

# Hard Spectator Interactions in $B \rightarrow \pi\pi$

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# Outline

Phenomenological Importance of  $B \rightarrow \pi\pi$

QCD-Factorization

Feynman Diagrams for  $B \rightarrow \pi\pi$

Hard Spectator Interactions

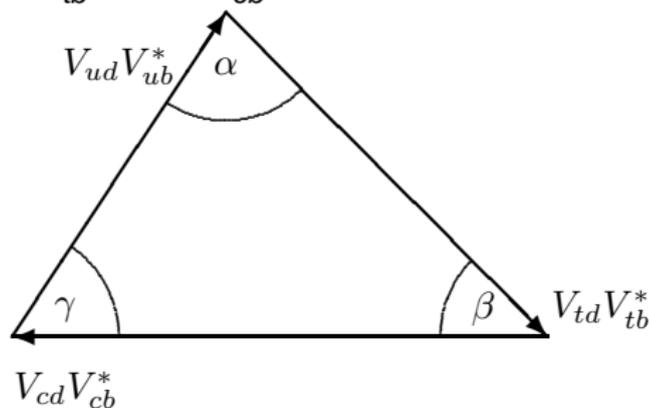
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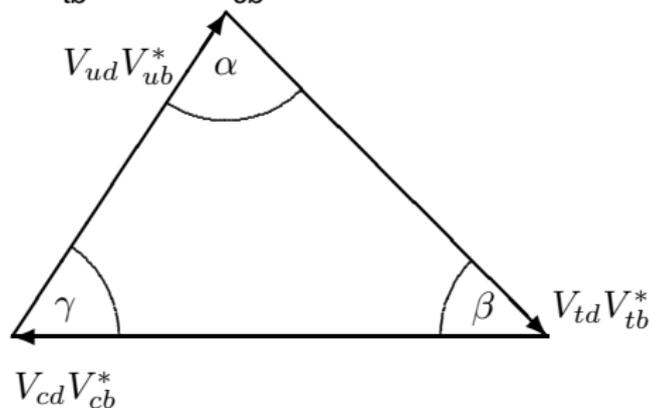


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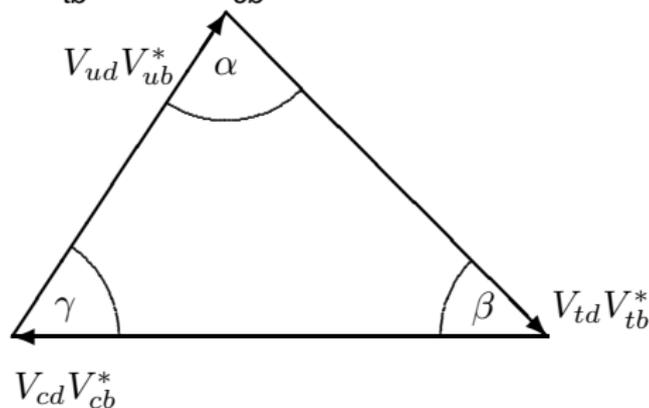
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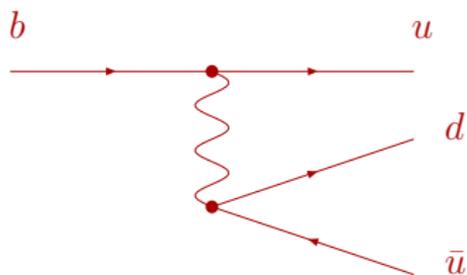


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- ▶ The angle  $\beta$  is well determined by  $B \rightarrow J/\psi K_S$
- ▶ For the angles  $\alpha, \gamma$  controlling of hadronic quantities in decays like  $B \rightarrow \pi\pi$  is needed

# B-decays and CKM-angles

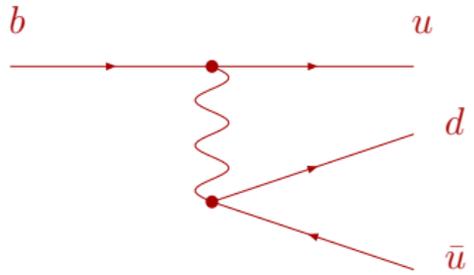
## ► Tree diagram



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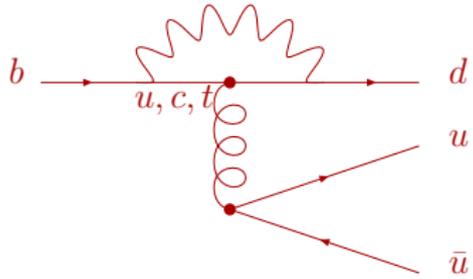
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$$\begin{aligned}\langle \pi\pi | \hat{H} | B \rangle &= F^{B \rightarrow \pi}(0) \int_0^1 du T^I(u) \phi_\pi(u) \\ &+ \int_0^1 d\xi dv du \phi_B(\xi) \phi_\pi(u) \phi_\pi(v) T^{II}(\xi, u, v)\end{aligned}$$

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⇒ Separation of short- and long-distance effects

# When does QCD-Factorization not work?

- ▶ Infrared (IR) singularities remain from loop integrals
  - ▶ Formally: IR singularities cannot be absorbed into counter terms
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- ▶ Endpoint singularities occur from integrals over meson wave functions i.e.

$$\int_0^1 d\xi dv du \phi_\pi(u) \phi_\pi(v) \phi_B(\xi) T(\xi, u, v)$$

diverges for  $u, v \rightarrow 0, 1$  or  $\xi \rightarrow 0$

⇒ Contribution of soft constituent quarks is dominant

# How does QCD-Factorization work?

- ▶ Perturbative expansion of the matrix element:

$$\begin{aligned}\langle \pi\pi | \hat{H} | B \rangle &= A^{(0)} + \alpha_s A^{(1)} + \dots \\ &= T \otimes \phi \\ &= T^{(0)} \otimes \phi^{(0)} + \\ &\quad \alpha_s [T^{(1)} \otimes \phi^{(0)} + T^{(0)} \otimes \phi^{(1)}] + \dots\end{aligned}$$

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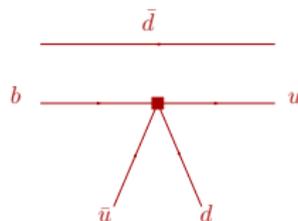
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- ▶ IR and endpoint singularities of  $A^{(1)}$  and  $\phi^{(1)}$  cancel each other  
 $\Rightarrow T$  stays finite

# Feynman diagrams for $B \rightarrow \pi\pi$

► LO:

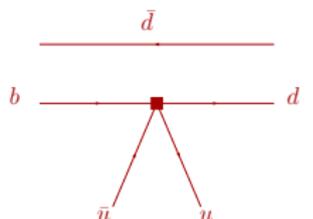
►  $\bar{B}^0 \rightarrow \pi^+\pi^-$ :



A Feynman diagram for the decay  $\bar{B}^0 \rightarrow \pi^+\pi^-$ . It shows a  $b$  quark line entering from the left and a  $u$  quark line entering from the right. They meet at a central vertex (represented by a red square). From this vertex, two lines go downwards to  $\bar{u}$  and  $d$  quarks. A  $\bar{d}$  quark line is shown above the vertex, connected to the  $b$  quark line.

$$\sim C_1 + \frac{C_2}{N_C} \approx 1$$

►  $\bar{B}^0 \rightarrow \pi^0\pi^0$ :



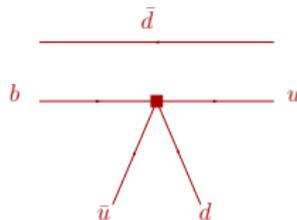
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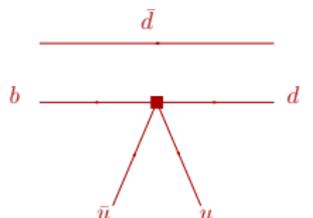
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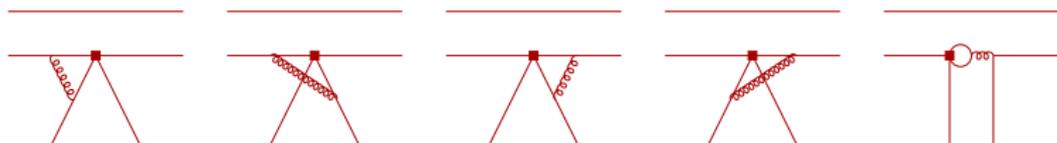


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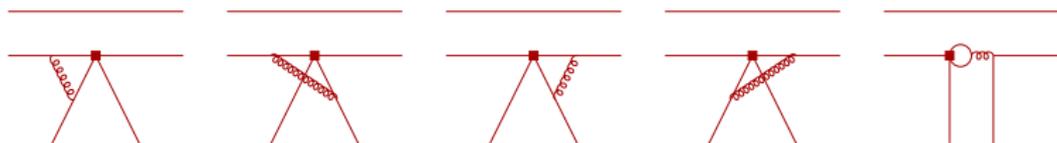
- ▶ In the case  $\bar{B}^0 \rightarrow \pi^0\pi^0$  LO is color suppressed and NLO is expected to be more important

► Form factor contributions:



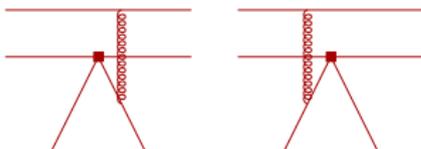
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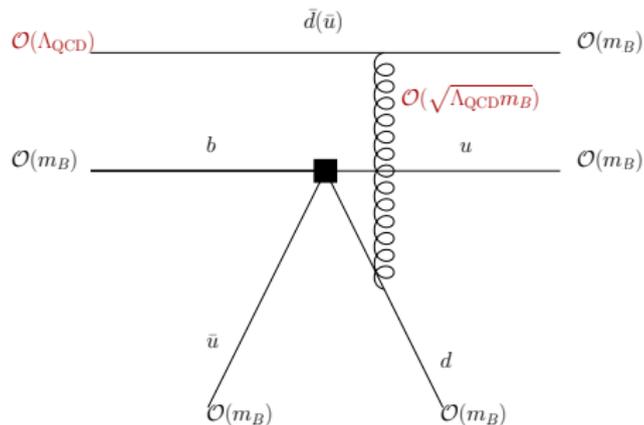
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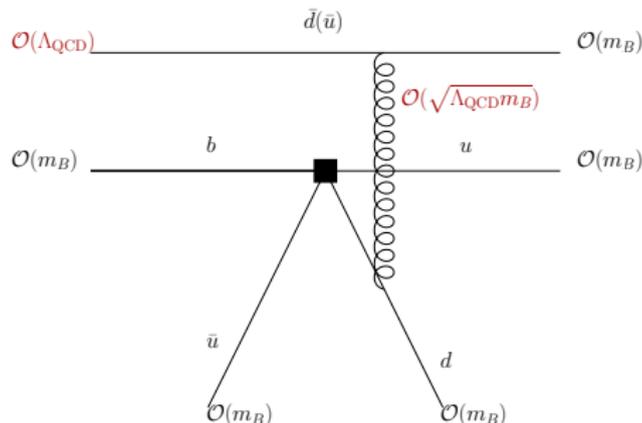
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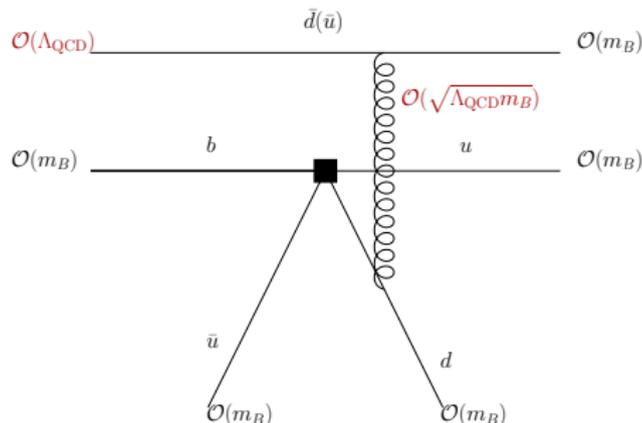
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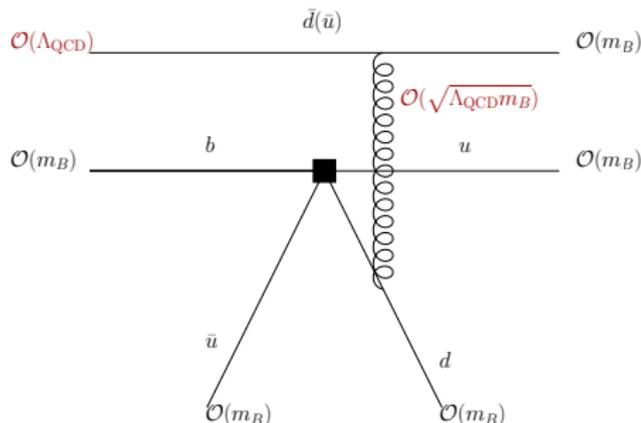
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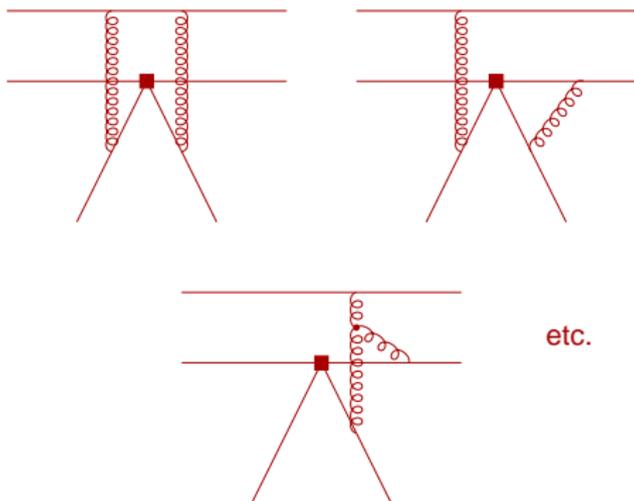
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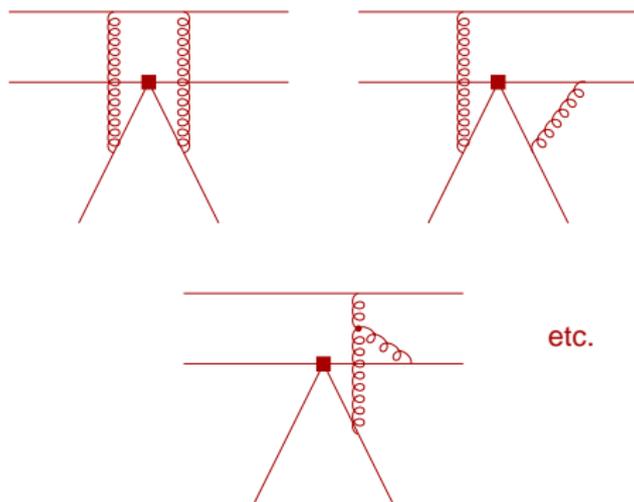
⇒ At NNLO hard spectator interactions are expected to be important

# Hard spectator interactions in NNLO



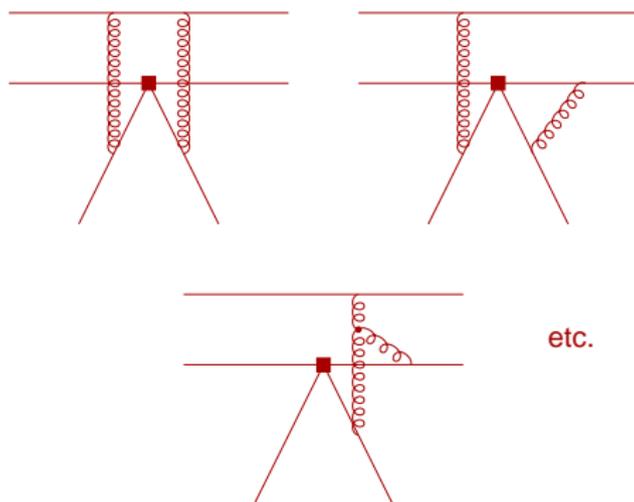
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# Hard spectator interactions in NNLO



- ▶ Altogether about fifty diagrams
- ▶ Only one loop calculations despite  $\mathcal{O}(\alpha_s^2)$
- ▶ Feynman integrals contain up to five propagators and three linearly independent momenta  
⇒ Expansion in  $\Lambda_{\text{QCD}}/m_B$  on the level of integrands rather than of Feynman integrals

- ▶ CP averaged branching ratios

Mode	NLO	NNLO	Experiment
$B^- \rightarrow \pi^- \pi^0$	5.1	$5.5 \pm 1.0$	$5.5 \pm 0.6$
$\bar{B}^0 \rightarrow \pi^+ \pi^-$	5.2	$5.0 \pm 1.3$	$5.0 \pm 0.4$
$\bar{B}^0 \rightarrow \pi^0 \pi^0$	0.7	$0.73^{+0.8}_{-0.6}$	$1.45 \pm 0.3$

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- ▶ NNLO taken from Beneke and Jäger. My work is still in progress

# Summary

- ▶  $B$ -decays play an important role for determining the CKM angles
- ▶ In leading power QCD-effects in  $B$ -decays can be handled in the framework of QCD-factorization
- ▶ Hard spectator interactions are important because they introduce large logarithms  $\ln \frac{\Lambda_{\text{QCD}}}{m_B}$