Interplay between the LHC and the ILC

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Some of the open questions in particle physics

What is the mechanism of electroweak symmetry breaking?

• SM, SUSY, . . . :

Higgs mechanism, elementary scalar particle(s)

- Strong electroweak symmetry breaking (technicolour, ...): new strong interaction, non-perturbative effects, resonances, ...
- Higgsless models in extra dimensions: boundary conditions for SM gauge bosons and fermions on Planck and TeV branes in higher-dimensional space

 \Rightarrow New phenomena required at the TeV scale

Open questions, cont'd

Hierarchy problem: how can the Planck scale be so much larger than the weak scale?

Weak-scale physics is affected by high (cutoff-) scale Λ via quantum corrections

If $\Lambda \approx M_{Planck} \Rightarrow$ would expect to drive up all physics to the Planck scale

Nature has found a way to prevent this; no explanation in SM \Rightarrow Expect new physics at the TeV scale

Supersymmetry: large corrections cancel out because of symmetry fermions \Leftrightarrow bosons

Extra-dim. models: fundamental Planck scale is $\sim TeV$ (large extra dim.), hierarchy of scales related to "warp factor" (RS)

Do the gauge interactions unify? SM: no, SUSY: yes, Extra-dim. models: ??

What is dark matter?

SM: —, SUSY: lightest SUSY particle (LSP): neutralino, gravitino, axino, ...,

Extra-dim. models: lightest massive KK mode (LKP), ...

What is dark energy? SM: ???, SUSY: ???, Extra-dim. models: ???

LHC and ILC will explore a new territory: TeV scale (1 TeV $\Leftrightarrow 2 \times 10^{-19}$ m)



How can we probe the TeV scale? — LHC

The Large Hadron Collider (LHC)

under construction at CERN (Geneva), scheduled to take first data in 2007, expected to run for about 15–20 years

Proton–proton scattering at 14 TeV: composite objects of quarks and gluons, bound together by strong interaction



Complicated scattering process, difficult to interpret 10^9 scattering events/ $s \Rightarrow$ only 1 event in 10^7 will be recorded

How can we probe the TeV scale? — ILC

The International Linear Collider (ILC)

world-wide project, construction could start as early as ≈ 2009 first data in $\gtrsim 2015$?

Electron–positron scattering at \approx 0.5–1 TeV:

fundamental particles, point-like, electroweak interaction well-defined initial state, full collision energy usable, tunable





Results are easy to interpret, all events can be recorded \Rightarrow high-precision physics

Physics at the LHC and ILC in a nutshell

LHC: pp scattering at 14 TeV



Scattering process of proton constituents with energy up to several TeV, strongly interacting

⇒ huge QCD backgrounds, low signal–to–background ratios ILC: e^+e^- scattering at \approx 0.5–1 TeV



Clean exp. environment:
well-defined initial state,
tunable energy,
beam polarization, GigaZ,
γγ, eγ, e⁻e⁻ options, ...
⇒ rel. small backgrounds
high-precision physics

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LHC / ILC complementarity

The results of LHC and ILC will be highly complementary

LHC: good prospects for producing new heavy states (in particular strongly interacting new particles)

ILC: direct production (in particular colour-neutral new particles)

 high sensitivity to effects of new physics via precision measurements (cf.: WMAP vs. COBE)

LHC / ILC synergy during concurrent running

LHC: $\gtrsim 2007$, expected to run for about 15–20 years ILC: $\gtrsim 2015$?

⇒ period of concurrent running

During concurrent running: LHC \otimes ILC

- ⇒ Information obtained at the ILC can be used to improve analyses at the LHC and vice versa
- ⇒ Enable improved strategies, dedicated searches

Interplay between lepton and hadron colliders: some examples from the past

LEP + SLC + Tevatron led to many success stories:

SM at quantum level, top quark, prediction of Higgs mass

HERA observation of high Q^2 events \Rightarrow dedicated leptoquark searches at the Tevatron, results fed back to HERA analyses

Belle discovery of X(3872) \Rightarrow dedicated search at CDF & D0

 \Rightarrow independent confirmation

LHC and ILC will explore a new energy domain

- ⇒ expect ground-breaking discoveries
- \Rightarrow large potential for synergy

What is the physics gain of LHC / ILC synergy? What is the added value of concurrent running?

Exploring physics gain from LHC / ILC interplay requires:

- Detailed information on how well LHC and ILC can measure wide variety of observables in different scenarios
- Close collaboration of experts from LHC and ILC as well as from theorists and experimentalists

\Rightarrow LHC / ILC Study Group

www.ippp.dur.ac.uk/~georg/lhclc

World-wide working group, started in spring 2002

Collaborative effort of Hadron Collider and Linear Collider experimental communities and theorists

First report: hep-ph/0410364, to appear in Physics Reports

Main focus of the LHC studies so far was to investigate whether a signal can be detected

Less results available on expected precisions, measurable properties of new physics, etc.

For scenarios where detailed experimental simulations are available for both the LHC and the ILC

⇒ Quantitative results on physics gain from LHC / ILC interplay have been worked out

The issue of LHC / ILC synergy has found a lot of attention within and outside our field, of funding agencies, etc.

Examples from the U.S.

Presentation from M. Turner (NSF) to HEPAP, Sep. 23, 2004:

Complementarity

Inevitably, the question will arise of why we need a second, *less* powerful accelerator to explore the energy frontier. To educate us and to clarify this issue more generally, we would like HEPAP to form a subpanel to address complementarity, paying particular attention to the following aspects of LC/LHC complementarity:

In the context of physics discoveries (e.g., low-energy supersymmetry) made at the Tevatron or early at the LHC, what is the role of a subTeV Collider?
In the context of physics discoveries made an LC, what is the role of the LHC
In the context of "known physics" (e.g., electroweak physics), what are the synergies and complementaries of these two machines?

You should assume that the LC and LHC (with possible upgrades) will have a significant period of overlapping operation.

We are looking for a short document (20 pages), accessible to knowledgeable nonexperts (e.g., members of the EPP2010 Study, OSTP Staff and ourselves). We ask that the report be completed by April 2005.

Finally, to further educate us as well as giving us an opportunity to refine and discuss the charge with you in more detail, we suggest a half-day session at the next HEPAP meeting devoted to Complementarity.

HEPAP subpanel on LHC / ILC complementarities

- Official request from NSF (R. Staffin, M. Turner) to HEPAP on March 21, 2005: form subpanel, provide report by summer 2005
- Panel members:

J. Lykken (Co-Chair), J. Siegrist (Co-Chair), J. Bagger, B. Barrish, N. Calder, A. De Roeck, J. Feng, F. Gilman, J. Hewett, J. Huth, J. Jackson, Y.-K. Kim, R. Kolb, K. Matchev, H. Murayama, R. Weiss

Report to the Elementary Particle Physics (EPP) 2010 Committee from HEPAP (July 27, 2005):

"Discovering the Quantum Universe"

www.science.doe.gov/hep/hepap.shtm

EPP 2010 Decadal Survey

U.S. National Academy of Science reviews each field of physics every ten years (last survey of Particle Physics was completed in 1998)

EPP 2010 charge:

- Identify, articulate, and prioritize the scientific questions and opportunities that define elementary particle physics
- Recommend a 15-year implementation plan with realistic, ordered priorities to realize these opportunities
- \Rightarrow emphasis on ranking science priorities

Some of the EPP questions on the ILC:

- What are the physics arguments for operating a Linear Collider during the same time frame as the LHC?
- How would the combination of the LHC and a Linear Collider answer questions that could not be addressed by either machine alone?
- What physics would a Linear Collider address that would be impossible to probe at the LHC?

 \Rightarrow The LHC / ILC Study Group was approached by the EPP, asked to provide a response to these questions

Response of the LHC / ILC Study Group to the EPP questions

Document prepared, writing team: J. Conway, J. Gunion, H. Haber, S. Heinemeyer, G. Moortgat-Pick, G. W.

The EPP Questions

Response from the LHC/ILC Study Group

Ground-breaking discoveries are expected from the experiments under construction at the Large Hadron Collider (LHC) and those planned for the International Linear Collider (ILC). These high-energy particle accelerators will open up a new energy domain that will allow us to examine the very fabric of matter, energy, space and time. The experimental results should reveal how particles obtain the property of mass, whether the different forces that we experience in nature are in fact different manifestations of only one fundamental force, whether space and time are embedded into a wider framework of "supersymmetric" coordinates, and whether dark matter can be produced on Earth.

The LHC and ILC will probe this new TeV energy regime (roughly equivalent to 1000 proton masses) in very different ways, as a consequence of the distinct features of the two machines. Due to its high collision energy and luminosity, the LHC has a large mass reach for the discovery of new heavy particles. The striking advantages of the ILC are its clean experimental environment, polarized beams, and tunable collision energy. The ILC can thus perform precision measurements and detailed studies of directly accessible new particles, and also has exquisite sensitivity to quantum effects of unknown physics. Indeed, the fingerprints of very high-scale new physics (e.g. very high mass particles) will often only be manifest in small effects whose measurement requires the greatest possible precision.

The need for instruments that are optimized in different ways is typical in all branches of natural sciences, for example earth- and space-based telescopes in astronomy. In high-energy physics there has historically been a great synergy between hadron colliders, which can reach the highest energies, and lepton colliders, at which high-precision measurements are possible. As an example, the precise knowledge of the Z boson mass from LEP and precise measurement of its decay properties led to the prediction of a heavy top quark. Its mass was well beyond the energy reach of LEP but accessible to the Tevatron. Following the Tevatron's discovery of the top quark, its mass was determined. Subsequently the Tevatron and LEP measured the W boson mass with high precision. In combination, these measurements point tantalizingly toward a light Higgs boson.



Precise measurements from concurrent running of LEP and the Tevatron experiments have brought us to the threshold of discovering the Higgs boson.

We expect an even greater synergy between the LHC and ILC. Discoveries made at the LHC will guide the operation of the ILC, and the precision ILC measurements can make it possible for the LHC to extract subtle signals for new physics and particles that may have escaped detection. Ultimately both machines will be needed to definitively connect TeV-scale measurements with the underlying theoretical structure.

In general, the LHC can most readily discover the heavy states of new physics that are "strongly coupled" (that is, produced via the strong interaction). These strongly coupled states typically decay via complicated cascades into new "weakly coupled" particles. The ILC is ideal for directly producing and detecting these weakly coupled particles.

How would the combination of the LHC and a Linear Collider answer questions that could not be addressed by either machine alone?

Precision ILC measurement of the properties of these particles are essential in understanding the strongly coupled ones and their decay patterns. Moreover, ILC measurements of quantum effects can be combined with direct LHC and ILC measurements to infer the existence and properties of additional heavy states at first missed by the LHC and too massive to be directly produced

at the ILC. In many cases, these could then be directly discovered using modified LHC procedures.

As an example, the existence and properties of heavy Higgs bosons and/or difficult-to-detect scalar Higgslike particles associated with extra dimensions can be inferred from precision ILC Higgs measurements. A dedicated LHC search can then confirm their existence. In supersymmetry and extra-dimension theories, the LHC and the ILC will typically access different parts of the spectrum of new states.

Summary

There will be a profound synergy between the physics results from the LHC and those from the ILC. The two machines complement and supplement one another in many ways, and concurrent operation will maximize Together the ILC and LHC can measure the unified supersymmetry masses much more precisely than either machine alone.

the impact of both. Understanding the physics of the TeV scale will have an important impact on cosmology and other fields, as well as give timely guidance regarding future facilities. The sconer the ILC can be brought into operation, the sconer these benefits can be exploited. Optimal use of the capabilities of both machines will greatly improve our knowledge of the fundamental nature of matter, energy, space and time.

We urge the international high energy physics community and the governments of all the countries involved to strive to make the ILC a reality in the coming decade.

See the full report of the LHC/ILC Study Group at http://arxiv.org/abs/hep-ph/0410364

John Conway, Jack Gunion, Howard Haber, Sven Heinemeyer, Gudrid Moortgat-Pick, and Georg Weiglein

May 16, 2005



Released to the EPP members at the EPP meeting in May '05 at Fermilab, see www.ippp.dur.ac.uk/~georg/lhcilc

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Some examples of recent results

In the following: some examples of LHC / ILC synergy

Focus on:

- Electroweak symmetry breaking
- Supersymmetry

Many more results on electroweak symmetry breaking, electroweak and QCD precision physics, SUSY, new gauge theories, extra dimensions, ... in hep-ph/0410364

Electroweak symmetry breaking

ILC will determine electroweak symmetry breaking mechanism regardless of its nature

Higgs discovery possible independent of decay modes "Golden" production channel: $e^+e^- \rightarrow ZH$, $Z \rightarrow e^+e^-$, $\mu^+\mu^-$

ILC is a "Higgs factory"

- e.g.: $E_{\rm CM} = 800$ GeV, 1000 fb⁻¹, $M_{\rm H} = 120$ GeV:
- $\Rightarrow \approx 160000$ Higgs events in "clean" experimental environment
- ⇒ Precise measurement of Higgs mass, couplings, determination of Higgs spin and quantum numbers,
- ⇒ Verification of Higgs mechanism in model-independent way distinction between different possible manifestations: extended Higgs sector, invisible decays, Higgs—radion mix., Interplay between the LHC and the ILC, G. Weiglein, MPI München 02/2006 – p.24

Higgs coupling determination at LHC \oplus ILC

LHC: no absolute measurement of total production cross section (no recoil method like LEP, ILC: $e^+e^- \rightarrow ZH$, $Z \rightarrow e^+e^-, \mu^+\mu^-$)

Measurement of $\sigma \times BR$: narrow-width approximation

$$\Rightarrow \quad \sigma(H) \times \mathrm{BR}(H \to a + b) = \frac{\sigma(H)^{\mathrm{SM}}}{\Gamma_{\mathrm{prod}}^{\mathrm{SM}}} \cdot \frac{\Gamma_{\mathrm{prod}}\Gamma_{\mathrm{decay}}}{\Gamma_{\mathrm{tot}}}$$

Observation of different channels (or upper bound from non-observation) provides information on combinations of $\Gamma_g, \Gamma_W, \Gamma_Z, \Gamma_\gamma, \Gamma_\tau, \Gamma_b, Y_t^2$

Large uncertainty on $H \rightarrow b\overline{b}$, ...

⇒ LHC can directly determine only ratios of couplings

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Higgs coupling determination at LHC \oplus ILC

Need additional (mild) theory assumption to obtain absolute values of the couplings at the LHC:

[M. Dührssen, S. Heinemeyer, H. Logan, D. Rainwater, G. W., D. Zeppenfeld '04]

Assume: *HVV* couplings bounded by SM values: $\Gamma_V \leq \Gamma_V^{\text{SM}}$ (V = W, Z), valid in wide class of models: SM, MSSM, ... \Rightarrow Upper bound on Γ_V

Observation of Higgs production \Rightarrow Lower bound on production couplings and Γ_{tot}

Observation of $H \rightarrow VV$ in WBF \Rightarrow Determines $\Gamma_V^2/\Gamma_{tot} \Rightarrow$ Upper bound on Γ_{tot}

 \Rightarrow Absolute determination of Γ_{tot} and Higgs couplings

Use LHC / ILC interplay use ILC input instead of theory assumption: top Yukawa coupling

Only crude measurement of *tth* coupl. at 500 GeV ILC (light Higgs)

Precision measurement requires ILC with 800–1000 GeV

LHC measures ($\sigma \times BR$)

⇒ Yukawa coupling can
 be extracted if precise
 measurement of Higgs
 BR's from ILC are used

LHC \oplus ILC (500 GeV): [K. Desch, M. Schumacher '04]



LHC / ILC interplay: fit of Higgs couplings

Fit of Higgs couplings with input from LHC and ILC

[K. Desch, M. Dührssen, S. Heinemeyer, H. Logan, D. Rainwater, G. W., D. Zeppenfeld '05, preliminary]

Combined fit for all LHC channels with ILC input:

 $M_{\rm H}, \, \sigma(e^+e^- \to HZ), \, {\rm BR}(H \to b\bar{b}, \tau^+\tau^-, gg, WW^*), \\ \sigma(e^+e^- \to \nu\bar{\nu}H) \times {\rm BR}(H \to b\bar{b})$

Comparison: LHC only vs. LHC \oplus ILC



 $\Rightarrow \text{ higher accuracy on } g_{\text{Ht}\bar{\text{t}}} \text{ (and also } g_{\text{H}\gamma\gamma} \text{) than for LHC alone} \\ \text{(+ theory) and ILC}_{500} \text{ alone: } \Delta g_{\text{Ht}\bar{\text{t}}} / g_{\text{Ht}\bar{\text{t}}} \approx 11-14\% \\ \text{Interplay between the LHC and the ILC, G. Weiglein, MPI München 02/2006 - p.2:} \end{cases}$

Determination of M_A from heavy Higgs decays into SUSY particles at the LHC

[F. Moortgat '04]

 $H, A \rightarrow \tilde{\chi}_2^0 \tilde{\chi}_2^0$: Four lepton invariant mass distribution for $M_{\rm A} = 393 \pm 20 \; {\rm GeV}$ (left) and $M_1 = 100 \pm 10 \; {\rm GeV}$ (right)



 \Rightarrow Precise knowledge of LSP mass from ILC crucial for determination of $M_{\rm A}$

Indirect constraints on M_A from Higgs BR measurements at

the ILC using LHC / ILC input



 $\Rightarrow \textbf{Sensitive indirect bounds on } M_{A} \textbf{ only with high-precision} \\ \textbf{measurements, LHC} \oplus \textbf{ILC information}_{\text{Interplay between the LHC and the ILC, G. Weiglein, MPI München 02/2006 – p.2}}$

Precision Higgs physics

Large coupling of Higgs to top quark



One-loop correction $\sim G_{\mu} m_{\rm t}^4$

 $\Rightarrow M_{\rm H}$ depends sensitively on $m_{\rm t}$ in all models where $M_{\rm H}$ can be predicted (SM: $M_{\rm H}$ is free parameter)

SUSY as an example: $\Delta m_{\rm t} \approx \pm 3 \text{ GeV} \Rightarrow \Delta m_{\rm h} \approx \pm 3 \text{ GeV}$

⇒ Precision Higgs physics needs precision top physics LHC: $\Delta m_{\rm h} \approx 0.2 \text{ GeV}$, $\Delta m_{\rm t} \gtrsim 1 \text{ GeV}$, ILC: $\Delta m_{\rm t} \lesssim 0.1 \text{ GeV}$

Higgs-radion mixing

[M. Battaglia, S. De Curtis, A. De Roeck, D. Dominici, J. Gunion '03]

Models with 3-branes in extra dimensions predict radion ϕ , can mix with the Higgs

 \Rightarrow Higgs properties modified, can be difficult to detect at the LHC

LHC may observe radion instead

ILC guarantees Higgs observation over full parameter space

⇒ precision measurements at ILC crucial to disentangle the nature of the observed state

LHC: large sensitivity to production of Kaluza-Klein excitations

Electroweak symmetry breaking without Higgs

If no light Higgs boson exists

⇒ dynamics of electroweak symmetry breaking can be probed in quasi-elastic scattering processes of W and Z at high energies

LHC / ILC sensitive to different scattering channels, yield complementary information

- LHC: direct sensitivity to resonances
- ILC: detailed measurements of cross sections and angular distributions
- ⇒ combination of LHC results with ILC data on cross-section rise essential for disentangling new states

Strong electroweak symmetry breaking

Sensitivity of LHC and ILC measurements to signals of strong electroweak symmetry breaking:

[American LC WG '01]

Signal significance in σ for various masses M_{ρ} of vector resonance in $W_{\rm L}W_{\rm L}$ scattering:



 $\Rightarrow Strong electroweak symmetry breaking scenarios can be probed in detail at LHC \oplus ILC$

The Minimal Supersymmetric Standard Model (MSSM)

Superpartners for Standard Model particles:

$$\begin{bmatrix} u, d, c, s, t, b \end{bmatrix}_{L,R} \begin{bmatrix} e, \mu, \tau \end{bmatrix}_{L,R} \begin{bmatrix} \nu_{e,\mu,\tau} \end{bmatrix}_{L} \qquad \text{Spin } \frac{1}{2}$$

$$\begin{bmatrix} \tilde{u}, \tilde{d}, \tilde{c}, \tilde{s}, \tilde{t}, \tilde{b} \end{bmatrix}_{L,R} \begin{bmatrix} \tilde{e}, \tilde{\mu}, \tilde{\tau} \end{bmatrix}_{L,R} \begin{bmatrix} \tilde{\nu}_{e,\mu,\tau} \end{bmatrix}_{L} \qquad \text{Spin } 0$$

$$g \qquad \underbrace{W^{\pm}, H^{\pm}}_{\tilde{1},2} \qquad \underbrace{\gamma, Z, H_{1}^{0}, H_{2}^{0}}_{1,2,3,4} \qquad \text{Spin } 1 \text{ / Spin } 0$$

$$\underbrace{\tilde{g}} \qquad \widetilde{\chi}_{1,2}^{\pm} \qquad \underbrace{\tilde{\chi}_{1,2,3,4}^{0}} \qquad \underbrace{Spin } \frac{1}{2}$$

Enlarged Higgs sector: two Higgs doublets, physical states: h^0, H^0, A^0, H^{\pm}

SUSY breaking

MSSM: no particular SUSY breaking mechanism assumed, parametrization of possible soft SUSY-breaking terms

 \Rightarrow relations between dimensionless couplings unchanged \Rightarrow no quadratic divergencies (hierarchy problem)

Most general case: 105 new parameters Good phenomenological description for universal breaking terms

Specific models for soft-SUSY breaking:

"Hidden sector": → Visible sector: SUSY breaking MSSM "Gravity-mediated": mSUGRA

"Gauge-mediated": GMSB

"Anomaly-mediated": AMSB

SUSY at LHC and ILC

LHC: good prospects for strongly interacting new particles long decay chains \Rightarrow complicated final states,

e.g.: $\tilde{g} \to \bar{q}\tilde{q} \to \bar{q}q\tilde{\chi}_2^0 \to \bar{q}q\tilde{\tau}\tau \to \bar{q}q\tau\tau\tilde{\chi}_1^0$

Many states are produced at once, difficult to disentangle

 \Rightarrow It quacks like SUSY!

But ist it really SUSY? Which particles are actually produced?

Main background for SUSY is SUSY itself!

SUSY phenomenology investigated in detail for SPS 1a benchmark point: "best case scenario"

more results needed for less favourable points (in progress at ATLAS & CMS)

It quacks like SUSY, but ...

- Joes every SM particle really have a superpartner?
- do their spins differ by 1/2?
- are their gauge quantum numbers the same?
- are their couplings identical?
- do the SUSY predictions for mass relations hold, ...?

Even when we are sure that it is actually SUSY, we will still want to know:

- is the lightest SUSY particle really the neutralino, or the stau or the sneutrino, or the gravitino or ...?
- is it the MSSM, or the NMSSM, or the mNSSM, or the N²MSSM, or …?
- what are the experimental values of the 105 (or more) SUSY parameters?
- Joes SUSY give the right amount of dark matter?
- what is the mechanism of SUSY breaking?

We will ask similar questions for other kinds of new physics

When and how will we find out?

- How much will we learn from the LHC alone?
- How much more will we know once we have ILC data?
- What is the added value of having the LHC and the ILC run concurrently?

SUSY at the ILC: clean signatures, small backgrounds

⇒ precise determination of masses, spin, couplings, mixing angles, complex phases ...,

good prospects for weakly interacting SUSY particles precision measurement of mass of lightest SUSY particle (factor 100 improvement)

Production of new particles at the ILC

Tunable energy \Rightarrow can run directly at threshold

Example: Determination of mass and spin of SUSY particle $\tilde{\mu}_R$ from production at threshold: [TESLA TDR '01]

$$\Rightarrow \quad \frac{\Delta m_{\tilde{\mu}_R}}{m_{\tilde{\mu}_R}} < 1 \times 10^{-3}$$

 \Rightarrow test of J = 0 hypothesis





SUSY analyses at LHC / ILC

[M. Chiorboli, B.K. Gjelsten, J. Hisano, K. Kawagoe, E. Lytken, U. Martyn, D. Miller, M. Nojiri, P. Osland, G. Polesello, A. Tricomi '03]

Cascade decays: complicated decay chains for squarks and



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LHC analysis with ILC input: mass of lightest SUSY particle (LSP)

Reconstruction of the states in decay chain requires precise knowledge of LSP mass



 $\Rightarrow Precision measurement of m_{LSP} at ILC leads to significant$ improvement in determination of slepton, squark and gluino masses at the LHC Interplay between the LHC and the ILC, G. Weiglein, MPI München 02/2006 – p.4

Detailed analysis for SPS1a benchmark scenario: potential

of LHC (300 fb $^{-1}$) alone and LHC \oplus ILC

	LHC	LHC⊕ILC	LHC⊕ILC accuracy
$\Delta m_{\tilde{\chi}^0_1}$	4.8	0.05 (input)	limited by LHC jet en-
$\Delta m_{\tilde{l}_R}$	4.8	0.05 (input)	ergy scale resolution
$\Delta m_{ ilde{\chi}^0_2}$	4.7	0.08	
$\Delta m_{\tilde{q}_L}$	8.7	4.9	SPS 1a benchmark
$\Delta m_{\tilde{q}_R}$	11.8	10.9	scenario:
$\Delta m_{ ilde{g}}$	8.0	6.4	
$\Delta m_{\tilde{b}_1}$	7.5	5.7	tavourable scenario for
$\Delta m_{\tilde{b}_2}$	7.9	6.2	
$\Delta m_{\tilde{l}_L}$	5.0	0.2 (input)	
$\Delta m_{ ilde{\chi}_4^0}$	5.1	2.23	

 \Rightarrow ILC input improves accuracy significantly

Determination of the gluino mass: using ILC input to resolve ambiguities at the LHC

[B. Gjelsten, D. Miller, P. Osland '05]

Mass determination from cascade decays: invert endpoint formulas, fit masses

⇒ yields correct minimum
+ false minima (can be off by 10–20 GeV)

Determination of the gluino mass: using ILC input to resolve ambiguities at the LHC

Problem due to compositeness of formulas:

If masses are close to border of 'region', may find a similar-quality or better minimum in 'other' region



ILC input on LSP mass

⇒ correct minimum can be identified, ambiguities resolved

Determination of GUT-scale parameters m_0 , $m_{1/2}$ in the mSUGRA scenario

LHC only vs. LHC

HC



⇒ Combined information from LHC and ILC yields drastic improvement

LHC / ILC interplay in SUSY searches

- Precise determination of the properties of the SUSY particles accessible at the ILC
 - ⇒ identify whether or not these particles appear in the decay cascades at the LHC
- Precision measurement of the LSP mass at the ILC as input for LHC analyses
 - ⇒ significantly improves precision of mass determination of heavier SUSY particles at the LHC
- From part of the SUSY spectrum accessible at the ILC
 can predict the properties of heavier particles
 tell the LHC where to look

"Telling the LHC where to look"

ILC prediction transforms search for edge in di-lepton mass spectrum into single hypothesis test

- \Rightarrow Increase of LHC statistical sensitivity!
- \Rightarrow crucial for extracting statistically marginal signal at LHC
- ⇒ Optimised searches at the LHC: Improved selection criteria, modified triggers, different running strategy, ...

Compare the situation at LEP, where we had a statistically marginal excess of Higgs-like events

Suppose a collider running concurrently with LEP had predicted a Higgs boson with $M_{\rm H} = 115 \pm 1 \text{ GeV}$

this would have certainly affected the running strategy of LEP Interplay between the LHC and the ILC, G. Weiglein, MPI München 02/2006 - p.4

Example of LHC / ILC interplay

SUSY case study where the lightest neutralino and chargino states $(\chi_1^0, \chi_2^0, \chi_1^{\pm})$ are precisely measured at the ILC

[K. Desch, J. Kalinowski, G. Moortgat-Pick, M. Nojiri, G. Polesello '04]

- \Rightarrow Identification of $(\chi_1^0, \chi_2^0, \chi_1^{\pm})$ in LHC decay chains
- ⇒ Determination of all parameters in neutralino/chargino sector
- \Rightarrow Prediction of masses, decay prop. of all neutralinos, charginos
- ⇒ Prediction of masses of particles that are too heavy to be produced at the ILC but are produced with low statistics at the LHC, e.g. heaviest neutralino: $m_{\tilde{\chi}_{4}^{0}} = 378.3 \pm 8.8$ GeV
- ⇒ With this information the heaviest neutralino can be identified at the LHC using a dilepton "edge"

Search for the heaviest neutralino at LHC following the prediction from ILC



The new particle can be identified at the LHC via this "edge"

- $\Rightarrow \mbox{Determination of } m(\tilde{\chi}_4^0) \\ \mbox{with high precision}$
- \Rightarrow Crucial test of the model

Feeding information on $m(\tilde{\chi}_4^0)$ back into ILC analysis \Rightarrow Improved accuracy of parameter determination at ILC

ILC analysis with LHC input

Determination of neutralino parameter M_1 and chargino mixing angles $\cos \phi_{\rm L}$, $\cos \phi_{\rm R}$:





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Distinguishing between MSSM and NMSSM

Case study of scenario where Higgs sector and light neutralino / chargino spectra and cross sections are almost identical in the two models

[G. Moortgat-Pick, S. Hesselbach, F. Franke, H. Fraas '05]

Parameter determination as in MSSM \Rightarrow no contradiction

ILC input \Rightarrow prediction of masses and mixing character of heavy neutralinos

⇒ Detection of $\tilde{\chi}_3^0$ at LHC yields contradiction with MSSM prediction

 \Rightarrow Evidence for NMSSM

Mixing character of $\tilde{\chi}_3^0$, $\tilde{\chi}_4^0$: MSSM vs. NMSSM

[G. Moortgat-Pick, S. Hesselbach, F. Franke, H. Fraas '05]



 $\Rightarrow ILC prediction of mixing character yields distiction of NMSSM from MSSM Interplay between the LHC and the ILC, G. Weiglein, MPI München 02/2006 – p.5$

Prospects for SUSY parameter determination at LHC and ILC investigated in detail for SPS 1a benchmark point: "bulk" region of mSUGRA scenario ('best case scenario')

 $m_0 = 100 \,\mathrm{GeV}$, $m_{1/2} = 250 \,\mathrm{GeV}$, $A_0 = -100 \,\mathrm{GeV}$, $\tan \beta = 10$, $\mu > 0$

Most observables depend on variety of SUSY parameters

⇒ Need global fit to large set of observables
[R. Lafaye, T. Plehn, D. Zerwas '04] [P. Bechtle, K. Desch, P. Wienemann '04]

 \Rightarrow Reliable determination of SUSY parameters only possible from combined LHC \oplus ILC data

Determination of SUSY parameters: global fit

Use *Fittino* to compare the ability of LHC only and LHC \oplus ILC for SPS1a' point

[P. Bechtle, K. Desch, P. Wienemann '05]

LHC input:

mass measurements and precisions as above

- + assumption on $\tilde{t}_{1,2}$ mass measurement
- + ratios of Higgs branching ratios (see above)

Fittino: LHC only vs. LHC \oplus ILC

Parameter	"True" value	ILC Fit value	Uncertainty	Uncertainty			
			(ILC+LHC)	(LHC only)			
$\tan\beta$	10.00	10.00	0.11	6.7			
μ	400.4 GeV	400.4 GeV	1.2 GeV	811. GeV			
X_{τ}	-4449. GeV	-4449. GeV	20.GeV	6368. GeV			
$M_{\tilde{e}_R}$	115.60 GeV	115.60 GeV	0.27 GeV	39. GeV			
$M_{\tilde{\tau}_R}$	109.89 GeV	109.89 GeV	0.41 GeV	1056. GeV			
$M_{\tilde{e}_L}$	181.30 GeV	181.30 GeV	0.10 GeV	12.9 GeV			
$M_{\tilde{\tau}_L}$	179.54 GeV	179.54 GeV	0.14 GeV	1369. GeV			
X_t	-565.7 GeV	-565.7 GeV	3.1 GeV	548. GeV			
X_b	-4935. GeV	-4935. GeV	1284. GeV	6703. GeV			
$M_{\tilde{u}_R}$	503. GeV	503. GeV	24. GeV	25. GeV			
$M_{\tilde{b}_B}$	497. GeV	497. GeV	8. GeV	1269. GeV			
$M_{\tilde{t}_B}$	380.9 GeV	380.9 GeV	2.5 GeV	753. GeV			
$M_{ ilde{u}_L}$	523. GeV	523. GeV	10. GeV	19. GeV			
$M_{\tilde{t}_L}$	467.7 GeV	467.7 GeV	3.1 GeV	424. GeV			
M_1	103.27 GeV	103.27 GeV	0.06 GeV	8.0 GeV			
M_2	193.45 GeV	193.45 GeV	0.10 GeV	132. GeV			
M_3	569. GeV	569. GeV	7. GeV	10.1 GeV			
m_{Arun}	312.0 GeV	311.9 GeV	4.6 GeV	1272. GeV			
m_t	178.00 GeV	178.00 GeV	0.050 GeV	0.27 GeV			
χ^2 for unsmeared observables: 5.3×10^{-5}							

⇒ most of the Lagrangian parameters can hardly be constrained by LHC data alone

How precisely do we need to know the SUSY parameters?

Dark matter relic density: measurement vs. prediction

Aim:

match the precision of the relic density measurement with the prediction based on collider data

⇒ sensitive test of SUSY dark matter hypothesis

Relic density measurement:

current (WMAP): $\approx 10\%$

future (Planck): $\approx 2\%$

Prediction of the dark matter density

LHC analysis for SPS 1a point: [G. Polesello, D. Tovey '04] [P. Janot '04]

- LHC input for SPS 1a point
- Theoretical assumption: mSUGRA is correct SUSY-breaking mechanism (4 parameters and a sign)
- ⇒ Prediction of relic density with $\approx 3\%$ accuracy for full LHC-design integrated luminosity

Significantly better accuracy with ILC input (same th. assumption)

Much worse prospects at the LHC for less favourable parameter points [*P. Janot '04*]

Prediction of the dark matter density

Ultimately cannot assume SUSY-breaking scenario, have to test it

Even if SUSY-breaking scenario is assumed to be known:

RGE running from high scale introduces dependence on SM input parameters and uncertainties from unknown higher orders

 \Rightarrow need precise measurement of $m_{\rm t}$

extreme case: mSUGRA focus point scenario

 \rightarrow need $m_{\rm t}$ with accuracy of $\mathcal{O}(20 \ {\rm MeV})$

 \rightarrow tough to achieve even at the ILC, hopeless at LHC

ILC accuracy on $m_{\rm t}$ also needed in mSUGRA Higgs funnel region in order to match Planck precision

[B. Allanach, G. Belanger, F. Boudjema, A. Pukhov '04]

Prediction of the dark matter density

Precision measurements at ILC crucial for precise and model-independent prediction of relic density

Example: mSUGRA coannihilation region

Precise determination of masses of LSP and lightest slepton at ILC for general MSSM (no mSUGRA assumption) [*P. Bambade, M. Berggren, F. Richard, Z. Zhang '04*]

Planck accuracy ⇔ NLSP–LSP mass difference needed with precision of 0.2 GeV

Need also precise knowledge of cross section: neutralino mixing matrix, stau sector, masses of heavy Higgs bosons [*B. Allanach, G. Belanger, F. Boudjema, A. Pukhov '04*]

Conclusions

LHC / ILC interplay is a very rich field

 \Rightarrow promising potential for important physics gain

- ATLAS and CMS are actively preparing for the start of data taking: CMS writes physics TDR, many new studies in ATLAS (full simulations, new scenarios)
 - + ongoing ILC studies
 - \Rightarrow Many new results, ideal input for studying LHC / ILC interplay
- LHC will start to take data very soon
 We are looking forward to finding out what the real data will tell us