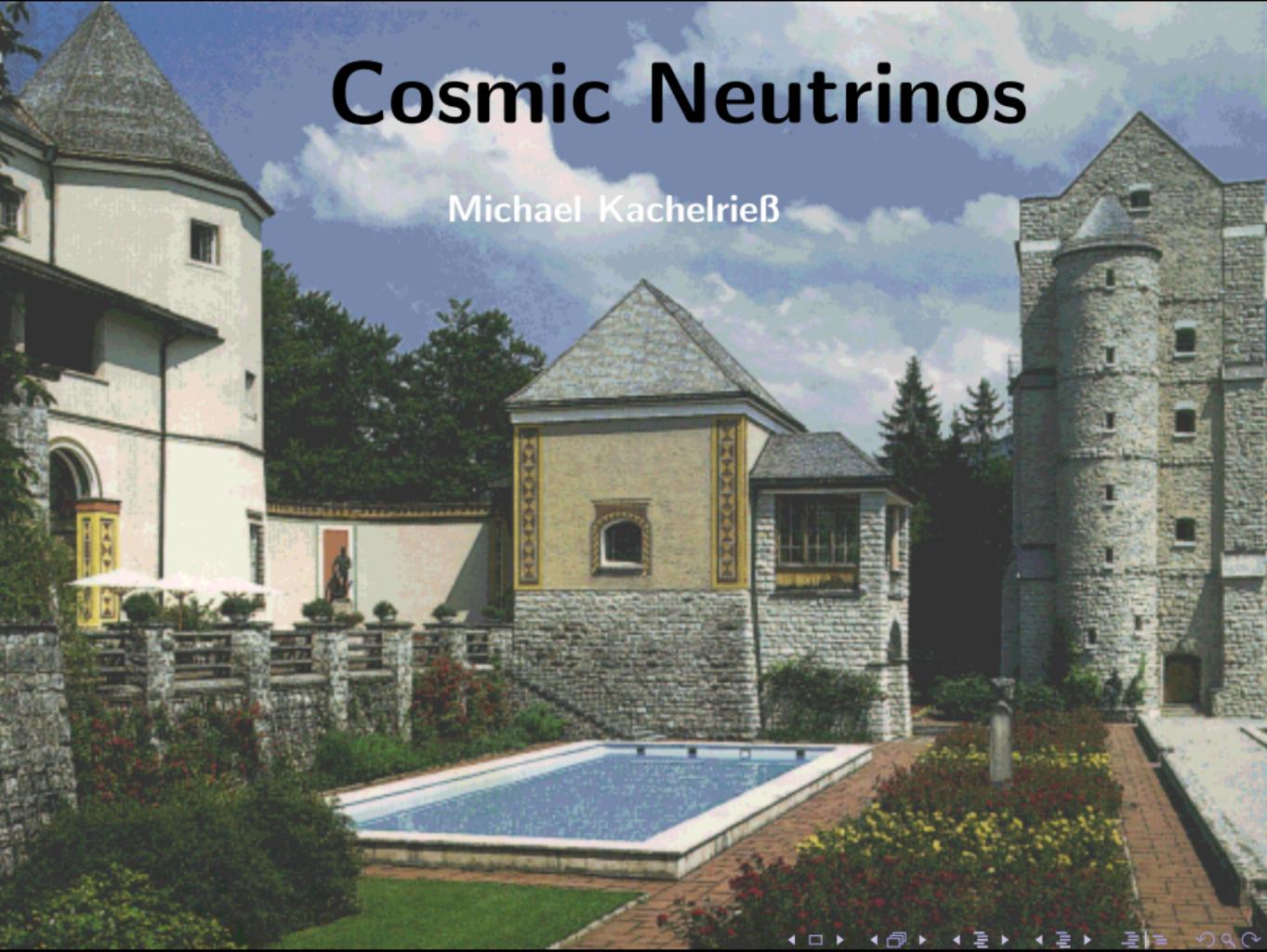
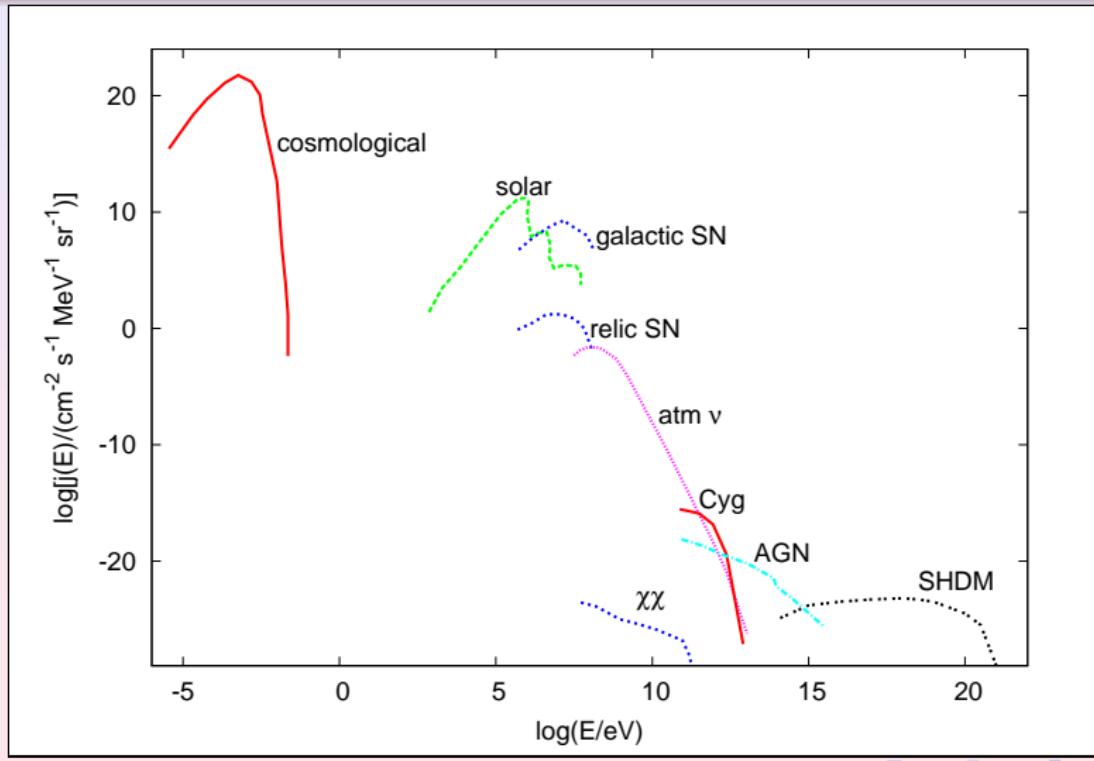


Cosmic Neutrinos

Michael Kachelrieß



Neutrino opportunities:



Neutrinos: what we do know

- neutrino oscillations are solution to
 - solar neutrino problem
 - atmospheric neutrino anomaly

⇒ neutrinos have mass and mix

$$\begin{bmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{bmatrix} = U_{\alpha i} \begin{bmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{bmatrix}$$

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$$U = \underbrace{\begin{bmatrix} 1 & 0 & 0 \\ 0 & c_a & s_a \\ 0 & -s_a & c_a \end{bmatrix}}_{\text{atm. osc.}} \underbrace{\begin{bmatrix} c_x & 0 & s_x e^{-i\delta} \\ 0 & 1 & 0 \\ -s_x e^{i\delta} & 0 & c_x \end{bmatrix}}_{\theta_x \equiv \theta_{13}} \underbrace{\begin{bmatrix} c_s & s_s & 0 \\ -s_s & c_s & 0 \\ 0 & 0 & 1 \end{bmatrix}}_{\text{solar osc.}} \underbrace{\overbrace{\quad}^{\text{diag}(1, \phi_2, \phi_3)}}_{\text{Majorana phases}}$$

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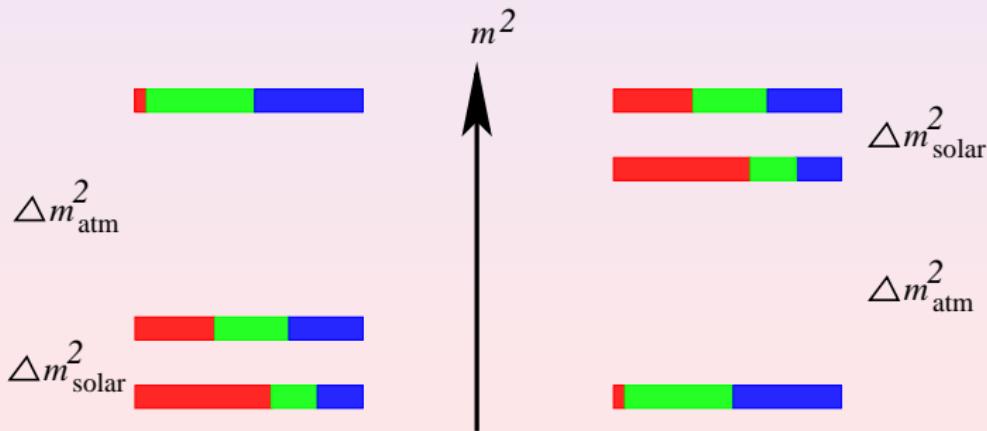
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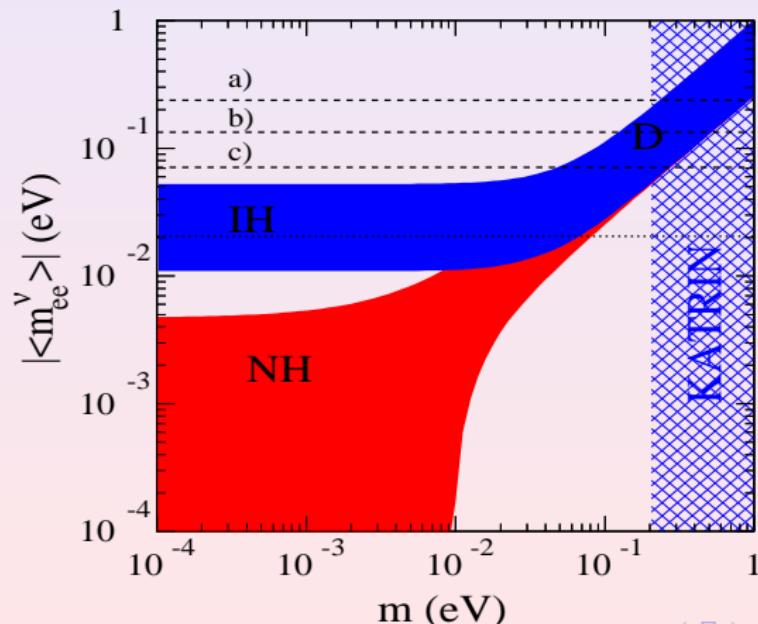
Neutrinos: what we do not know

- Dirac or Majorana particles?
- absolute scale of neutrino masses
- CP violation
- value of θ_{13}
- hierarchy of masses: normal or inverted? (ν_e , ν_μ , ν_τ)



Impact of mass hierarchy:

- constraint for **models** of neutrino mass
- neutrinoless double-beta decay:**



Neutrinos: what we do not know

only 2 confirmed neutrino sources

- the Sun
- SN 1987A

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- complete understanding of SN dynamics

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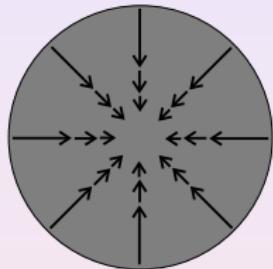
- complete understanding of SN dynamics

open: sources for HE neutrinos

- Galaxy, AGN's, ...
- decay/annihilations of cosmic relics

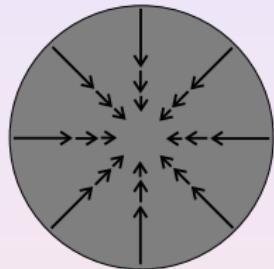
Core collapse supernovae:

- gravitational core collapse:

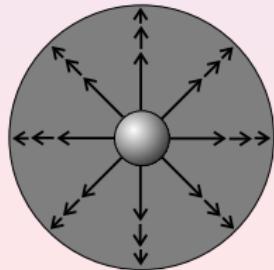


Core collapse supernovae:

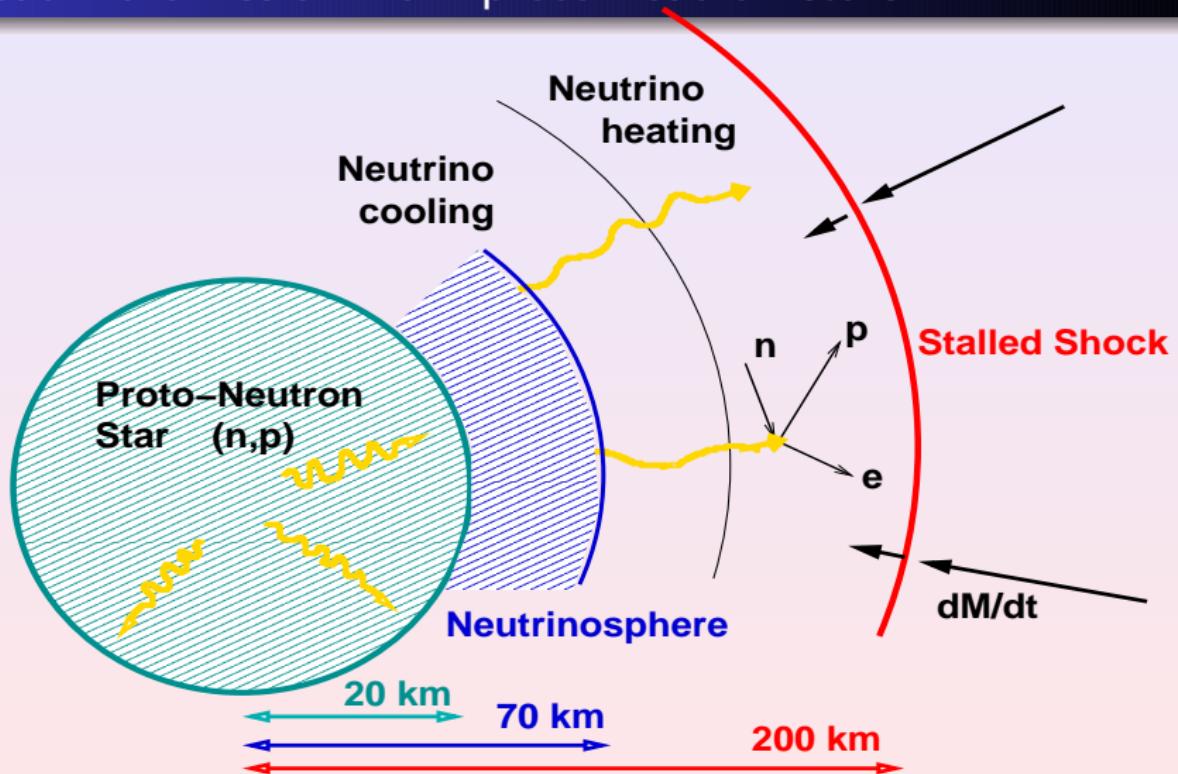
- gravitational core collapse:



- reflection for $\rho > \rho_{\text{nuc}}$, generation of out-going shock wave



Neutrino emission from proto-neutron stars



problem

SN neutrino spectra are model dependent:
how can mixing parameters be determined reliably?

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how can mixing parameters be determined reliably?

solution

use only things in which you are sure:

- shock wave
- Earth matter effect
- neutronization burst

Neutrino oscillations and matter effects

- SN core → envelope → interstellar medium → Earth

Neutrino oscillations and matter effects

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- consider **effective Hamiltonian** in (ν_e, ν_μ) basis:

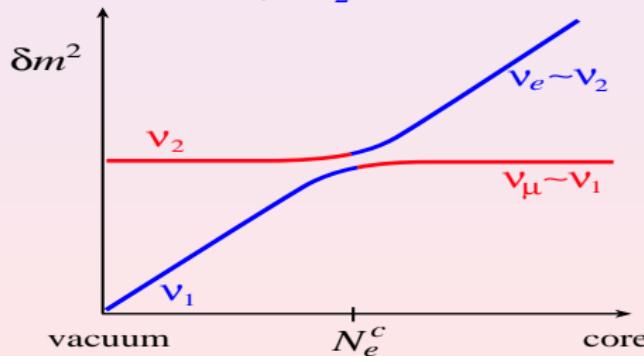
$$H = \frac{1}{4E} \begin{pmatrix} -\Delta m^2 \cos 2\theta + A & \Delta m^2 \sin 2\theta \\ \Delta m^2 \sin 2\theta & \Delta m^2 \cos 2\theta \end{pmatrix}, \quad A \propto \pm 2E\rho$$

Neutrino oscillations and matter effects

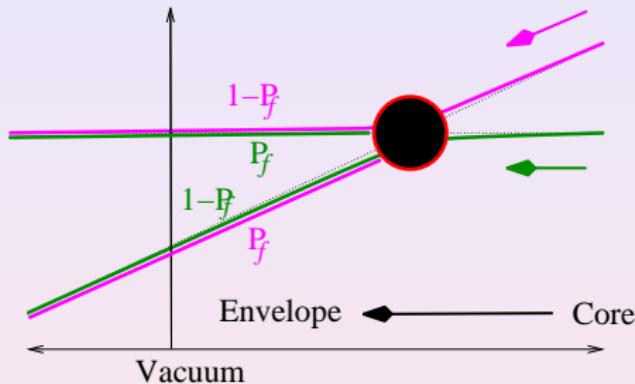
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- eigenvalues: $m_i^2 = \frac{A}{2} \mp \sqrt{(\Delta m^2 \cos 2\theta - A)^2 + (\Delta m^2 \sin 2\theta)^2}$



Adiabaticity at resonance:



$$P_c \approx \exp\left(-\frac{\pi}{2}\gamma\right), \quad \gamma \equiv \frac{\Delta m^2}{2E} \frac{\sin^2 2\theta}{\cos 2\theta} \left(\frac{1}{n_e} \frac{dn_e}{dr}\right)_{\text{res}}^{-1}$$

$\gamma \gg 1 \Rightarrow P_c \ll 1 \Rightarrow$ adiabatic resonance

depends on $\Delta m^2/E$, mixing angle, density profile

more

Observable 1: SN shock wave propagation

- **inverted hierarchy:** “atm. resonance” in **anti-neutrino** sector

Observable 1: SN shock wave propagation

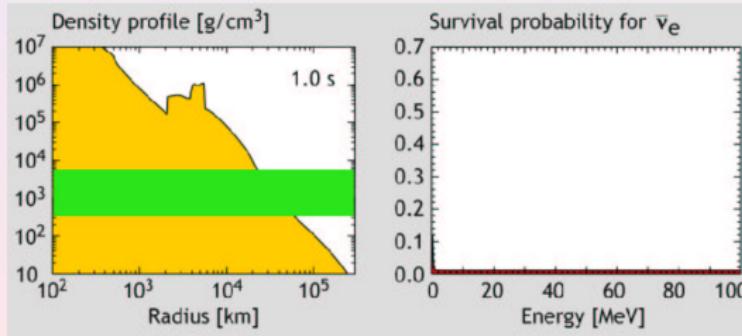
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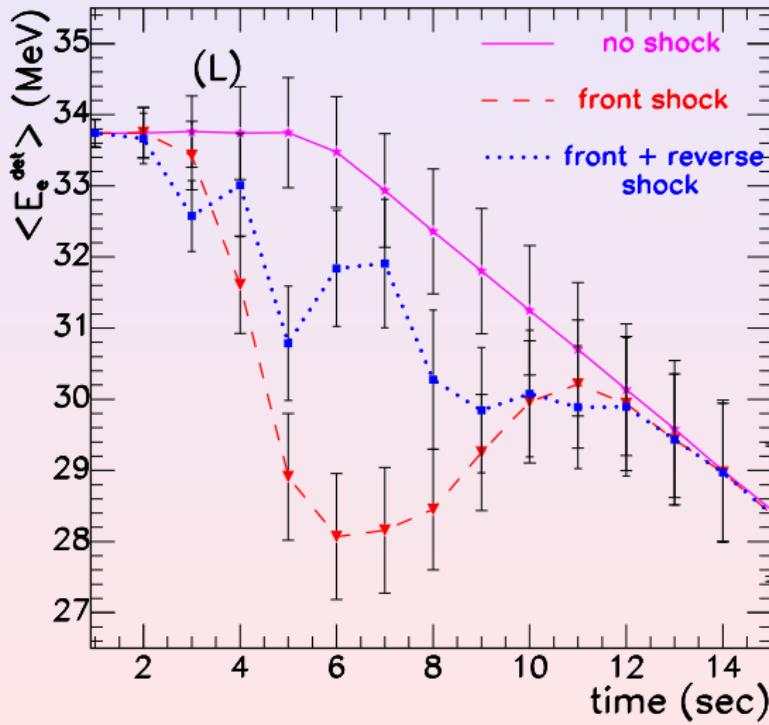
Observable 1: SN shock wave propagation

- inverted hierarchy: “atm. resonance” in anti-neutrino sector
- large θ_{13} : “atm. resonance” is adiabatic for progenitor profile
- shock waves passing through resonance break adiabaticity
- position of **resonance is energy dependent**
 \Rightarrow energy binned $\bar{\nu}_e$ spectra allows **tomography of SN**

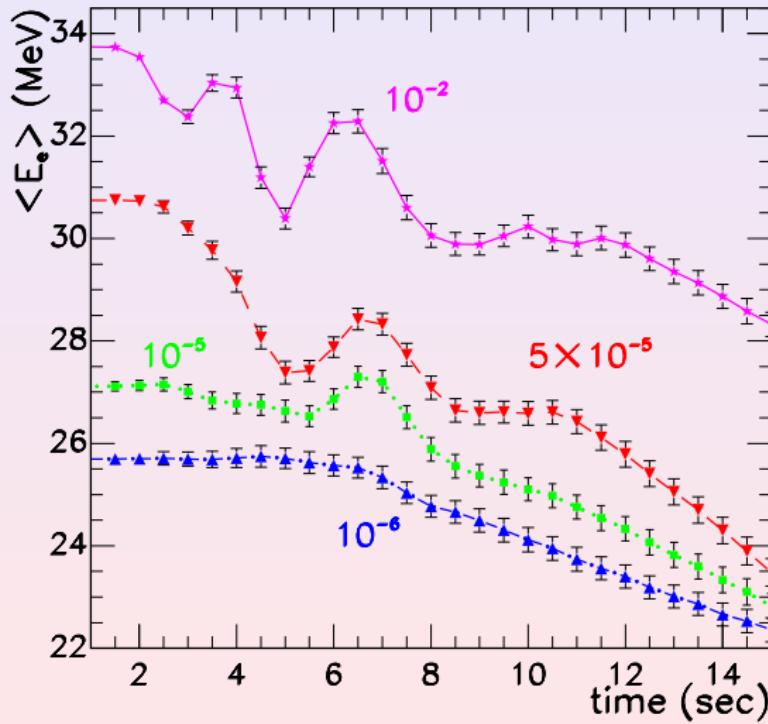


[Tomàs, MK, Raffelt, Dighe, Scheck, Janka '04]

Statistical significance: SuperKamiokande

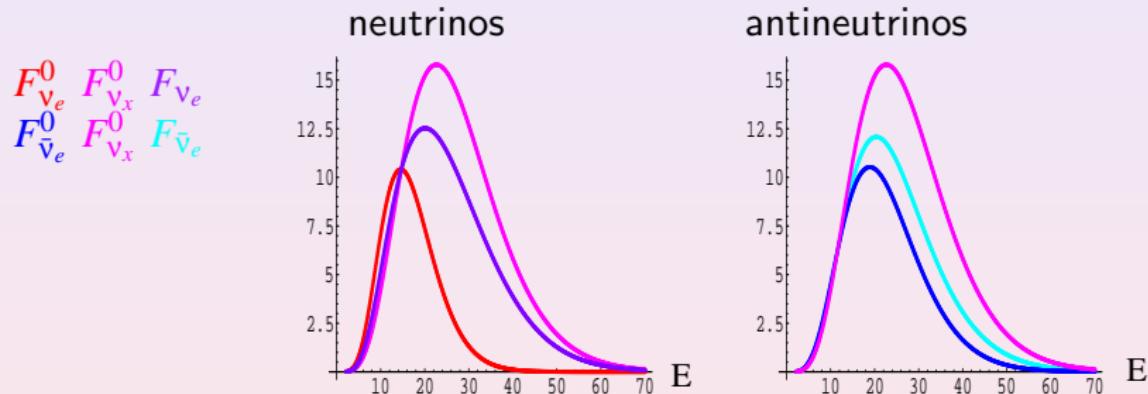


Statistical significance: HyperKamiokande



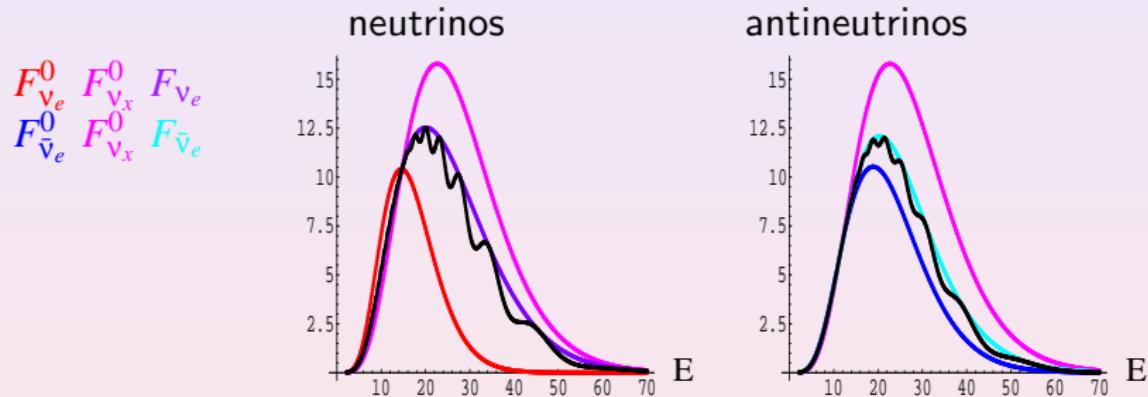
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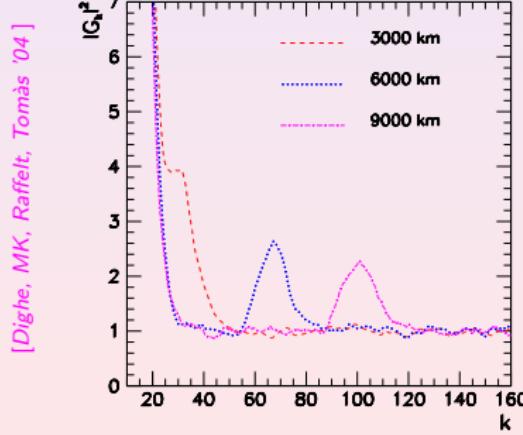
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[Dighe, MK, Raffelt, Tomàs '04]

Observable 2: Earth matter effects

- oscillations in the Earth superimpose wiggles on spectra
- frequencies are
 - analytically known
 - independent of the primary neutrino spectra
- Fourier transform of “inverse” energy spectrum $y = 1/E$



$$G(k) = \frac{1}{\sqrt{N}} \sum_{\text{events}} e^{iky}$$

SN neutrino summary

Hierarchy	$\sin^2 \theta_{13}$	Earth effects	Shocks	ν_e burst
Normal	$\gtrsim 10^{-3}$	$\bar{\nu}_e$	ν_e	absent
Inverted	$\gtrsim 10^{-3}$	ν_e	$\bar{\nu}_e$	present
Any	$\lesssim 10^{-5}$	ν_e and $\bar{\nu}_e$	—	present

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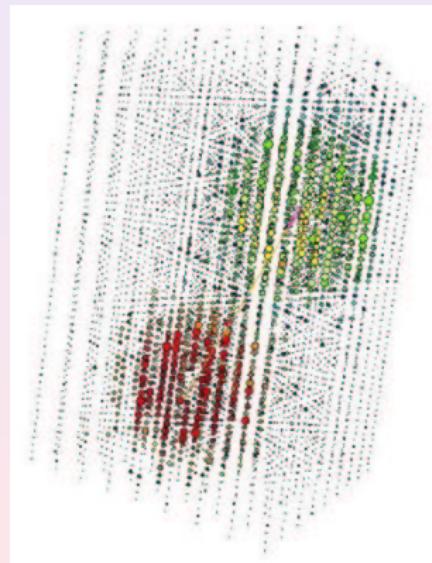
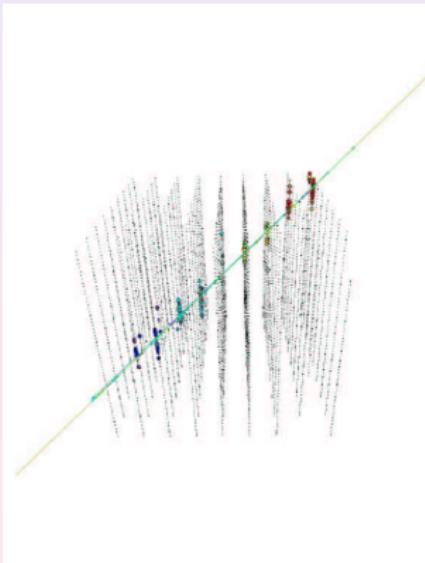
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galactic SN & water Cherenkov/scintillation detector allows

- identification of neutrino mixing scenario
- a lot of astrophysics

Neutrino telescopes and neutrino mixing

- neutrino telescopes can distinguish muon neutrinos from electron and tau neutrino events:



Neutrino telescopes and neutrino mixing

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$$\phi_e : \phi_\mu : \phi_\tau = 1 : 2 : 0 \quad \Rightarrow \quad \phi_e : \phi_\mu : \phi_\tau = 1 : 1 : 1$$

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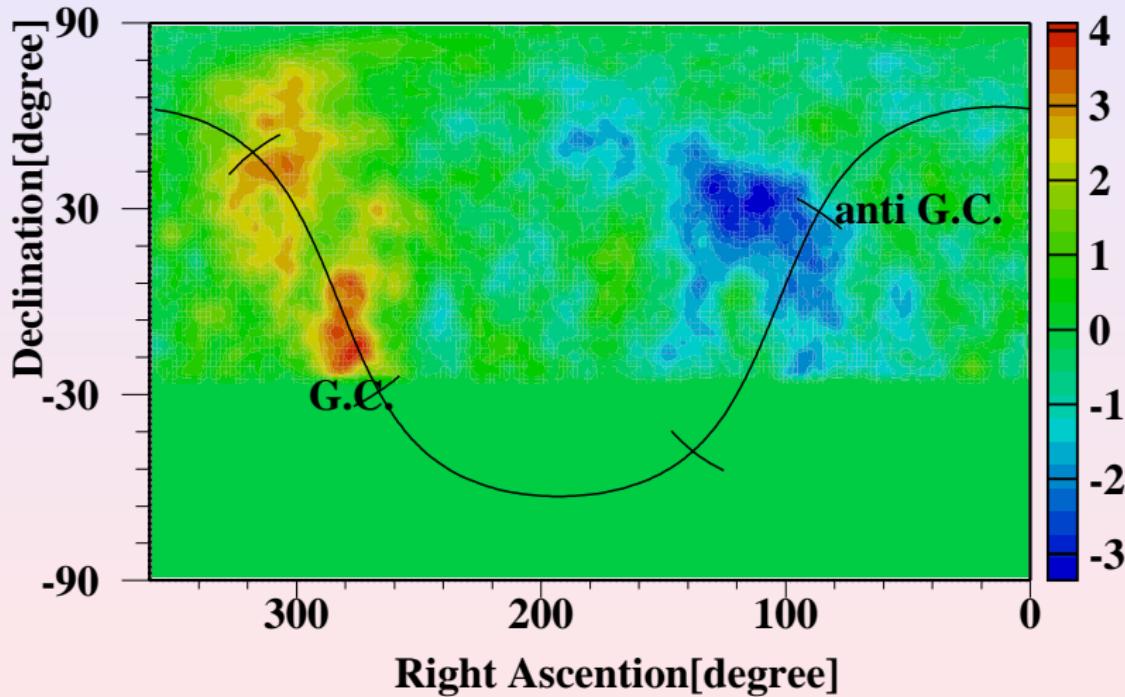
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- exception: e.g. beta beam from neutron decay
- example: galactic CR source near Cygnus region, if nuclei are accelerated

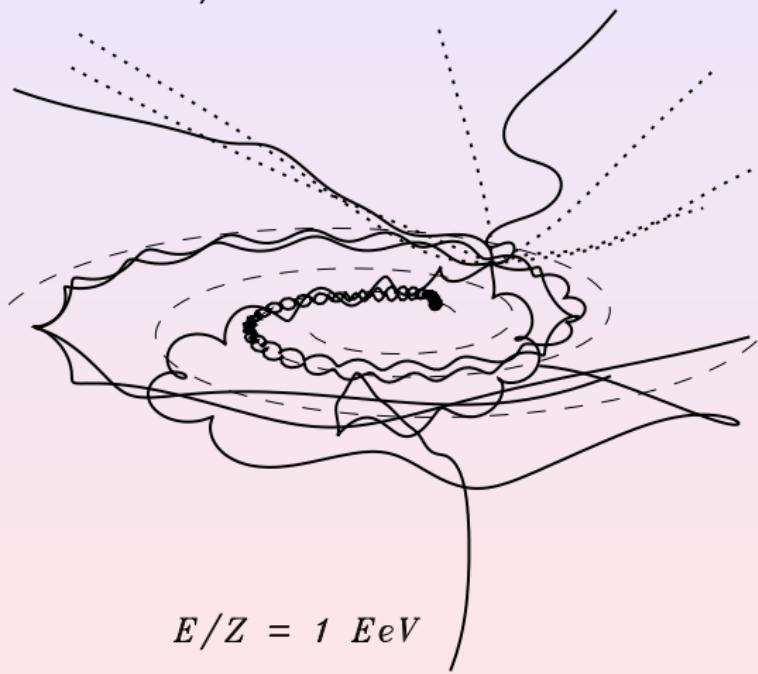
[MK, P. Serpico '05]

Galactic anisotropy around $E = 10^{18}$ eV: significance

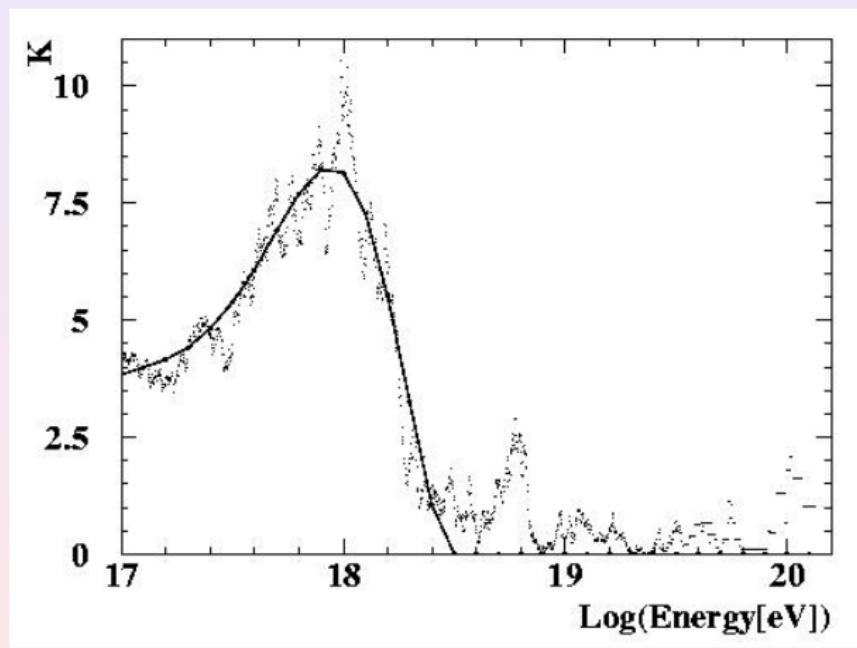


Deflection of charged particles in Galactic B -field:

$$E/Z = 10 \text{ EeV}$$



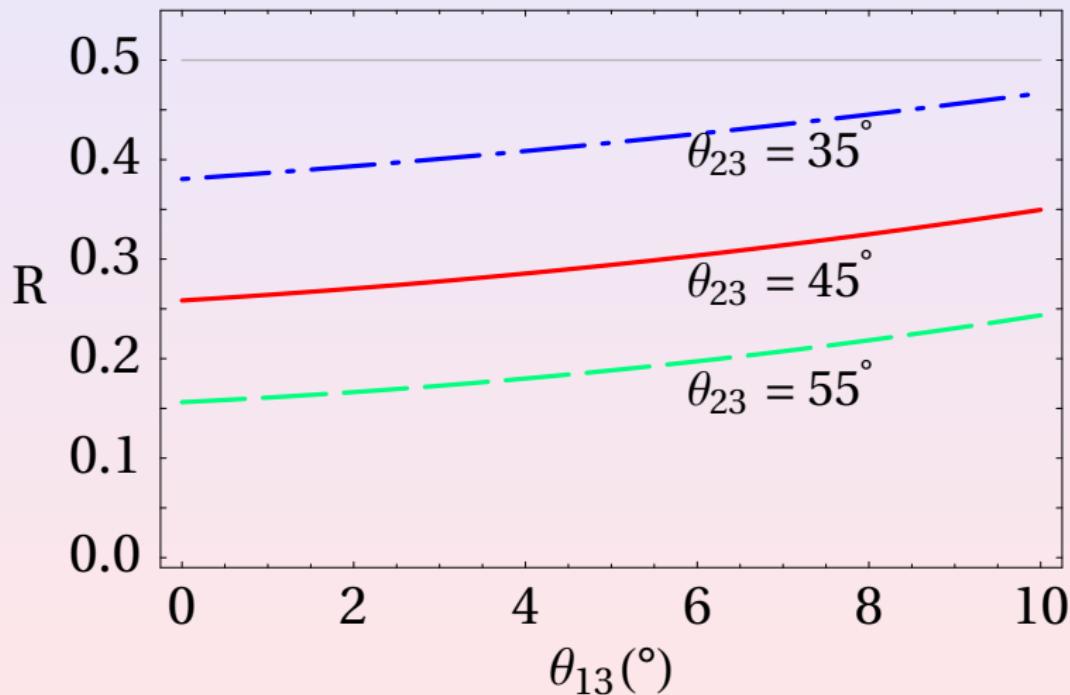
Fit of neutron source at GC in Leaky Box model:



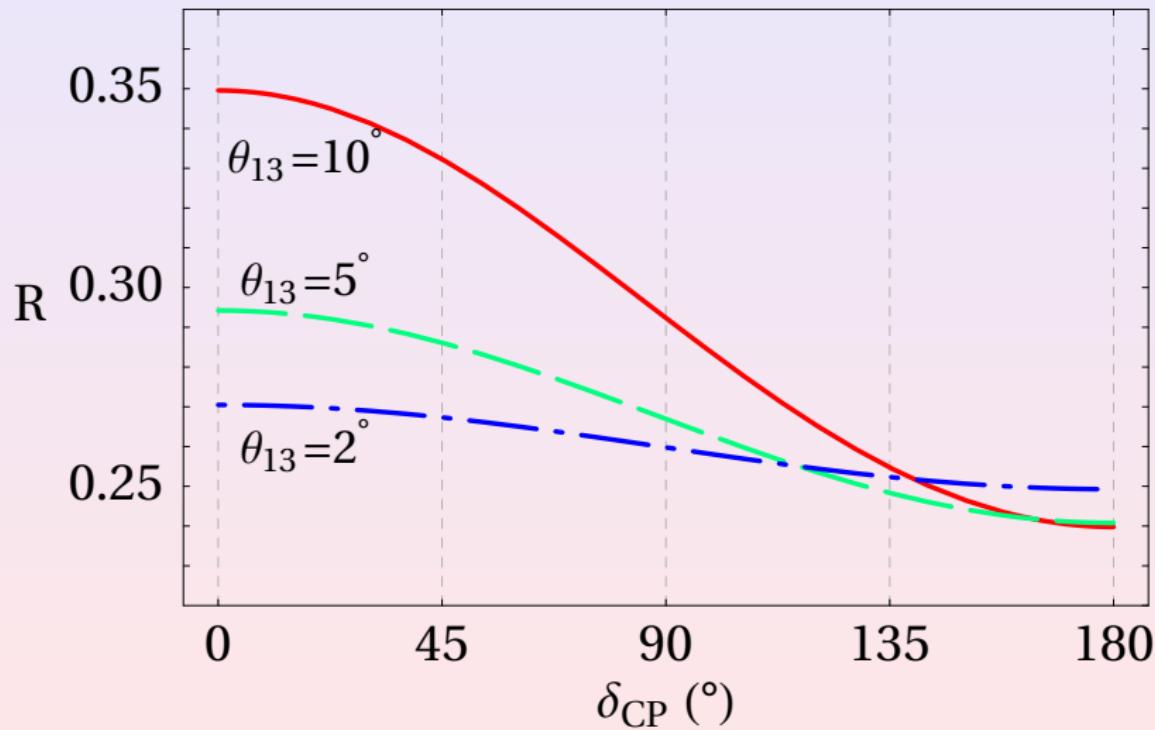
best fit with
 $dN/dE \propto E^{-2.5}$
distance 10kpc
 $E_{\text{max}} = 10^{18.5} \text{ eV}$

[AGASA '98]

Neutrino telescopes and neutrino mixing: $R = \phi^\mu / \phi^{e+\tau}$



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Sensitivity of ICECUBE:

statistical significance of Cyg detection:

- with $\ell = 25^\circ$ angular resolution for ν_e and ν_τ :
4.2 σ for 10 years measurement

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- with $\ell = 25^\circ$ angular resolution for ν_e and ν_τ :
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- with $\ell = 10^\circ$ angular resolution for ν_e and ν_τ :
3.3 σ for 1 year measurement

⇒ end

Top-Down Models

UHECR primaries are produced by **decays of supermassive particle X** with $M_X \gtrsim 10^{12}$ GeV.

- topological defects: monopoles, strings, ...

[Hill '83; Ostriker, Thompson, Witten '86]

- superheavy metastable particles

[Berezinsky, MK, Vilenkin '97; Kuzmin, Rubakov '97]

Advantages:

- no acceleration problem
- no visible sources
- if $X \in \text{CDM}$, no GZK-cutoff
- theoretically motivated; testable predictions

Fragmentation of heavy particles

- Consider Bremsstrahlung, $X \rightarrow \bar{f}fV$:

soft and collinear singularities generate $\ln^2(m_V^2/m_X^2)$ for $m_X^2 \gg m_V^2$

\Rightarrow they can compensate the small couplings g^2 ,

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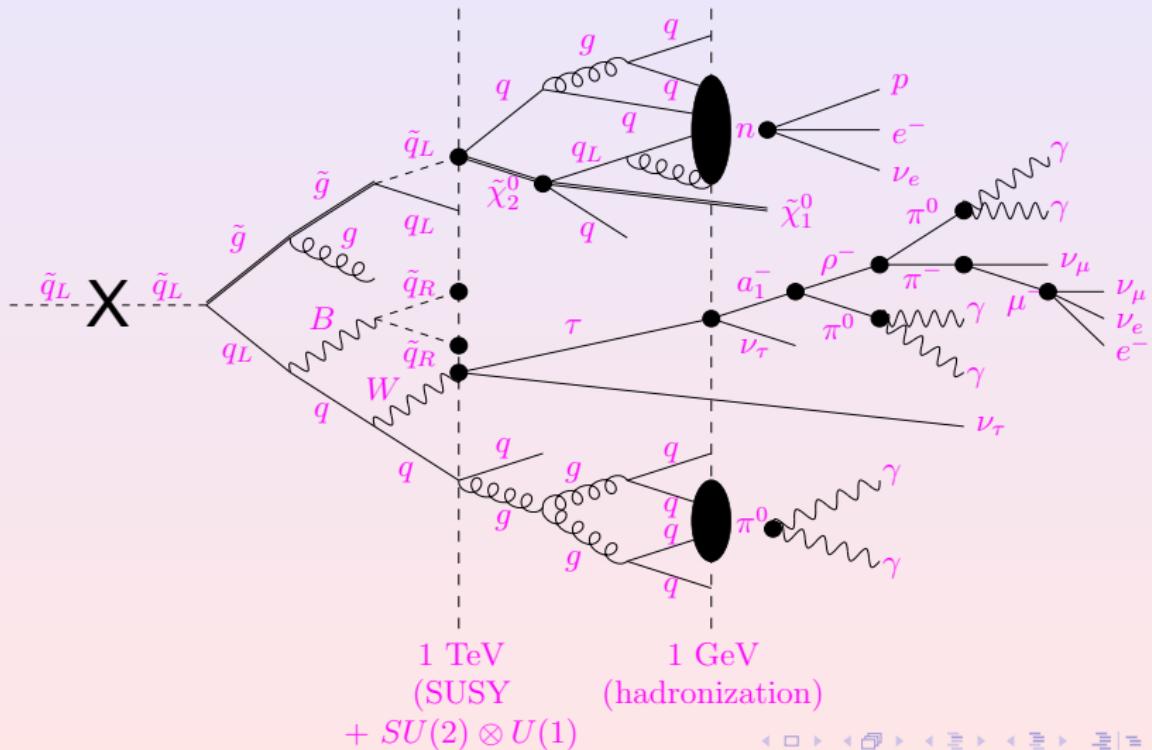
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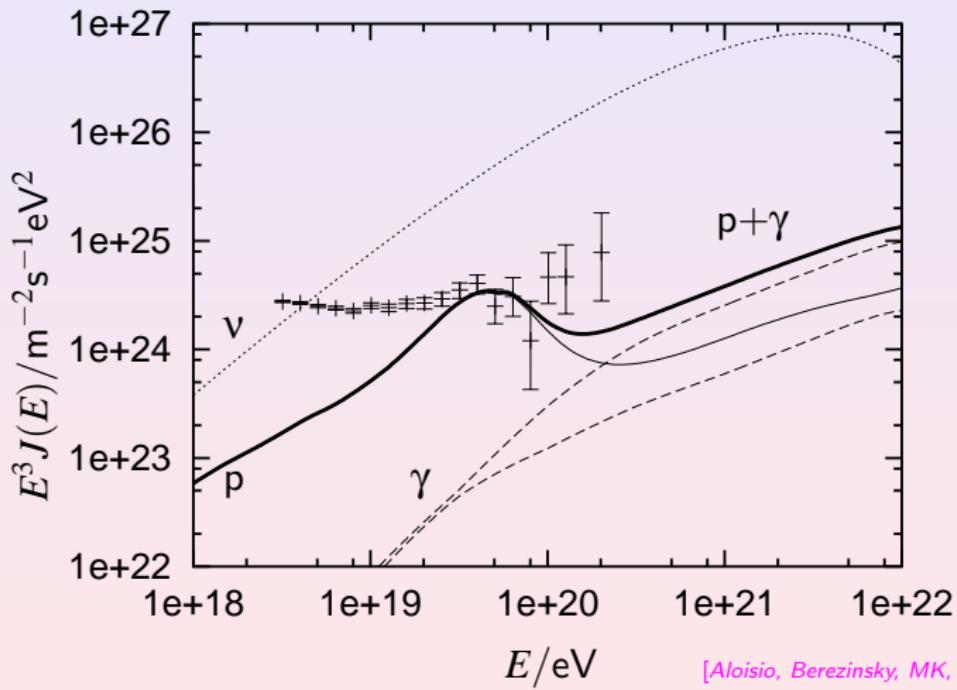
- $M_X \gtrsim 10^6$ GeV, \Rightarrow naive perturbation theory breaks down:
electroweak and SUSY sector have a QCD-like behavior
("jets")

[Berezinsky, MK '98, Berezinsky, MK, Ostapchenko '02]

Fragmentation of heavy particles



Neutrino fluxes from topological defects:



[Aloisio, Berezhinsky, MK, '03]

Summary

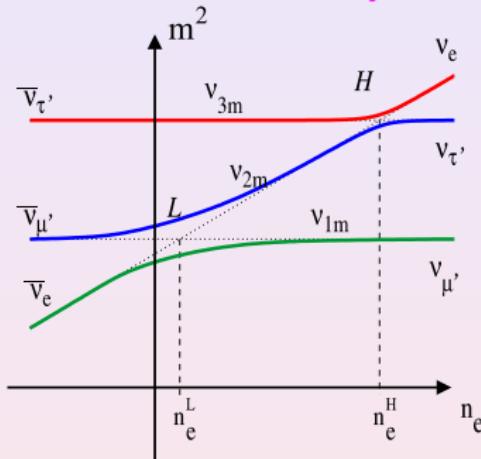
- Cosmic neutrinos encompass **24 orders in energy** from 2.7 K relic background to UHE neutrinos with 10^{20} eV
- Relic neutrinos allow to test the **absolute neutrino masses** (LSS, Z-Burst model)
- SN neutrinos may identify **mass hierarchy** and θ_{13}
- Cygnus neutrinos as test for δ_{CP} and θ_{13}
- HE neutrinos as test for **SUSY DM** and new interactions
- UHE neutrinos are one of best tests for **Lorentz invariance**
- UHE neutrinos test **superheavy dark matter or topological defects**

...

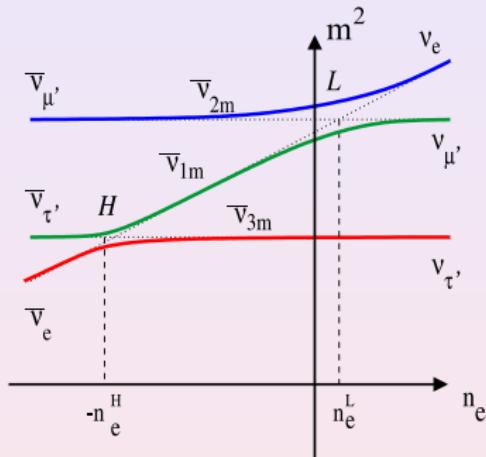
essential: interplay between **particle physics, astrophysics and cosmology**

3-v level crossing schemes

normal mass hierarchy

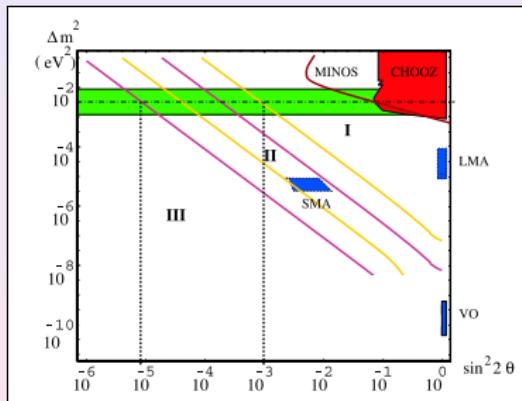


inverted mass hierarchy



- H -resonance: $(\Delta m_{\text{atm}}^2, \theta_{13})$, $\rho \sim 10^3 \text{ g/cm}^3$
- L -resonance: $(\Delta m_{\odot}^2, \theta_{\odot})$, $\rho \sim 10 \text{ g/cm}^3$
- Δm^2 hierarchy \Rightarrow resonances are independent

Adiabaticity at H and L resonances



- L always adiabatic
- H adiabatic for $|U_{e3}|^2 \gtrsim 10^{-3}$, non-adiabatic for $|U_{e3}|^2 \lesssim 10^{-3}$
- H in ν_e channel for normal hierarchy
 H in $\bar{\nu}_e$ channel for inverted hierarchy

SN neutrino spectra sensitive for

- $|U_{e3}|^2 \gtrsim 10^{-3}$ or $|U_{e3}|^2 \lesssim 10^{-3}$
- normal or inverted mass hierarchy

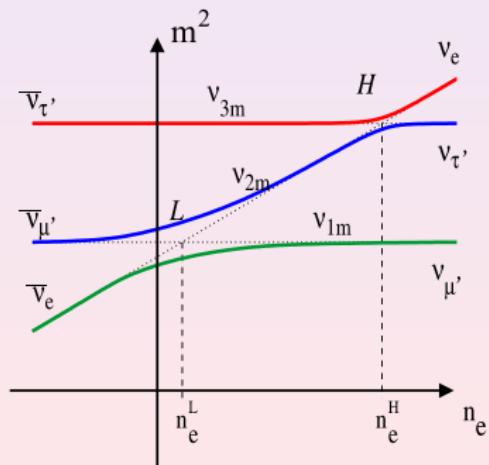
Fluxes arriving at the Earth

Mixing of the emitted $\bar{\nu}_e$ flux:

$$F_{\bar{\mathbf{v}}_e} = \bar{p} F_{\bar{\mathbf{v}}_e}^0 + (1 - \bar{p}) F_{\mathbf{v}_x}^0$$

Survival probability for different mixing scenarios:

case	hierarchy	$\sin^2 \Theta_{13}$	\bar{p}
A	normal	$\gtrsim 10^{-3}$	$\cos^2 \Theta_\odot$
B	inverted	$\gtrsim 10^{-3}$	0
C	both	$\lesssim 10^{-3}$	$\cos^2 \Theta_\odot$



→

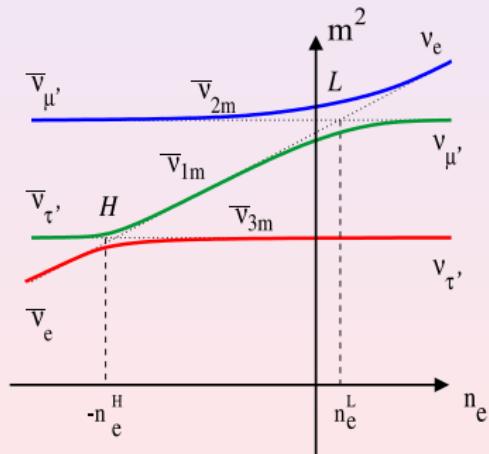
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C	both	$\lessgtr 10^{-3}$	$\cos^2 \Theta_\odot$



→

Earth matter effects II

⇒

$$F_{\bar{v}_e} = \sin^2 \Theta_{12} F_{\bar{v}_x}^0 + \cos^2 \Theta_{12} F_{\bar{v}_e}^0 + \Delta F^0 \bar{A}_{\oplus} \sin^2(\overline{\Delta m_{\oplus}^2} L_y)$$

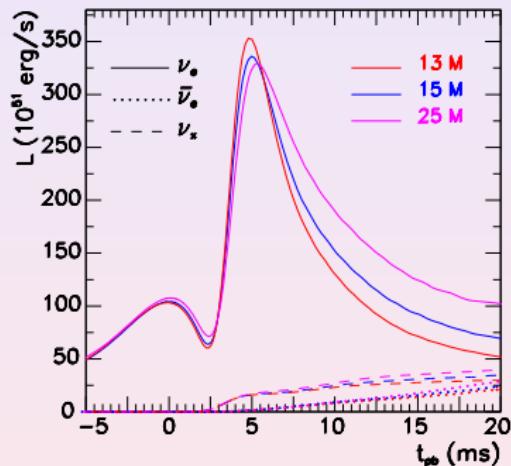
$$(F_{\bar{v}_e}^0 - F_{\bar{v}_x}^0) \quad \sin 2\bar{\Theta}_{12}^\oplus \sin(2\bar{\Theta}_{12}^\oplus - 2\Theta_{12}) \quad (12.5/E)$$

- coefficient $\Delta F^0 \bar{A}_{\oplus}$ varies slowly with E
- neutrino spectrum as function of $y \equiv 12.5/E$ oscillate with frequency k_{\oplus} fast around primary spectrum

Oscillation frequency: $k_{\oplus} \equiv 2\overline{\Delta m_{\oplus}^2} L$

- completely independent of the primary neutrino spectra
- depends only on solar oscillation parameters, Earth density and the distance traveled through the Earth

Observable 3: ν_e neutronization burst

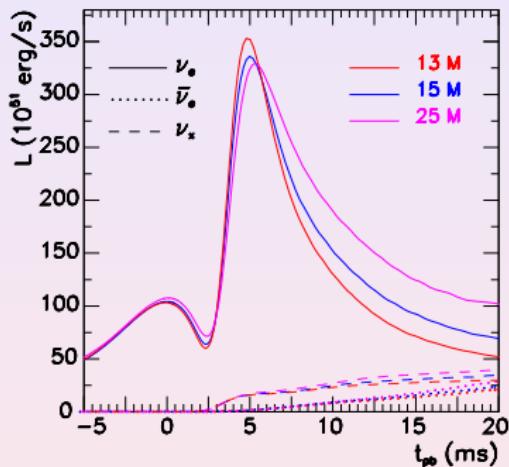


independent from

- progenitor mass
- EoS
- simulation

[MK, Tomàs, Buras, Janka, Marek, Rampp '04]

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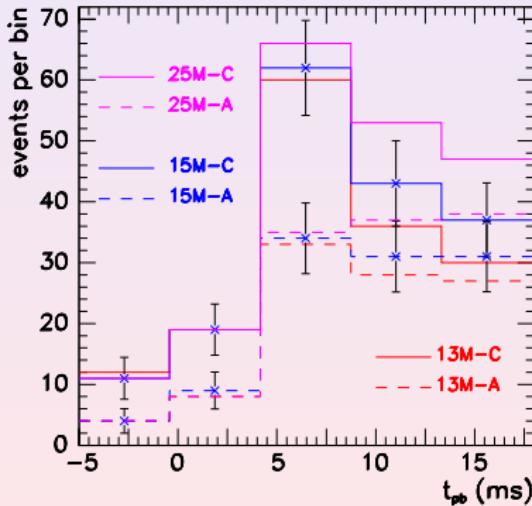
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[MK, Tomàs, Buras, Janka, Marek, Rampp '04]

- allows to measure SN distance with 10% precision

Observable 3: ν_e neutronization burst

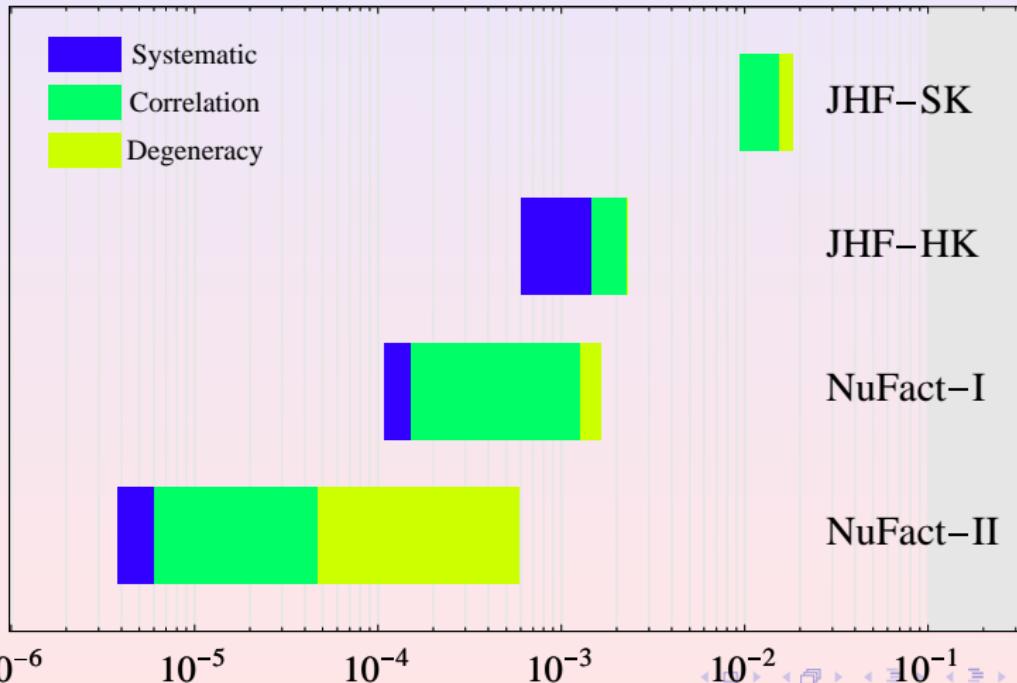
- peak in time-binned spectra disappears for normal hierarchy and “large θ_{13} ”



[MK, Tomàs, Buras, Janka, Marek, Rampp '04]

Sensitivity of future LBL experiments:

Sensitivity to $\sin^2 2\theta_{13}$



Scintillation vs. water Cherenkov detectors:

