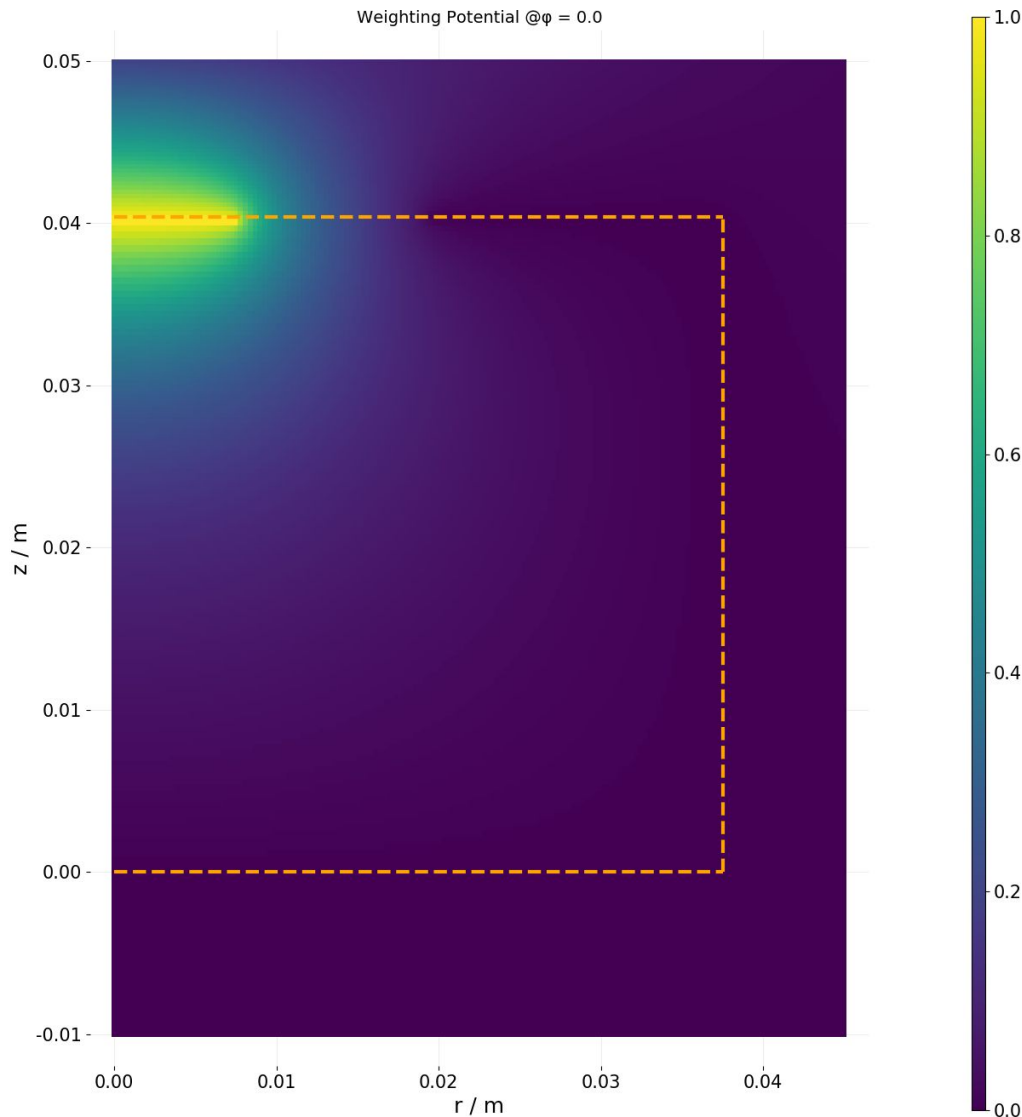
$$Q = q\Delta\phi_0$$

q = moving charge,
 ϕ_0 = Weighting Potential

Solve:

$$\nabla(\epsilon_r(\vec{r})\nabla\phi_0(\vec{r})) = 0$$

Weighting Potential & Signal Formation



Shockley-Ramo Theorem:
Charge induced on given electrode by motion of charge carriers according to:

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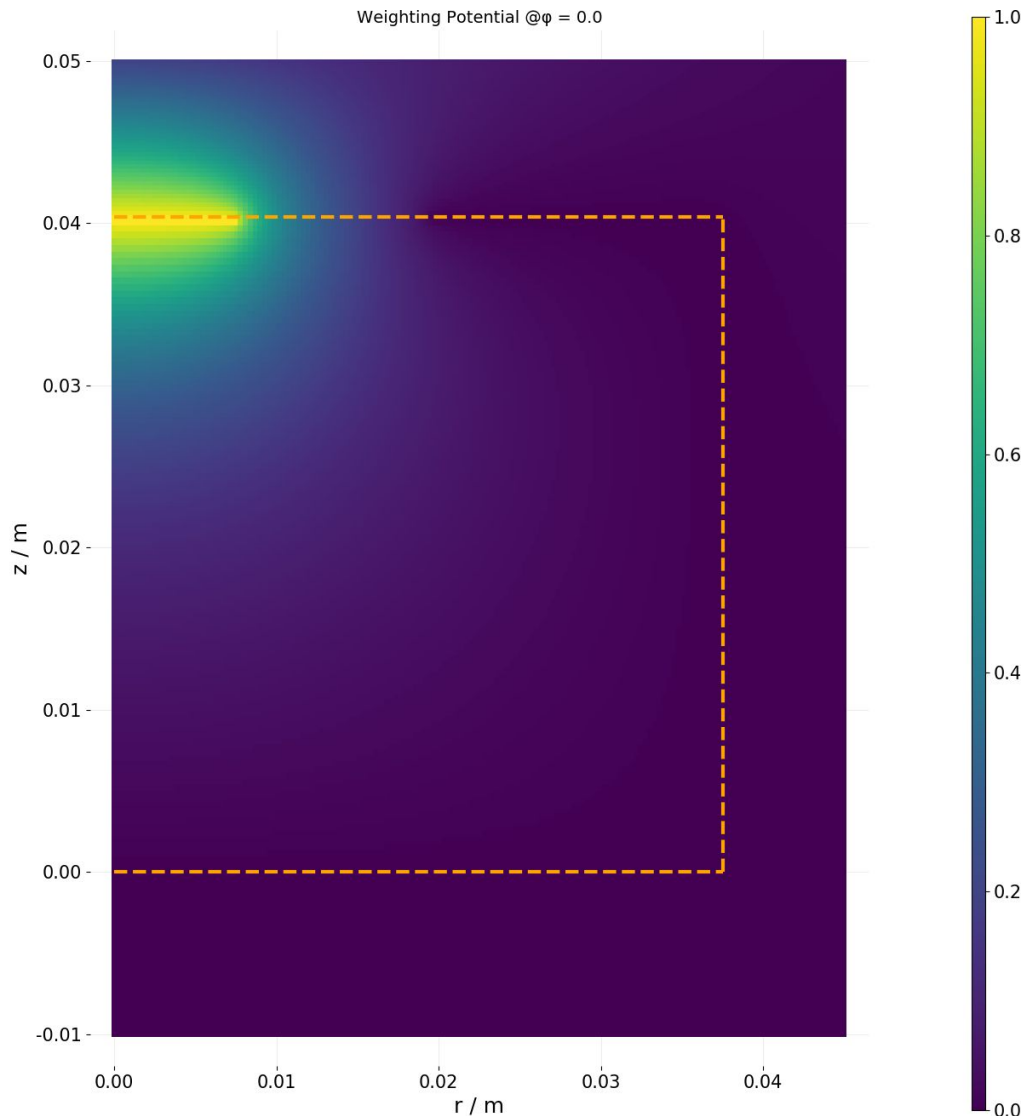
Solve:

$$\nabla(\epsilon_r(\vec{r})\nabla\phi_0(\vec{r})) = 0$$

$$Q_e(t) = n_e \cdot (-1 \cdot e) \cdot (\phi_0(\vec{r}_e(t)) - \phi_0(\vec{r}_e(t_0)))$$

$$Q_h(t) = n_h \cdot e \cdot (\phi_0(\vec{r}_h(t)) - \phi_0(\vec{r}_h(t_0)))$$

Weighting Potential & Signal Formation



Shockley-Ramo Theorem:
 Charge induced on given electrode by
 motion of charge carriers according to:

$$Q = q \Delta \varphi_0$$

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 φ_0 = Weighting Potential

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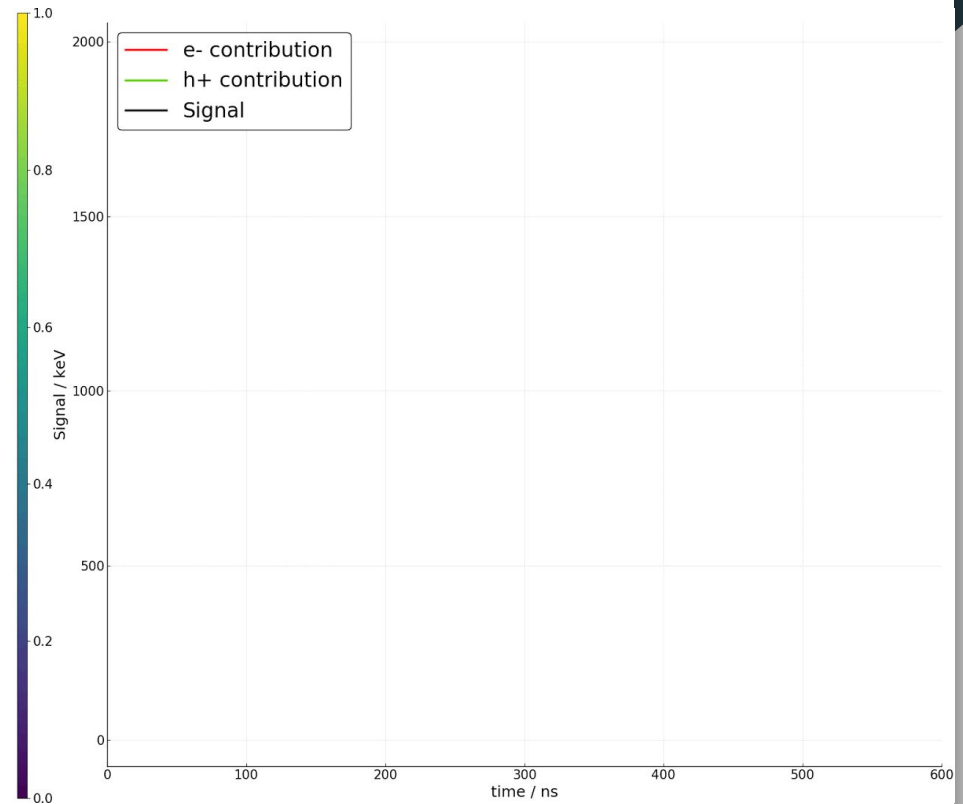
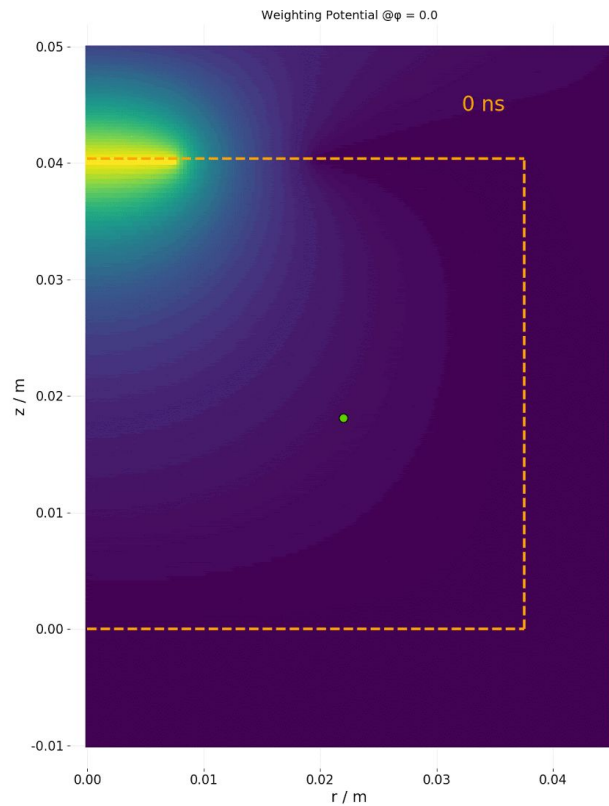
$$n_e = n_h, \quad \vec{r}_e(t_0) = \vec{r}_h(t_0)$$

$$Q_{tot} = Q_e + Q_h$$

$$Q_{tot}(t) = n_{e/h} \cdot e \cdot (\varphi_0(\vec{r}_h(t)) - \varphi_0(\vec{r}_e(t)))$$

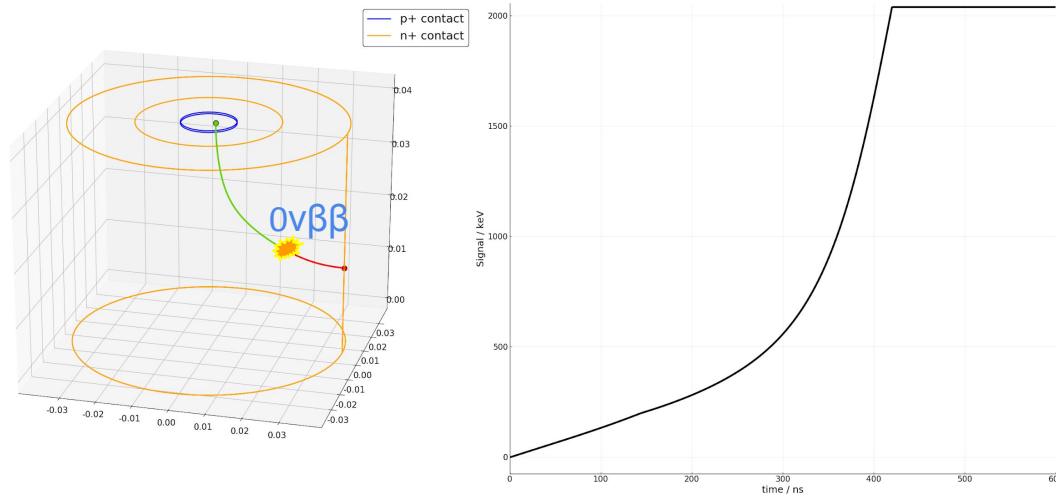
Weighting Potential & Signal Formation

$$Q_{tot}(t) = n_{e/h} \cdot e \cdot (\varphi_0(\vec{r}_h(t)) - \varphi_0(\vec{r}_e(t)))$$

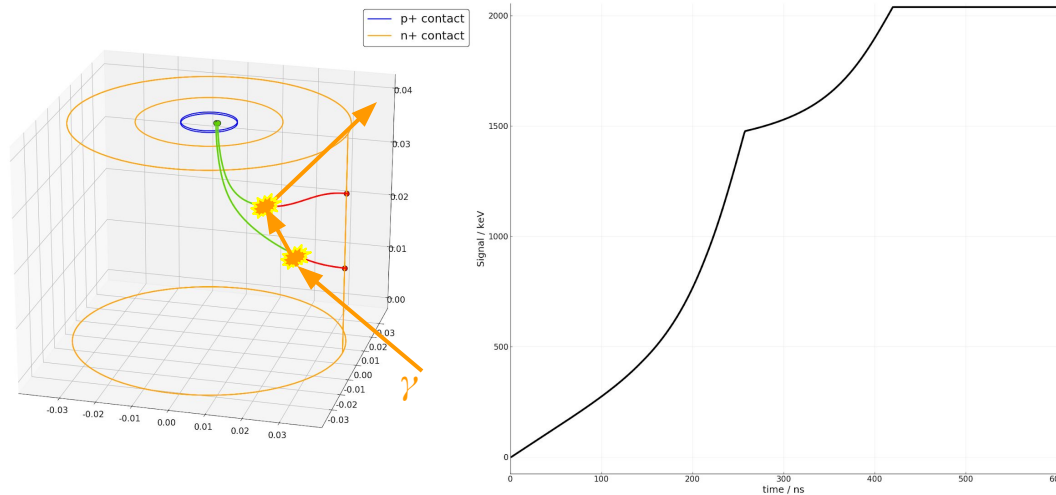


Pulse Shape Analysis(PSA)

PSA in GERDA / LEGEND (example)

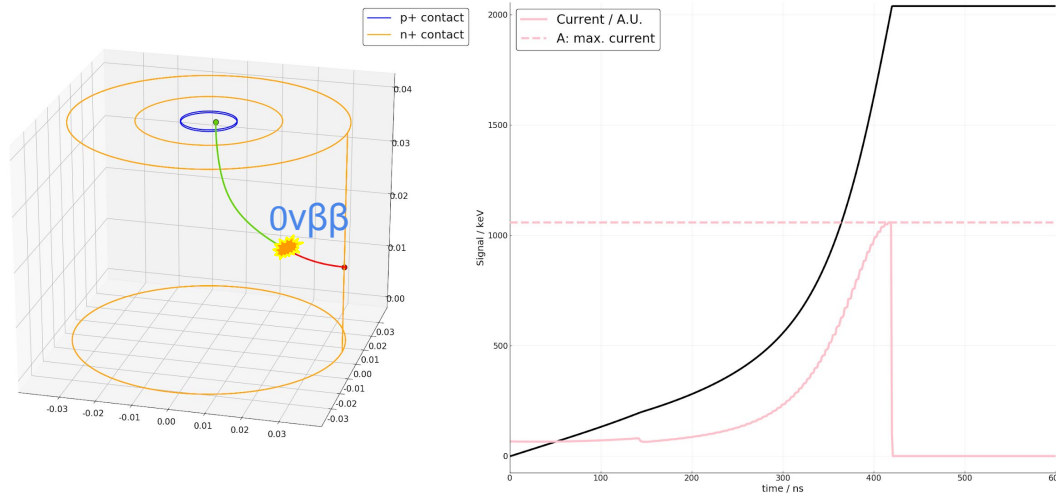


Signature of $0\nu\beta\beta$ in Ge:
 2 e^- depositing their energy
 (2.039 MeV) very locally
 → Single - Site Event

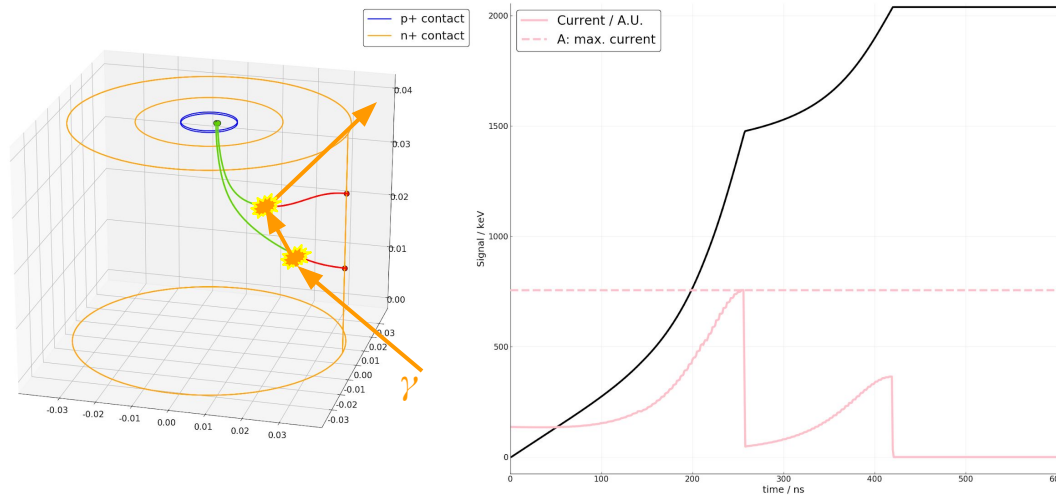


Background in the Region of
 Interest:
 A high energy γ energy Compton
 scatters within the crystal
 depositing 2.039 MeV in total
 → Multi - Site Event

PSA in GERDA / LEGEND (example)



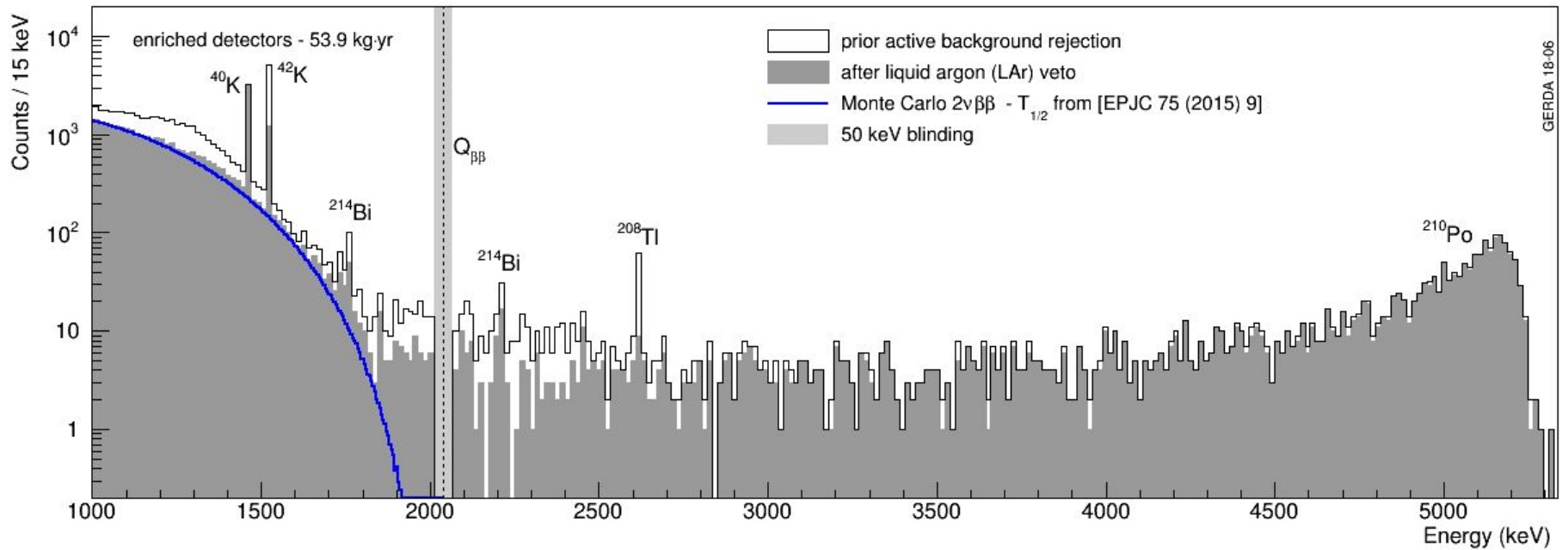
Signal energy: 2039 keV (E)
High max current value (A)



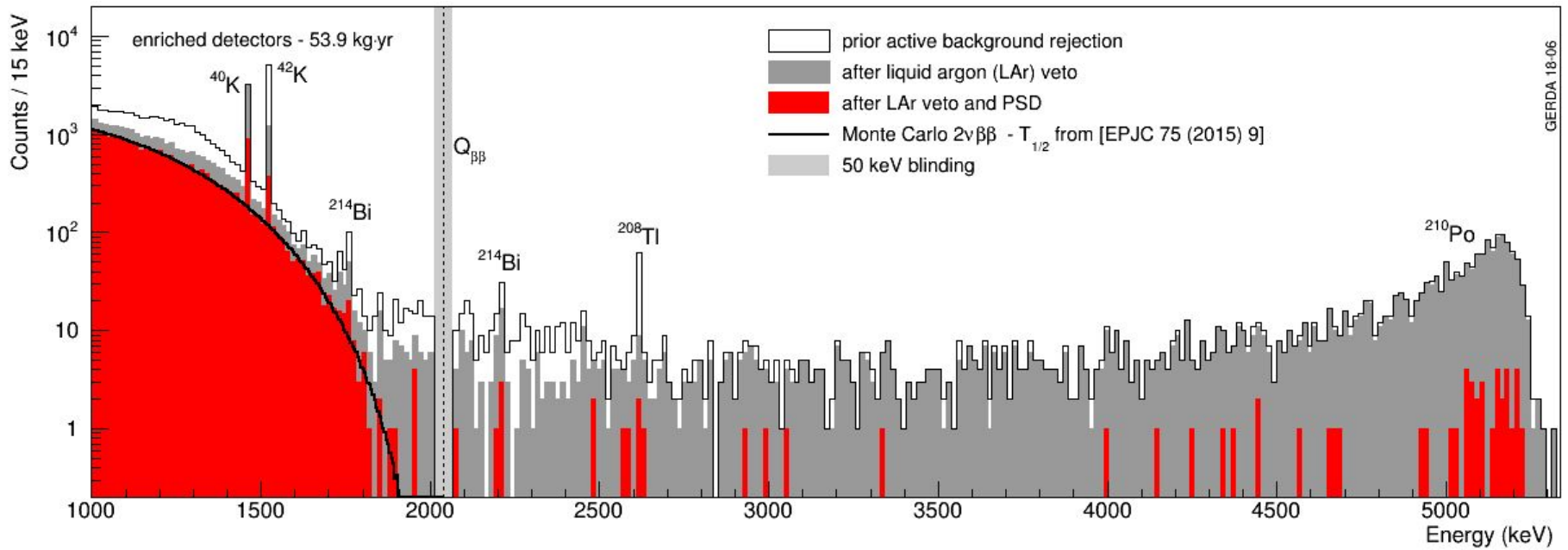
Signal Energy: 2039 keV (E)
Lower max current value (A)

→ Discriminate on A / E ratio

Pulse Shape Analysis Cuts on GERDA

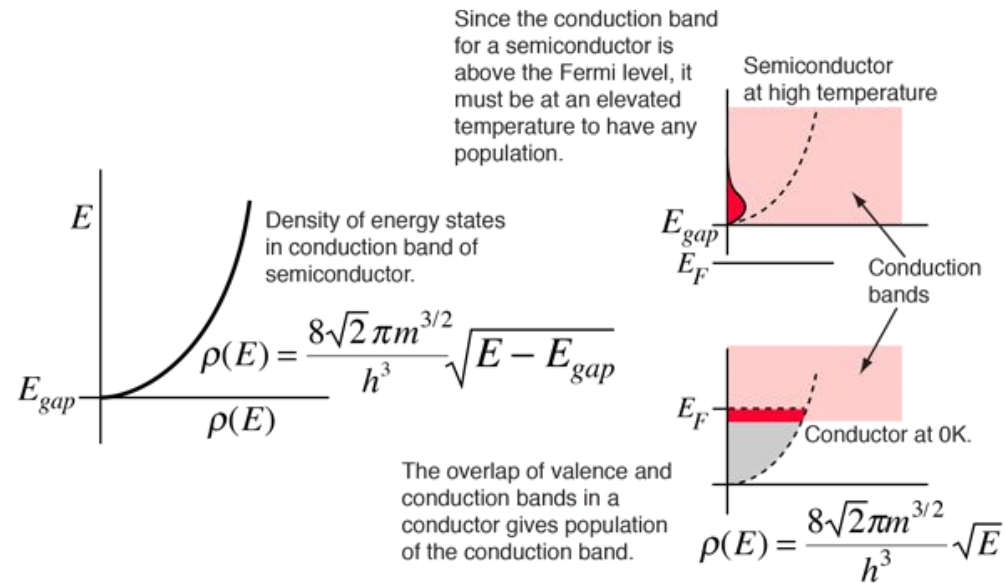


Pulse Shape Analysis Cuts on GERDA

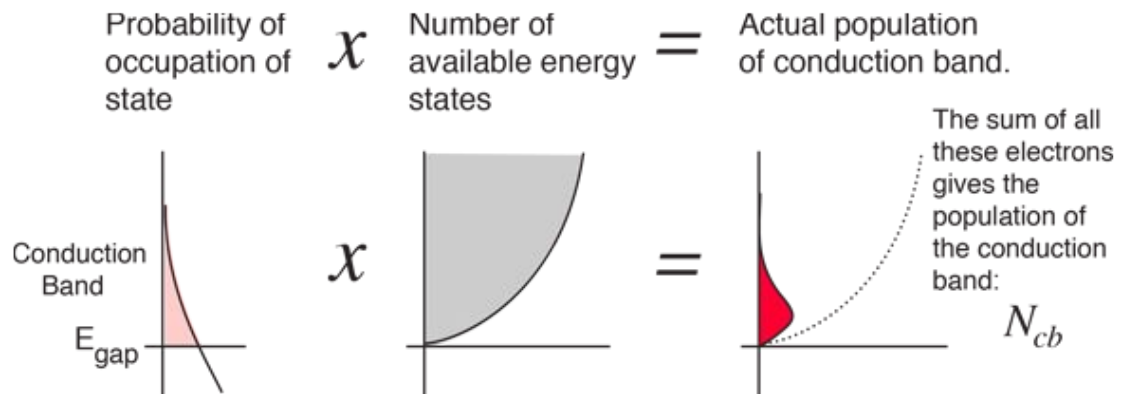


BACKUP

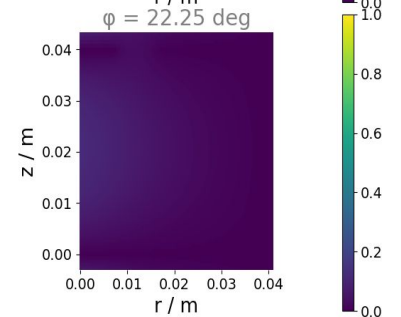
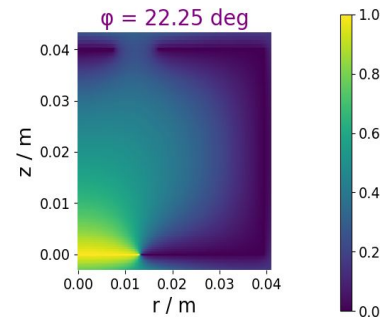
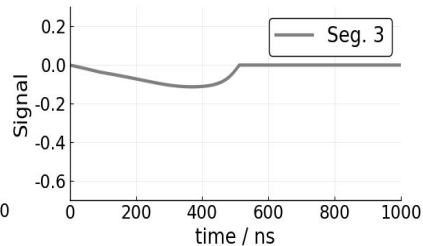
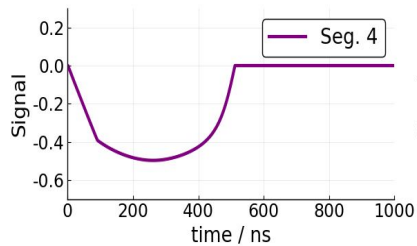
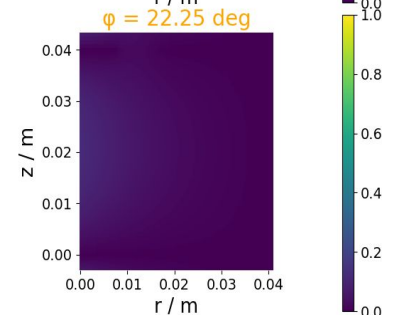
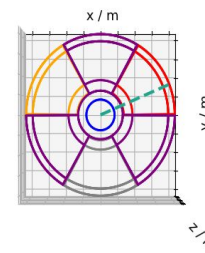
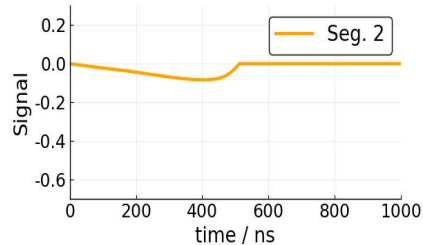
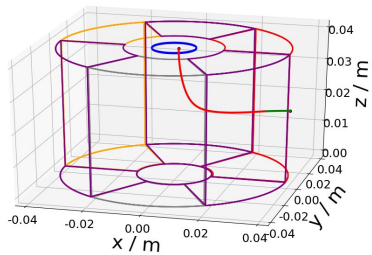
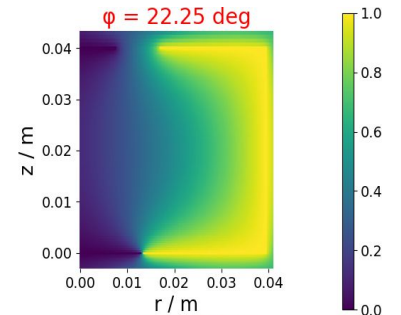
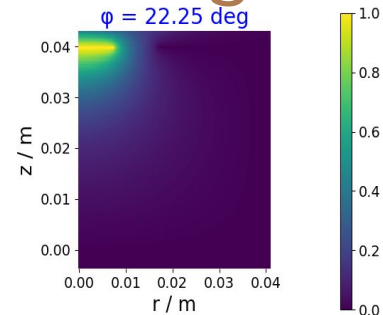
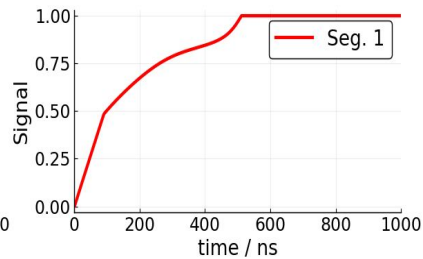
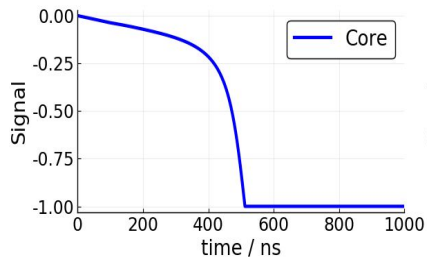
Density of States



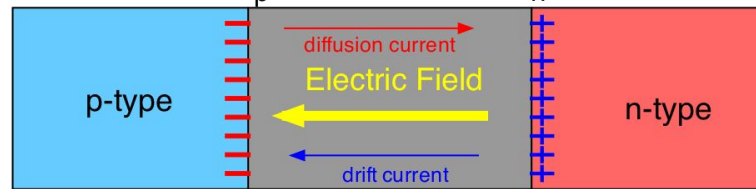
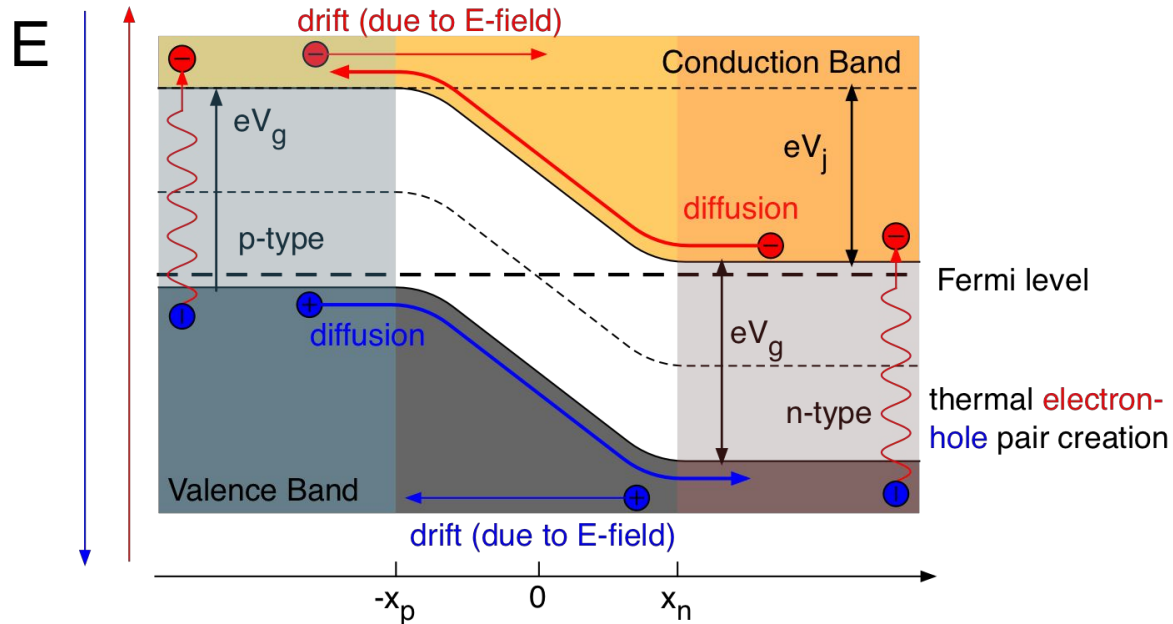
Population of Conduction Band



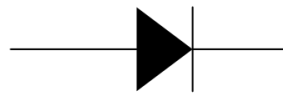
Signal and mirror pulses in segments



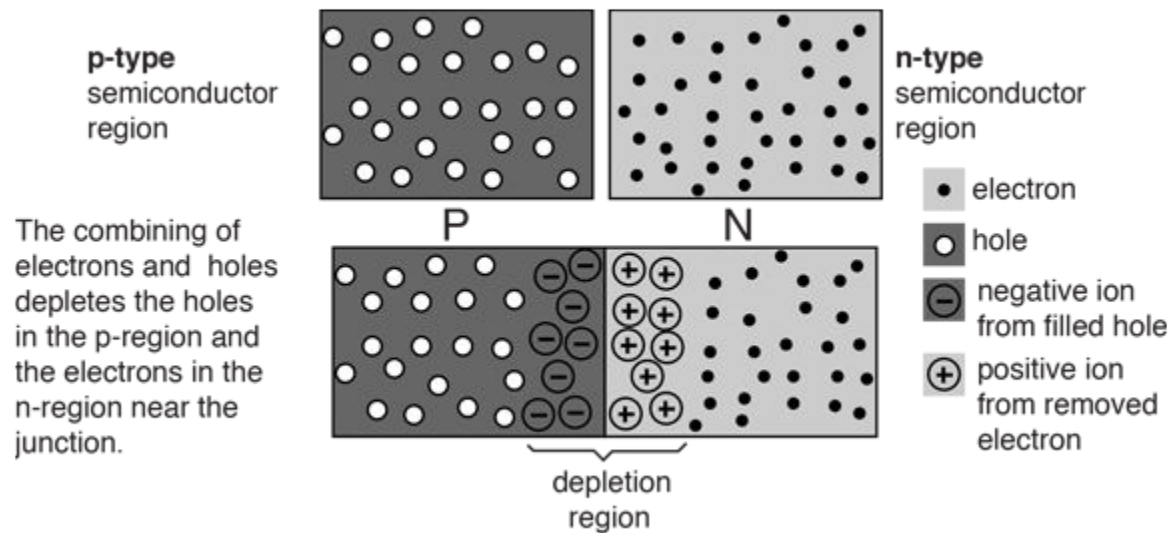
p-n-Junction



Schematic of pn-junction



Depletion Region & Reverse Bias



Field calculation

$$\nabla(\epsilon_r(\vec{\mathbf{r}})\nabla\varphi(\vec{\mathbf{r}})) = \frac{\rho(\vec{\mathbf{r}})}{\epsilon_0}, \quad \vec{\mathbf{r}} = (r, \phi, z)$$

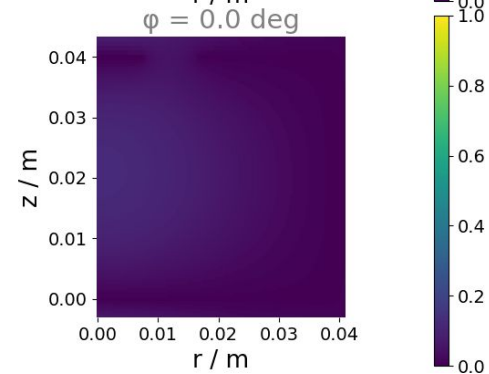
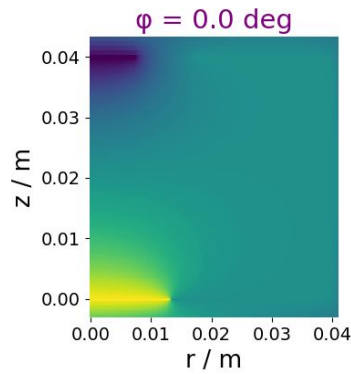
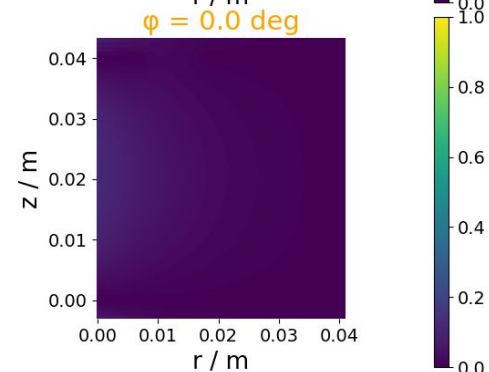
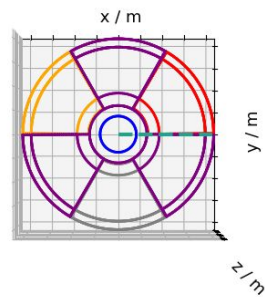
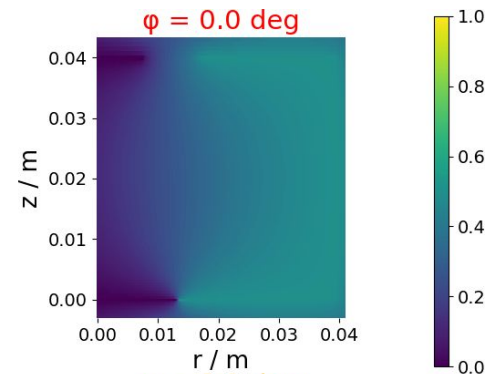
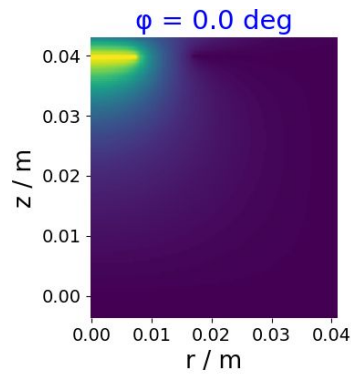
- Electric potential calculated by Gauss' law

$\varphi(\mathbf{r})$ electric potential, ϵ_0 vacuum permittivity, $\epsilon_r = 16$ dielectric constant of germanium, $\rho(\mathbf{r})$ impurity density

- Weighting potential calculation $\nabla(\epsilon_r(\vec{\mathbf{r}})\nabla\varphi_W(\vec{\mathbf{r}})) = 0$

with the boundary conditions that the weighting potential equals unity on the considered electrode and zero otherwise

Weighting Potentials



Drift velocity model

$$v_l = \frac{\mu_0 E}{(1 + (E/E_0)^\beta)^{1/\beta}} - \mu_n E$$

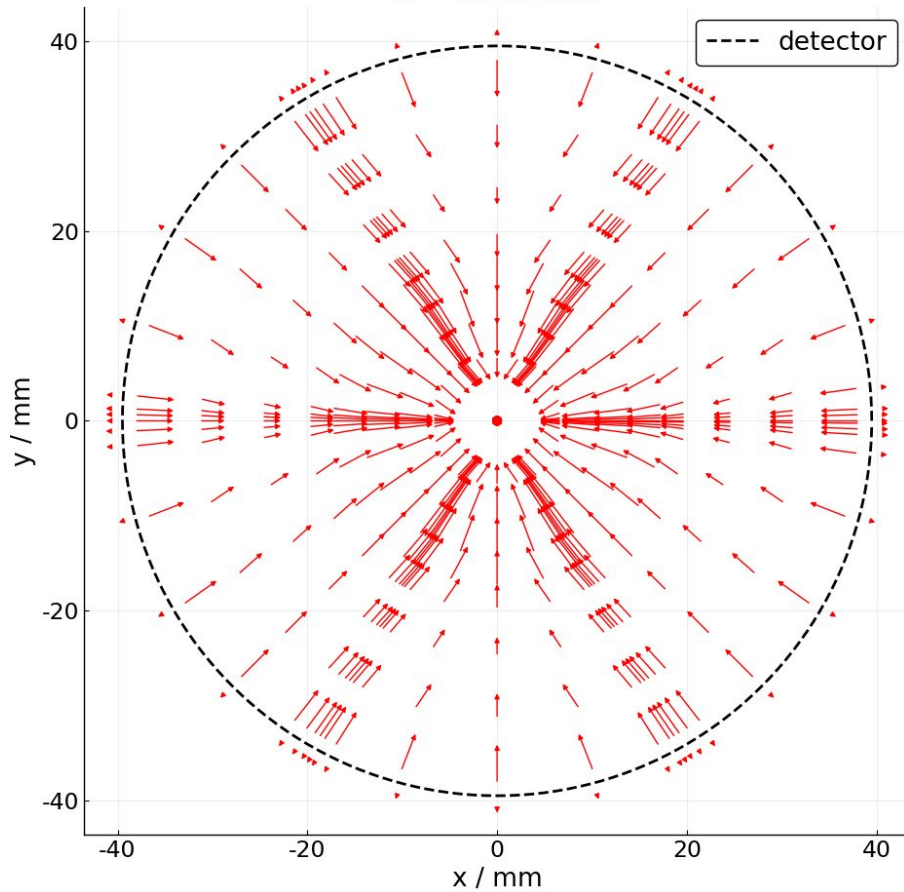
- Parallel to the principal crystal axis the drift velocity is parallel to the electric field
- μ_0 , E_0 , β and μ_n parameters different for electrons and holes and for $\langle 100 \rangle$ and $\langle 111 \rangle$ axes
- Different models for electrons and holes implemented from [1] B. Bruyneel et al., [NIM A 569 \(2006\) 764](#)
also used by e.g. AGATA

Drift velocity model

Electron Drift Velocity Field

$r\phi$ - cut

$z = 29.97$ mm



rz - cut

$\phi = 0.0$ deg

