# A Different Perspective on Strong CP & Axions

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Goal: Present alternative language to understand...

... how axions solve strong CP problem.

... which kind of new physics can spoil the solution.

[Gia Dvali, hep-th/0507215 (2005)]

QCD 0 0 Conclusion



## Maxwell and the $\theta$ -Term

Maxwell:  $\mathscr{L}_{MW} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu}$ ,  $F_{\mu\nu} = \partial_{\mu}A_{\nu} - \partial_{\nu}A_{\mu}$ 

Why these terms?

- Gauge invariant under U(1)
- Up to energy dimension 4

Wait! There is one more:  $\mathscr{L}_{\theta} \sim \theta F_{\mu\nu} \tilde{F}^{\mu\nu}$ ,  $\sim \theta \vec{E} \cdot \vec{B}$   $\tilde{F}^{\mu\nu} = \epsilon^{\mu\nu\rho\sigma} F_{\rho\sigma}$  $\theta = \text{const.}$ 

CP violation!

QCD 0 0 Conclusion



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$$\theta = \text{const}.$$

CP violation!

Wait again! Can write as:  $\mathscr{L}_{\theta} \sim \theta \partial_{\mu} K^{\mu}$ 

Boundary term!

# QED is not CP violating

QCD 0 0 Conclusion

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# Why QCD makes Life complicated...

QCD:  $\mathscr{L}_{QCD} = -\frac{1}{4}F^a_{\mu\nu}F^{\mu\nu a}$ 

Why these terms?

- Gauge invariant under SU(3)
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$$\mathscr{L}_{\theta} \sim \theta F^{a}_{\mu\nu} \tilde{F}^{\mu\nu a}$$
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 $\tilde{F}^{\mu\nu a} = \epsilon^{\mu\nu\rho\sigma} F^a_{\rho\sigma}$  $\theta = \text{const}.$ CP violation!

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QCD 0 0 Conclusion

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# Why QCD makes Life complicated...

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$$\theta = \text{const}.$$
$$CP \text{ violation!}$$

Wait again! Can write as:  $\mathscr{L}_{\theta} \sim \theta \partial_{\mu} K^{\mu}$ 

Boundary term! Instantons do not care

# QCD is CP violating!

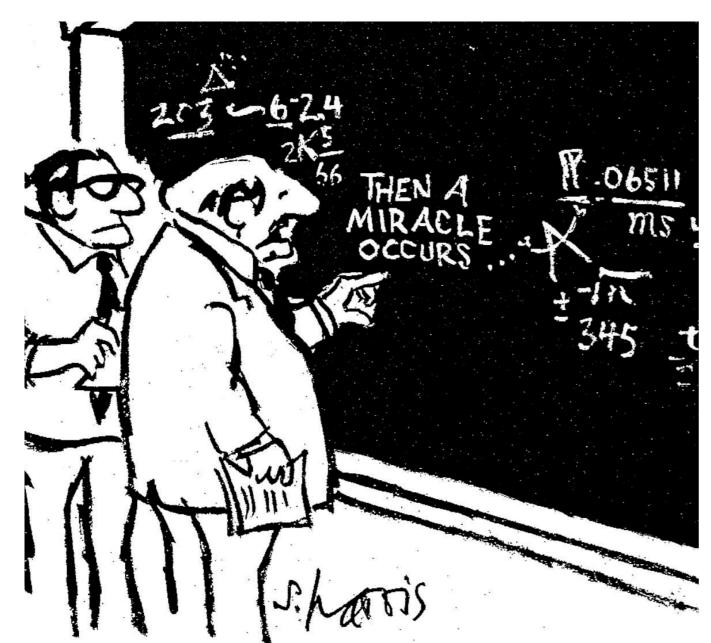
The Strong CP Problem

### Schwinger Model

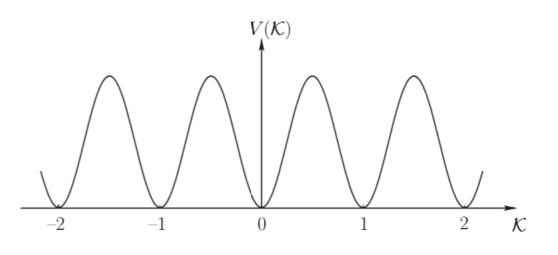
Conclusion



### Instantons = Tunneling Processes



"I THINK YOU SHOULD BE MORE EXPLICIT HERE IN STEP TWO." Lead to non-trivial QCD vacuum structure



Main point for now: "Generate correlation to boundary"



- Contributes to electric dipole moment of e.g. neutron
- $\Rightarrow$  Calculate contribution and compare to experiment to find  $\theta$ !

- We do not measure CP violation!  $\implies \theta < 10^{-10}$  or so
- Strong CP problem: Why is  $\theta$  so small?

The Strong CP Problem 0000

Schwinger Model

**OCD** 

Conclusion

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#### Higg Solution: Axion lived pseudoscalar boson. [The SCI @ indicates quar that this paper has been cited in more than 605 been publications.] and t-qua The Birth of Axions locat Fol tions Frank Wilczek disco Institute for Advanced Study Princeton, NJ 08540 was whic by fi The basic idea of axions occurred to me as a was result of my wife's ear infection. This was in thing the summer of 1975, when we were visiting ular Fermilab with a small baby daughter in tow. mass My wife was ill and confined to bed, and I had of th a rather difficult day trying to cope. Finally, stanc both wife and daughter were safely asleep and even it was a beautiful midwestern night with a gor-But geous clear sky, and I decided to take a long by C walk. been Turning with relief from the cares of the day, rame I decided to think about the Higgs sector. At be vi that time, what is now called the standard that model of particle physics, although less than cons three years old, was already established as far



Peccei-Quinn Mechanism:

- Propose new global axial symmetry, U(1)<sub>PQ</sub>
- New Higgs particle to spontaneously break U(1)<sub>PQ</sub>
- Axion *a* is the appearing Goldstone boson

$$\mathscr{L} = \frac{1}{2} \partial_{\mu} a \partial^{\mu} a - \frac{1}{4} F^{a}_{\mu\nu} F^{\mu\nu a} + (a + \theta) F^{a}_{\mu\nu} \tilde{F}^{\mu\nu a}$$

 $\theta$  now dynamical!

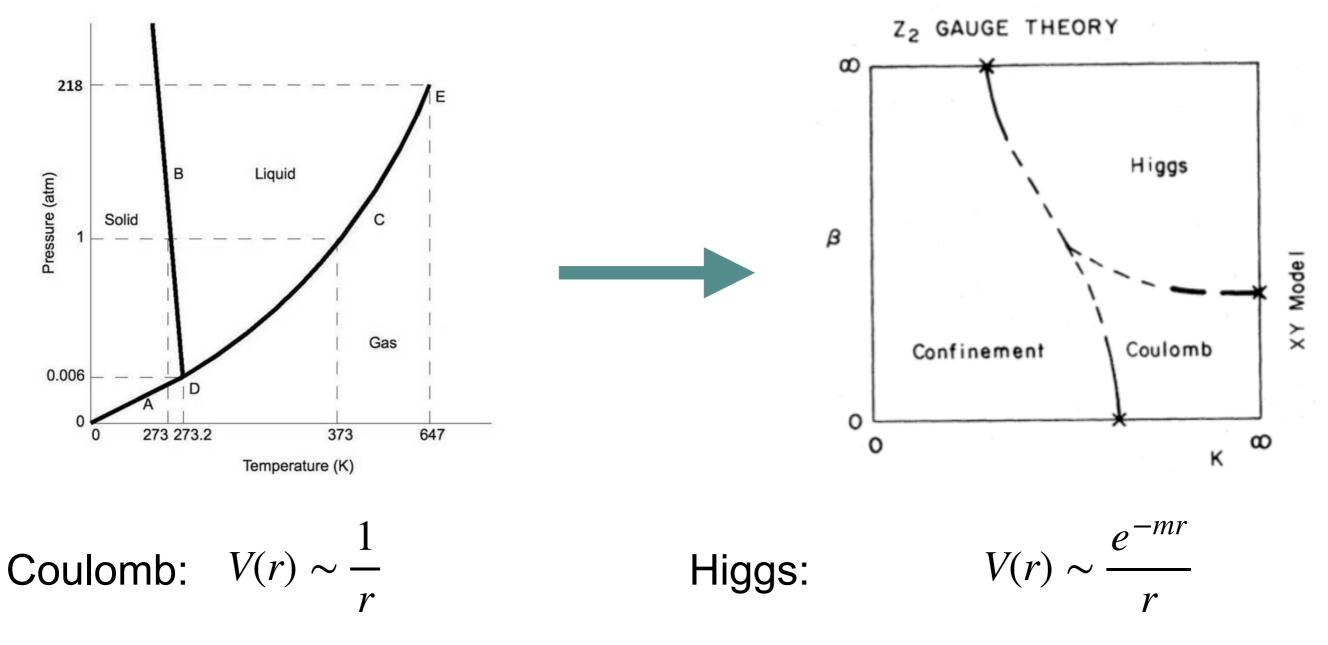
Vafa-Witten: Vacuum where  $\theta = 0$ 

QCD

Conclusion



## **Recap: Phases of Gauge Theories**



Confinement:  $V(r) \sim r$ 



QCD 0 0

Conclusion



# QED in 1+1 Dimensions

Maxwell's Theory: 
$$\mathscr{L}_{MW} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu}$$
  $\mu, \nu = 0,1$ 

Solution to Maxwells equation: E(x) = const.

- Constant electric field everywhere in space
- Can not be changed

 $\mathscr{L}$ 

Add  $\theta$ -term:

$$= -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \frac{1}{2}\theta \epsilon_{\mu\nu}F^{\mu\nu} , \qquad \theta = \text{const.}$$

- Total derivative
- Just redefines the constant value of the electric field

Main Point:  $\theta$  equivalent to constant electric field in the vacuum

QCD 0 0 Conclusion



# <u>Getting Rid of Global Electric Field</u>

Idea: Bring to Higgs phase

$$\mathscr{L}_{\text{Proca}} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \frac{m^2}{2}(A_{\mu} - \partial_{\mu}B)^2 + \frac{1}{2}\theta\epsilon_{\mu\nu}F^{\mu\nu}$$

Let's Hodge dualize just for fun:  $B \leftrightarrow a$ 

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \frac{1}{2}\partial_{\mu}a\partial^{\mu}a + \frac{1}{2}\Big(a+\theta\Big)\epsilon_{\mu\nu}F^{\mu\nu}$$

- Constant electric field no longer a solution, now dynamical!
- Unique vacuum:  $E = 0, a = -\theta$

# Axion making $\theta$ dynamical = Putting global electric field in Higgs Phase



- Generalizable to arbitrary kinetic term K(F),  $F = F_{\mu\nu} \epsilon^{\mu\nu}$
- Massive fermions give additional contribution to  $\theta$ , but do not spoil conclusions

The Strong CP Problem	Schwinger Model	QCD • •	Conclusion o	$\Delta_{p} \cdot \Delta_{g} \ge \frac{1}{2} t$
Back to QCD				
Maxwell in 1+	1 d	(	QCD in 3+1 d	
$\mathscr{L}_{\theta} \sim \theta \epsilon^{\mu \nu} F_{\mu \nu}$	$\leftarrow$	$\rightarrow \mathscr{L}_{\theta}$	$\sim \theta F^a_{\mu u} \tilde{F}^{\mu u a} \sim \theta \epsilon^{lpha eta \mu u} F_{lpha}$	βμν
$F_{\mu\nu} = \partial_{[\mu} A_{\nu]}$	$\leftarrow$	$ \qquad \qquad$	$_{\nu} = \partial_{[\alpha} C_{\beta\mu\nu]}$	
$\mathscr{L} \sim \frac{m^2}{2} (A_{\mu} - \partial_{\mu} E)$	$(\mathbf{S})^2$	$\rightarrow$ $\mathscr{L} \sim$	$\frac{m^2}{2}(C_{\beta\mu\nu}-\partial_{\beta}B_{\mu\nu})^2$	
$B \longleftrightarrow a$	$\leftarrow$	$\rightarrow B_{\mu\nu}$	$\longleftrightarrow a$	
Wait! QCD in strong coupling regime at low energies				
$\longrightarrow \mathcal{C} = \frac{1}{\rho_{\alpha}} \partial_{\alpha} \partial_{\mu} \nu_{F} \qquad \forall K(F) \qquad \longrightarrow \qquad ($				

$$\Longrightarrow \mathscr{L} = \frac{1}{2} \theta \epsilon^{\alpha \beta \mu \nu} F_{\alpha \beta \mu \nu} + K(F) \qquad \Longrightarrow \checkmark$$



# How to spoil axion solution?

- •Single axion cannot Higgs more than one 3-form field.
- $\implies$  New physics must provide independently-massless 3-form field
- Example: Gravity, Chern-Simons spin connection 3-form  $R\tilde{R} = \epsilon^{\alpha\beta\mu\nu}R^{i}_{j\alpha\beta}R^{j}_{i\mu\nu} = \frac{1}{3}\epsilon^{\alpha\beta\mu\nu}\partial_{[\alpha}G_{\beta\mu\nu]}$



# **Conclusion**

- Hodge duality is a nice tool
- $\theta$  equivalent to constant "electric" field in vacuum
- Axion making θ dynamical
   = Putting global "electric" field in Higgs phase
- New physics must provide independently-massless 3-form field to spoil axion solution

The Strong CP Problem

Schwinger Model

Conclusion



# Thank you for your attention!

