Top Instability and Electroweak Effects at the ttroduction Threshold

#### Christoph Reißer

Max-Planck-Institut für Physik, Munich

A.H.Hoang & C.R.: Phys. Rev. **D** 71, 074022 (2005)



IMPRS seminar, MPI Munich, 11 Nov 2005

## Outline

- Motivation
- Status for  $\sigma_{tot}(e^+e^- \rightarrow t\bar{t})$  at threshold
- Non-relativistic QCD (vNRQCD)
- Effects of the top instability
- (Usual) electroweak effects
- Numerical analysis
- Summary and Outlook



## tt Production Threshold

<u>e<sup>+</sup>e<sup>-</sup></u> collisions: c.m. energy  $\sqrt{s} \approx 340 - 360$  GeV

Top quarks are non-relativistic

 $\sim 1$ 



⇒ Perturbation theory in  $\alpha_s$  breaks down  $v \sim \alpha_s$ ⇒ Non-relativistic QCD  $\simeq$  Schrödinger theory at LO

Ŧ

 $\sim \frac{\alpha_s}{v}$ 

## tt Production Threshold

<u>e<sup>+</sup>e<sup>-</sup></u> collisions: c.m. energy  $\sqrt{s} \approx 340 - 360$  GeV

Top quarks are non-relativistic

Top quarks decay fast 
$$\Gamma$$

- $\Rightarrow$  No bound states
- ⇒ Smooth line-shape
- ⇒ Non-perturbative effects suppressed [Fadin,Khoze]

$$\Gamma_t \approx 1.5 \text{ GeV} \gg \Lambda_{\rm QCD}$$



 $\frac{4m_t^2}{s} \ll 1$ 



IMPRS seminar, MPI Munich, 11 Nov 2005

## tt Production Threshold

<u>e<sup>+</sup>e<sup>-</sup></u> collisions: c.m. energy  $\sqrt{s} \approx 340 - 360$  GeV

Top quarks are non-relativistic

$$\mathsf{v}=\sqrt{1-\frac{4\mathsf{m}_t^2}{\mathsf{s}}}\ll 1$$

Top quarks decay fast

$$\Gamma_{t} \approx 1.5 \text{ GeV} \gg \Lambda_{QCD}$$

 $\sigma^{\rm obs}(s) = \int_0^1 dx \mathcal{L}(x) \sigma^{\rm theo}(x^2 s)$ 

- Measured cross section contains
  - beam spread
  - beamstrahlung
  - ISR



Christoph Reißer – p.3



#### Measurements

Simulations of Threshold Scan ( $\mathcal{L} \sim 300 \, {\rm fb}^{-1}$ ): [Martinez, Miquel]

Top quark mass

 $(\delta m_t)^{exp} \sim 50 \text{ MeV}$ 

Top Yukawa coupling

 $(\delta y_t/y_t)^{\rm exp} \sim 0.35$  (light Higgs)

Top decay width

 $(\delta\Gamma_{\rm t})^{
m exp}\sim$  50 MeV



#### Measurements

Simulations of Threshold Scan ( $\mathcal{L} \sim 300 \, {\rm fb}^{-1}$ ): [Martinez, Miquel]

Top quark mass

 $(\delta m_t)^{exp} \sim 50 \text{ MeV}$ 

Top Yukawa coupling

 $(\delta y_t/y_t)^{\rm exp} \sim 0.35$  (light Higgs)

Top decay width

 $(\delta\Gamma_{t})^{\mathrm{exp}}\sim$  50 MeV

 $\Rightarrow$  Theory needs



## **Theory Status**

RGE methods to sum large logs  $(\alpha_s \ln v)^m$ 

QCD effects

$$\begin{array}{lllllllllll} \text{LL }\checkmark & \qquad & \text{LL } &\sim \left(\frac{\alpha_{\text{s}}}{\nu}\right)^{\text{n}} \\ \\ \text{NLL }\checkmark & \qquad & \text{NLL } &\sim \alpha_{\text{s}} \left(\frac{\alpha_{\text{s}}}{\nu}\right)^{\text{n}} \\ \\ \text{NNLL (nearly)} \rightarrow \left(\frac{\delta\sigma_{\text{tot}}}{\sigma_{\text{tot}}}\right)^{\text{NNLL}} \sim \pm 6 \% & \qquad & \text{NNLL } &\sim \alpha_{\text{s}}^2 \left(\frac{\alpha_{\text{s}}}{\nu}\right)^{\text{n}} \end{array}$$



## **Theory Status**

RGE methods to sum large logs  $(\alpha_{\rm s}\ln{\rm v})^{\rm m}$ 

- QCD effects
- Electroweak effects

LL √

NLL (partly)

NNLL?



IMPRS seminar, MPI Munich, 11 Nov 2005

# **Theory Status**

RGE methods to sum large logs  $(\alpha_{\rm s}\ln{\rm v})^{\rm m}$ 

- QCD effects
- Electroweak effects

#### $LL \checkmark$

 $\begin{array}{ll} \mbox{NLL (partly)} & \rightarrow \mbox{New parametric NLL corrections} \\ \mbox{NNLL ?} & \rightarrow \mbox{NNLL top decay corrections} \end{array}$ 



IMPRS seminar, MPI Munich, 11 Nov 2005

[Hoang, Luke, Manohar, Rothstein, Stewart]

#### Relevant scales

 $m_t \text{ (hard)} \quad \gg \quad \mathbf{p} \sim m_t v \text{ (soft)} \quad \gg \quad E \sim m_t v^2 \text{ (ultrasoft)}$ 



IMPRS seminar, MPI Munich, 11 Nov 2005

[Hoang, Luke, Manohar, Rothstein, Stewart]

#### Relevant scales

 $\begin{array}{ll} \mathsf{m}_{t} \ (\mathsf{hard}) & \gg & \mathbf{p} \sim \mathsf{m}_{t} \mathsf{v} \ (\mathsf{soft}) & \gg & \mathsf{E} \sim \mathsf{m}_{t} \mathsf{v}^{2} \ (\mathsf{ultrasoft}) \end{array}$   $\begin{array}{ll} \underline{\mathsf{Lagrangian}} \ (\mathsf{LL}) & \mathcal{L} = \mathcal{L}_{\mathrm{bilinear}} + \mathcal{L}_{\mathrm{potential}} \\ \mathcal{L}_{\mathrm{bilinear}} = \psi^{\dagger}_{\mathbf{p}} (\mathsf{x}) \ \left\{ \mathsf{i} \mathsf{D}^{0} - \frac{\mathbf{p}^{2}}{2\mathsf{m}_{t}} - \delta\mathsf{m}_{t} \right\} \psi_{\mathbf{p}} (\mathsf{x}) + \dots \\ \mathcal{L}_{\mathrm{potential}} = - \frac{\mathcal{V}_{c}(\mu)}{(\mathbf{p} - \mathbf{p}')^{2}} \psi^{\dagger}_{\mathbf{p}'} \psi_{\mathbf{p}} \chi^{\dagger}_{-\mathbf{p}'} \chi_{-\mathbf{p}} + \dots \end{array}$ 



[Hoang, Luke, Manohar, Rothstein, Stewart]

#### Relevant scales

 $m_t$  (hard)  $\gg p \sim m_t v$  (soft)  $\gg E \sim m_t v^2$  (ultrasoft) Lagrangian (LL)  $\mathcal{L} = \mathcal{L}_{\text{bilinear}} + \mathcal{L}_{\text{potential}}$  $\mathcal{L}_{\text{bilinear}} = \psi_{\mathbf{p}}^{\dagger}(\mathbf{x}) \left\{ \mathsf{i}\mathsf{D}^{0} - \frac{\mathbf{p}^{2}}{2\mathsf{m}_{t}} - \delta\mathsf{m}_{t} \right\} \psi_{\mathbf{p}}(\mathbf{x}) + \dots$  $\mathcal{L}_{\text{potential}} = -\frac{\mathcal{V}_{c}(\mu)}{(\mathbf{p}-\mathbf{p}')^{2}}\psi^{\dagger}_{\mathbf{p}'}\psi_{\mathbf{p}}\chi^{\dagger}_{-\mathbf{p}'}\chi_{-\mathbf{p}} + \dots$  $\Rightarrow$  Schrödinger equation with Coulomb potential  $G^{LL}(0,0,\sqrt{s}) = \frac{m_t^2}{4\pi} \left\{ i v - C_F \alpha_s \left| \frac{1}{4\epsilon} + \ln\left(\frac{-i m_t v}{\mu}\right) + \dots \right| \right\}$ 

IMPRS seminar, MPI Munich, 11 Nov 2005

<u>Currents</u> for  $t\overline{t}$  production & annihilation

$$\mathbf{O}_{\mathbf{p}} = \mathsf{C}(\mu) \cdot (\psi_{\mathbf{p}}^{\dagger} \,\boldsymbol{\sigma} \, \tilde{\chi}_{-\mathbf{p}}^{*}) \qquad (^{3}\mathsf{S}_{1})$$



IMPRS seminar, MPI Munich, 11 Nov 2005

<u>Currents</u> for  $t\bar{t}$  production & annihilation

$$\mathbf{O}_{\mathbf{p}} = \mathsf{C}(\mu) \cdot (\psi_{\mathbf{p}}^{\dagger} \,\boldsymbol{\sigma} \, \tilde{\chi}_{-\mathbf{p}}^{*}) \qquad (^{3}\mathsf{S}_{1})$$

<u>Total cross section</u> from  $e^+e^- \rightarrow e^+e^-$  using the Optical Theorem [Strassler, Peskin]

$$\sigma_{\text{tot}} \propto \text{Im} \left[ \int d^4 x \, e^{-i\hat{q} \cdot x} \left\langle 0 \, \left| \, T \, \mathbf{O}_{\mathbf{p}}^{\dagger}(0) \, \mathbf{O}_{\mathbf{p}'}(x) \, \right| \, 0 \right\rangle \right]$$

$$\propto \text{Im} \left[ \mathsf{C}(\mu)^2 \, \mathsf{G}_{(}0, 0, \sqrt{\mathsf{s}}) \right]$$



IMPRS seminar, MPI Munich, 11 Nov 2005



**Optical Theory of absorptive medium** 

- Complex refractive indices in the Maxwell equations
- Can be derived systematically from microscopic processes



**Optical Theory of absorptive medium** 

- Complex refractive indices in the Maxwell equations
- Can be derived systematically from microscopic processes

vNRQCD accounting for  $t \rightarrow Wb$ 

- Wilson coefficients become complex
- Obtained by matching effective theory to SM
- Effective Lagrangian <u>non-Hermitian</u>

 $\rightarrow$  Total cross section through the optical theorem



Matching conditions accounting for  $t \rightarrow \mathsf{Wb}$ 

Bilinear quark operators





IMPRS seminar, MPI Munich, 11 Nov 2005

Matching conditions accounting for  $t \to \mathsf{Wb}$ 

Bilinear quark operators



Power counting for ew. effects:  $\Gamma_t \sim m_t g^2 \sim m_t \alpha_s^2 \Rightarrow g \sim g' \sim v \sim \alpha_s$ 

- $E \rightarrow E + i\Gamma_t$  replacement to include finite lifetime at LL [Fadin, Khoze] (not sufficient beyond LL order)
- Time dilatation is NNLL

IMPRS seminar, MPI Munich, 11 Nov 2005

Matching conditions accounting for  $t \rightarrow Wb$ <u>Currents</u> for  $t\overline{t}$  production & annihilation

$$\mathbf{O}_{\mathbf{p}} = \left[ \mathsf{C}^{\mathrm{LL}} + \mathsf{C}^{\mathrm{NLL}} + \mathsf{C}^{\mathrm{NNLL}} + \mathsf{i} \, \mathsf{C}^{\mathrm{NNLL}}_{\mathrm{abs}} + \dots \right] \cdot \left( \begin{array}{c} e^{+} & t \\ e^{-} & \overline{t} \end{array} \right) + \dots$$



IMPRS seminar, MPI Munich, 11 Nov 2005

Matching conditions accounting for  $t \rightarrow Wb$ <u>Currents</u> for  $t\overline{t}$  production & annihilation

$$\mathbf{O}_{\mathbf{p}} = \begin{bmatrix} \mathsf{C}^{\mathrm{LL}} + \mathsf{C}^{\mathrm{NLL}} + \mathsf{C}^{\mathrm{NNLL}} + \mathsf{i} \, \mathsf{C}^{\mathrm{NNLL}}_{\mathrm{abs}} + \dots \end{bmatrix} \cdot \begin{bmatrix} e^{-\overline{t}} & \overline{t} \\ e^{-\overline{t}} & \overline{t} \end{bmatrix} + \dots$$



IMPRS seminar, MPI Munich, 11 Nov 2005

Christoph Reißer – p.1

 $\left(\begin{array}{c} e^{+} \\ t \end{array}\right)$ 

Matching conditions accounting for  $t \rightarrow Wb$ <u>Currents</u> for  $t\overline{t}$  production & annihilation





IMPRS seminar, MPI Munich, 11 Nov 2005

Interactions in vNRQCD constrained by symmetries





IMPRS seminar, MPI Munich, 11 Nov 2005

Interactions in vNRQCD constrained by symmetries



#### Potentials:

 $\rightarrow$  no NNLL electroweak corrections

Soft, ultrasoft interactions:

- $\rightarrow$  no NNLL electroweak corrections
- → also: no non-factorizable effects from ultrasoft gluons [Melnikov, Yakovlev]





IMPRS seminar, MPI Munich, 11 Nov 2005

Optical Theorem  $\Rightarrow \sigma_{tot} = 2 N_c \operatorname{Im} \left[ C(\mu)^2 G(0, 0, \sqrt{s}) \right]$ 

NNLL finite lifetime correction:

$$\Delta^{\Gamma,1}\sigma_{\rm tot} = 2\,\mathsf{N}_{\mathsf{c}}\,\mathrm{Im}\left[2\,\mathsf{C}^{\mathrm{LL}}\,\mathsf{i}\mathsf{C}^{\mathrm{NNLL}}_{\mathrm{abs}}\,\mathsf{G}^{\mathrm{LL}}(0,0,\sqrt{\mathsf{s}}) + \left(\mathsf{C}^{\mathrm{LL}}\right)^{2}\delta\mathsf{G}^{\mathrm{NNLL}}(0,0,\sqrt{\mathsf{s}})\right]$$



Optical Theorem  $\Rightarrow \sigma_{tot} = 2 N_c \operatorname{Im} \left[ C(\mu)^2 G(0, 0, \sqrt{s}) \right]$ 

NNLL finite lifetime correction:

$$\Delta^{\Gamma,1}\sigma_{\rm tot} = 2\,\mathsf{N}_{\mathsf{c}}\,\mathrm{Im}\left[\underbrace{2\,\mathsf{C}^{\rm LL}\,\mathsf{i}\mathsf{C}^{\rm NNLL}_{\rm abs}\,\mathsf{G}^{\rm LL}(0,0,\sqrt{\mathsf{s}})}_{\mathsf{C}^{\rm LL}} + \underbrace{\left(\mathsf{C}^{\rm LL}\right)^2\delta\mathsf{G}^{\rm NNLL}(0,0,\sqrt{\mathsf{s}})}_{\mathsf{C}^{\rm LL}}\right]$$







IMPRS seminar, MPI Munich, 11 Nov 2005

Optical Theorem  $\Rightarrow \sigma_{tot} = 2 N_c \operatorname{Im} \left[ C(\mu)^2 G(0, 0, \sqrt{s}) \right]$ 

NNLL finite lifetime correction:

$$\Delta^{\Gamma,1}\sigma_{\rm tot} = 2\,\mathsf{N}_{\mathsf{c}}\,\mathrm{Im}\left[\,2\,\mathsf{C}^{\rm LL}\,\mathsf{i}\mathsf{C}^{\rm NNLL}_{\rm abs}\,\mathsf{G}^{\rm LL}(0,0,\sqrt{\mathsf{s}}) + \left(\mathsf{C}^{\rm LL}\right)^2\delta\mathsf{G}^{\rm NNLL}(0,0,\sqrt{\mathsf{s}})\right]$$

■ NLL mixing effect:  $\Delta^{\Gamma,1}\sigma_{tot}$  contains

 $\Gamma_{t} \alpha_{s\frac{1}{\epsilon}}$  (logarithmic UV phase space divergence due to finite lifetime)



Christoph Reißer – p.1

Optical Theorem  $\Rightarrow \sigma_{tot} = 2 N_c \operatorname{Im} \left[ C(\mu)^2 G(0, 0, \sqrt{s}) \right]$ 

NNLL finite lifetime correction:

$$\Delta^{\Gamma,1}\sigma_{\rm tot} = 2\,\mathsf{N}_{\mathsf{c}}\,\mathrm{Im}\left[\,2\,\mathsf{C}^{\rm LL}\,\mathsf{i}\mathsf{C}^{\rm NNLL}_{\rm abs}\,\mathsf{G}^{\rm LL}(0,0,\sqrt{\mathsf{s}}) + \left(\mathsf{C}^{\rm LL}\right)^2\delta\mathsf{G}^{\rm NNLL}(0,0,\sqrt{\mathsf{s}})\right]$$

● NLL mixing effect:  $\Delta^{\Gamma,1}\sigma_{tot}$  contains

 $\Gamma_{t} \alpha_{s\frac{1}{\epsilon}}$  (logarithmic UV phase space divergence due to finite lifetime)

 $\Rightarrow$  Anomalous dimension for  $(e^+e^-)(e^+e^-)$  operator:

 $\mathsf{i}\, \tilde{\mathsf{C}}^{\mathsf{eeee}}(\mu) \cdot \left( \begin{smallmatrix} e^+ & e^- \\ e^- & e^+ \end{smallmatrix} 
ight)$ 

- $\Rightarrow$  Correction  $\Delta^{\Gamma,2}\sigma_{\rm tot}$ 
  - $\sqrt{s}$ -independent
  - scale-dependent

\* Matching coefficient  $\tilde{C}^{eeee}(\mu = m_t)$  not yet determined  $\rightarrow$  w.i.p.

IMPRS seminar, MPI Munich, 11 Nov 2005



## **Absorptive Parts**



- $\Rightarrow$  Comparable to NNLL QCD corrections
- $\Rightarrow$  LL peak position shifted by 30 50 MeV

IMPRS seminar, MPI Munich, 11 Nov 2005

## **Absorptive Parts**



- $\Rightarrow$  Comparable to NNLL QCD corrections
- $\Rightarrow$  LL peak position shifted by 30 50 MeV

IMPRS seminar, MPI Munich, 11 Nov 2005

## **Electroweak Corrections** (preliminary)

 $\int_{\bar{t}}^{t} = \left[ C^{\text{LL}} + C^{\text{NLL}} + C^{\text{NNLL}} + i C^{\text{NNLL}}_{\text{abs}} + \dots \right] \cdot \left( \int_{e^{-t}}^{e^{+t}} \int_{\bar{t}}^{t} \right)$ 

Electroweak short distance corrections: Real parts C<sup>NNLL</sup>

- > Full electroweak theory one-loop diagrams  $O(\alpha_{em}) \rightarrow NNLL$
- >> Sizable shift of total cross section normalization
- >> Pure QED & ISR diagrams require separate treatment (SCET) → w.i.p.



### **Electroweak Corrections** (preliminary)

Correction to the total cross section normalization:  $(1 + 2\delta c_{ew}) \sigma_{tot}$ 





IMPRS seminar, MPI Munich, 11 Nov 2005

# **Summary and Outlook**

Summary

- **•** Threshold scan allows for precise  $m_t, y_t, \Gamma_t$  determination
- Effective theory approach (unstable particles)
  - UV divergencies from instablility
  - Interference of resonant and non-resonant diagrams
  - Symmetries restrict interactions
  - Short distance electroweak corrections
- Corrections comparable to QCD

<u>Outlook</u>

- Phase space matching  $\tilde{C}^{\text{eeee}}(\mu = m_t)$
- NNLL running of  $\tilde{C}^{eeee}(\mu)$
- QED contributions: ISR, beamstrahlung, Coulomb singularities

IMPRS seminar, MPI Munich, 11 Nov 2005