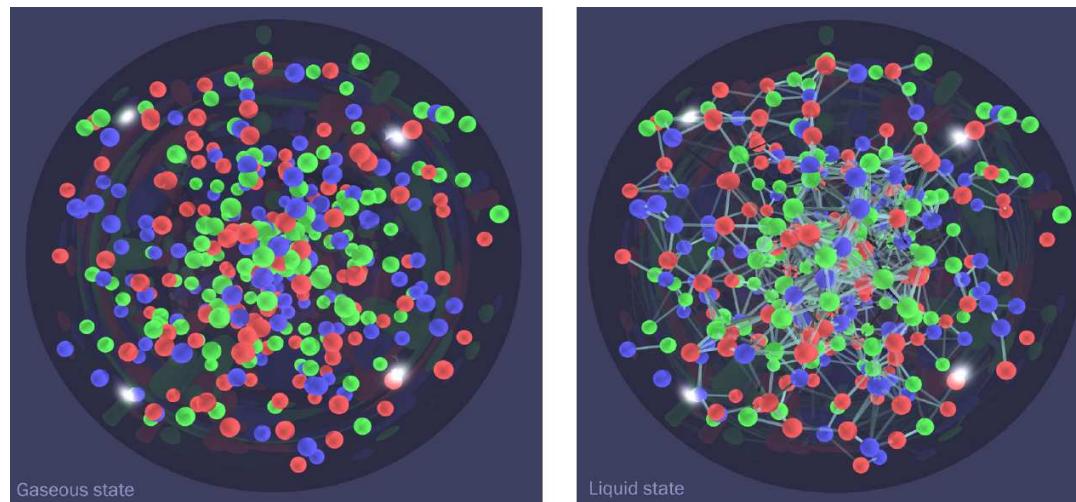


Quark Gluon Plasma and the AdS/CFT-correspondence

by Matthias Kaminski

IMPRS Particle Physics School Munich Seminar, 12th of May 2006



[Berkley Lab webpages: <http://www.lbl.gov/Science-Articles/Archive/sabl/2005/May/02-RHIC.html>]

Collision movie

Overview

1. Quark-Gluon-Plasma at the RHIC collider
2. Strings and the AdS/CFT-correspondence
3. Kubo formulae
4. Application to different theories
5. Summary and Outlook

2. Quark-Gluon-Plasma at RHIC

[BNL 73847-2005 Formal Report]

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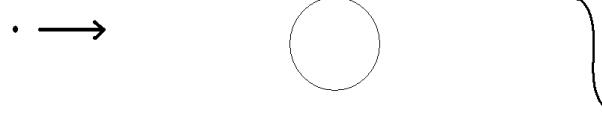
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- ☞ comparing with hydrodynamical model calculations suggests that local equilibrium is established (no evidence)
- ☞ PHENIX collaboration claims that they will be able to '... address color and *thermal transport properties* ...'

3. Strings and AdS/CFT

Question: How is String Theory related to Quantum Field Theory?

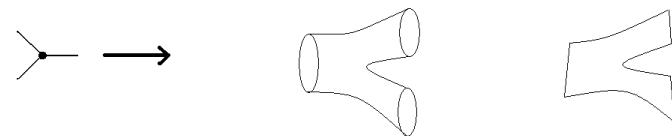
3.1 – String Theory

point particle



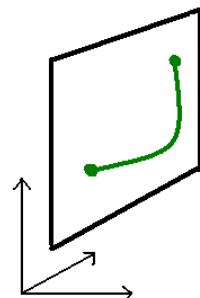
extended strings

pointlike vertex



worldsheets

Also contains D branes e.g. D3-brane



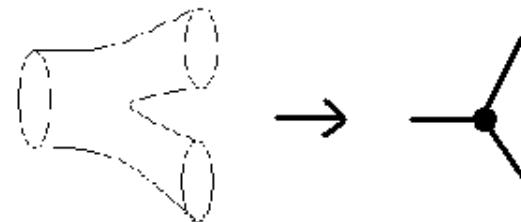
- open strings end on D-branes (determines dynamics)
- hypersurfaces embedded in 10-dimensional spacetime
- field theory 'lives' on branes (at low-energy string → ·)

3.2 – SUGRA-limit of String Theory

$$\lambda = g_s N \rightarrow \infty$$

(later $N \rightarrow \infty$)

string length: $l_{\text{string}}^4 = \frac{L^4}{4\pi(g_s N)} \rightarrow 0$



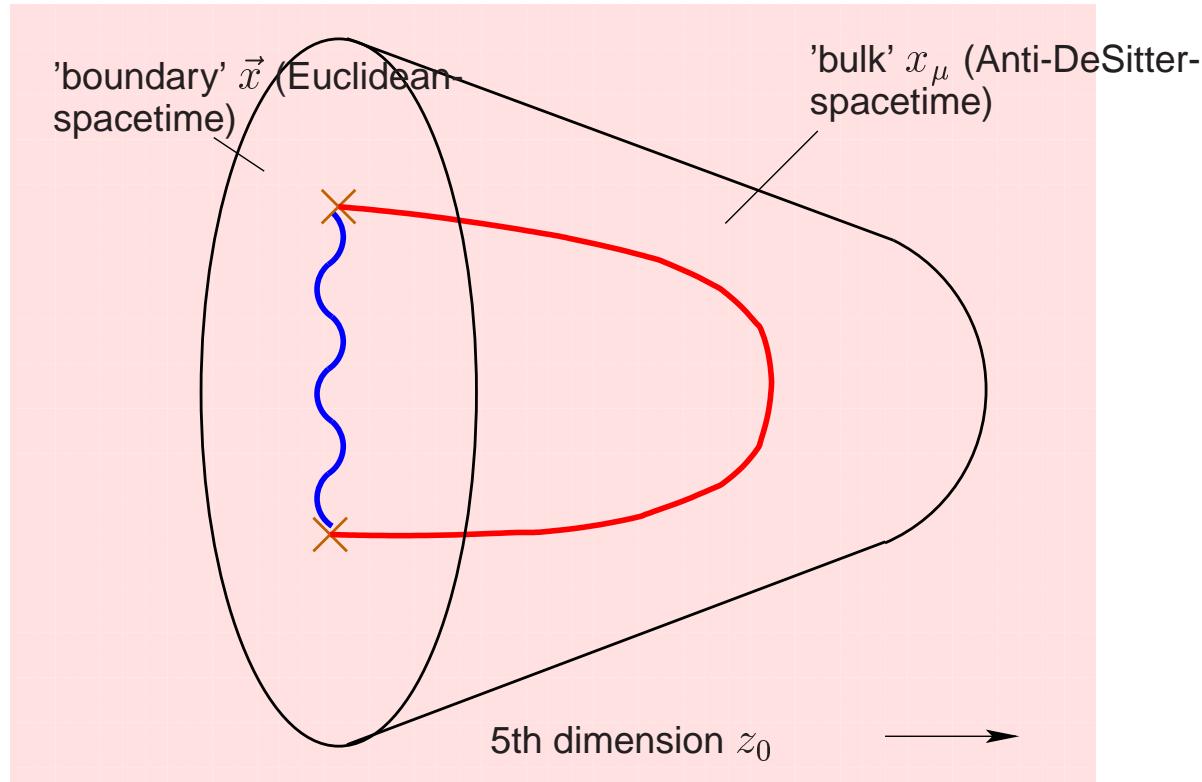
\implies SUGRA is a classical perturbative field theory expanded in $\alpha' = l_{\text{string}}^2$.

Possible solutions to SUGRA: solitons \equiv D-branes

3.3 – AdS/CFT-conjecture

[Maldacena, '97]

Using the two 'faces' of D-branes



Dictionary (by limit $\lambda \rightarrow \infty$):

operator \mathcal{O} in Conformal Field Theory (CFT) \longleftrightarrow SUGRA-field ϕ in Anti-deSitter space $\text{AdS}_5 \times \text{S}^5$

3.4 – Correlators from AdS/CFT

Formally stated for Euclidean metric:

$$\langle e^{\int d^4x \phi_0(\vec{x}) \mathcal{O}(\vec{x})} \rangle_{CFT} = Z_{\text{SUGRA}}[\phi(z, \vec{x})] \Big|_{\phi(0, \vec{x}) = \phi_0(\vec{x})} \quad (1)$$

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$$\langle \mathcal{O}(\vec{x}) \mathcal{O}(\vec{y}) \rangle_{CFT} = \frac{\delta}{\delta \phi_0(\vec{x})} \frac{\delta}{\delta \phi_0(\vec{y})} Z_{\text{SUGRA}} \Big|_{\phi = \phi_0} \quad (2)$$

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Conjectured prescription for Minkowskian metric: [Son & Starinets, '02]

$$\langle \mathcal{O} \mathcal{O} \rangle_{CFT} = -2 \mathcal{F}(k, z) \Big|_{z_B}$$

where \mathcal{F} comes from the quadratic part of the SUGRA **on-shell** action

$$S_{\text{SUGRA}}^{\text{on-shell}} = \int \frac{d^4 k}{(2\pi)^4} \phi_0(-k) \mathcal{F}(k, z) \phi_0(k) \Big|_{z=z_B}^{z=z_H} \quad (3)$$

4. Kubo formulae

[Hosoya et al., '83]

Question: How are 2-point correlators connected to physical observables?

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For the operator $\mathcal{O} \equiv T_{\mu\nu}$: [e.g. Zubarev]

$$\varrho_{\text{non-eq}} = \underbrace{\exp\left\{-\int d^3x F^\nu T_{0\nu}\right\}}_{=\beta H \text{ equilibrium}} + \int d^3x \int_{-\infty}^t dt_1 e^{\epsilon(t_1-t)} \underbrace{T_{\mu\nu}}_{\text{current}} \underbrace{\partial^\mu F^\nu}_{\text{thermo force/gradient}} \quad (6)$$

where

$$F^\nu = \beta u^\nu \quad (7)$$

$\beta \equiv$ inverse temperature, $u^\nu \equiv$ four-velocity component of fluid

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$$\langle \textcolor{blue}{T}_{\mu\nu} \rangle_{\text{non-eq}} \approx \langle T_{\mu\nu} \rangle_{\text{eq}} + \int d^3x' \int_{-\infty}^t dt' e^{\epsilon(t'-t)} \underbrace{\langle T_{\mu\nu}(\vec{x}, t), T_{\rho\sigma}(\vec{x}', t') \rangle}_{\propto \langle T_{\mu\nu}, T_{\rho\sigma} \rangle \text{ retarded}} \partial^\rho F^\sigma(\vec{x}', t') \quad (9)$$

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To find coefficient, separate tensor, vector and scalar processes [Curie's Theorem, 1894]:

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Separating e.g. the tensor process from others:

$$\langle \pi_{\mu\nu} \rangle_{\text{non-eq}} \approx \underbrace{\langle \pi_{\mu\nu} \rangle_{\text{eq}}}_{\equiv 0} + \int d^3x' \int_{-\infty}^t dt' e^{\epsilon(t'-t)} \underbrace{(\pi_{\mu\nu}(\vec{x}, t), \pi_{\rho\sigma}(\vec{x}', t'))}_{\propto (\text{tensor structures}_{\mu\nu\rho\sigma}) \times (\pi_{\alpha\beta}, \pi^{\alpha\beta})} \beta \partial^\rho u^\sigma \quad (12)$$

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- ☞ interpret coefficient as shear viscosity (by definition):

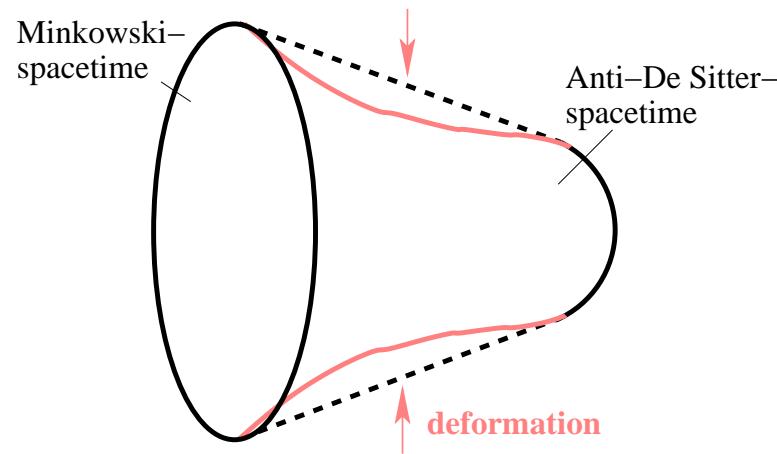
$$\eta \equiv \frac{\beta}{5} \int d^3x' \int_{-\infty}^t dt' e^{\epsilon(t'-t)} \underbrace{(\pi_{\alpha\beta}(\vec{x}, t), \pi^{\alpha\beta}(\vec{x}', t'))}_{\langle \pi_{\alpha\beta}, \pi^{\alpha\beta} \rangle \text{ retarded}}$$

Kubo Formula

5. Applications

Using different SUGRA-theories (Actions and metrics) we can learn much about the strongly coupled dual field theories.

5.1 – Towards QCD



☞ for a R-charged $\mathcal{N} = 4$ SUSY $SU(N)$ Yang-Mills plasma

$$\frac{\eta}{s} = \frac{1}{4\pi}$$

[Son, Starinets, '06]

5.2 – Deriving a coupling

[M.K., Erdmenger, '06]

D7-brane worldvolume action [Erdmenger,Große,Guralnik,'05]

$$S_{D7} = \frac{(2\pi\alpha')^2 T_7}{4} \int d^4x d^4y \operatorname{tr} \left[H(r) F_{\mu\nu} F_{\mu\nu} + 2F_{m\nu} F_{m\nu} + H^{-1}(r) \frac{1}{2} (F_{mn} - {}^* F_{mn})(F_{mn} - {}^* F_{mn}) \right], \quad (13)$$

and the metric

$$ds_{D7}^2 = H^{-1/2}(r) \eta_{\mu\nu} dx^\mu dx^\nu + H^{1/2}(r) d\vec{y}^2 \quad (14)$$

with $r^2 = y^2 + (2\pi\alpha')^2 m^2$, $H(r) = \frac{L^4}{r^4}$

Coordinates									
0	1	2	3	4	5	6	7	8	9
D3									
	D7								
$x^{\mu,\nu,\dots}$		$y^{m,n,\dots}$		$z^{i,j,\dots}$					
			r						
		y							

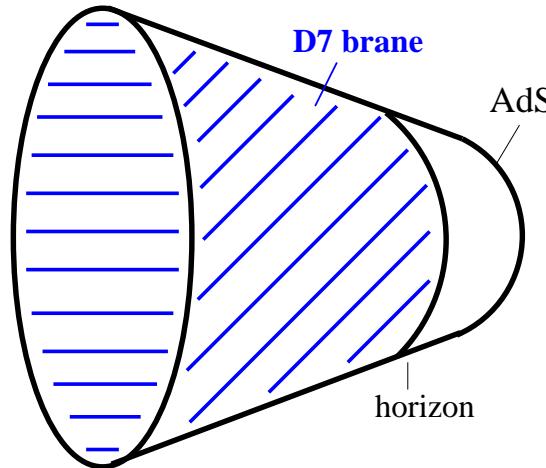
Two-point correlators from this action are proportional to $g_s N$.

Summary and Outlook

☞ back to 'roadmap'

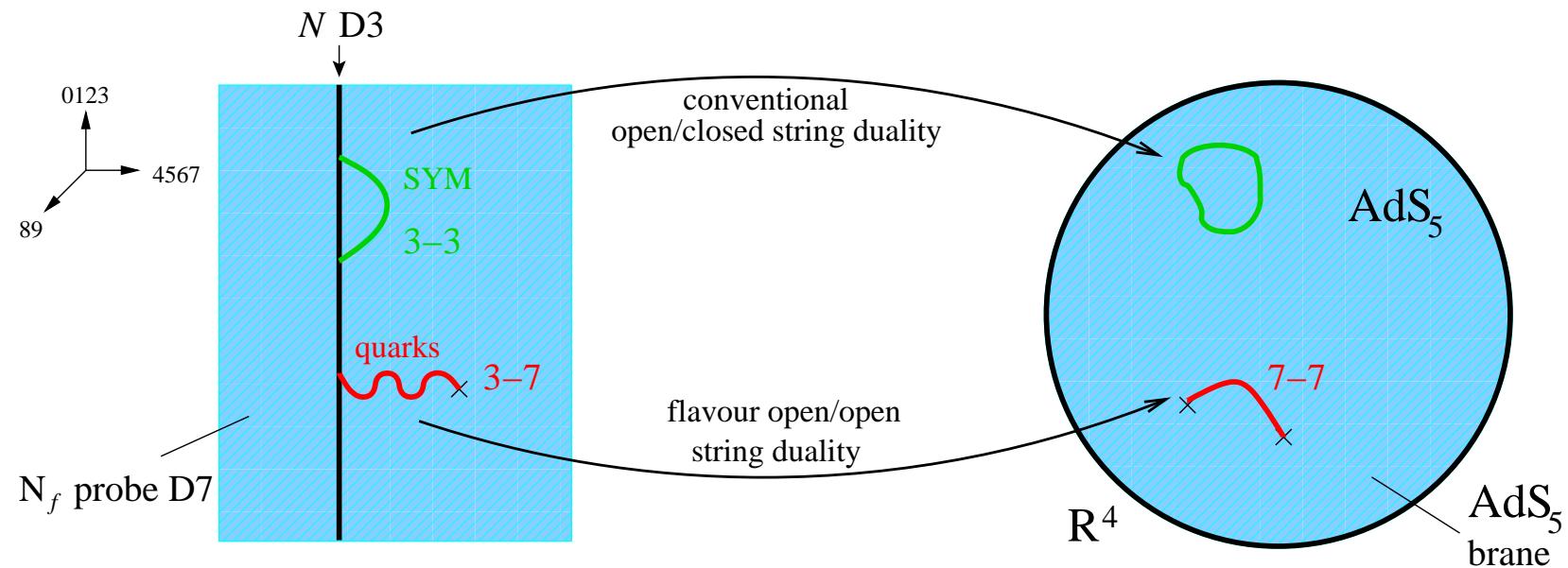
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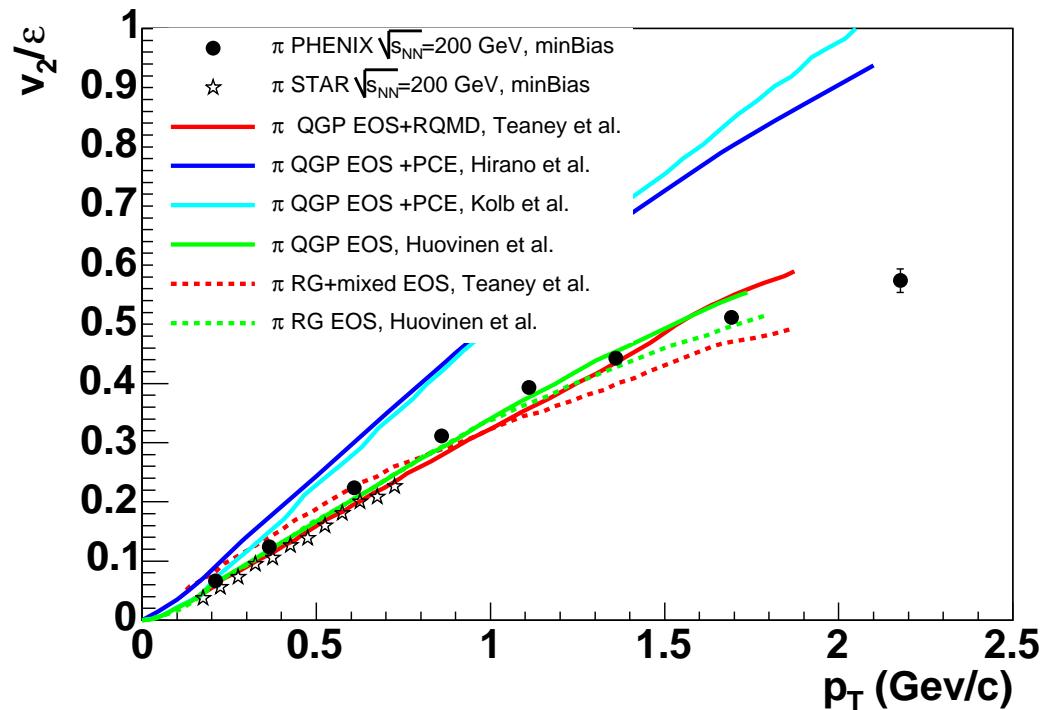
- ☞ incorporate massive fields by embedding D7-branes
- ☞ include finite chemical potential (\equiv R-charge)
- ☞ transport coefficients for this theory

APPENDIX



Comparision: pure hydrodynamics calculations \leftrightarrow RHIC data

[PHENIX coll.]



Eccentricity defined with Cartesian x and y : $\varepsilon = \frac{\langle y^2 \rangle - \langle x^2 \rangle}{\langle y^2 \rangle + \langle x^2 \rangle}$.

Measure for 'elliptic flow':

Second Fourier coefficient v_2 where ϕ is the azimuthal angle: $\frac{d^2 N}{d\phi dp_T} = N_0 (1 + 2v_2(p_T) \cos(2\phi))$.