



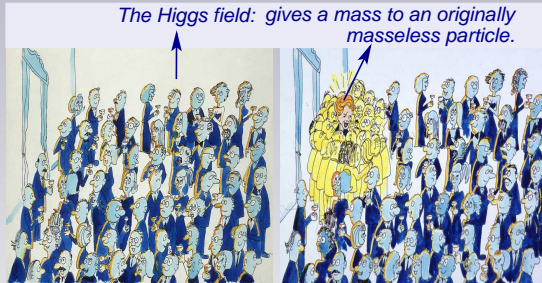
Sandra Horvat

Max-Planck-Institut für Physik, München

Standard Model Higgs searches at the LHC

3rd MPI Young Scientists Workshop
Ringberg Castle • October 27, 2004

The Higgs Mechanism...



...and the Higgs Boson

The Higgs Mechanism

The invariant Lagrangian \mathcal{L} of the $SU(2)_L \times U(1)$ electroweak theory contains **massless** gauge bosons (**W**, **B**) and fermions (ψ).

$$\begin{aligned} \mathcal{L} = & -\frac{1}{4} (\partial_\nu W_\mu^i - \partial_\mu W_\nu^i + g\epsilon^{ijk} W_\mu^j W_\nu^k) (\partial^\nu W^{i\mu} - \partial^\mu W^{i\nu} + g\epsilon^{ijk} W^{j\mu} W^{k\nu}) \\ & -\frac{1}{4} (\partial_\nu B_\mu - \partial_\mu B_\nu) (\partial^\nu B^\mu - \partial^\mu B^\nu) \\ & + i\bar{\psi}_R \gamma^\mu (\partial_\mu + i\frac{g'}{2} Y_R B_\mu) \psi_R + i\bar{\psi}_L \gamma^\mu (\partial_\mu + i\frac{g}{2} \tau_i W_\mu^i + i\frac{g'}{2} Y_L B_\mu) \psi_L. \end{aligned}$$

P.W.Higgs 1964; F.Englert and R.Brout 1964;

Introduction of a scalar field Φ provides for the **particle masses**.

$$\Phi = \begin{pmatrix} \phi^\dagger \\ \phi^0 \end{pmatrix} ; V(\Phi) = -\lambda v^2 |\Phi^\dagger \Phi| + \lambda (|\Phi^\dagger \Phi|)^2$$

$$\mathcal{L}_\Phi = (D^\mu \Phi)^\dagger (D_\mu \Phi) - g_f (\bar{\psi}_L \Phi \psi_R + \bar{\psi}_R \Phi^\dagger \psi_L) - V(\Phi)$$

The Higgs Mechanism

The invariant Lagrangian \mathcal{L} of the $SU(2)_L \times U(1)$ electroweak theory contains **massless** gauge bosons (**W**, **B**) and fermions (ψ).

$$\begin{aligned} \mathcal{L} = & -\frac{1}{4} (\partial_\nu W_\mu^i - \partial_\mu W_\nu^i + g\epsilon^{ijk} W_\mu^j W_\nu^k) (\partial^\nu W^{i\mu} - \partial^\mu W^{i\nu} + g\epsilon^{ijk} W^{j\mu} W^{k\nu}) \\ & -\frac{1}{4} (\partial_\nu B_\mu - \partial_\mu B_\nu) (\partial^\nu B^\mu - \partial^\mu B^\nu) \\ & + i\bar{\psi}_R \gamma^\mu (\partial_\mu + i\frac{g'}{2} Y_R B_\mu) \psi_R + i\bar{\psi}_L \gamma^\mu (\partial_\mu + i\frac{g}{2} \tau_i W_\mu^i + i\frac{g'}{2} Y_L B_\mu) \psi_L. \end{aligned}$$

P.W.Higgs 1964; F.Englert and R.Brout 1964;

Introduction of a scalar field Φ provides for the **particle masses**.

$$\Phi = \begin{pmatrix} \phi^\dagger \\ \phi^0 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v + H \end{pmatrix}; \quad V(\Phi) = -\lambda v^2 |\Phi^\dagger \Phi| + \lambda (|\Phi^\dagger \Phi|)^2$$

$$\mathcal{L}_\Phi = (D^\mu \Phi)^\dagger (D_\mu \Phi) - g_f (\bar{\psi}_L \Phi \psi_R + \bar{\psi}_R \Phi^\dagger \psi_L) - V(\Phi)$$

The Higgs Mechanism

The invariant Lagrangian \mathcal{L} of the $SU(2)_L \times U(1)$ electroweak theory contains **massless** gauge bosons (**W**, **B**) and fermions (ψ).

$$\begin{aligned} \mathcal{L} = & -\frac{1}{4} (\partial_\nu W_\mu^i - \partial_\mu W_\nu^i + g \epsilon^{ijk} W_\mu^j W_\nu^k) (\partial^\nu W^{i\mu} - \partial^\mu W^{i\nu} + g \epsilon^{ijk} W^{j\mu} W^{k\nu}) \\ & -\frac{1}{4} (\partial_\nu B_\mu - \partial_\mu B_\nu) (\partial^\nu B^\mu - \partial^\mu B^\nu) \\ & + i \bar{\psi}_R \gamma^\mu (\partial_\mu + i \frac{g'}{2} Y_R B_\mu) \psi_R + i \bar{\psi}_L \gamma^\mu (\partial_\mu + i \frac{g}{2} \tau_i W_\mu^i + i \frac{g'}{2} Y_L B_\mu) \psi_L. \end{aligned}$$

P.W.Higgs 1964; F.Englert and R.Brout 1964;

Introduction of a scalar field Φ provides for the **particle masses**.

$$\Phi = \begin{pmatrix} \phi^\dagger \\ \phi^0 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v + H \end{pmatrix}; \quad V(\Phi) = -\lambda v^2 |\Phi^\dagger \Phi| + \lambda (|\Phi^\dagger \Phi|)^2$$

$$\mathcal{L}_\Phi = (D^\mu \Phi)^\dagger (D_\mu \Phi) - g_f (\bar{\psi}_L \Phi \psi_R + \bar{\psi}_R \Phi^\dagger \psi_L) - V(\Phi)$$



interaction with gauge bosons

$$m_{W^\pm} = \frac{g v}{2}, \quad m_Z = \frac{v \sqrt{(g^2 + g'^2)}}{2}$$

The Higgs Mechanism

The invariant Lagrangian \mathcal{L} of the $SU(2)_L \times U(1)$ electroweak theory contains **massless** gauge bosons (**W**, **B**) and fermions (**ψ**).

$$\begin{aligned} \mathcal{L} = & -\frac{1}{4}(\partial_\nu W_\mu^i - \partial_\mu W_\nu^i + g\epsilon^{ijk}W_\mu^jW_\nu^k) (\partial^\nu W^{i\mu} - \partial^\mu W^{i\nu} + g\epsilon^{ijk}W^{j\mu}W^{k\nu}) \\ & -\frac{1}{4}(\partial_\nu B_\mu - \partial_\mu B_\nu) (\partial^\nu B^\mu - \partial^\mu B^\nu) \\ & + i\bar{\psi}_R\gamma^\mu (\partial_\mu + i\frac{g'}{2}Y_R B_\mu) \psi_R + i\bar{\psi}_L\gamma^\mu (\partial_\mu + i\frac{g}{2}\tau_i W_\mu^i + i\frac{g'}{2}Y_L B_\mu) \psi_L. \end{aligned}$$

P.W.Higgs 1964; F.Englert and R.Brout 1964;

Introduction of a scalar field Φ provides for the **particle masses**.

$$\Phi = \begin{pmatrix} \phi^\dagger \\ \phi^0 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v + H \end{pmatrix}; \quad V(\Phi) = -\lambda v^2 |\Phi^\dagger \Phi| + \lambda (|\Phi^\dagger \Phi|)^2$$

$$\mathcal{L}_\Phi = (D^\mu \Phi)^\dagger (D_\mu \Phi) - g_f (\bar{\psi}_L \Phi \psi_R + \bar{\psi}_R \Phi^\dagger \psi_L) - V(\Phi)$$



interaction with fermions

$$m_f = \frac{g_f v}{\sqrt{2}}$$

The Higgs Mechanism

The invariant Lagrangian \mathcal{L} of the $SU(2)_L \times U(1)$ electroweak theory contains **massless** gauge bosons (**W**, **B**) and fermions (**ψ**).

$$\mathcal{L} = -\frac{1}{4}(\partial_\nu W_\mu^i - \partial_\mu W_\nu^i + g\epsilon^{ijk}W_\mu^jW_\nu^k)(\partial^\nu W^{i\mu} - \partial^\mu W^{i\nu} + g\epsilon^{ijk}W^{j\mu}W^{k\nu}) - \frac{1}{4}(\partial_\nu B_\mu - \partial_\mu B_\nu)(\partial^\nu B^\mu - \partial^\mu B^\nu) + i\bar{\psi}_R\gamma^\mu(\partial_\mu + i\frac{g'}{2}Y_R B_\mu)\psi_R + i\bar{\psi}_L\gamma^\mu(\partial_\mu + i\frac{g}{2}\tau_i W_\mu^i + i\frac{g'}{2}Y_L B_\mu)\psi_L.$$

P.W.Higgs 1964; F.Englert and R.Brout 1964;

Introduction of a scalar field Φ provides for the **particle masses**.

$$\Phi = \begin{pmatrix} \phi^\dagger \\ \phi^0 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v + H \end{pmatrix}; \quad V(\Phi) = -\lambda v^2 |\Phi^\dagger \Phi| + \lambda (|\Phi^\dagger \Phi|)^2$$

$$\mathcal{L}_\Phi = (D^\mu \Phi)^\dagger (D_\mu \Phi) - g_f (\bar{\psi}_L \Phi \psi_R + \bar{\psi}_R \Phi^\dagger \psi_L) - V(\Phi)$$



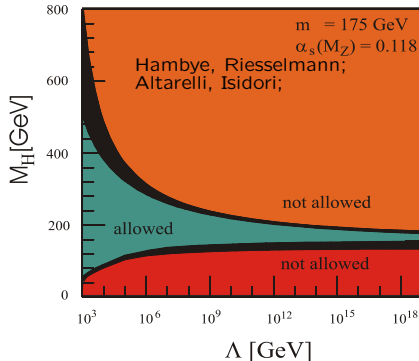
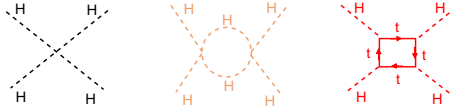
self-interaction

Existence of a spin -0 particle, the **Higgs boson** with a mass $m_H = \sqrt{2\lambda}v$.

The Higgs Boson Mass

Not specified explicitly, but...

Bounds derived from the evolution of the self-coupling λ with energy:



Upper bound:

SM valid up to the scale Λ

$$\lambda(\Lambda) < \infty$$

Lower bound:

vacuum stability
(finite minimum for $V(\Phi)$)

$$\lambda(\Lambda) > 0$$

The Higgs boson discovery can provide the constraints on the validity of the Standard Model.

Current experimental status



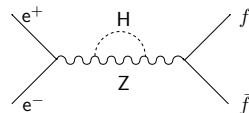
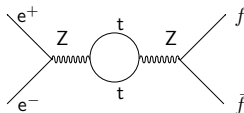
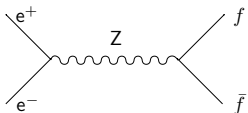
Experimental Data

Collider	Running Period	Beam	\sqrt{s} (GeV)
LEP1	1989-1995	e^+e^-	91
SLC	1991-1998	e^+e^-	91
LEP2	1995-2000	e^+e^-	~ 209
TEVATRON	2001-	$p\bar{p}$	1 800
LHC	2007-	pp	14 000
LINEAR COLLIDER?	???	e^+e^-	500

Experimental Data

Collider	Running Period	Beam	\sqrt{s} (GeV)
LEP1	1989-1995	e^+e^-	91
SLC	1991-1998	e^+e^-	91
LEP2	1995-2000	e^+e^-	~ 209
TEVATRON	2001-	$p\bar{p}$	1 800
LHC	2007-	pp	14 000
LINEAR COLLIDER?	???	e^+e^-	500

indirect tests, setting bounds on the Higgs mass
 from precision measurement of electroweak observables
 taking into account the radiative corrections

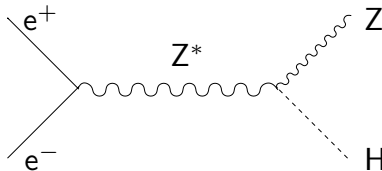


$$\propto \text{Log} \frac{m_H^2}{m_z^2}$$

Experimental Data

Collider	Running Period	Beam	\sqrt{s} (GeV)
LEP1	1989-1995	e^+e^-	91
SLC	1991-1998	e^+e^-	91
LEP2	1995-2000	e^+e^-	~ 209
TEVATRON	2001-	$p\bar{p}$	1 800
LHC	2007-	pp	14 000
LINEAR COLLIDER?	???	e^+e^-	500

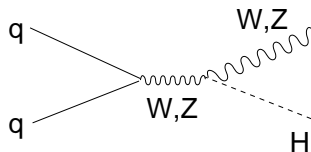
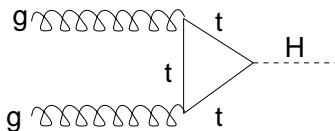
direct search, Higgs produced by Higgs-Strahlung;
large coupling to Z, all other processes negligible



Experimental Data

Collider	Running Period	Beam	\sqrt{s} (GeV)
LEP1	1989-1995	e^+e^-	91
SLC	1991-1998	e^+e^-	91
LEP2	1995-2000	e^+e^-	~ 209
TEVATRON	2001-	$p\bar{p}$	1 800
LHC	2007-	pp	14 000
LINEAR COLLIDER?	???	e^+e^-	500

direct searches; Higgs produced by
gluon fusion and Higgs-Strahlung



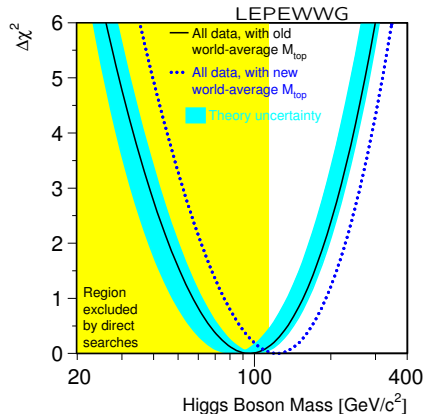
Precision measurements of electroweak observables:

- Z boson lineshape
- forward-backward asymmetries
- W boson mass

With the recent world average value
 $m_{top} = (178.0 \pm 4.3)$ GeV (April 2004):

$$m_H = 117_{-45}^{+67} \text{ GeV}$$

$m_H < 251$ GeV, 95% confidence level

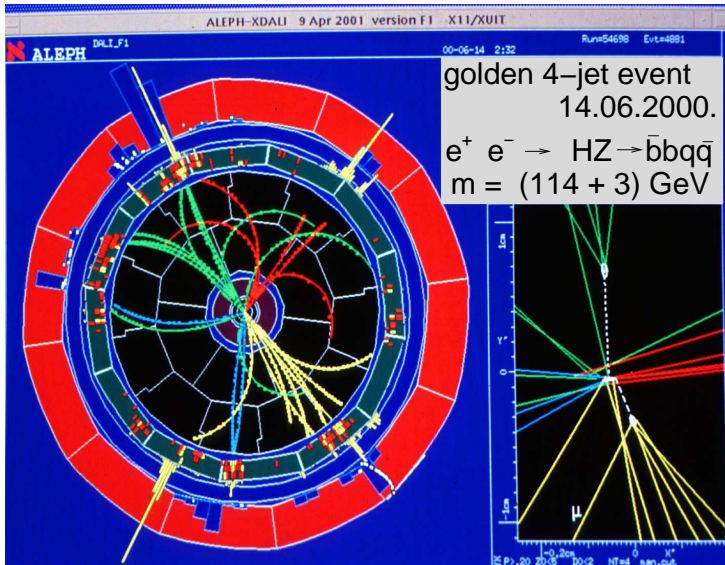


Direct searches

in the production channel $e^+e^- \rightarrow Z^* \rightarrow HZ$ give the lower bound

$$m_H > 114.4 \text{ GeV}$$

Observed Higgs event?



All other channels in all experiments are compatible with background!

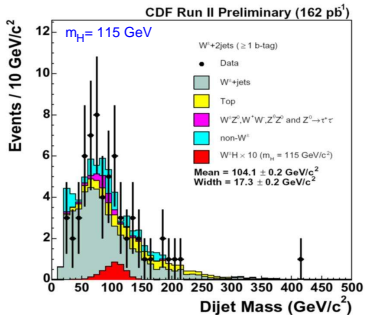
Higgs Search at the Tevatron

- collected integrated luminosity of $\sim 0.4 \text{ fb}^{-1}$,
 $\sim 0.2 \text{ fb}^{-1}$ analysed so far
- goal: $4\text{-}9 \text{ fb}^{-1}$ by the end of 2009
- still an early stage of data taking,
emphasis on understanding the detector performance
- the Higgs search has begun

Main signatures for the Standard Model Higgs boson:

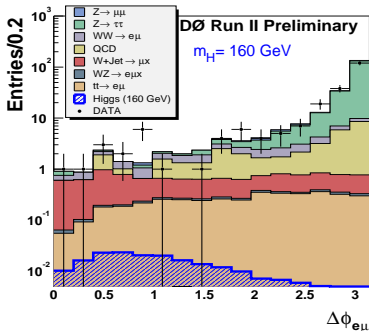
- For $m_H < 130 \text{ GeV}$: $q\bar{q} \rightarrow (W/Z)H \rightarrow (W/Z)b\bar{b}$
- For $m_H > 130 \text{ GeV}$: $gg \rightarrow H \rightarrow WW^* \rightarrow 2l2\nu$

Recent Tevatron Results (July 2004)



$WH \rightarrow Wb\bar{b}$

- leptonic W-decay
 - 2 jets required, ≤ 1 b-tagged
- $\sigma(Wb\bar{b}) < 20.3 \text{ pb}$
 $\sigma(WH) \times B(H \rightarrow b\bar{b}) < 12.4 \text{ pb}$



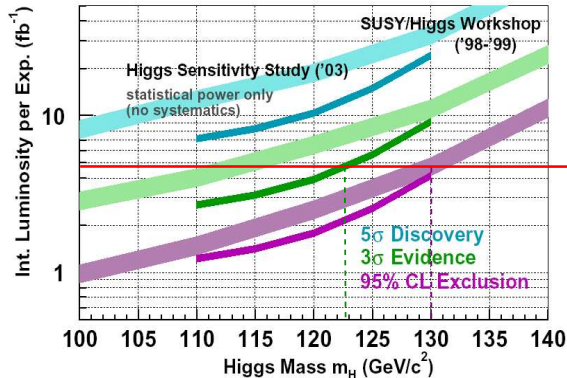
$H \rightarrow WW^* \rightarrow 2l2\nu$

- 2 high- p_T leptons
 - large missing energy
 - spin correlation
(charged leptons are collinear)
- $\sigma \times B(H \rightarrow WW) < 5.7 \text{ pb}$

Tevatron Discovery Potential

Higgs discovery in a single channel is not possible.

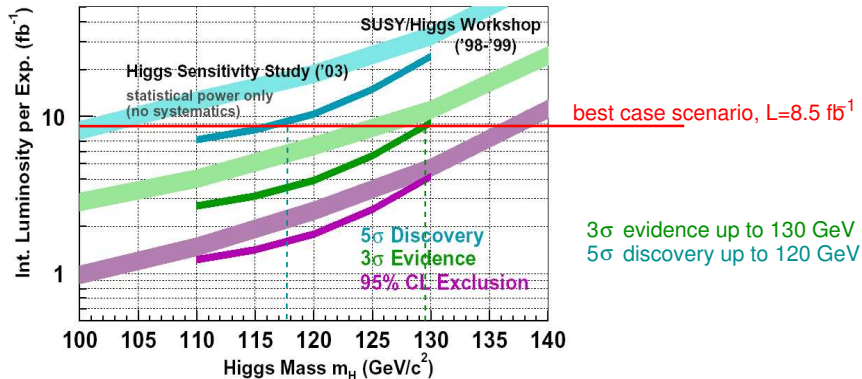
Combined CDF/D0 discovery reach up to 2009:



Tevatron Discovery Potential

Higgs discovery in a single channel is not possible.

Combined CDF/D0 discovery reach up to 2009:



Large Hadron Collider at CERN

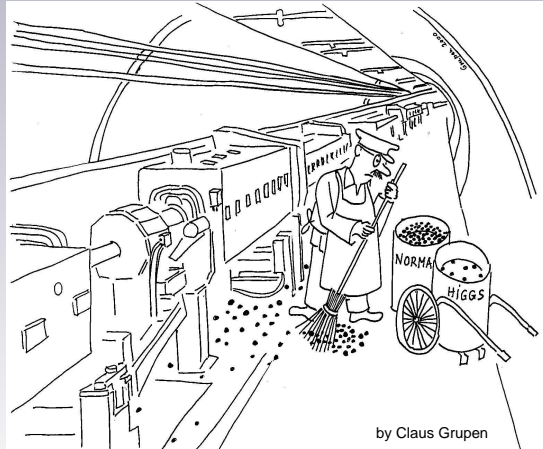
first 3 years at low luminosity: $L = 30 \text{ fb}^{-1}$
high luminosity afterwards: $L = 100 \text{ fb}^{-1} / \text{year}$

ATLAS

4.3 km

CMS

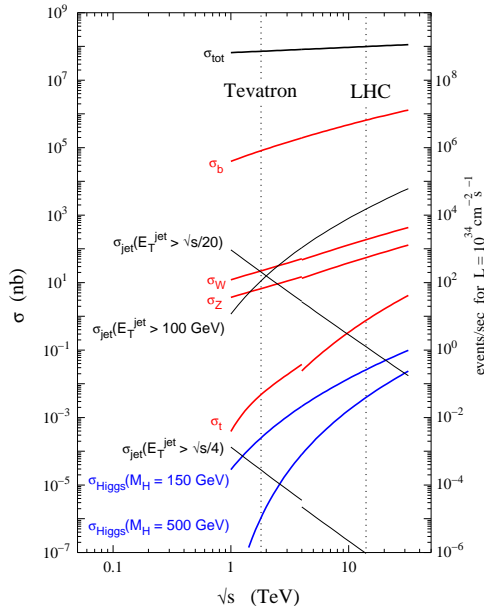
Higgs Signatures at the LHC



by Claus Grupen

Overall LHC Rates

proton – (anti)proton cross sections



The **Higgs signal** hidden behind the large **background** contributions.

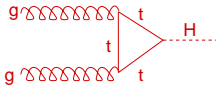


Challenging detector performance needed:

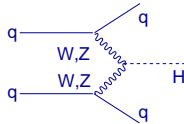
- powerful trigger
- high granularity
- radiation hardness
- high particle detection efficiency and resolution

Higgs Production Mechanisms

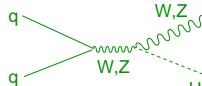
gluon–gluon fusion



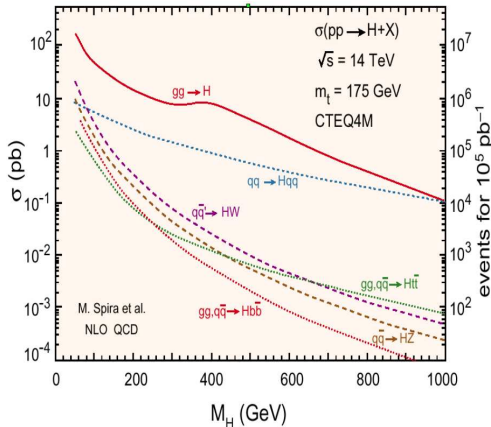
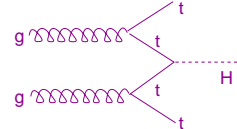
vector boson fusion, VBF



Higgs–Strahlung



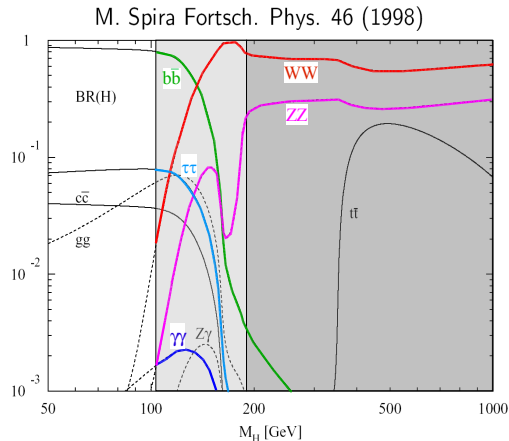
associated production



10 000 - 1 000 000
Higgs Bosons
produced per year
at high luminosity

Decay Channels

Branching ratios are completely determined by the Higgs mass.



Low mass $m_H < 2m_Z$:

- The most difficult region, combining several channels:
 $H \rightarrow b\bar{b}$, $H \rightarrow \gamma\gamma$,
 $H \rightarrow \tau^+\tau^-$ (via VBF),
 $H \rightarrow ZZ^* \rightarrow 4l$,
 $H \rightarrow WW^* \rightarrow 2l2\nu$ (via VBF)

Large mass $m_H > 180$ GeV:

- the gold-plated channel
 $H \rightarrow ZZ \rightarrow 4l$
- supplementary channel at very high masses (>800 GeV)
 $H \rightarrow WW \rightarrow l\nu jj$

Hadronic final states dominate for all values of m_H ,
but are difficult to distinguish from large QCD background.

We rather look for final states with **leptons, photons, missing energy**.

LHC Detectors

Ladbrokes

(the biggest betting company in the world):

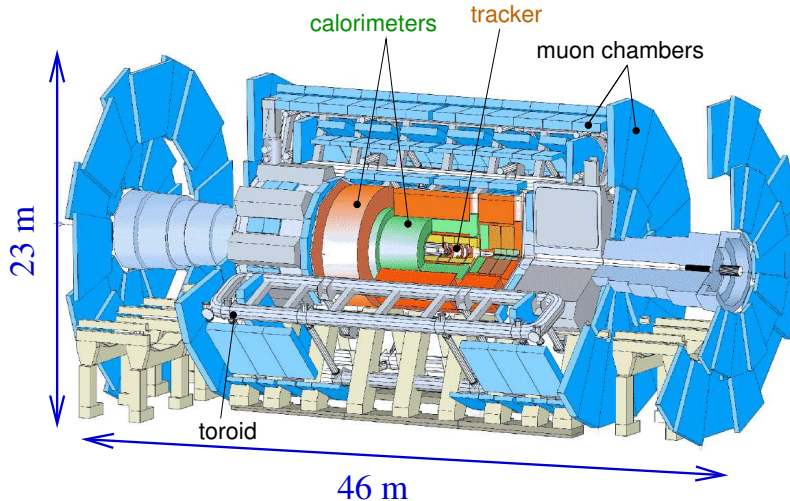
As of August 2004

accepts the bets on finding the Higgs boson at LHC before 2010,
the odds are set to **6 to 1** .

ATLAS (A Toroidal LHC ApparatuS)

Design optimized for the Higgs discovery in a wide mass range:

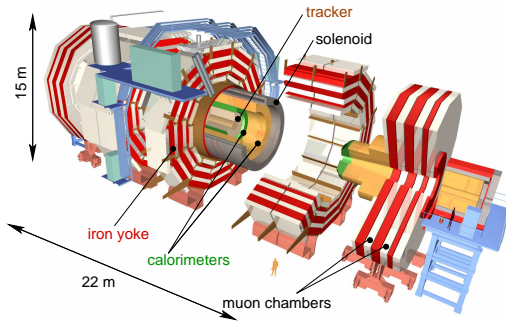
- large solid angle, efficient and precise particle detection, good calorimetry and missing energy measurement



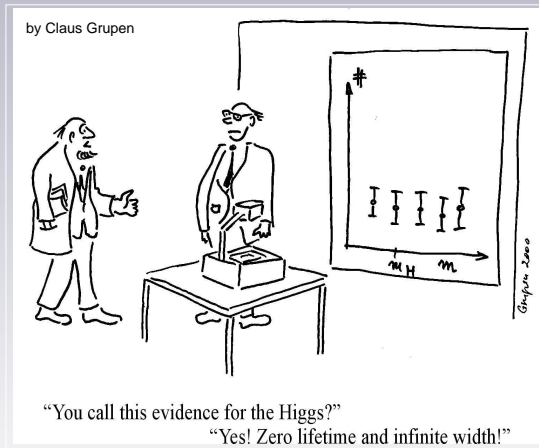
CMS (Compact Muon Solenoid)

Design optimized for the Higgs discovery in a wide mass range:

- large solid angle, efficient and precise particle detection, good calorimetry and missing energy measurement

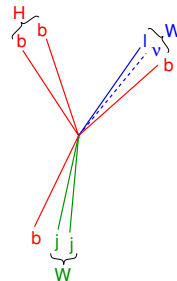


Discovery Potential

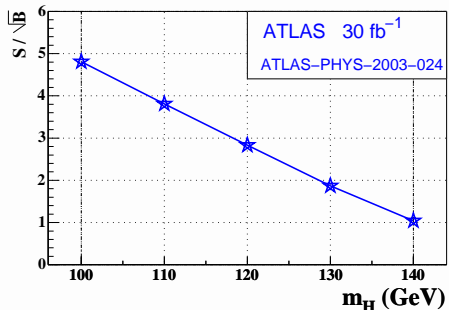
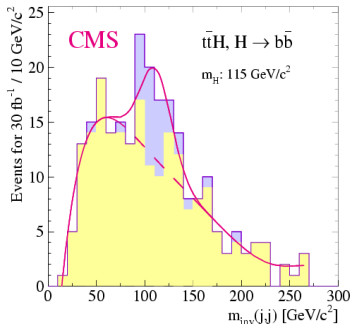


Low-mass Higgs: $t\bar{t}H, H \rightarrow b\bar{b}$

- identification via top-quark tagging
 - 1.) $t \rightarrow Wb \rightarrow l\nu b$: **lepton used for triggering**
 - 2.) $t \rightarrow Wb \rightarrow jjb$
- selecting final states with 4 b-tagged jets, **precise vertexing needed**
- need for a very good understanding of **background (ttbb, ttjj, Ztt)**



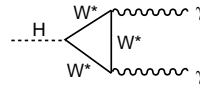
Signal significance (S/\sqrt{B}) at 30 fb^{-1} : 3.4σ for $m_H = 115 \text{ GeV}$



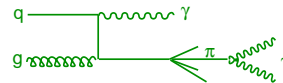
$B(H \rightarrow \gamma\gamma) = 10^{-3} B(H \rightarrow b\bar{b})$ only, but $H \rightarrow \gamma\gamma$ is a clean channel.

Main background processes:

- reducible: $qg \rightarrow q\gamma \rightarrow q\gamma\gamma$
 need excellent γ /jet separation,
 (in particular γ/π^0),
high granularity detectors
- irreducible: $gg \rightarrow \gamma\gamma, q\bar{q} \rightarrow \gamma\gamma$
 subtraction of a smooth
 continuous background,
excellent calorimetry needed
 for a narrow $m_{\gamma\gamma}$ -resonance



reducible background



irreducible background

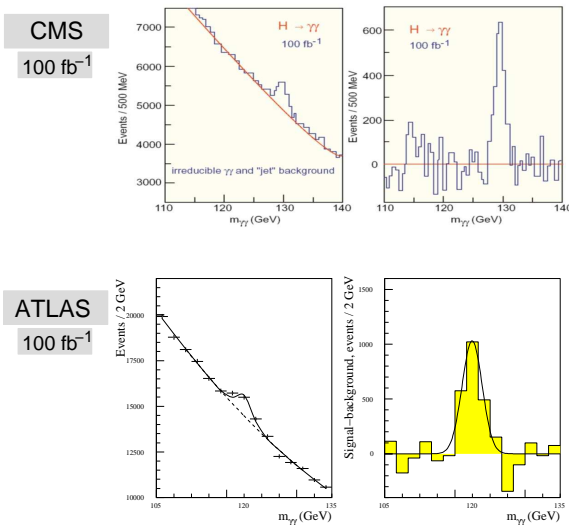


Additional improvement:

- associated Higgs production: WH, ZH and $t\bar{t}H$,
 process triggered by a lepton from W, Z or t

Signal significance (S/\sqrt{B}) at 100 fb^{-1} : $3-4\sigma$ for $m_H < 150 \text{ GeV}$

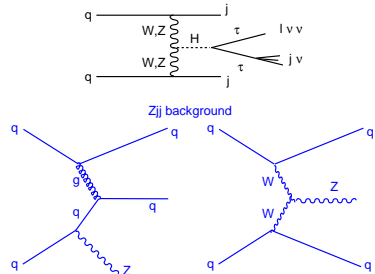
CMS has a 10% better calorimeter resolution than ATLAS.



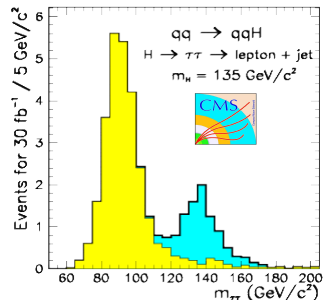
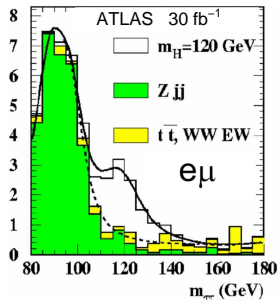
Low-mass Higgs:

$$qqH, H \rightarrow \tau^+\tau^-$$

- hadronic jets in forward-backward region, **forward jet-tagging, central jet veto**
- leptonic or hadronic decay in τ -direction, **missing p_T measurement, lepton identification**
- background: Zjj , $t\bar{t}$ +jets, $WWjj$



Signal significance (S/\sqrt{B}) at 40 fb^{-1} : $>5\sigma$ for $m_H < 140 \text{ GeV}$

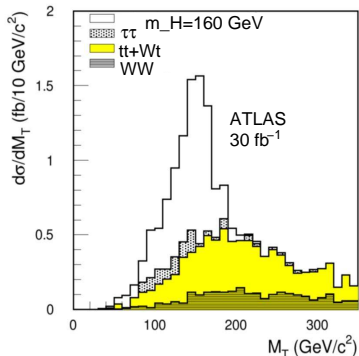


Low-mass Higgs: $(qq)H, H \rightarrow WW^{(*)} \rightarrow l\nu l\nu$

This is the main decay channel in the region around $m_H=170$ GeV. A good missing energy measurement is needed.

Main backgrounds:

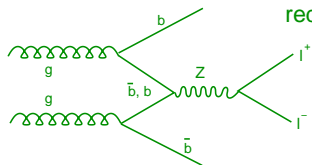
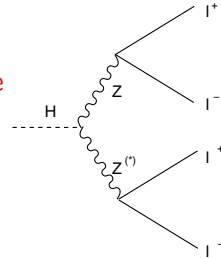
- $WW, WbW\bar{b}$ - suppressed by WW spin correlations in the signal
- Drell-Yan $ee, \mu\mu$ - lepton cuts ($p_T^{miss} > 30$ GeV, $m < 75$ GeV)
- $t\bar{t}, \tau\tau$ - b-jet and τ veto in the central region



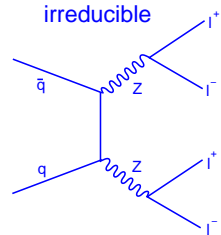
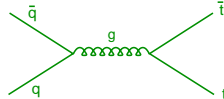
Signal significance (S/\sqrt{B})

of 5σ reachable at 10 fb^{-1}
for $m_H \approx 170$ GeV.

- relatively clean signature
one pair with m_{2l} -peak at the Z-resonance
- reducible background: $Zb\bar{b}, t\bar{t}$
lepton isolation, small impact parameter
- irreducible background: $ZZ \rightarrow 4l$

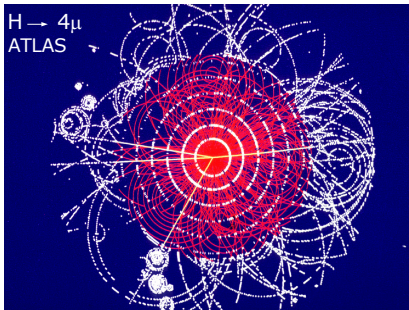


reducible

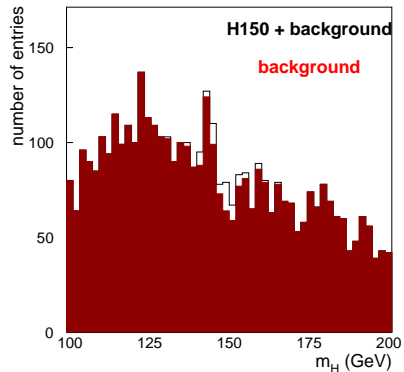


irreducible

Particularly clean signature from muons:

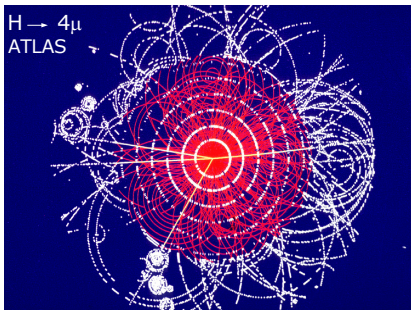


- no cuts applied on the muon tracks
- $m_{2\mu} = (m_Z \pm 2\Gamma_Z)$
- selecting isolated leptons
- cut on the impact parameter

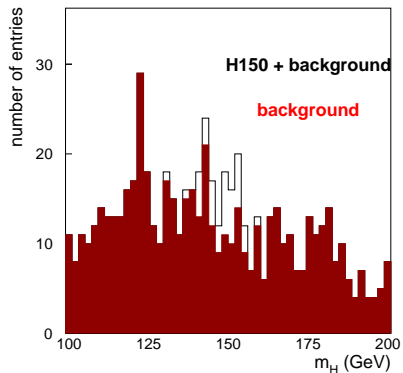


Signal significance for $H \rightarrow 4l$ at 30 fb^{-1} : **3-5 σ** for $m_H < 200 \text{ GeV}$.

Particularly clean signature from muons:

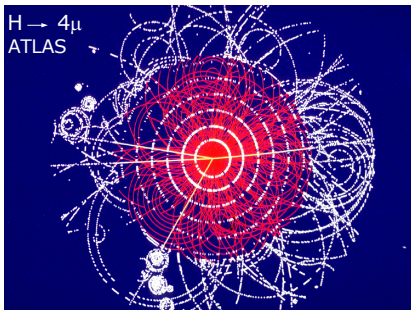


- no cuts applied on the muon tracks
- $m_{2\mu} = (m_Z \pm 2\Gamma_Z)$
- selecting isolated leptons
- cut on the impact parameter

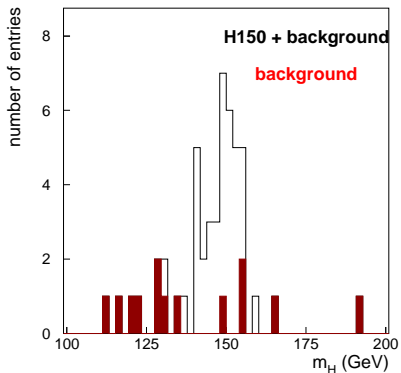


Signal significance for $H \rightarrow 4l$ at 30 fb^{-1} : **3-5 σ** for $m_H < 200 \text{ GeV}$.

Particularly clean signature from muons:

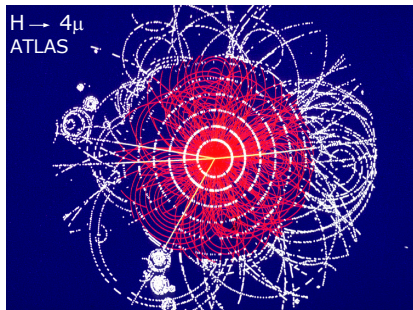


- no cuts applied on the muon tracks
- $m_{2\mu} = (m_Z \pm 2\Gamma_Z)$
- selecting isolated leptons
- cut on the impact parameter

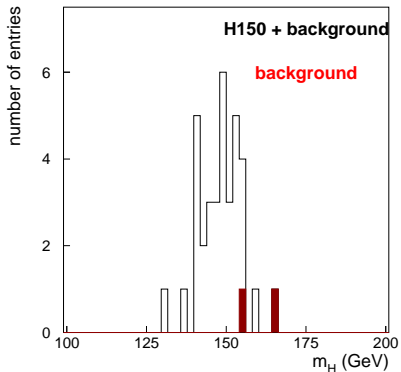


Signal significance for $H \rightarrow 4l$ at 30 fb^{-1} : **3-5 σ** for $m_H < 200 \text{ GeV}$.

Particularly clean signature from muons:



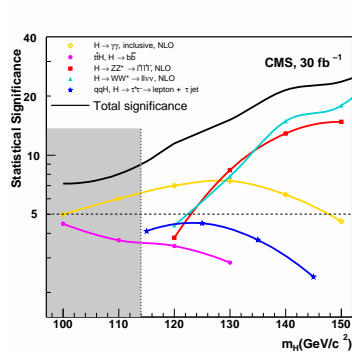
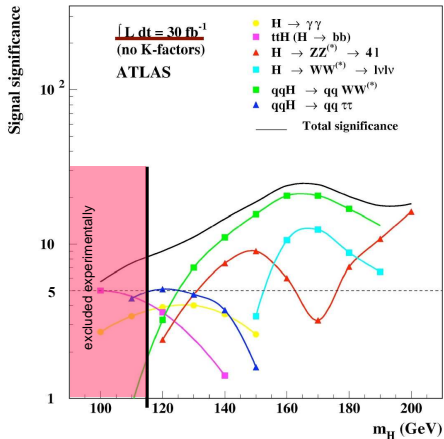
- no cuts applied on the muon tracks
- $m_{2\mu} = (m_Z \pm 2\Gamma_Z)$
- selecting isolated leptons
- cut on the impact parameter



Signal significance for $H \rightarrow 4l$ at 30 fb^{-1} : **3-5 σ** for $m_H < 200 \text{ GeV}$.

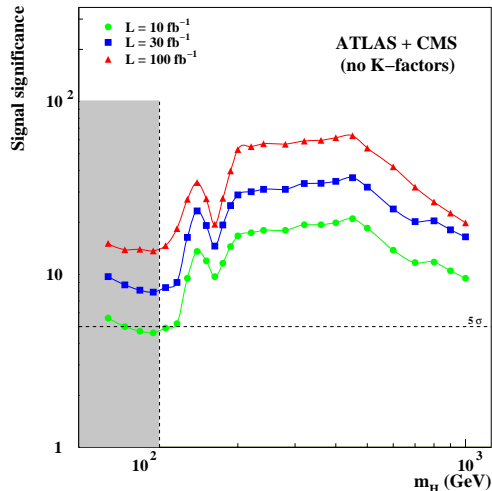
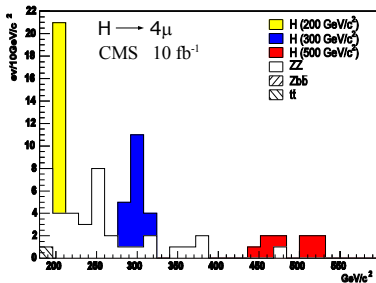
Discovery potential for a low Higgs mass

- combination of all channels is needed for the 5σ discovery in the first three years of running LHC

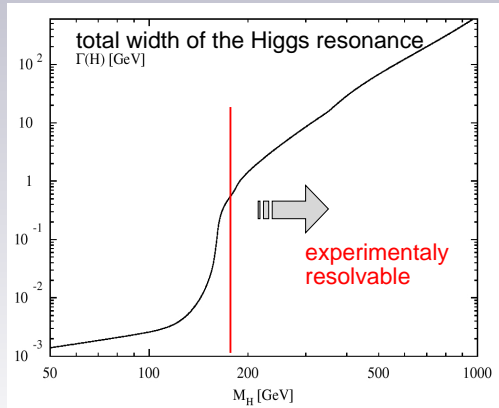


Discovery potential at the LHC

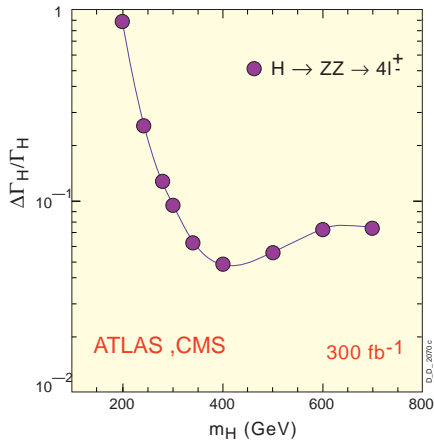
- $H \rightarrow ZZ \rightarrow 4l$ gives a clean signature for the region $m_H > 200$ GeV
- **complete coverage** of the theoretically allowed mass range m_H ;
 5σ discovery possible already **in the first year**



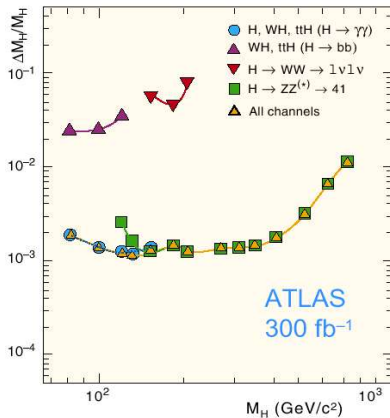
Precision Measurements of Higgs Properties



Mass and Width

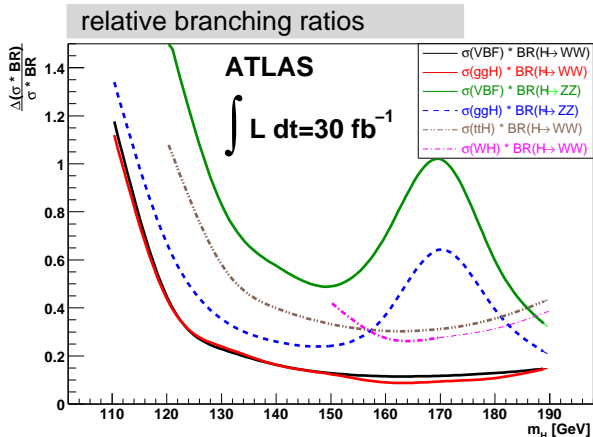
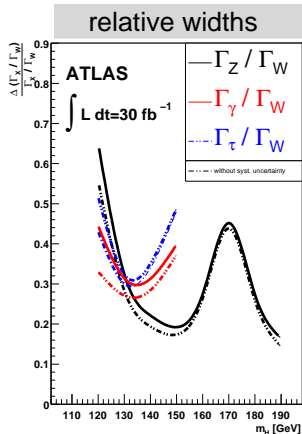


Width measurement:
6% for $m_H > 300$ GeV

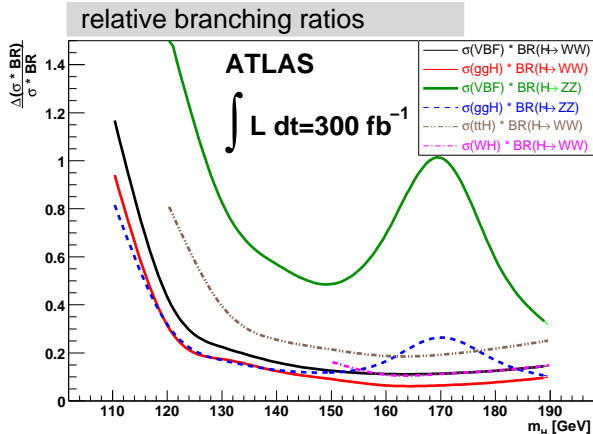
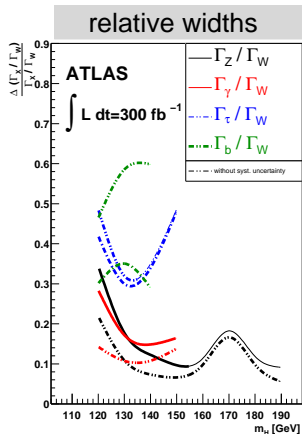


Mass measurement:
<1% for all m_H

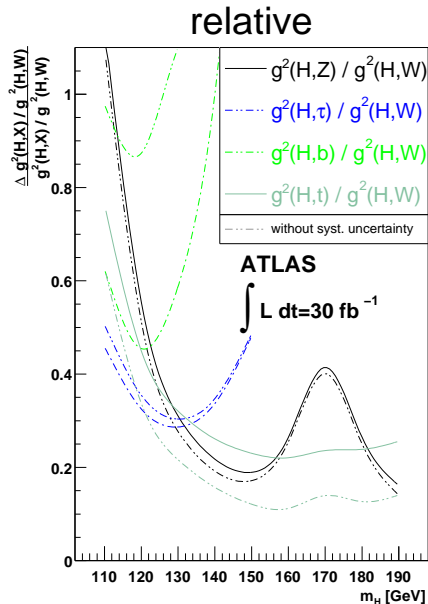
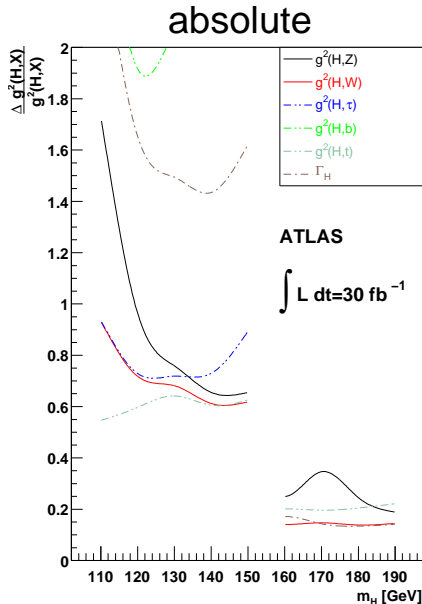
Relative widths and branching ratios



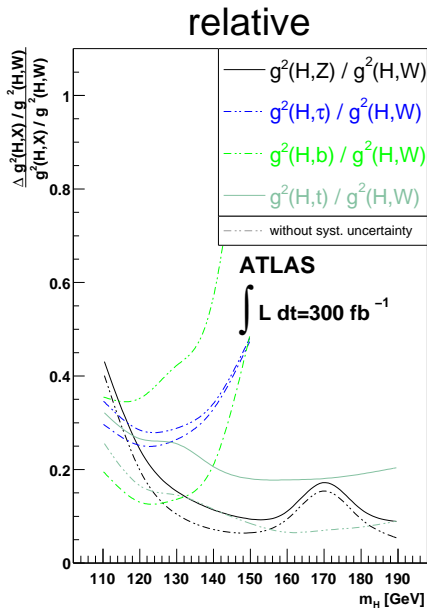
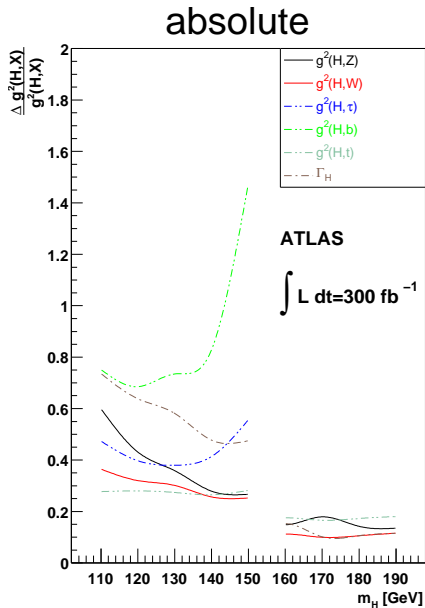
Relative widths and branching ratios



Couplings to fermions and gauge bosons



Couplings to fermions and gauge bosons



- The first pp-collisions at the LHC are expected **end of 2007**.
- The LHC detectors (ATLAS, CMS) are designed for the **Higgs search** covering the full mass spectrum.
- The 5σ -discovery limit can be reached in the whole mass range already after one year.
(However, a time will be needed to understand the detectors.)
- Physics studies are currently **concentrated on** the more difficult region of **lower Higgs masses**, which is favoured by the theoretical and experimental observations.
- **If** the Standard Model Higgs boson **exists**, **it will be seen** at the LHC.

Searching Beyond the Standard Model...

